

Running of the Charm-Quark Mass from

HERA Deep-Inelastic Scattering Data



α_s

m_b

m_c

Achim Geiser
DESY Hamburg

work partially done
within scope of

PROSA, ZEUS and H1
collaborations



DIS 2017, Birmingham, UK, 5 April 2017

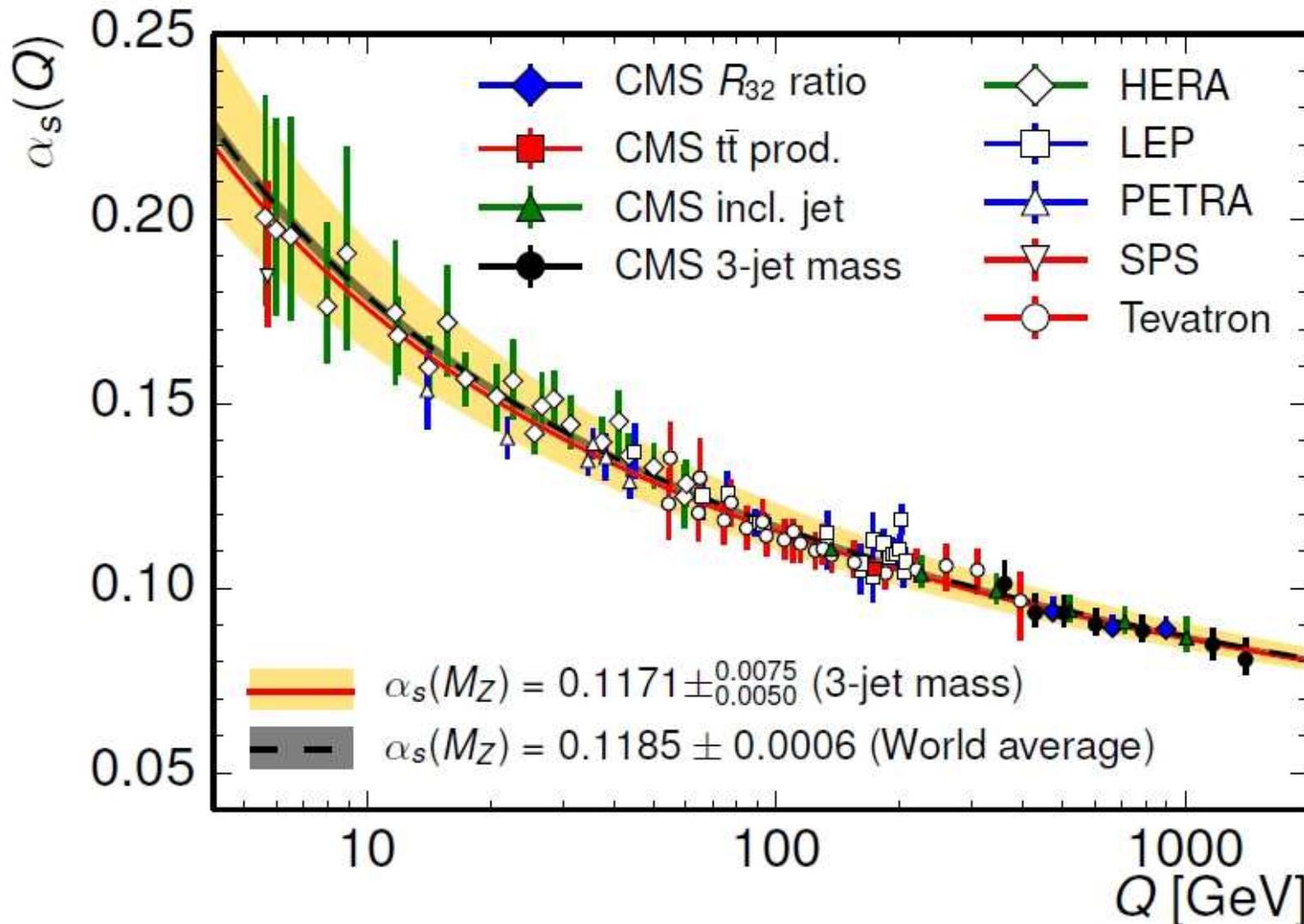
- Introduction: running of α_s and m_b
- Final results on charm mass running: **DESY-17-048** (on arXiv soon)
- Interpretation in terms of **Higgs Yukawa couplings**: PoS CHARM2016 (2017) 012

Running of strong coupling „constant“ α_s

reminder

EPJC 75 (2015) 186

e.g. from jet production at $e+e^-$, ep , and pp at DESY, Fermilab and CERN



updates see talks
K. Rabbertz
and D. Britzger

**Yes,
it runs!**

Running of α_s and quark masses m_Q

- α_s running depends on number of colours N_c and number of quark flavours n_f

$$\alpha_s(\mu) = \frac{\alpha_s(\mu_0)}{1 + \alpha_s \times (11N_c - 2n_f)/12\pi \ln(\mu^2/\mu_0^2)}$$

- quark mass running depends on α_s , e.g.

$$\begin{aligned} m_Q(\text{pole}) &= m_Q(m_Q) (1 + 4/3 \alpha_s/\pi) \\ &= m_Q(\mu) (1 + \alpha_s/\pi (4/3 + \ln(\mu^2/m_Q^2))) \end{aligned}$$

leading
order
QCD
formulae

or

$$m_Q(\mu) = m_Q(m_Q) \times \left(\frac{\alpha_s(\mu)}{\alpha_s(m_Q)} \right)^{c_0} \quad c_0 = 4/(11 - 2n_f/3) = 4/9 \quad \begin{matrix} N_c \\ n_f = 3 \end{matrix}$$

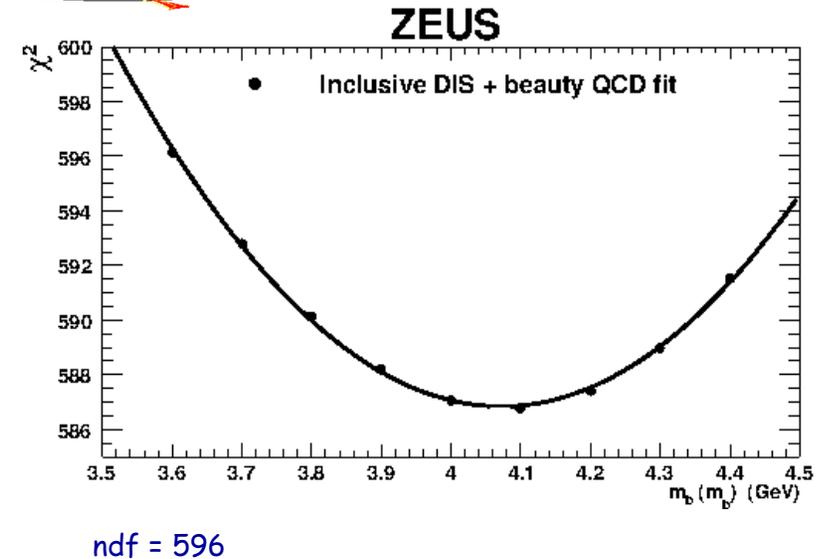
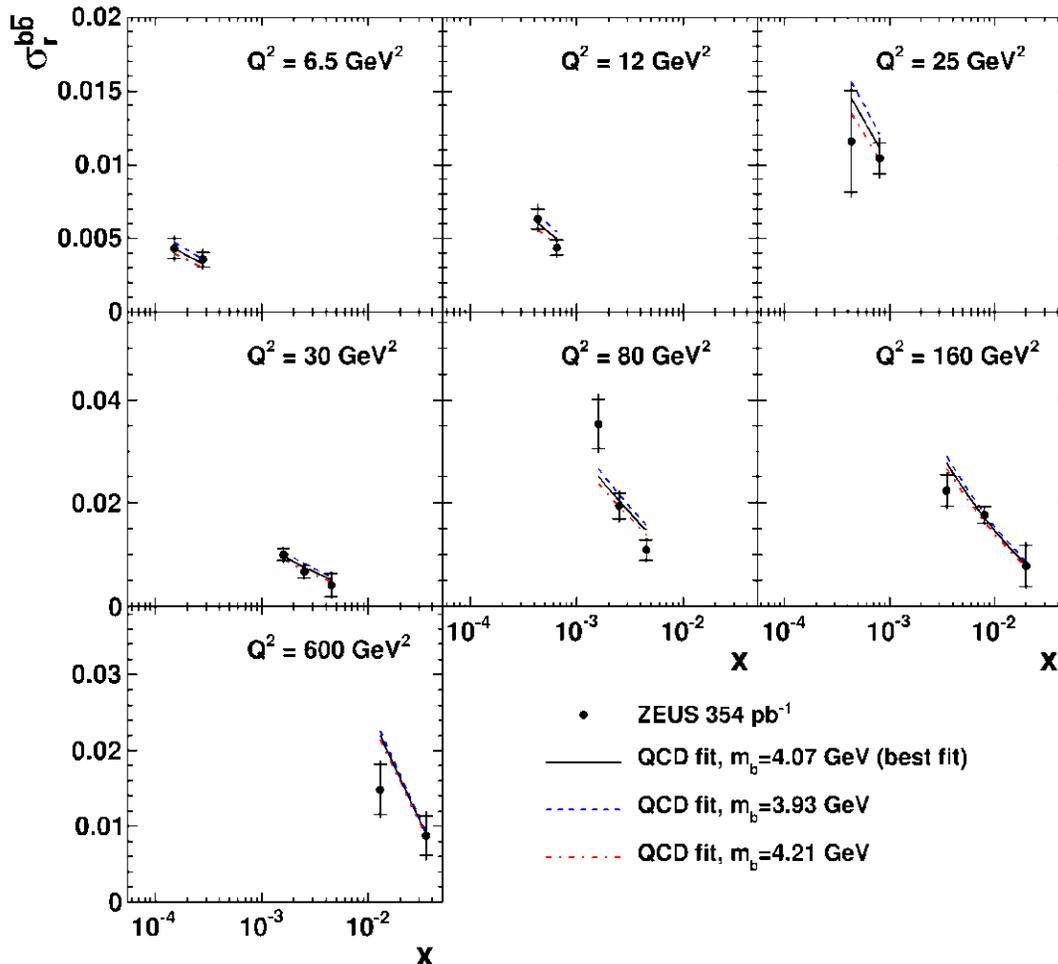
- part of gluon field around quark not 'visible' any more when 'looking' at smaller distances/larger energy scales
-> **effective quark mass decreases**

m_b from reduced beauty cross section in DIS

reminder

JHEP 1409 (2014) 127

ZEUS



update see talk
O. Zenaiev

$$m_b(m_b) = 4.07 \pm 0.14_{\text{fit}} \quad +0.01 \quad -0.07_{\text{mod}} \quad +0.05 \quad -0.00_{\text{par}} \quad +0.08 \quad -0.05_{\text{th}} \quad \text{GeV}$$

PDG: 4.18 ± 0.03 GeV (lattice QCD + time-like processes)

The running beauty quark mass



ZEUS, JHEP 1409 (2014) 7;

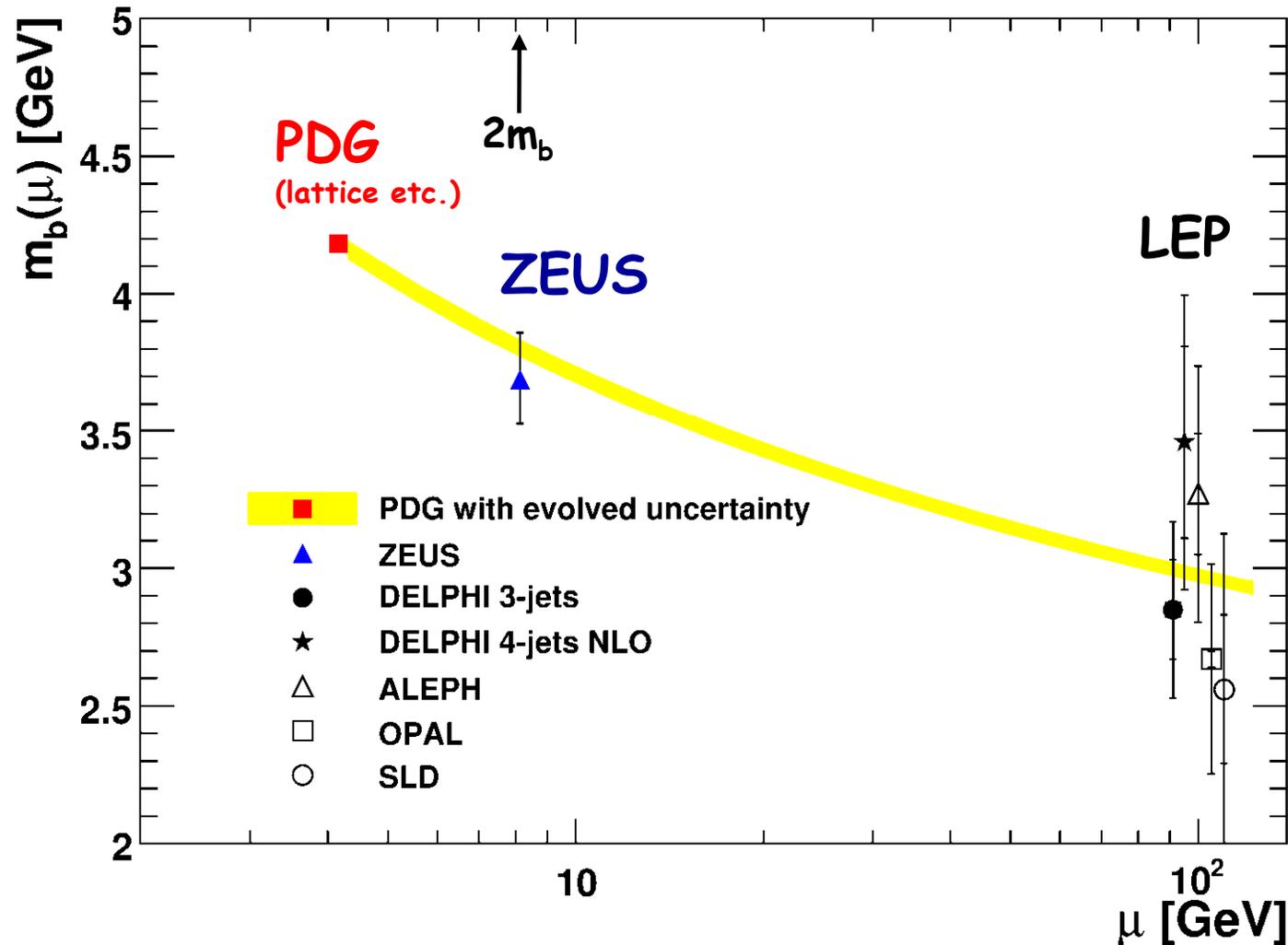
review, arXiv:1506.07519

LEP, Eur. Phys. J. C55 (2008) 525

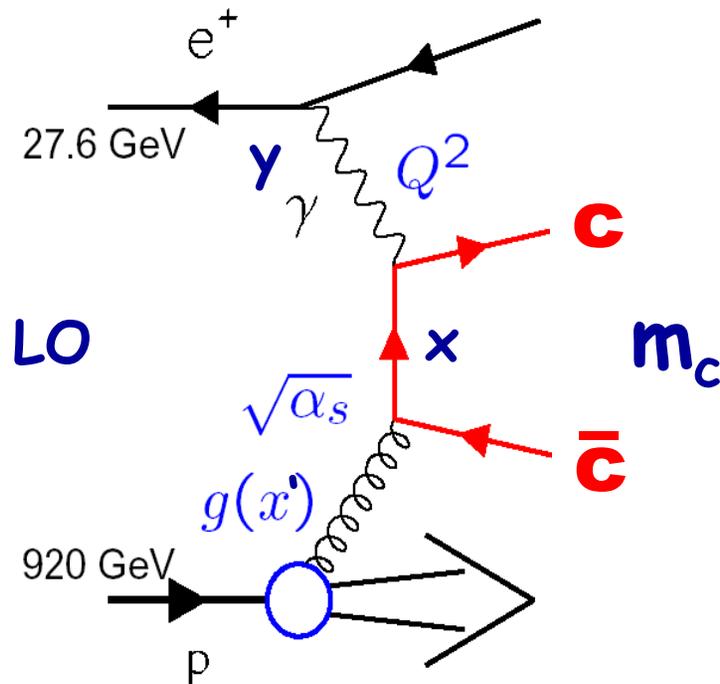
Prog. Part. Nucl. Phys. 84 (2015) 1

translate to $2m_b$

ZEUS



Fixed Flavour Number Scheme (FFNS)



+ NLO (+partial NNLO) corrections,

“natural” scale:
 $\mu^2 = Q^2 + 4m_c^2$

- no charm in proton

- full kinematical treatment of charm mass ☺

(multi-scale problem:
 $Q^2, p_T, m_c \rightarrow$ logs of ratios)

- no resummation of logs ☹

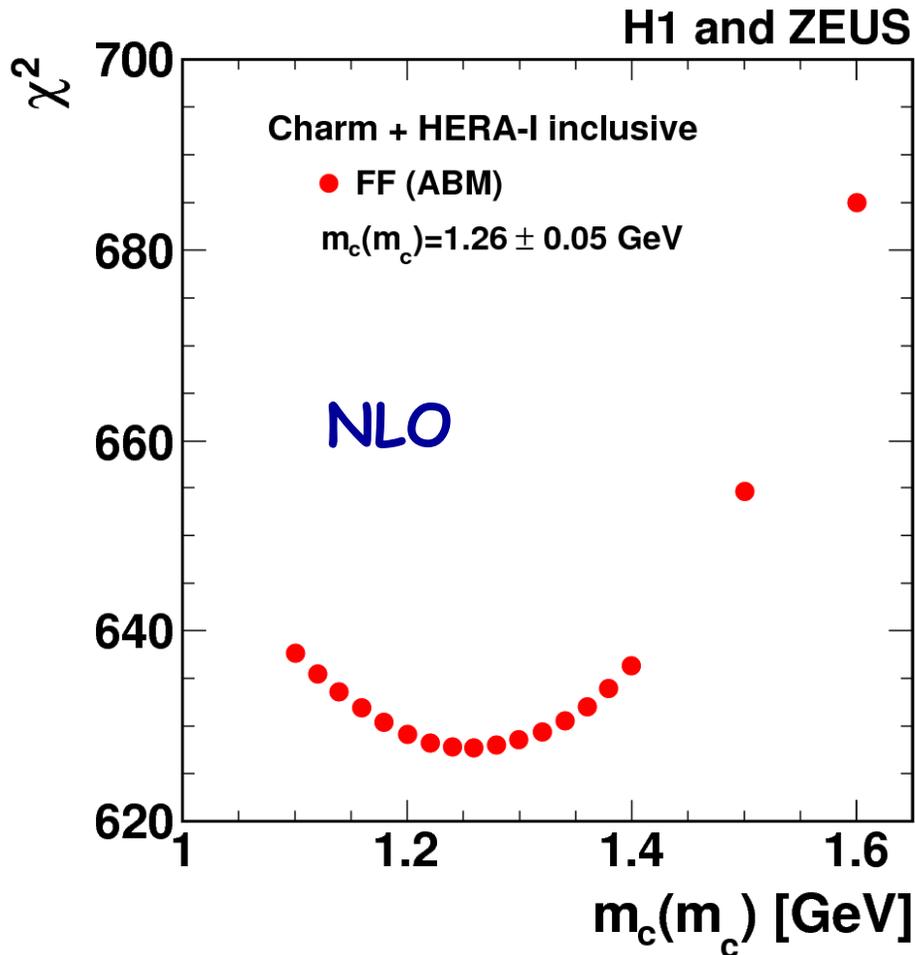
- no extra matching parameters ☺



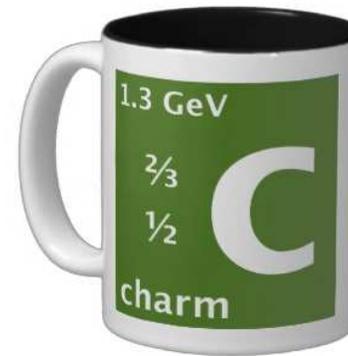
Measurement of \overline{MS} charm mass

reminder

EPJ C73 (2013) 2311



simultaneous mass + PDF fit of combined charm data and inclusive HERA I DIS data



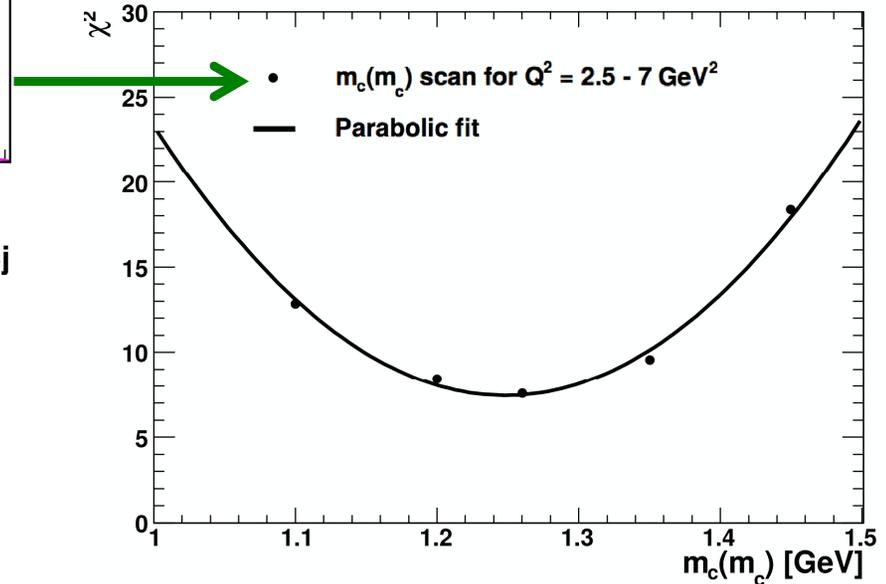
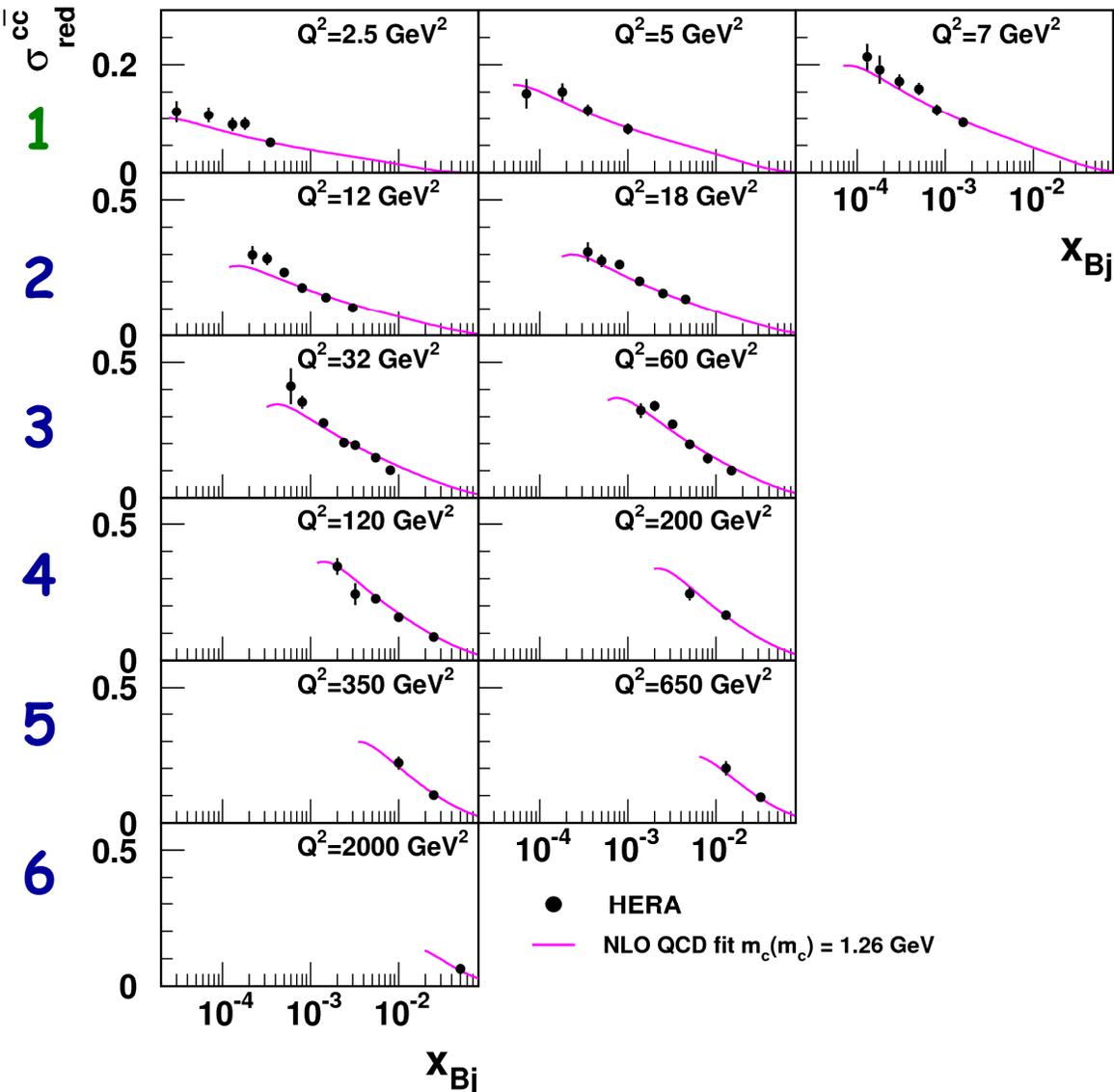
update see talk O. Zenaiev

$$m_c(m_c) = 1.26 \pm 0.05_{\text{exp}} \pm 0.03_{\text{mod}} \pm 0.02_{\alpha_s} \text{ GeV}$$

PDG: $1.275 \pm 0.025 \text{ GeV}$ (lattice QCD + time-like processes)

Measurement of m_c running

A. Gizhko et al., DESY-17-048



Step 1:
 extract $m_c(m_c)$ separately
 for 6 different kinematic
 ranges in $\mu^2 = Q^2 + 4m_c^2$

(take log average for central scale)

m_c fit and uncertainties

A. Gizhko et al., DESY-17-048

use appropriate PDF set for each mass
(from inclusive DIS data only),
fit charm data

Fit uncertainty

- Was estimated by taking $\Delta\chi^2 = 1$ (dominant uncertainty)

Parametrisation

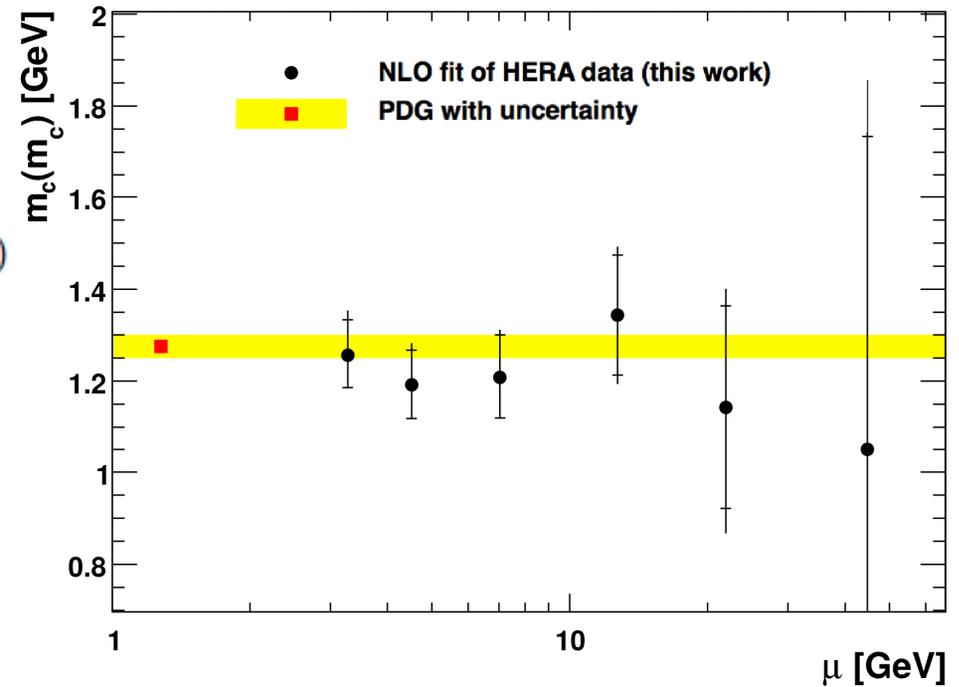
- Adding extra parameter in the PDF parametrisation

Model uncertainty

- Variation of the strangeness suppression factor
- Lower cut on Q^2 for inclusive data
- The evolution starting scale
- The b-quark mass

Theory

- Variation of α_s
- Variation of the factorisation and renormalization scales of heavy quarks by factor 2 → outer error bar



sensitivity to $m_c(m_c)$ decreases with increasing scale $\mu^2 = Q^2 + 4m_c^2$

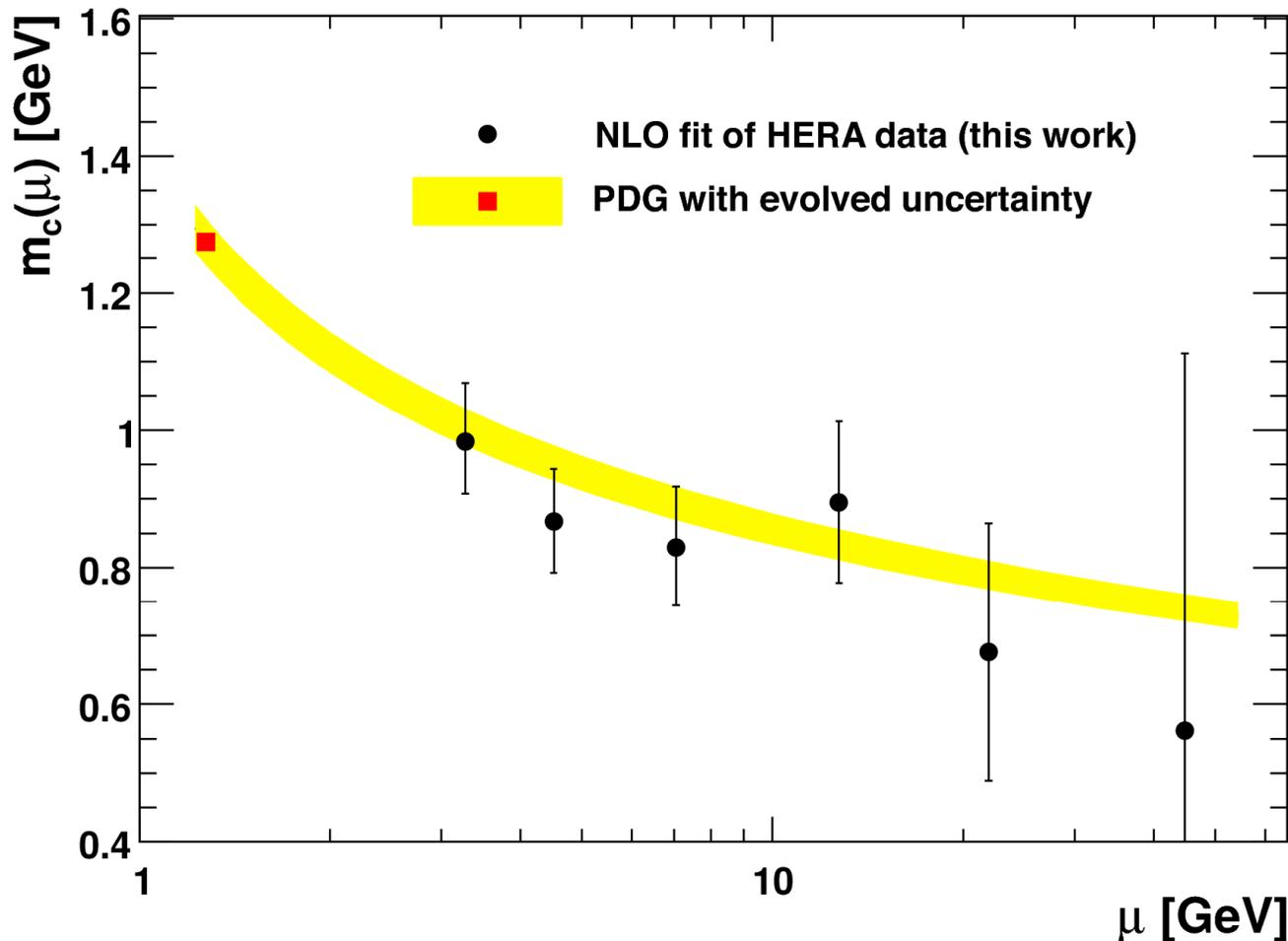
'in reality', have measured $m_c(\mu)$ at each scale

The running charm quark mass

A. Gizhko et al., DESY-17-048

Step 2: translate back to $m_c(\mu)$, which was actually measured, using LO formula consistent with NLO \overline{MS} QCD fit

(OpenQCDrad, Alekhin et al.)



running mass
concept in QCD
is self-consistent !

Numerical details

Subset	N_{dat}	Q^2 range [GeV ²]	μ [GeV]	$m_c(m_c)$		$m_c(\mu)$	
				[GeV] fit	scale	[GeV] fit	
1	15	2.5–7	3.3	1.256 ^{+0.078} _{-0.070}	^{+0.054} _{-0.000}	0.984 ± 0.061	
2	12	12–18	4.5	1.192 ^{+0.075} _{-0.073}	^{+0.043} _{-0.000}	0.867 ± 0.055	
3	13	32–60	7.0	1.208 ^{+0.092} _{-0.088}	^{+0.045} _{-0.000}	0.830 ± 0.063	
4	7	120–200	12.7	1.344 ^{+0.130} _{-0.131}	^{+0.073} _{-0.074}	0.895 ± 0.087	
5	4	350–650	21.9	1.143 ^{+0.222} _{-0.221}	^{+0.133} _{-0.163}	0.676 ± 0.132	
6	1	2000	44.8	1.050 ^{+0.684} _{-0.760}	^{+0.400} _{-0.149}	0.562 ± 0.412	

Table 1: Values of $m_c(m_c)$ at different scales μ , determined from six different subsets, and corresponding values of $m_c(\mu)$. The first uncertainty (fit) corresponds to the uncertainty $\delta_{\text{fit}}^{\text{exp}}$ added in quadrature with the symmetrised systematic uncertainties $\delta_1 - \delta_6$. The second uncertainty (scale) of $m_c(m_c)$ corresponds to the scale variation uncertainty δ_7 . No scale uncertainty is quoted for $m_c(\mu)$ (see text). The range of Q^2 values contributing to the six data subsets shown in Fig. 1 is given. Also given is the corresponding logarithmic average scale μ for each subset according to Eq. (2), and the number N_{dat} of charm data points contributing to each measurement.

Breakdown of uncertainties on m_c

A. Gizhko et al., DESY-17-048

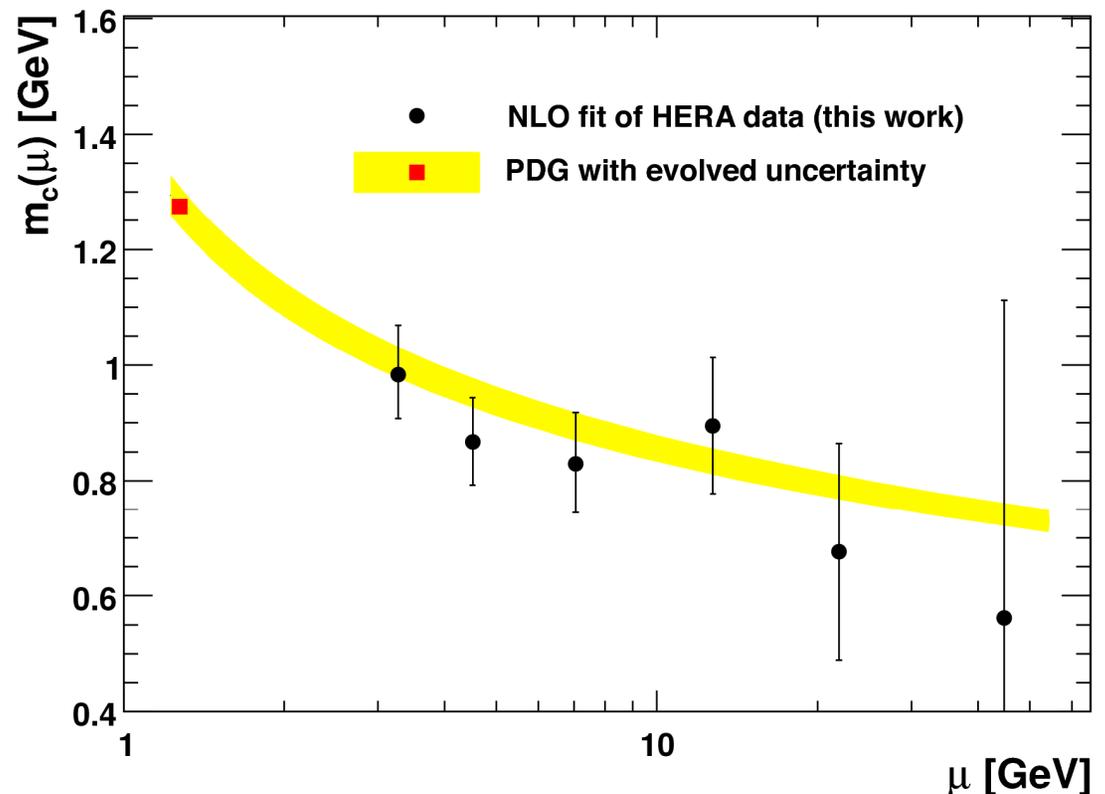
Subset	$\delta_{\text{fit}}^{\text{exp}}$	δ_1	δ_2	δ_3	δ_4	δ_5	δ_6	δ_7
		(m_b)	(α_s)	(f_s)	(Q_0)	(Q_{min}^2)	(param)	(scale)
	[%]	[%]	[%]	[%]	[%]	[%]	[%]	[%]
1	± 5.4	+0.1 -0.4	-1.2 +2.6	-0.4 +0.2	+0.5	+1.4	+0.5	+3.1 +4.3
2	± 6.0	+0.2 -0.5	-0.9 +0.7	-0.5 +0.2	+0.3	+1.0	+0.9	+2.4 +3.6
3	± 7.2	+0.3 -0.7	-0.4 +0.3	-0.8 +0.3	+1.7	+0.3	+1.8	+0.1 +3.7
4	± 9.6	+0.5 -0.8	+0.7 -0.6	-0.8 +0.5	+0.5	-1.2	+0.1	-5.5 +5.4
5	± 19.2	+0.5 -1.2	+1.6 -1.8	-1.2 +0.5	-0.5	+2.1	-1.7	-14.3 +11.6
6	± 63.8	-7.4 -2.9	+5.9 -5.7	-3.0 -7.6	+6.5	-33.3	+9.5	+38.1 -14.2

Table 2: Summary of the systematic uncertainties in the $m_c(m_c)$ determinations. The definitions of the uncertainty sources, the meaning of the symbols in the first and second row and related details are given in the text. In cases where opposite variations of a variable yield uncertainties with the same sign, only the larger one is considered for the uncertainty combination in Table 1. Except for δ_7 , these uncertainties also apply to $m_c(\mu)$.

Conclusions, part I

A. Gizhko et al., DESY-17-048

- Subdividing HERA DIS charm data into 6 kinematic intervals, running of charm-quark mass in \overline{MS} scheme has been determined for the first time (conceptually similar to running of α_s from jets)
- Interplay/treatment of correlations between mass and PDF fits nontrivial, details see DESY-17-048
- Charm-quark mass running consistent with QCD



Higgs Yukawa couplings from m_Q

PoS CHARM2016 (2017) 012

relate m_t, m_b, m_c to associated Higgs Yukawa couplings

LO EW (+NLO QCD) formula:

$$y_Q = \sqrt{2}m_Q/v$$

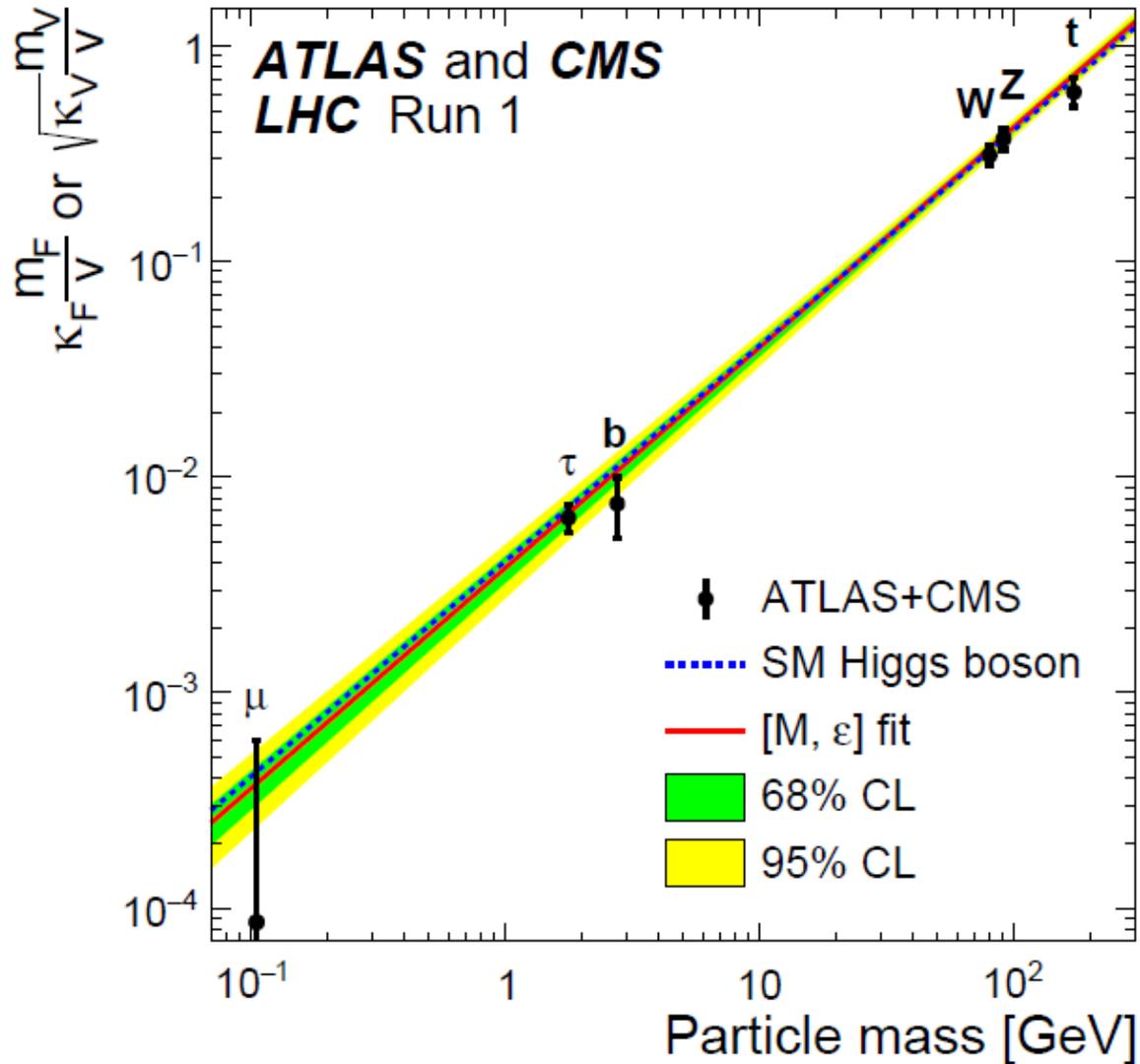
use Higgs/EW scheme in which this relation is exact !



source: viXra blog

Direct measurements of Higgs Yukawa couplings

ATLAS and CMS, JHEP08 (2016) 045



Running of α_s and quark Yukawa couplings

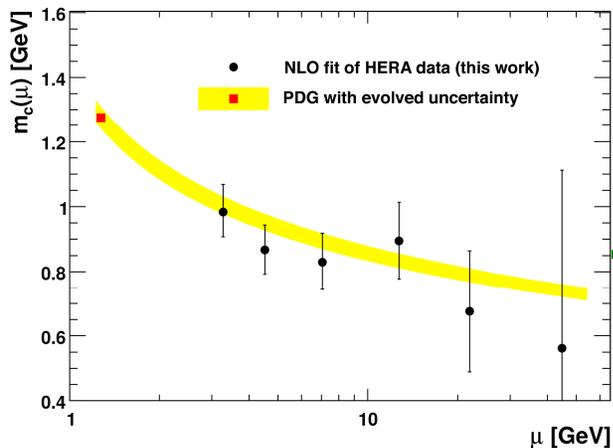
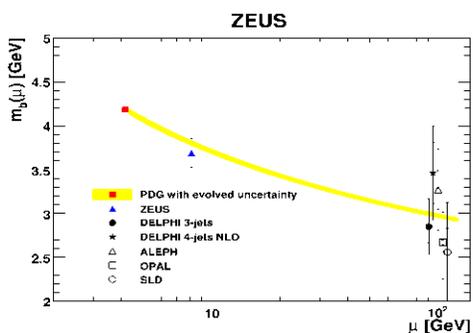
PoS CHARM2016 (2017) 012

relate m_t , m_b , m_c to associated Higgs Yukawa couplings

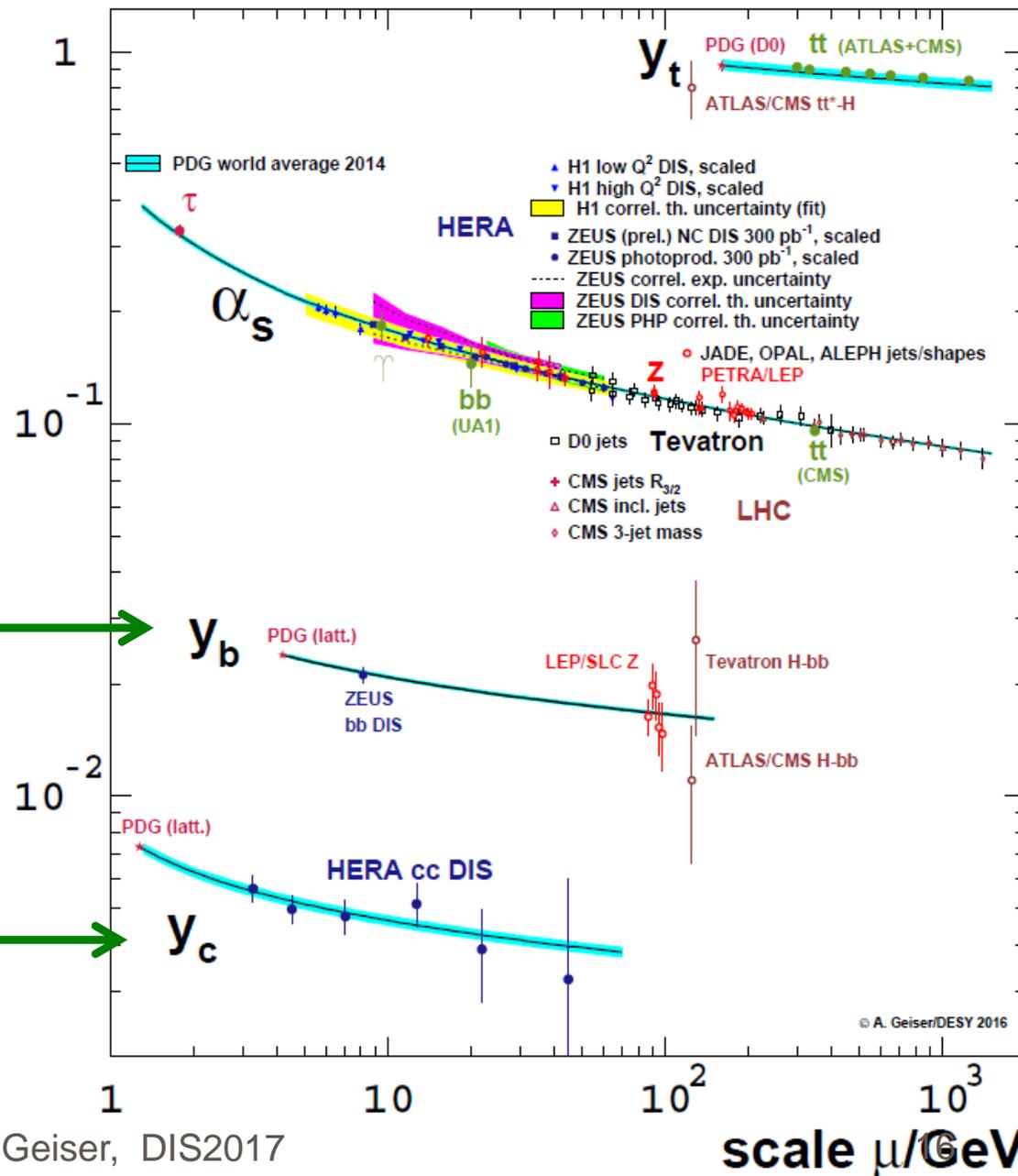
LO EW (+NLO QCD) formula:

$$y_Q = \sqrt{2}m_Q/v$$

for top see backup



running coupling



Conclusion

Part II

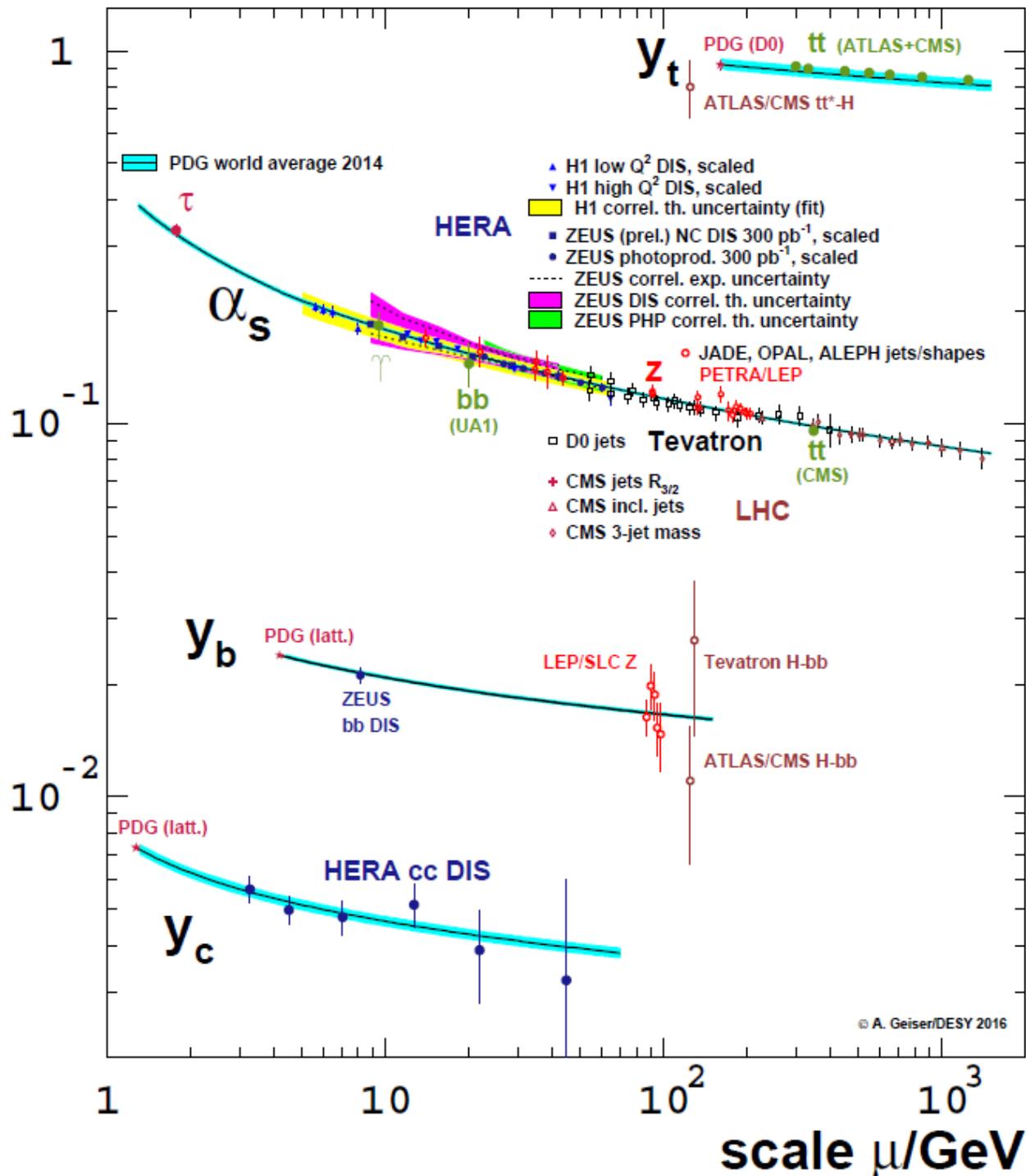
experimental
representation of
running Yukawa couplings
obtained
for the first time

heavy quark
physics is also
QCD + Higgs physics

so far, Higgs couplings
and their running
as obtained from quark
masses are consistent
with directly measured
Higgs couplings

05. 04. 17

running coupling



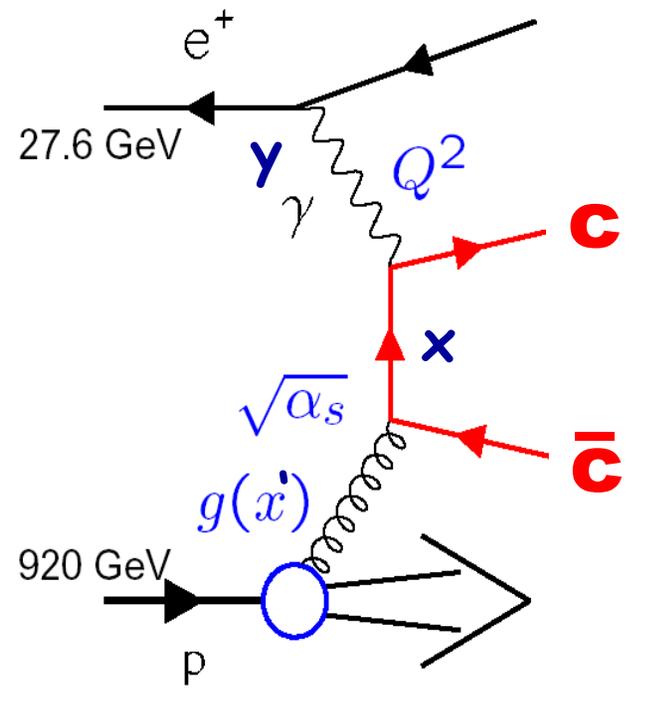
A. Geiser, DIS2017

17



Backup

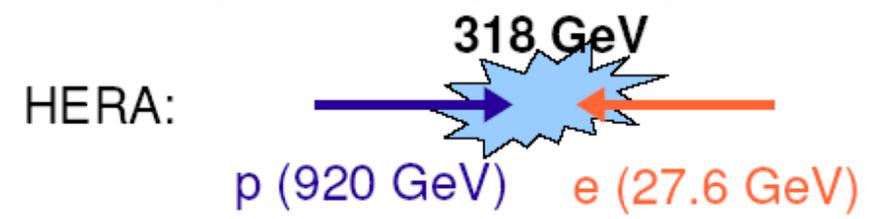
The HERA ep collider and experiments



up to 30%
of cross section

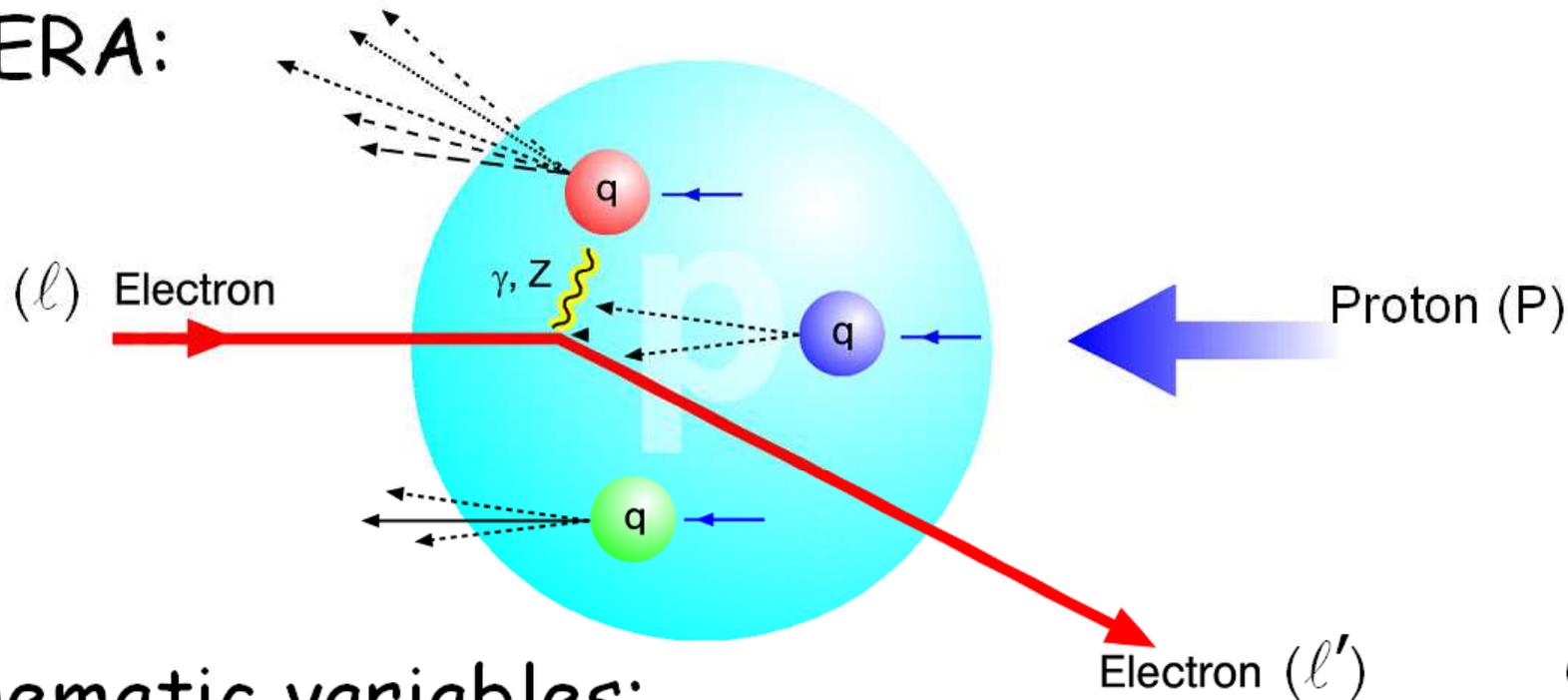


HERA I: $\sim 130 \text{ pb}^{-1}$ (physics)
HERA II: $\sim 380 \text{ pb}^{-1}$ (physics)
combined: $\sim 2 \times 0.5 \text{ fb}^{-1}$



Deep Inelastic ep Scattering at HERA

HERA:



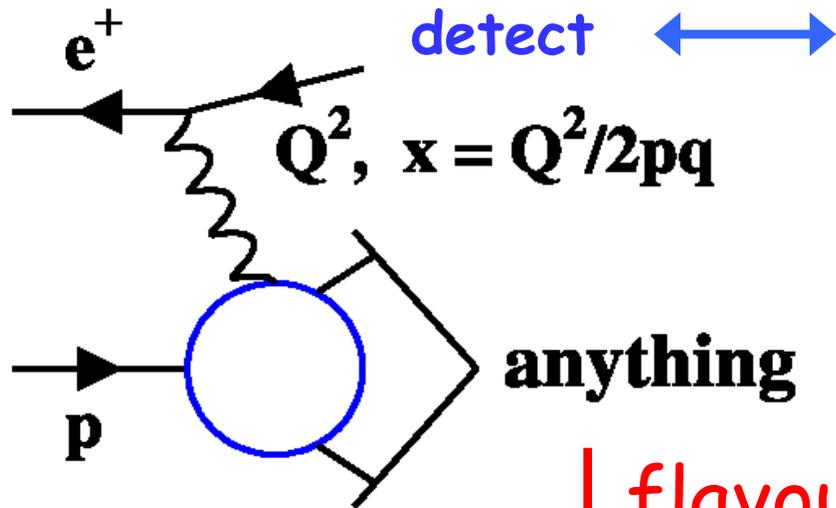
kinematic variables:

$Q^2 = -q^2$	photon (or Z) virtuality, squared momentum transfer
$x = \frac{Q^2}{2Pq}$	Bjorken scaling variable, for $Q^2 \gg (2m_q)^2$: momentum fraction of p constituent
$y = \frac{qP}{lP}$	inelasticity, γ momentum fraction (of e)

$$Q^2 \lesssim 1 \text{ GeV}^2: \text{ photoproduction}$$

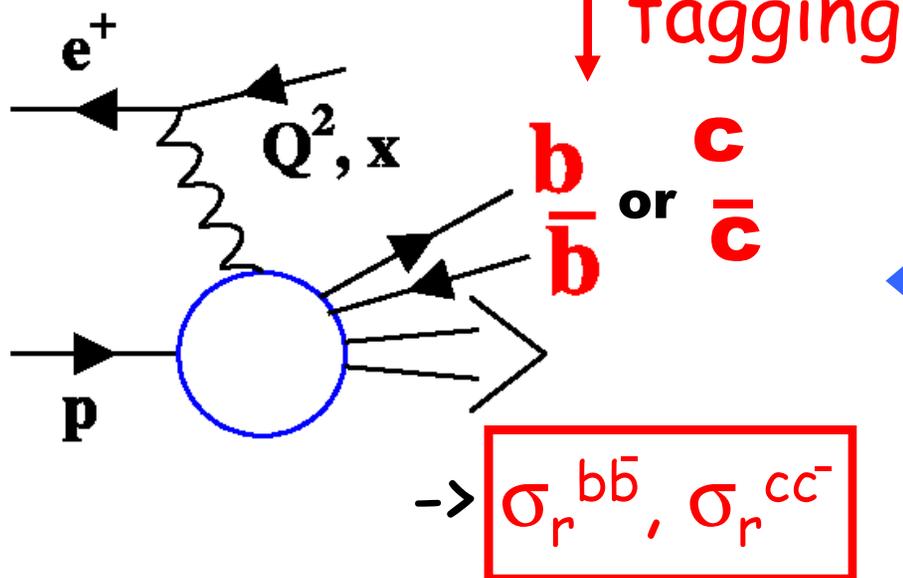
$$Q^2 \gtrsim 1 \text{ GeV}^2: \text{ DIS}$$

Heavy flavour contributions to F_2

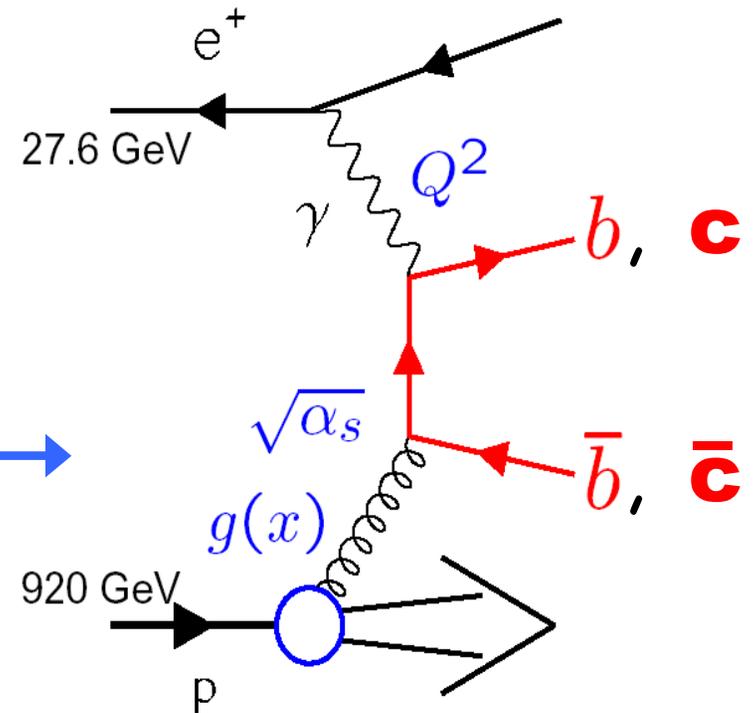


Measure cross section

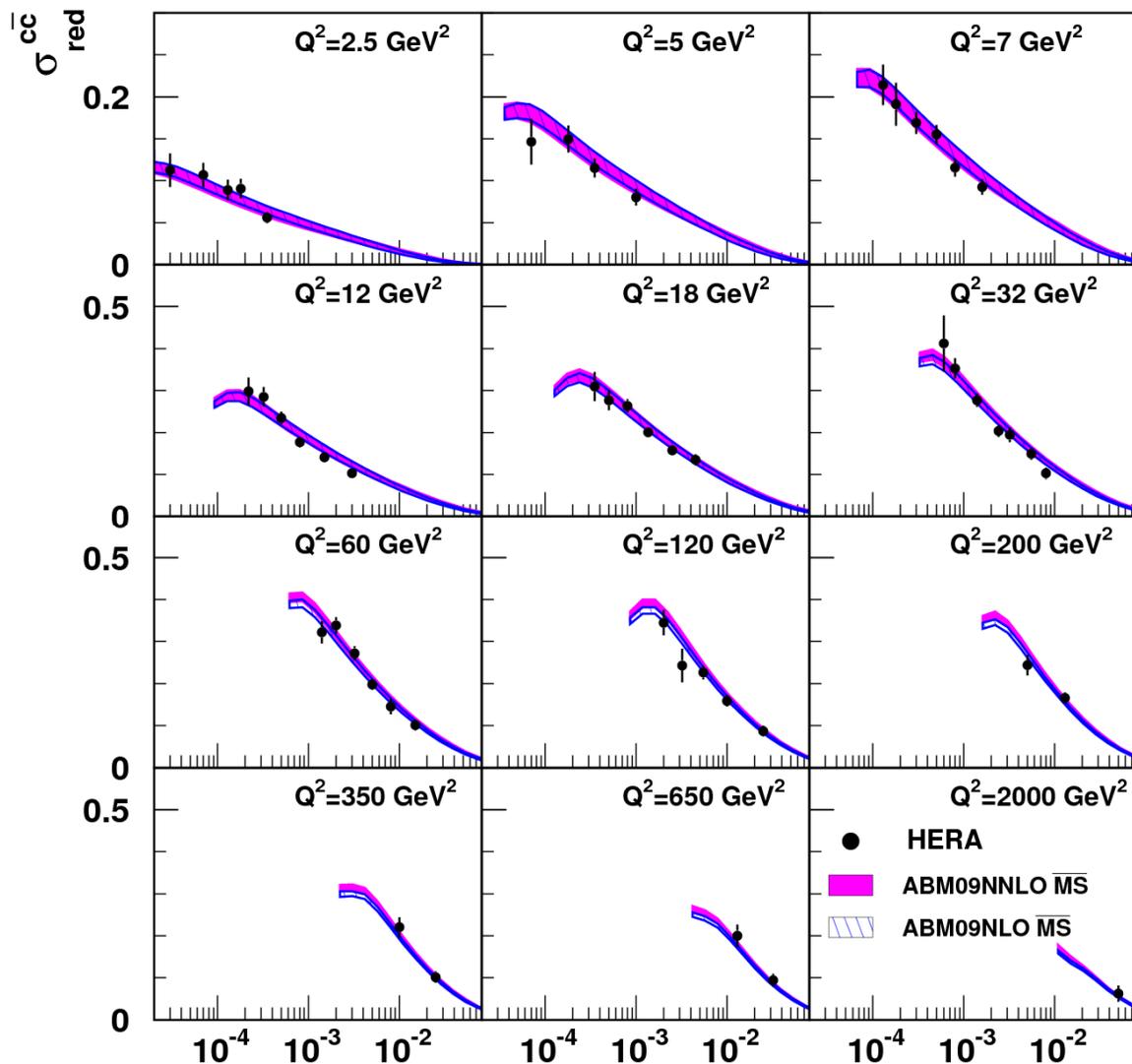
$$\frac{d^2\sigma}{dx dQ^2} \approx \frac{2\pi\alpha^2}{Q^4 x} \left\{ \left[1 + (1-y)^2 \right] \sigma_r(x, Q^2) \right\}$$



QCD



H1 and ZEUS



very good description
of data
in full kinematic range

unambiguous treatment
of m_c in all terms of
calculation

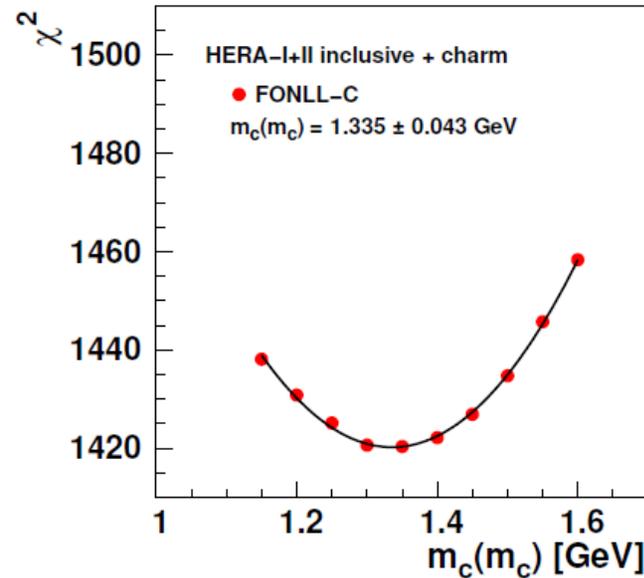
here: \overline{MS} running mass

(similar predictions for
pole mass)

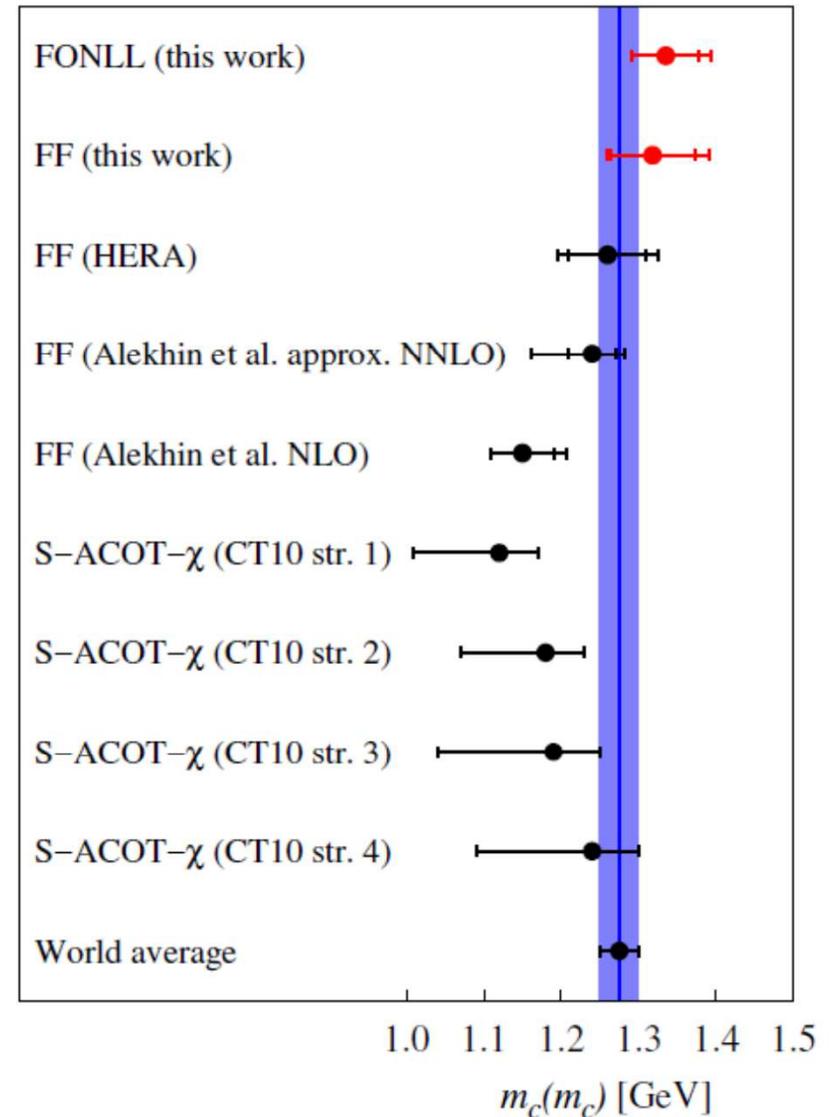
x

$m_c(m_c)$ from FONLL fit of HERA data

V. Bertone et al., arXiv 1605.01946, JHEP 1608 (2016) 050



scheme	$m_c(m_c)$ [GeV]
FONLL (this work)	$1.335 \pm 0.043(\text{exp})_{-0.000}^{+0.019}(\text{param})_{-0.008}^{+0.011}(\text{mod})_{-0.008}^{+0.033}(\text{th})$
FFN (this work)	$1.318 \pm 0.054(\text{exp})_{-0.010}^{+0.011}(\text{param})_{-0.019}^{+0.015}(\text{mod})_{-0.004}^{+0.045}(\text{th})$
FFN (HERA) [9]	$1.26 \pm 0.05(\text{exp}) \pm 0.03(\text{mod}) \pm 0.02(\text{param}) \pm 0.02(\alpha_s)$
FFN (Alekhin <i>et al.</i>) [24]	$1.24 \pm 0.03(\text{exp})_{-0.02}^{+0.03}(\text{scale})_{-0.07}^{+0.00}(\text{th})$ (approx. NNLO) $1.15 \pm 0.04(\text{exp})_{-0.00}^{+0.04}(\text{scale})$ (NLO)
S-ACOT- χ (CT10) [29]	$1.12_{-0.11}^{+0.05}$ (strategy 1) $1.18_{-0.11}^{+0.05}$ (strategy 2) $1.19_{-0.15}^{+0.06}$ (strategy 3) $1.24_{-0.15}^{+0.06}$ (strategy 4)
World average [53]	1.275 ± 0.025



top quark mass running

very preliminary procedure (with caveats, "cheated" a bit):

- use (conceptually constant) LO MC mass measured as function of scale-dependent quantity (e.g. $m_{\text{++}}^-$)
- check self-consistency of cross section measurements with data used for mass determination
- 'convert' LO MC mass to NLO pole mass by comparing MC and pole mass extractions from same data
- convert pole mass to $\overline{\text{MS}}$ mass using 3-loop QCD
- use 1-loop evolution for actual running (NLO QCD)

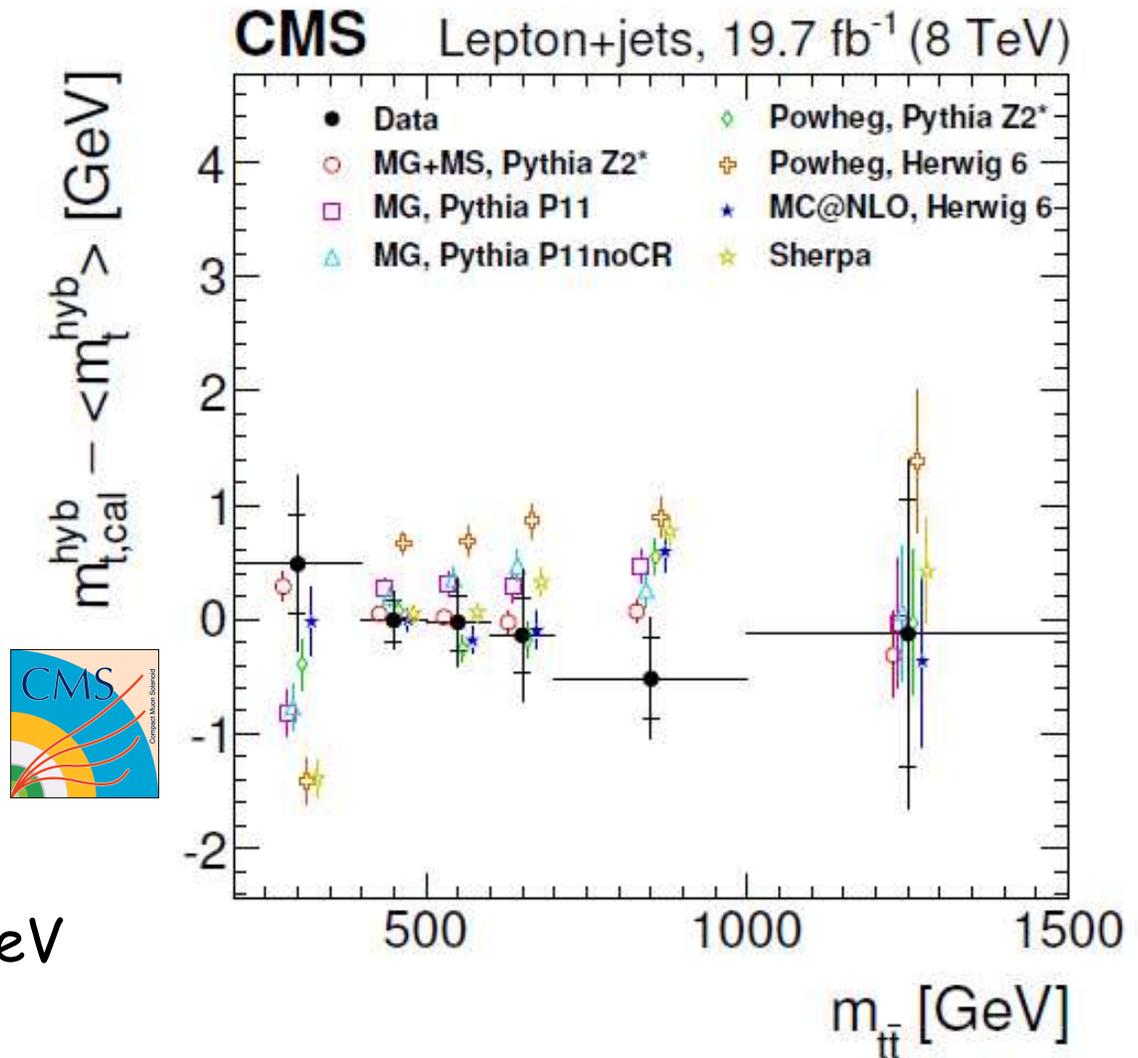
(in the future, like for m_c and m_b , extract NLO (or NNLO) running mass directly from data, e.g. cross section, in each kinematic bin)

top quark mass as function of $m_{t\bar{t}}$

CMS-TOP-14-022, Phys. Rev. D93 (2016) 072004

"MC mass"

deviation
from average of
 $172.35 \pm 0.16_{\text{stat}} \pm 0.48_{\text{sys}} \text{ GeV}$



differential top cross section shape consistent with NLO

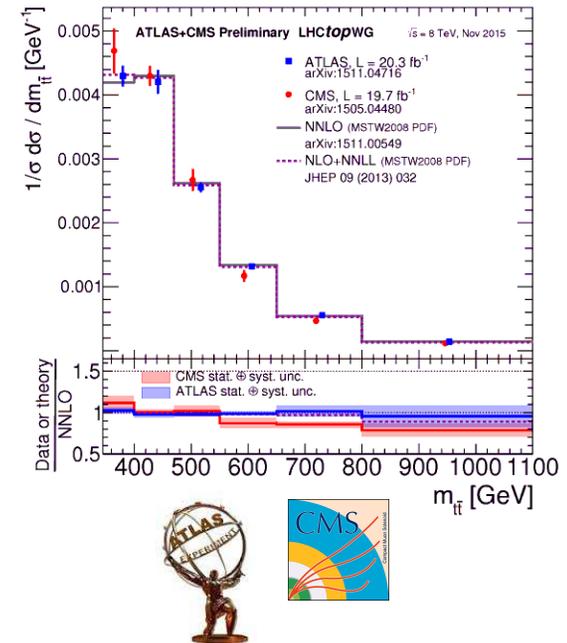
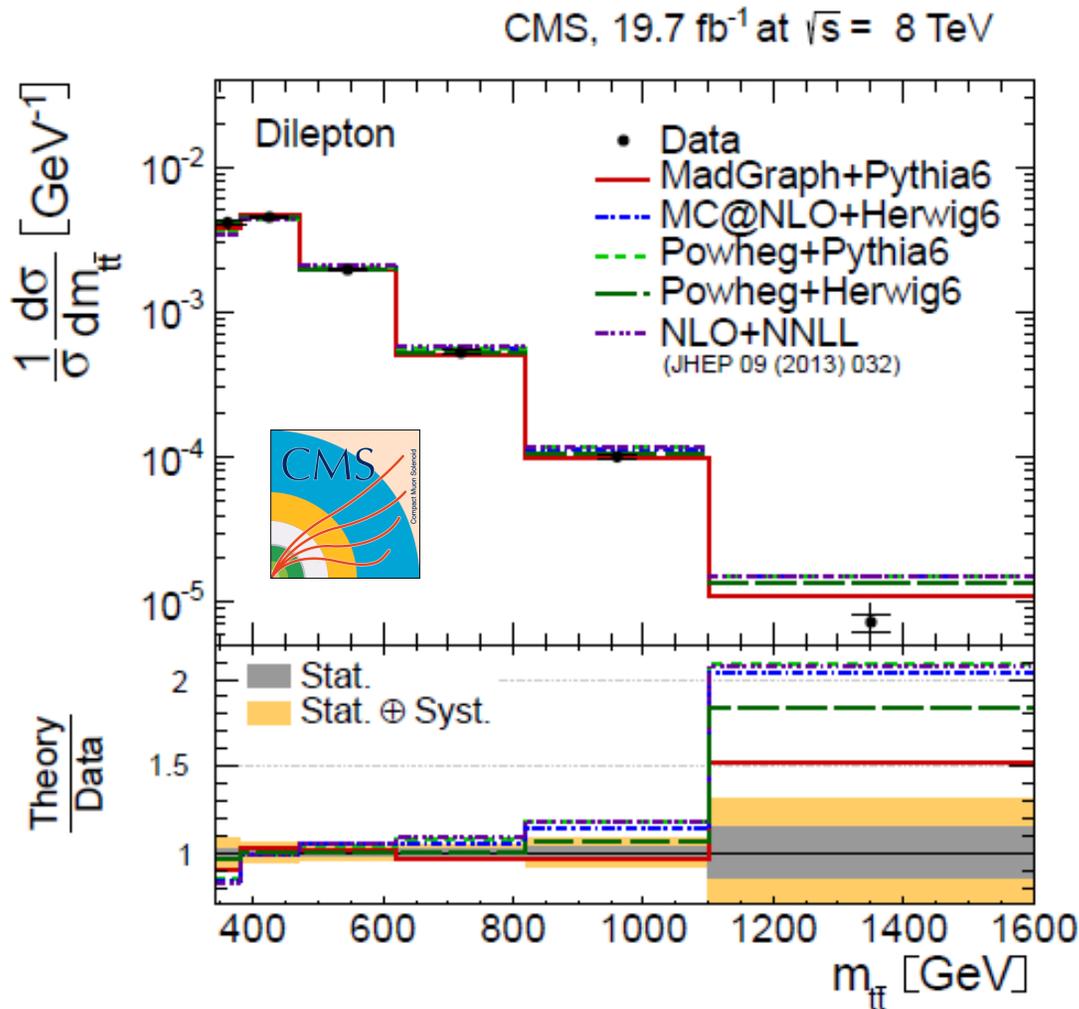
CMS-TOP-12-028, Eur. Phys. J. C75 (2015) 542

LHCtopWG

NLO theory
uses
pole mass
scheme

use CMS to
be consistent
with previous
slide

similar results
for lepton+jets
channel only



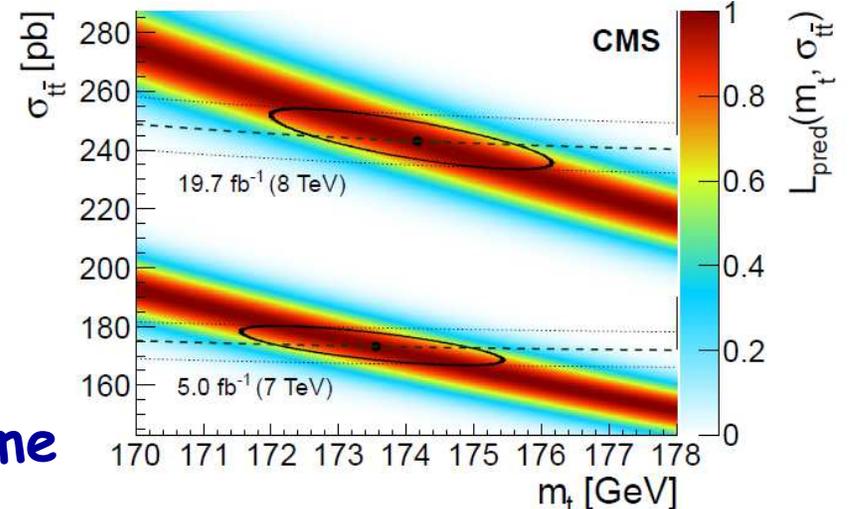
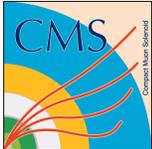
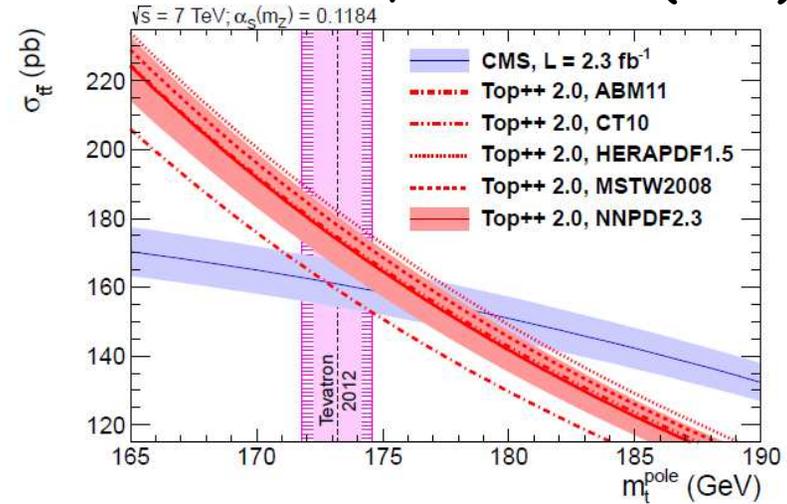
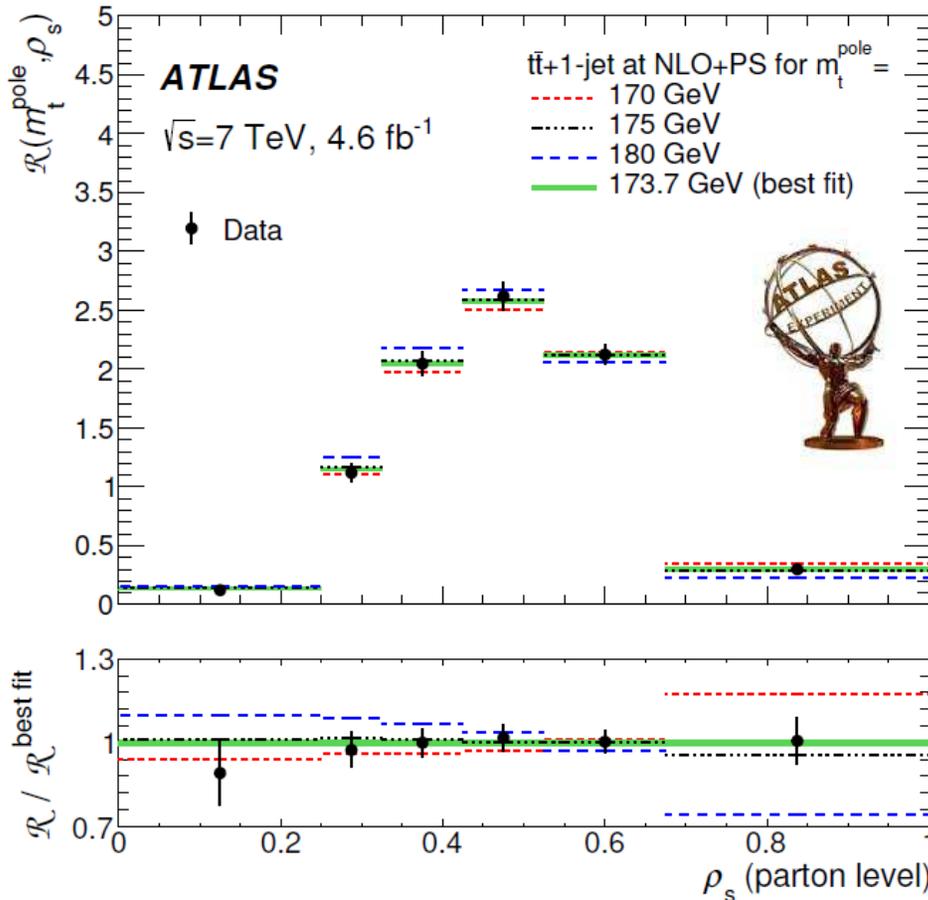
-> measurements and LO+PS/NLO theory are self-consistent and consistent with ATLAS and NNLO

convert top MC mass to pole mass

ATLAS, JHEP10 (2015) 121,

CMS-TOP-12-022, Phys. Lett. B728 (2014) 526

CMS-TOP-13-004, JHEP 1608 (2016) 029



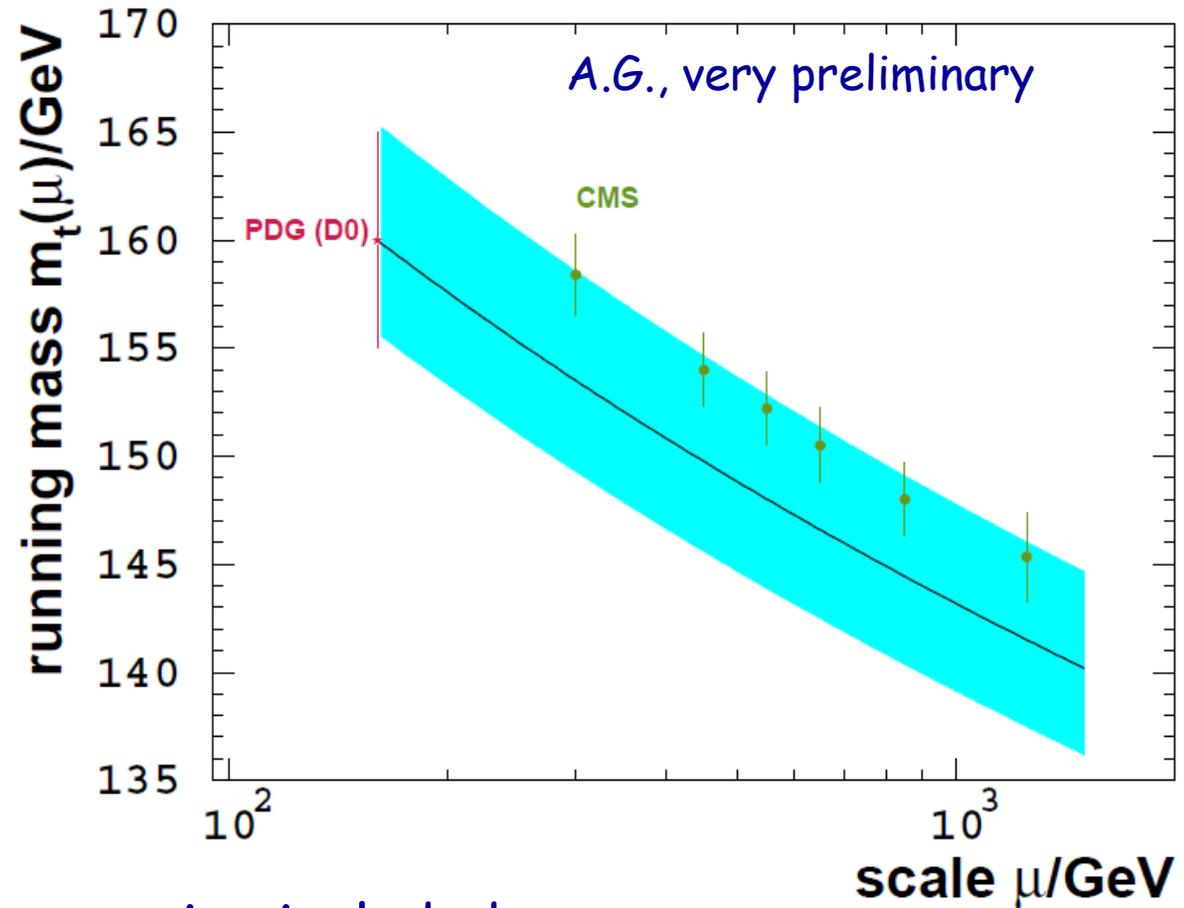
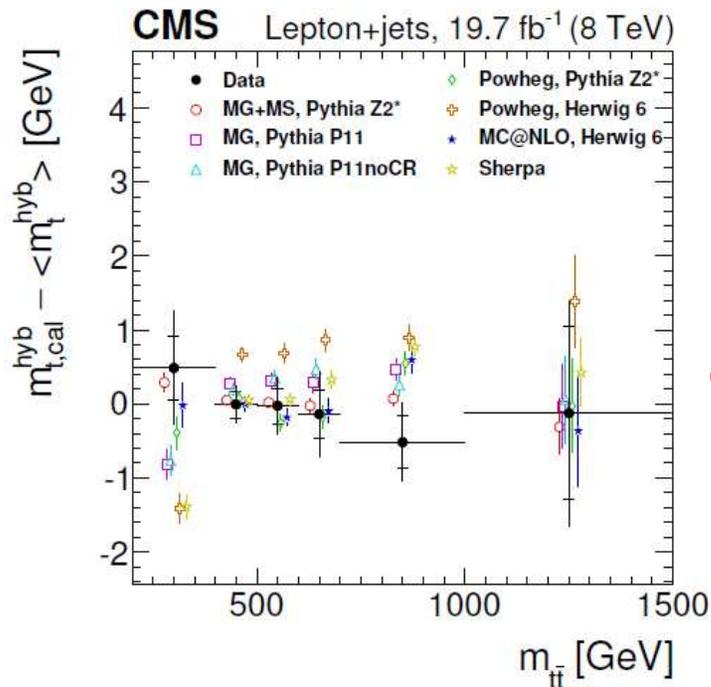
“MC” and pole masses almost the same

ATLAS: $m_t(\text{pole}) = 173.7 \pm 1.5_{\text{stat}} \pm 1.4_{\text{syst}} + 1.0 - 0.5_{\text{th}}\text{ GeV}$

CMS: $m_t(\text{pole}) = 173.8 + 1.7 - 1.8_{\text{total}}\text{ GeV} \leftrightarrow m_t(\text{MC}_{\text{CMS}, l+\text{jets}}) = 172.35 \pm 0.16_{\text{stat}} \pm 0.48_{\text{syst}}\text{ GeV}$

PDG: $m_t(\text{pole}) = 176.7 \pm 4.0 - 3.4\text{ GeV}$, “ $m_t(\text{MC})$ ” = $173.21 \pm 0.51_{\text{stat}} \pm 0.71_{\text{syst}}\text{ GeV}$

convert pole masses to running mass



caveat:

not all uncertainties from conversion included

(needs theoretically better defined procedure!)

-> take with grain of salt, **for illustration purposes**

Discussion

of future conceptual improvements

- avoid *MC mass* and *pole mass* intermediate steps for top
-> extract $m_+(\mu)$ directly from data, as already done for b,c
(e.g. from absolute m_{++} cross sections in CMS-TOP-16-008)
need *NLO QCD theory* for LHC using *running mass*
- extend *LO EW + NLO QCD* approach
(running of Higgs couplings is purely QCD-induced!)
to *NLO EW + NNLO QCD + interference*
highly non-trivial but eventually necessary
(Standard Model is not QCD only)