

Further studies of Isolated photon production in Deep Inelastic Scattering at HERA

Speaker:

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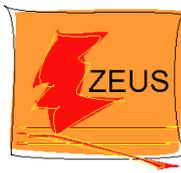
(For the ZEUS collaboration)

April 2017
Birmingham

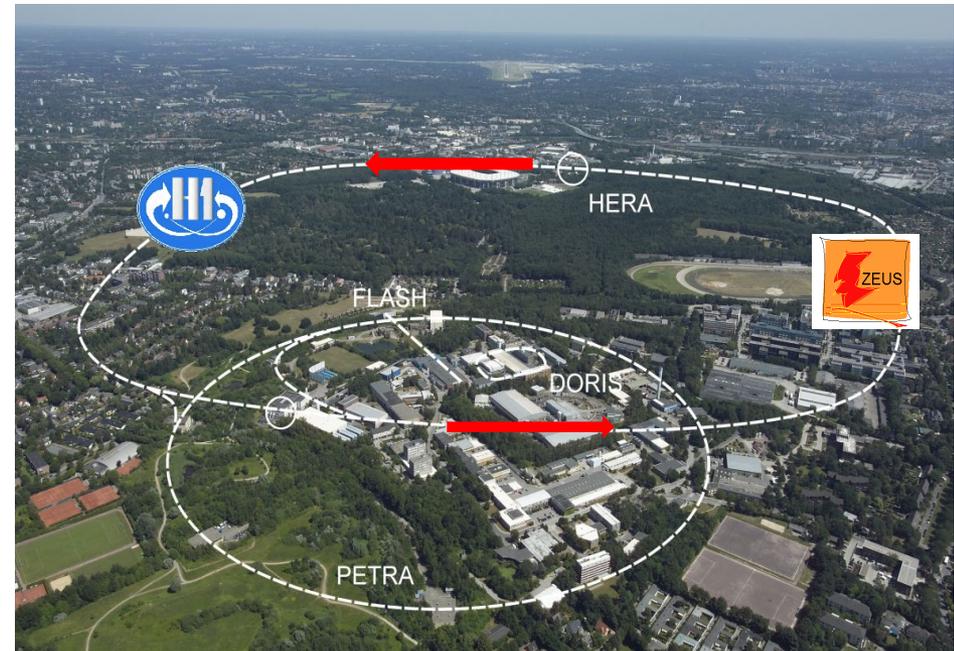


ZEUS-prel-16-001

The ZEUS Detector at HERA



- Protons: 920 GeV
- Electron/Positrons: 27.5 GeV
- Data
 - HERA II period (2004-2007)
 - Integrated luminosity: 326 pb^{-1}



Deep inelastic scattering

Kinematics:

→ $Q^2 = -q^2$ – virtuality
4-momentum transfer

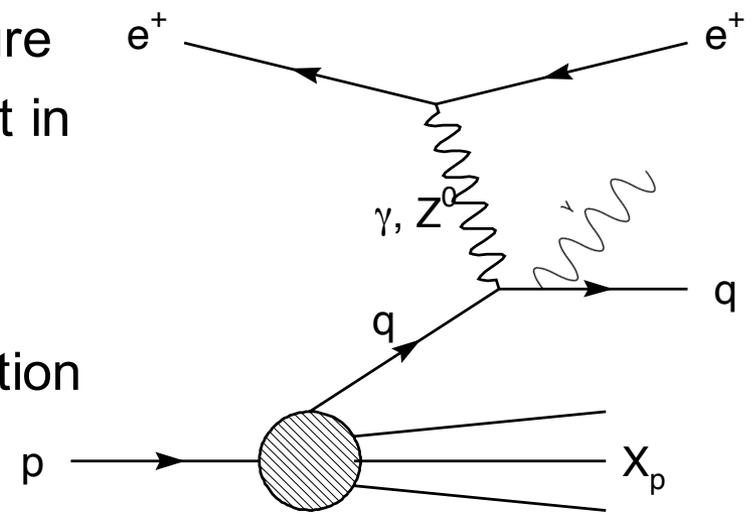
→ $y = \frac{P \cdot q}{P \cdot k}$ – inelasticity measure
fraction of the lepton energy lost in the interaction

→ $x = \frac{Q^2}{2P \cdot q}$ – Bjorken scaling
variable. QPM: momentum fraction carried by the incoming parton

DIS:

→ $Q^2 > 1 \text{ GeV}^2$
→ electron found

Neutral current



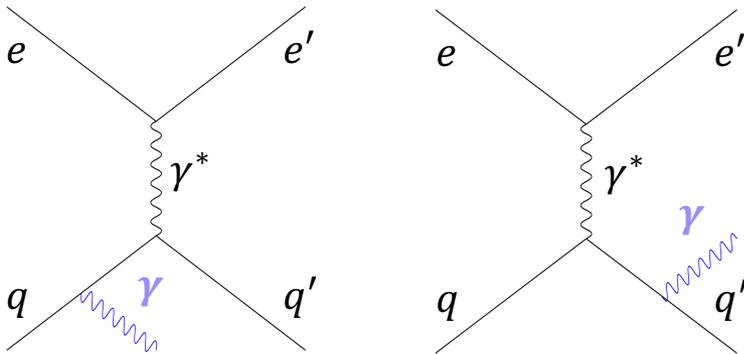
$$Q^2 = sxy$$

$$\sqrt{s} = 318 \text{ GeV}$$

Prompt photons

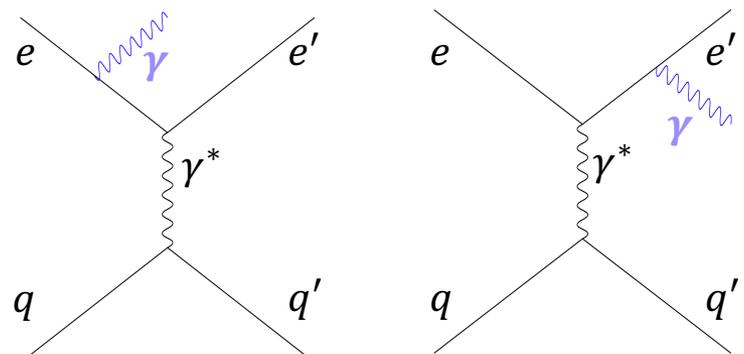
- Photons which are produced promptly in the collision - **before quarks and gluons form hadrons**

QQ - photons

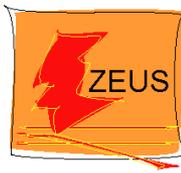


prompt photons are emitted from a quark as part of hard process

LL - photons



photon is radiated from an incoming or outgoing lepton



- A study of the dynamics of prompt photon emission can be used to probe different theoretical models such as the k_T -factorisation model and pQCD approaches
- It is interesting to know whether dynamics changes with virtuality scale
- A study of prompt photons can give a check of the proton's parton distribution functions
- Photons are a possible background to new physics processes: their production should be well understood
- Complements previous ZEUS publication of prompt photons in DIS *Physics Letters B* 715 (2012) 88–97

Event selection

■ Prompt photon selection

- $4 < E_T^\gamma < 15 \text{ GeV}$
- $-0.7 < \eta_\gamma < 0.9$ – in BCAL

Isolation
In CAL

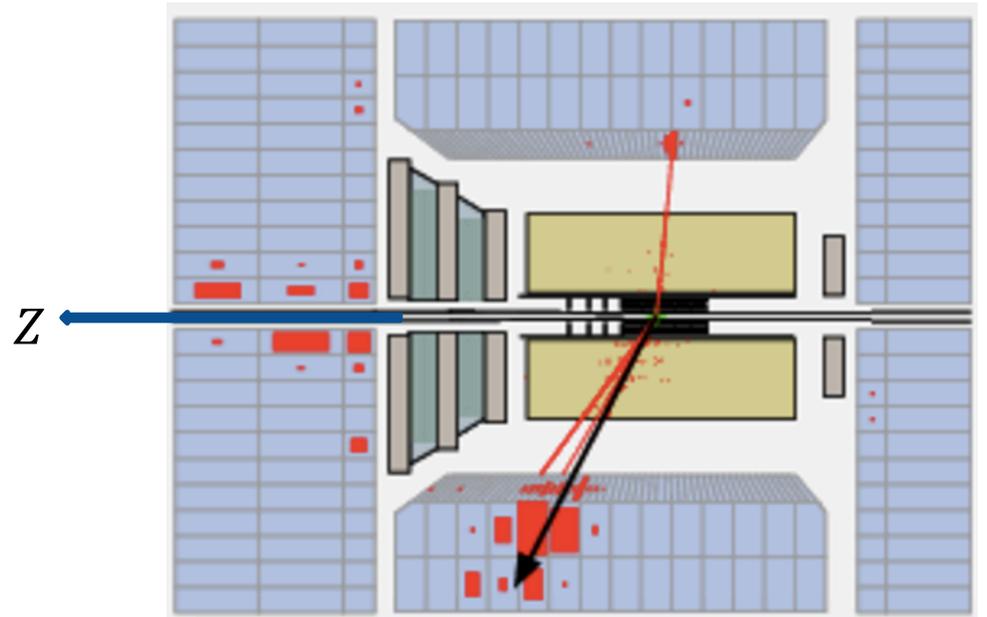
- $E_{EMC} / (E_{EMC} + E_{HAD}) > 0.9$
- $\Delta R(\eta, \varphi) < 0.2$
- $E_\gamma / E_{jet\ with\ \gamma} > 0.9$

■ Jet selection

- $E_T^{jet} > 2.5 \text{ GeV}$
- $-1.5 < \eta_{jet} < 1.8$
- Jet with $E_{T,max}^{jet}$

■ Some Kinematics:

- $10 < Q_{el}^2 < 350 \text{ GeV}^2$
- $E_{e,corr} > 10 \text{ GeV}$
- $140^\circ < \theta_{el} < 180^\circ$
- $35 < E - p_z < 65, \text{ GeV}$



BCAL is finely segmented in the Z direction

$$\bullet x_\gamma = \frac{\sum_{jet,\gamma}(E-p_z)}{2y_{JB}E_e}$$

$$\bullet x_p = \frac{\sum_{jet,\gamma}(E+p_z)}{2E_p}$$

$$\bullet \Delta\eta = \eta_{jet} - \eta_\gamma$$

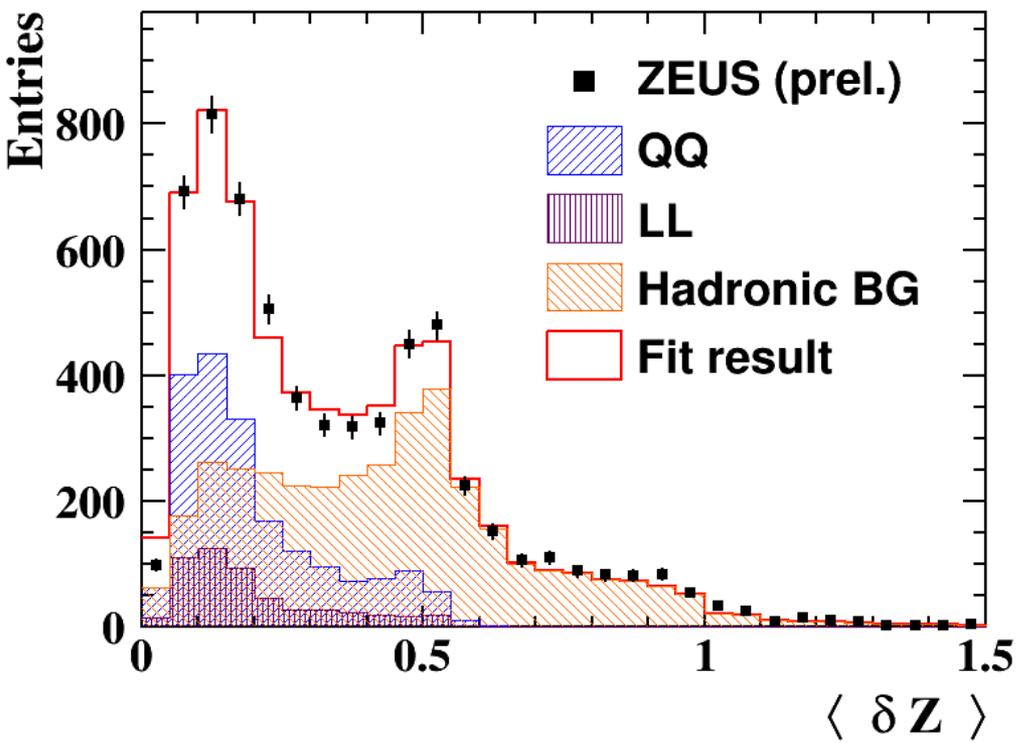
$$\bullet \Delta\varphi = \varphi_{jet} - \varphi_\gamma$$

$$\bullet \Delta\varphi_{e,\gamma} = \varphi_e - \varphi_\gamma$$

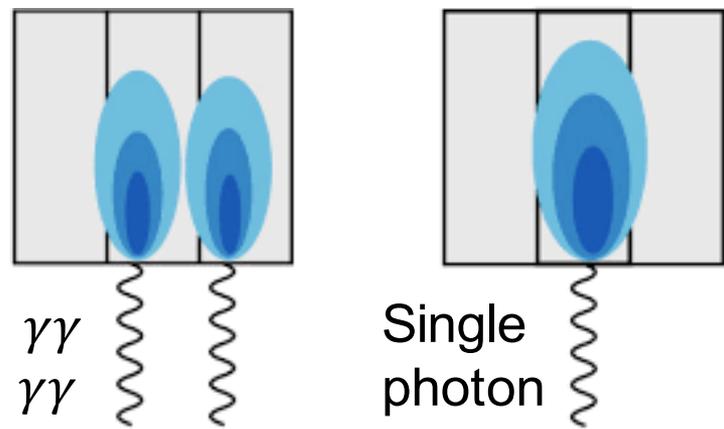
$$\bullet \Delta\eta_{e,\gamma} = \eta_e - \eta_\gamma$$

Signal extraction

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Energy-weighted mean width of the electromagnetic shower (cluster) in calorimeter relative to its centroid



Summary of uncertainties



- Typical statistical uncertainty is **13%**
 - ΔAcc – acceptance uncertainty, **~3-4%** effect
- Typical systematic uncertainty is **10%**
 - Dominated by the energy scale
- Fit of fraction of QQ in data
 - Δa – uncertainty of fit parameter, **~1%** effect
- $\Delta \mathcal{L}$ – 2%, but not included in the following plots

- For a given observable Y , the production cross section:

$$\frac{d\sigma}{dY} = \frac{N(\gamma_{QQ})}{A_{QQ} \cdot \mathcal{L} \cdot \Delta Y} + \frac{d\sigma_{LL}^{MC}}{dY}$$

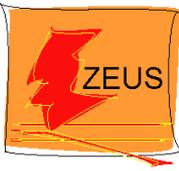
$N(\gamma_{QQ})$ - number of QQ photons extracted from the fit,

ΔY - bin width,

\mathcal{L} - total integrated luminosity,

$\frac{d\sigma_{LL}^{MC}}{dY}$ - cross section for LL photons

A_{QQ} - ratio of the number of events reconstructed to those generated in a given bin



Cross Sections compared to LO + Leading Log QQ (Pythia) and LL (Ariadne) MC

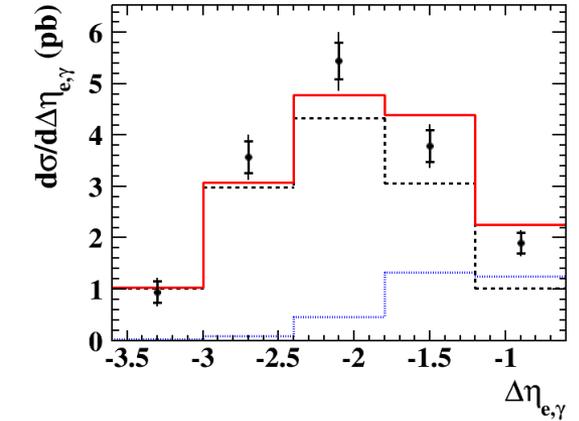
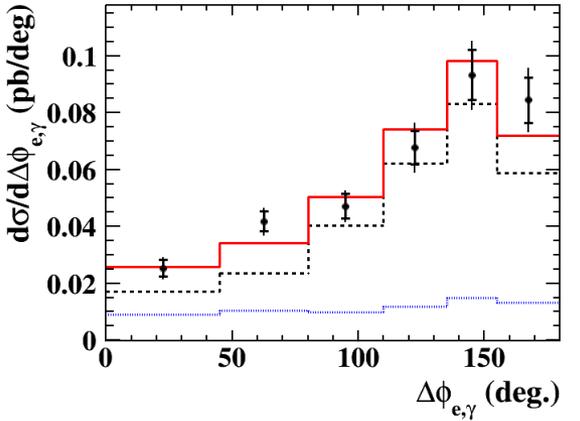
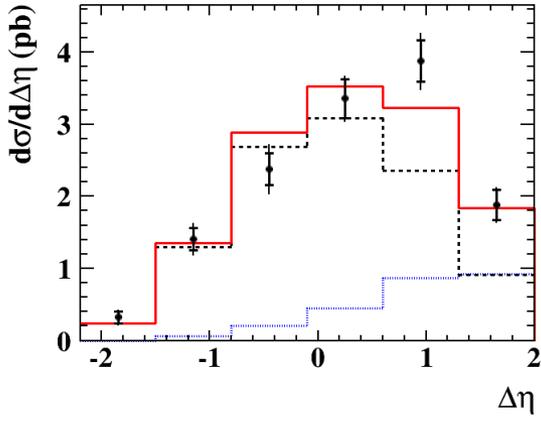
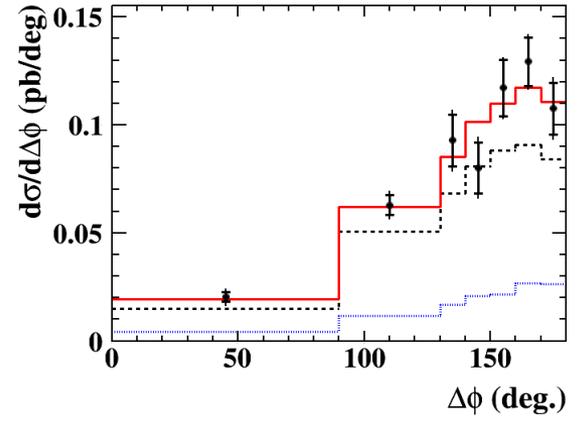
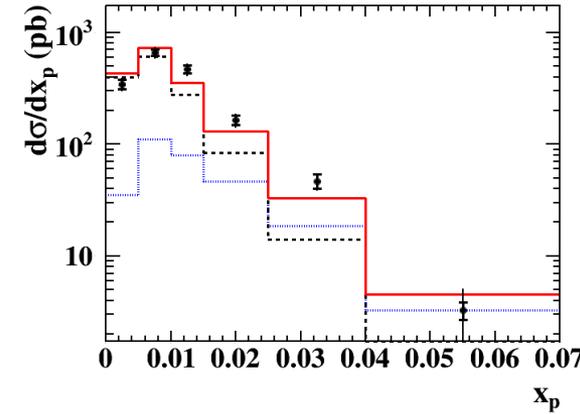
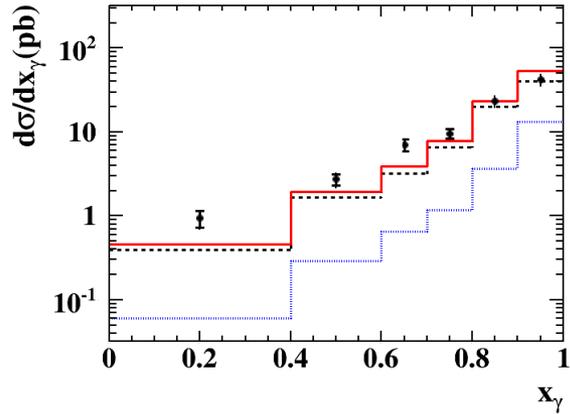
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16-001

● ZEUS (prel.)

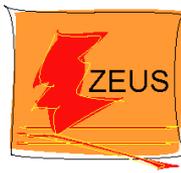
⋯ LL, MC

⋯ QQ*1.6, MC

▭ LL + QQ*1.6, MC



Two Q^2 kinematical regions on full statistics



- Cross sections for two separate kinematic regions: $10 < Q^2 < 30 \text{ GeV}^2$ and $30 < Q^2 < 350 \text{ GeV}^2$ were also calculated
- $Q^2 = 30 \text{ GeV}^2$ was chosen to divide available data sample into 2 with similar numbers of events

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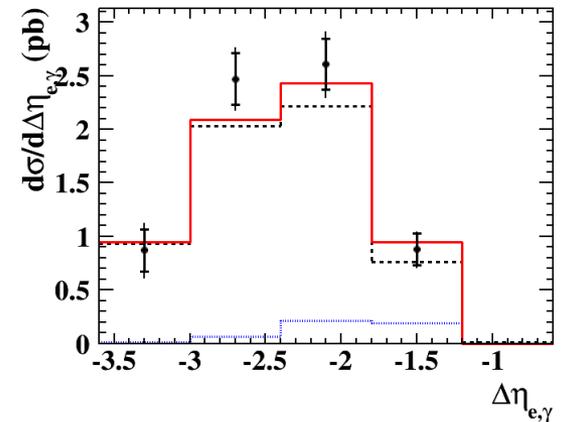
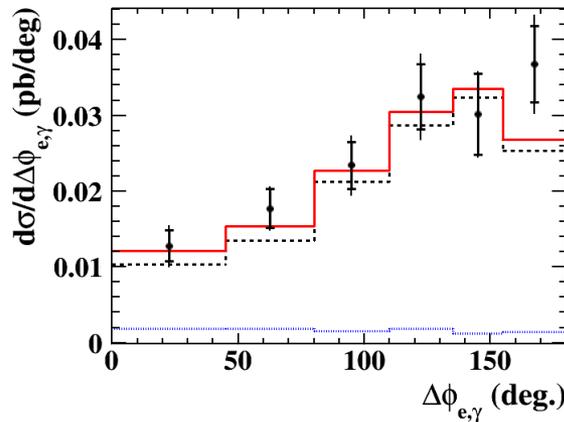
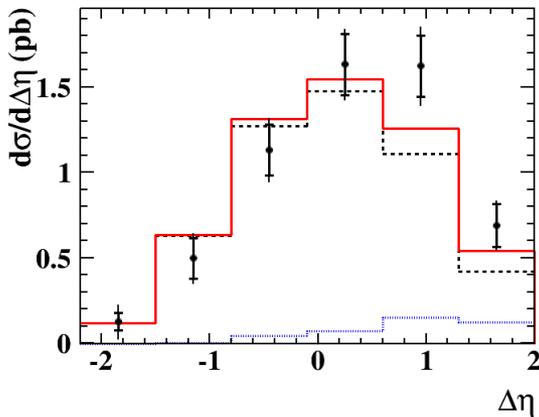
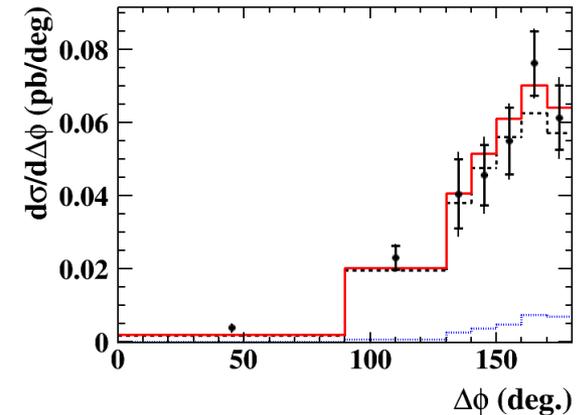
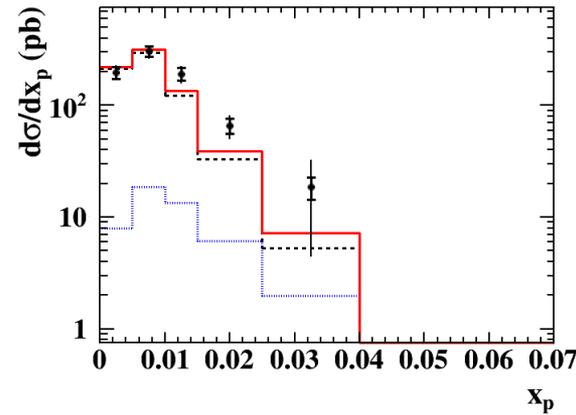
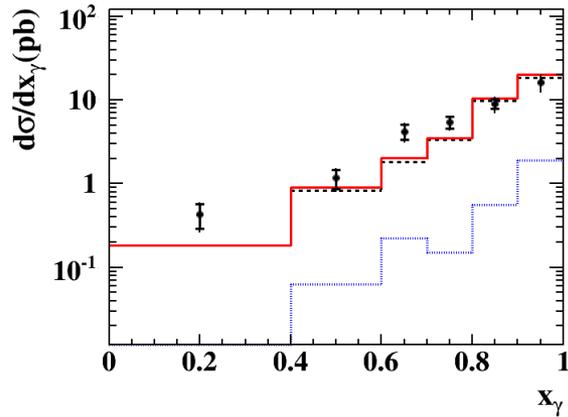
$10 < Q^2 < 30 \text{ GeV}^2$

● ZEUS (prel.)

○ QQ*1.6, MC

□ LL, MC

□ LL + QQ*1.6, MC



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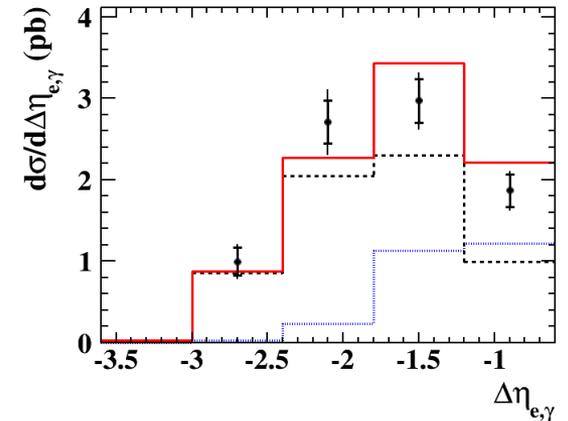
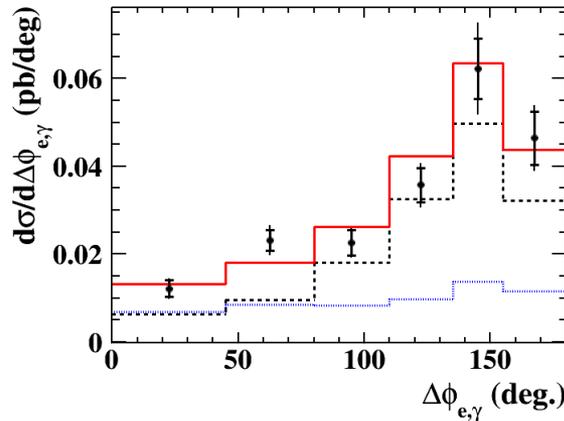
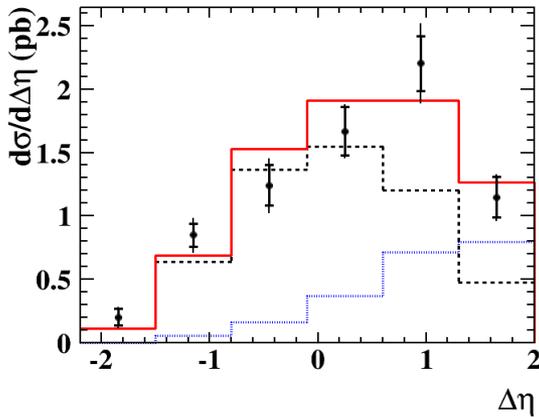
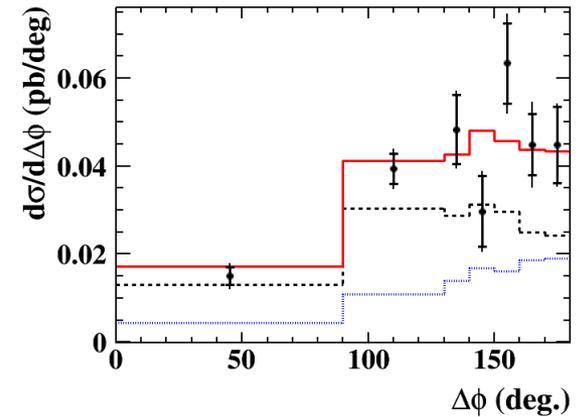
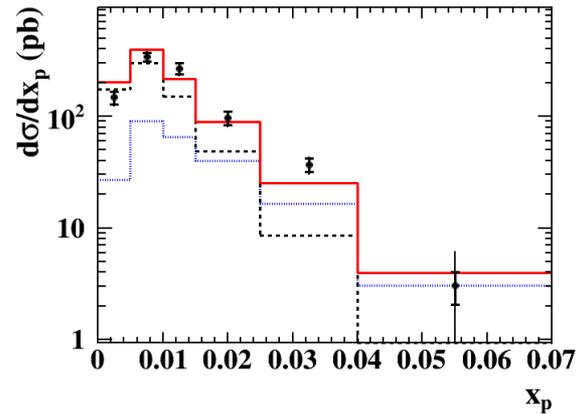
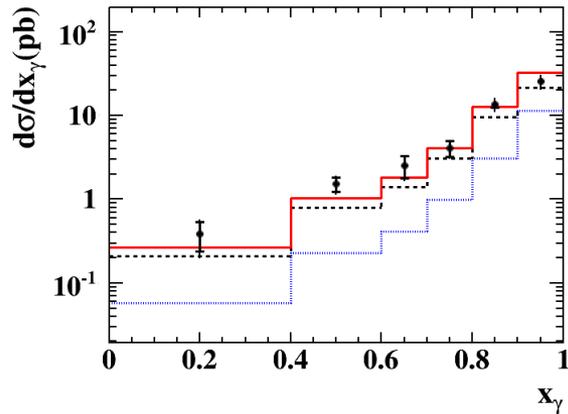
$30 < Q^2 < 350 \text{ GeV}^2$

● ZEUS (prel.)

○ QQ*1.6, MC

□ LL, MC

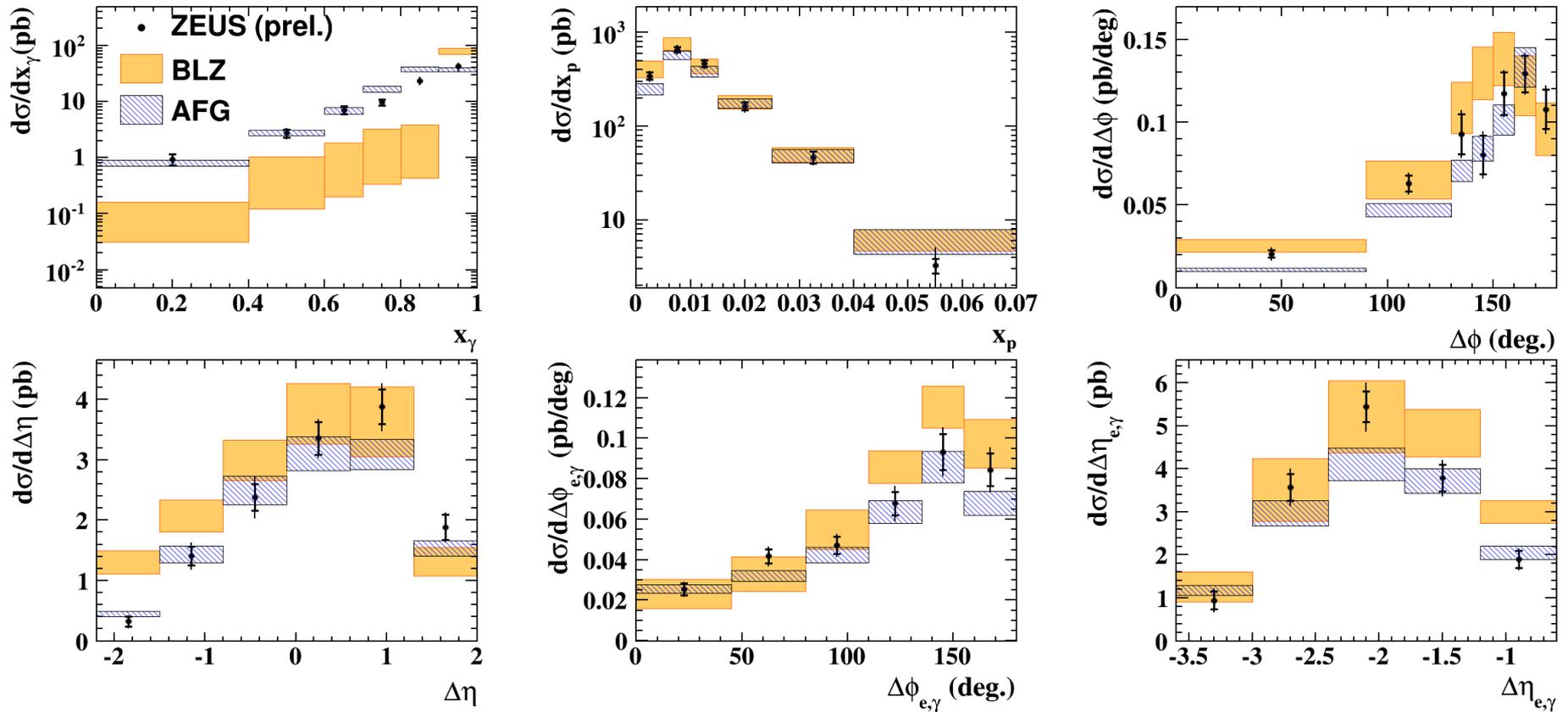
□ LL + QQ*1.6, MC





Comparison with theory

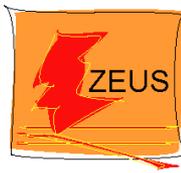
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Aurenche, Fontannaz and Guillet : LAPTH-005/17 LPT-Orsay 16-88

Baranov, Lipatov and Zotov: PRD81, 094034 (2010)

Conclusion



- Experimental differential cross sections have been obtained for x_γ , x_p , $\Delta\eta$, $\Delta\varphi$, $\Delta\eta_{e,\gamma}$, $\Delta\varphi_{e,\gamma}$ correlated observables
- Pythia describes the shape of the data reasonably well in both Q^2 ranges separately when rescaled by a factor 1.6, as in the previous ZEUS DIS publication
- AFG (NLO) calculations show an overall good agreement to the data
- k_t -factorisation (BLZ) predictions show a fair agreement with the data with the exception of x_γ and $\Delta\eta$
 - However x_γ and $\Delta\eta$ - variables sensitive to the gluon radiation - are reasonably well described by Pythia



Appendix

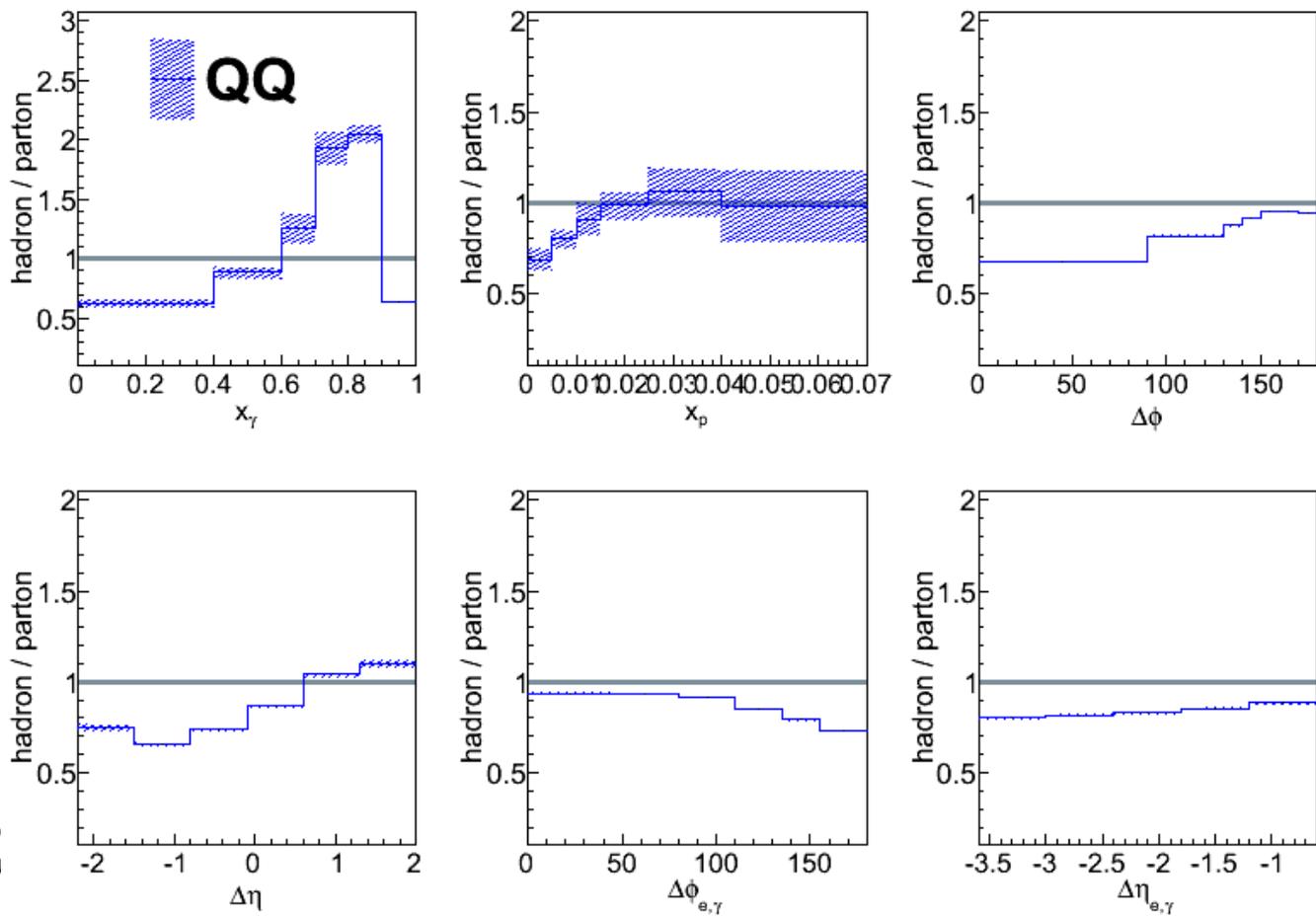
Systematics uncertainties



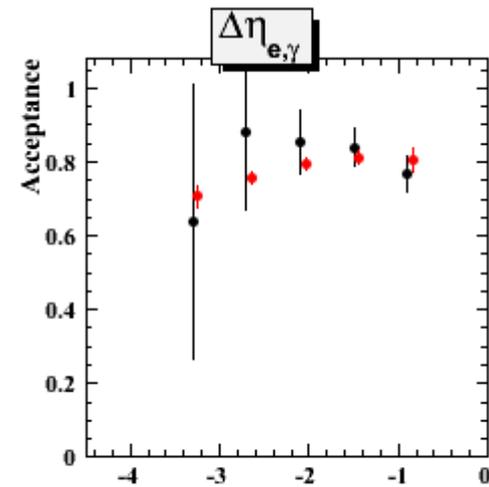
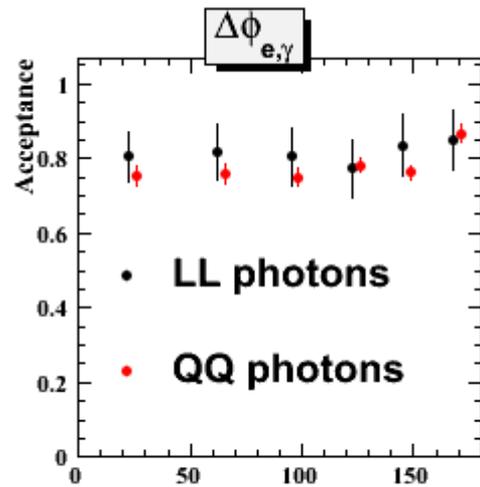
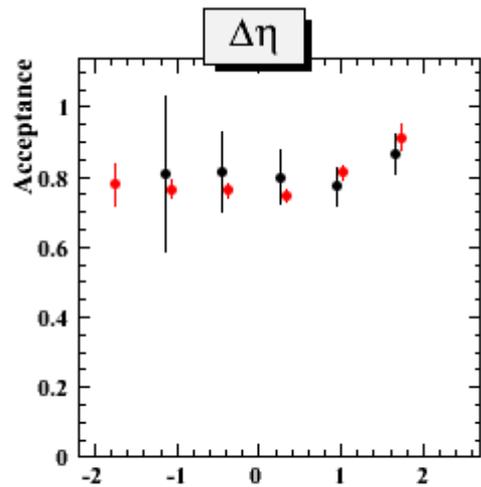
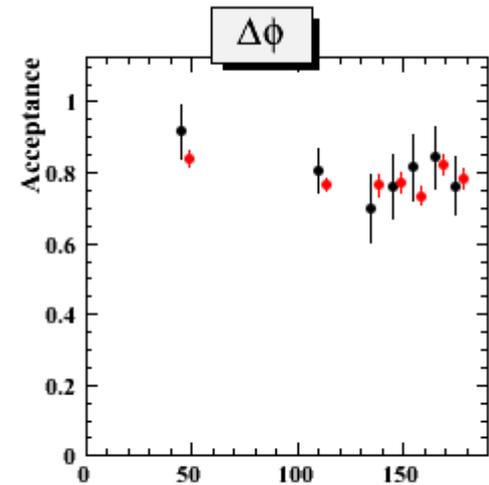
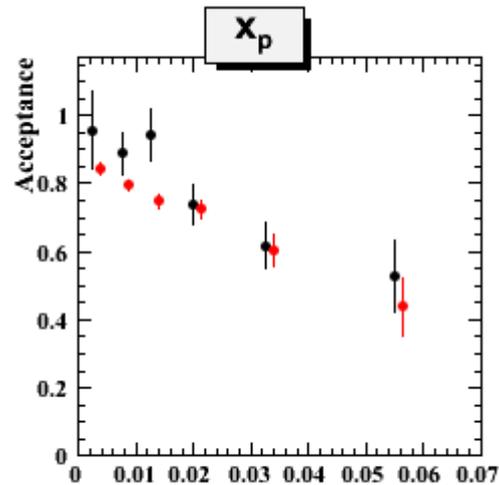
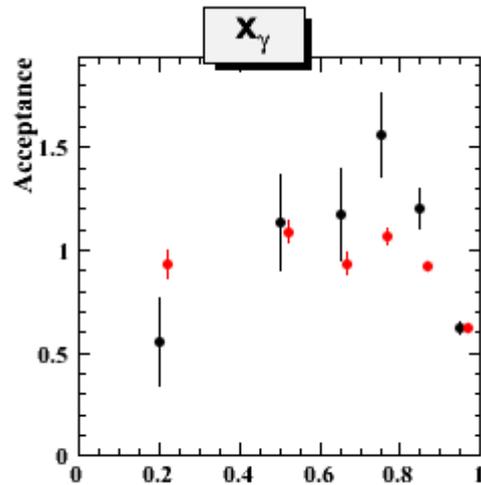
- Before applying cuts for the event selection the values of such variables as:
 - E_γ - energy of the photon varied by $\pm 2\%$ ($\sim 6\%$ effect)
 - E_{jet} - energy of the jet varied by
 - $\pm 1.5\%$ for $E_{jet} > 10$ GeV,
 - $\pm 2.5\%$ for $6 < E_{jet} < 10$ GeV,
 - $\pm 4\%$ for $E_{jet} < 6$ GeV(7% effect)
 - E_e - energy of the electron varied by $\pm 2\%$ ($\sim 1\%$ effect)
- The dependence on the modelling of hadronic background by Ariadne was investigated by changing the upper fit limit in range [0.6, 1.0] ($\sim 5\%$ effect)
- Uncertainty on the acceptance due to Pythia model is taking as half of the change attributable to the reweighting ($\sim 1\%$ effect)

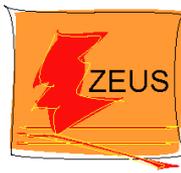
Hadronization corrections

- $\kappa = \frac{\text{had level}}{\text{part level}}$
 $\delta\kappa = \sqrt{\frac{C^2(D+E) + D^2(C+E) + E^2(C+D) + 2CDE}{(C+E)^4}}$
- Here: $C = \text{had} \ \& \ \text{part}$, $D = \text{had} \ \& \ \text{not part}$, $E = \text{not had} \ \& \ \text{part}$

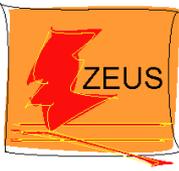


Acceptance

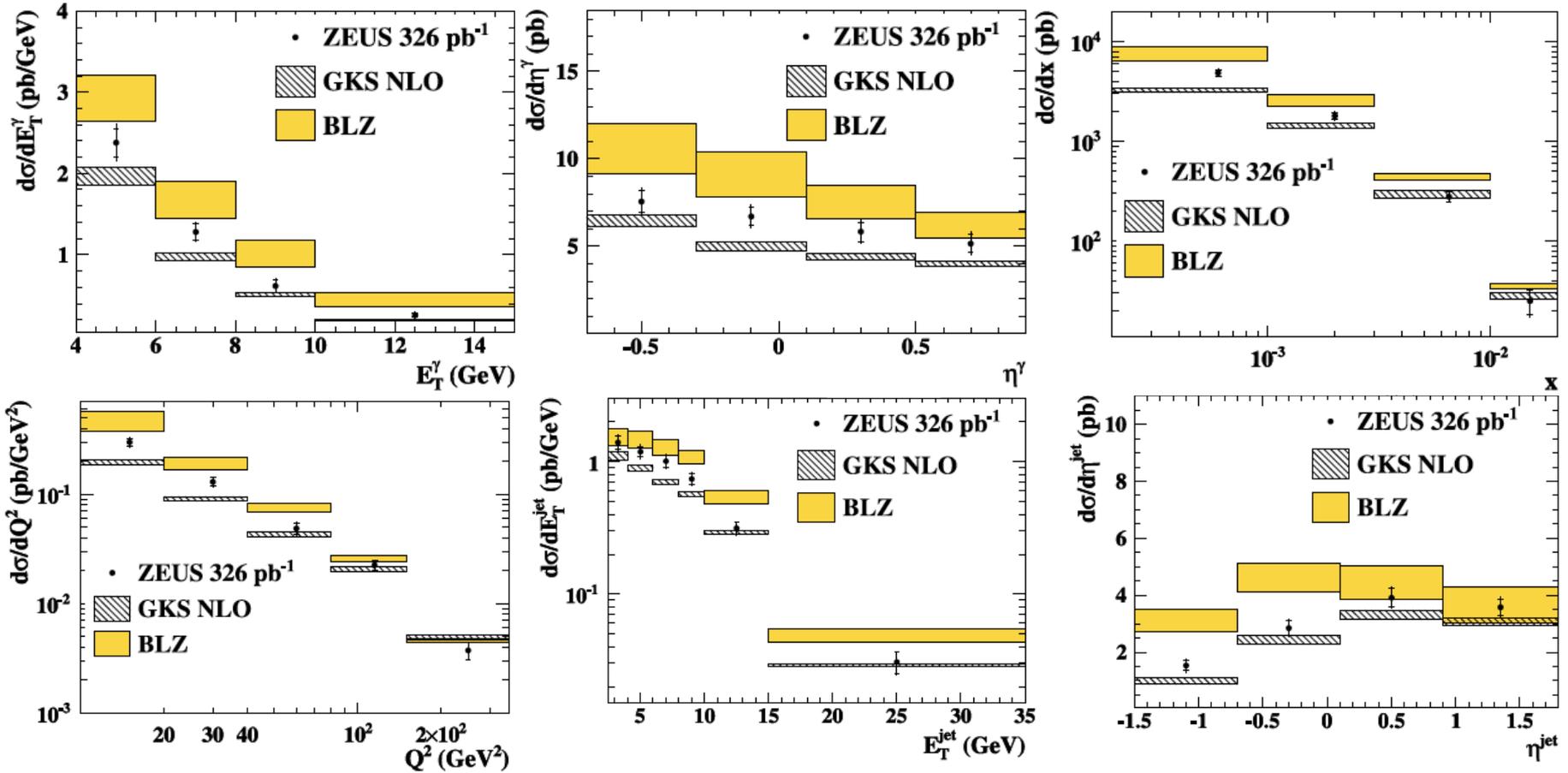




- Baranov, Lipatov and Zotov (BLZ) : Rev. D 81 (2010) 094034
 - approach is based on the calculation of the CS using the convolution of the off-shell matrix element and unintegrated quark-distribution (therefore the kt-factorisation).
 - In the kt-factorization theory some part of final state jets can originate not only from hard subprocess, but also from the parton evolution cascade in initial state.
 - To determine the 4-momenta of these jets (in particular, their rapidities) model approximation was used
 - $\Lambda_{\text{QCD}} = 200 \text{ MeV}$, number of flavours $N_f = 4$, scale $\mu_R^2 = \mu_F^2 = Q^2$; MSTW'2008 pdfs as an input for the KMR partons
- Aurenche, Fontannaz and Guillet (AFG) : Rev. D 81 (2010) 094034 – **NLO theory with conventional PDFs**
 - Predictions are available as CS with pt cut between p_{T_cut} in center-of-mass frame 2.5 GeV/c and .5 GeV/c



Previous study



The previous ZEUS DIS publication measured x , Q^2 , E_T^γ , η_γ , E_T^{jet} and η_{jet} .
(Physics Letters B 715 (2012) 88–97)

Study of photon-jet and photon-electron variables

$$\begin{aligned}
 \bullet x_\gamma &= \frac{\sum_{jet,\gamma}(E-p_z)}{2y_{JB}E_e} & \bullet \Delta\eta &= \eta_{jet} - \eta_\gamma \\
 \bullet x_p &= \frac{\sum_{jet,\gamma}(E+p_z)}{2E_p} & \bullet \Delta\varphi &= \varphi_{jet} - \varphi_\gamma \\
 & & \bullet \Delta\varphi_{e,\gamma} &= \varphi_e - \varphi_\gamma \\
 & & \bullet \Delta\eta_{e,\gamma} &= \eta_e - \eta_\gamma
 \end{aligned}$$

A similar kind of analysis was previously done for photoproduction ($Q^2 < 1 \text{ GeV}^2$)

