

Heavy flavour schemes, masses, scales


Achim Geiser, DESY Hamburg

Achim.Geiser@desy.de

QCD at LHC workshop

ECT Trento, 29 September 2010

- Measurement/fit of b, c masses
- Variants of FFNS scheme
- QCD scale choice



Can we measure the charm and bottom quark masses as physical parameters from F_2^c/F_2^b data?

heavy quark mass in different schemes

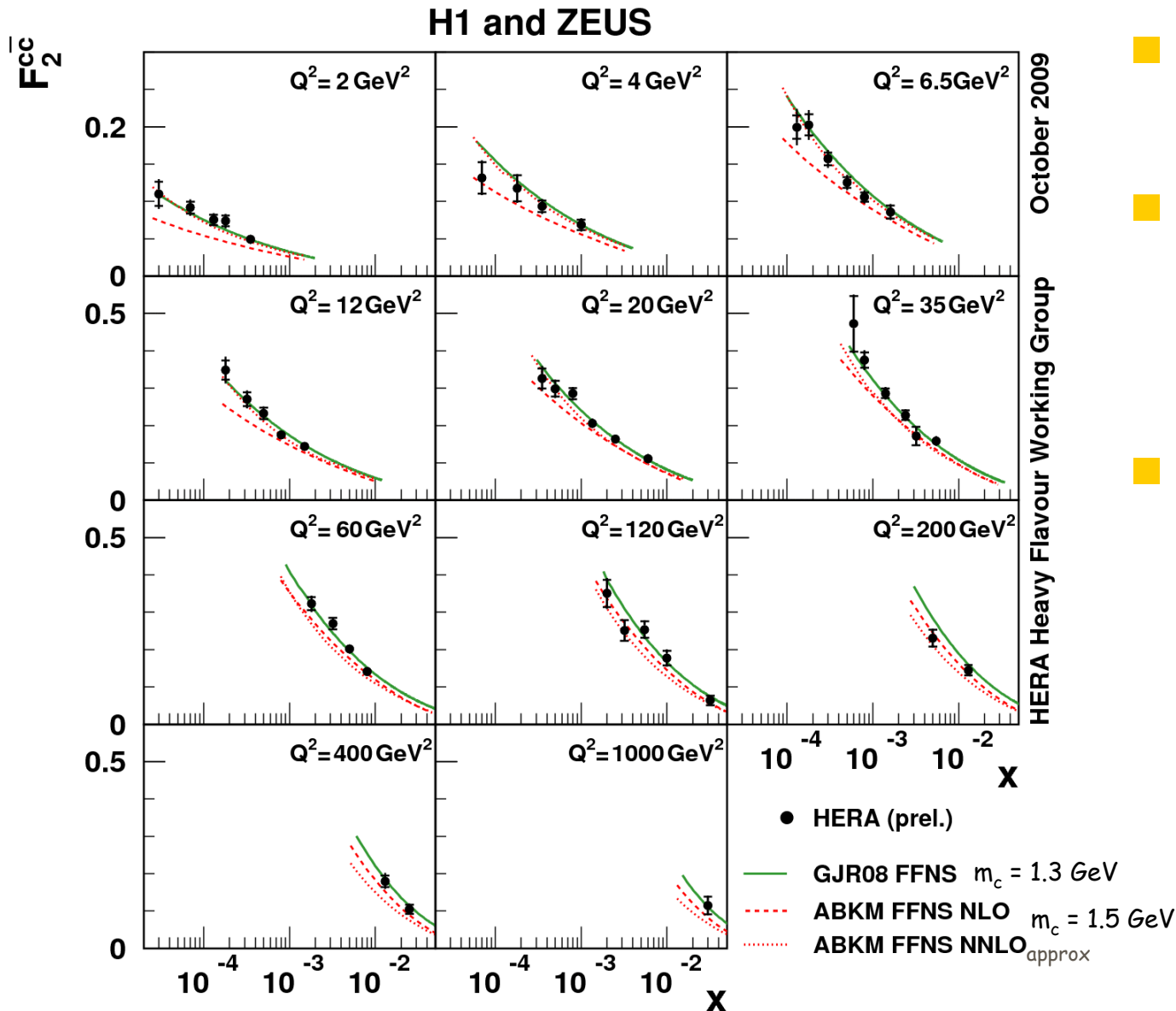
- **ZM-VFNS:** mass is just kinematic threshold parameter (onset of massless "heavy" quark PDF)
-> can not be used for meaningful mass measurement in QCD
- **FFNS:** until recently, all existing calculations for ep have used **pole mass** (on-shell mass renormalization scheme)
PDG09: $m_b \sim \underline{4.79+0.19-0.08} \text{ GeV}$, $m_c \sim \underline{1.65 \pm 0.18} \text{ GeV}$

but: pole mass definition has intrinsic uncertainty of order Λ_{QCD}
-> change scheme to $\overline{\text{MS}}$ running mass to predict and fit F_2^c ?
see talk S. Alekhin (ABKM FFNS)
-> $m_c(m_c) \sim 1.30 \pm 0.12 \text{ GeV}$ (PDG: $1.27+0.07-0.11 \text{ GeV}$)

might be relatively clean way to measure charm mass

but: beware of systematics from higher order QCD corrections !

FFNS predictions: NNLO vs. NLO



- both ABKM and GJR use pole mass
- ABKM: approximate NNLO correction (threshold resummation) very large at low Q^2
- GJR seems to compensate absence of threshold correction terms in matrix elements by choice of very low charm mass

-> systematic uncertainty on m_c from higher order corrections $\sim 0.2 \text{ GeV}$?

heavy quark mass in different schemes

- **GM-VFNS**: all existing calculations use **pole mass** for massive part of calculation, e.g. Martin, Stirling, Thorne, Watt, arXiv:1007.2624 [hep-ph]. Their pole mass evaluation from conversion of running masses: $m_b \sim \underline{4.9 \pm 0.2 \text{ GeV}}$, $m_c \sim \underline{1.5 \pm 0.2 \text{ GeV}}$

However, massive calculation is matched to massless part through various different matching schemes

-> sensitivity of cross sections to mass is modified through model assumptions in the matching

-> mass, although still based on pole mass concept, becomes **effective model parameter**

my conclusion: **GM-VFNS schemes can probably not be used for well-defined measurement of heavy quark masses** (unless unphysical zero mass effects can somehow be eliminated)

Example: variants of TR-VFNS

$$F_{2,c}^{\text{VFNS}} = \overset{\text{NLO mass effects}}{F_{2,c}^{\text{FFNS}}(N_f = 3)} + \overset{\text{NLL resummation of logs}}{F_{2,c}^{\text{ZMVFNS}}(N_f = 4)} - \overset{\text{interpolation}}{F_{2,c}^{\text{ASYMP}}(N_f = 3)}$$

pole mass "massless" !

Thorne: each term in the combination spoils mass definition
($F_2^{\text{ZMVFNS}} - F_2^{\text{asymp}}$) can be modified by corrections which fall like m_H^2/Q^2 .

-> 4 "free" parameters a,b,c,d (of range 0-1)

Variants of TR-VFNS ($m_c=1.4 \text{ GeV}$ fixed)

Thorne

+my comments

6 extreme variations tried.

GMVFNS1 – $b = -1, c = 1$.

GMVFNS2 – $b = -1, c = 0.5$.

GMVFNS3 – $a = 1$.

GMVFNS4 – $b = +0.3, c = 1$ – fit.

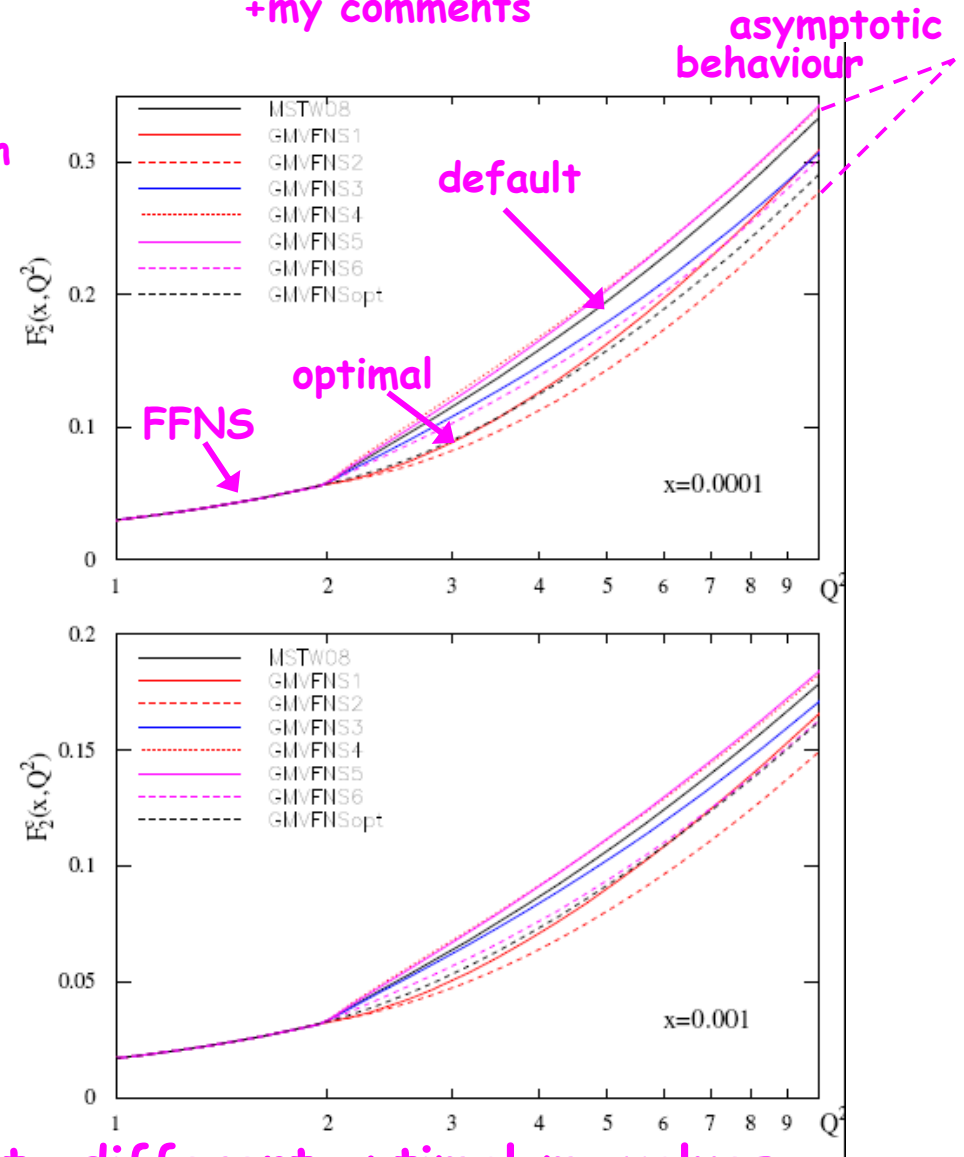
GMVFNS5 – $d = 0.1$ – fit.

GMVFNS6 – $d = -0.2$ – fit.

Variations in $F_2^c(x, Q^2)$ near the transition point at NLO due to different choices of GM-VFNS.

Optimal, $b = -2/3, c = -1, d = -1$, smooth behaviour.

can have noticeable effect even at LHC (3% on σ_Z)



fit to data with $Q^2 > 2 \text{ GeV}^2$ will lead to different optimal m_c values
 -> model dependence, not QCD! -> effective mass parameter

Variants of TR-VFNS, NNLO (approx)

Thorne

+my comments

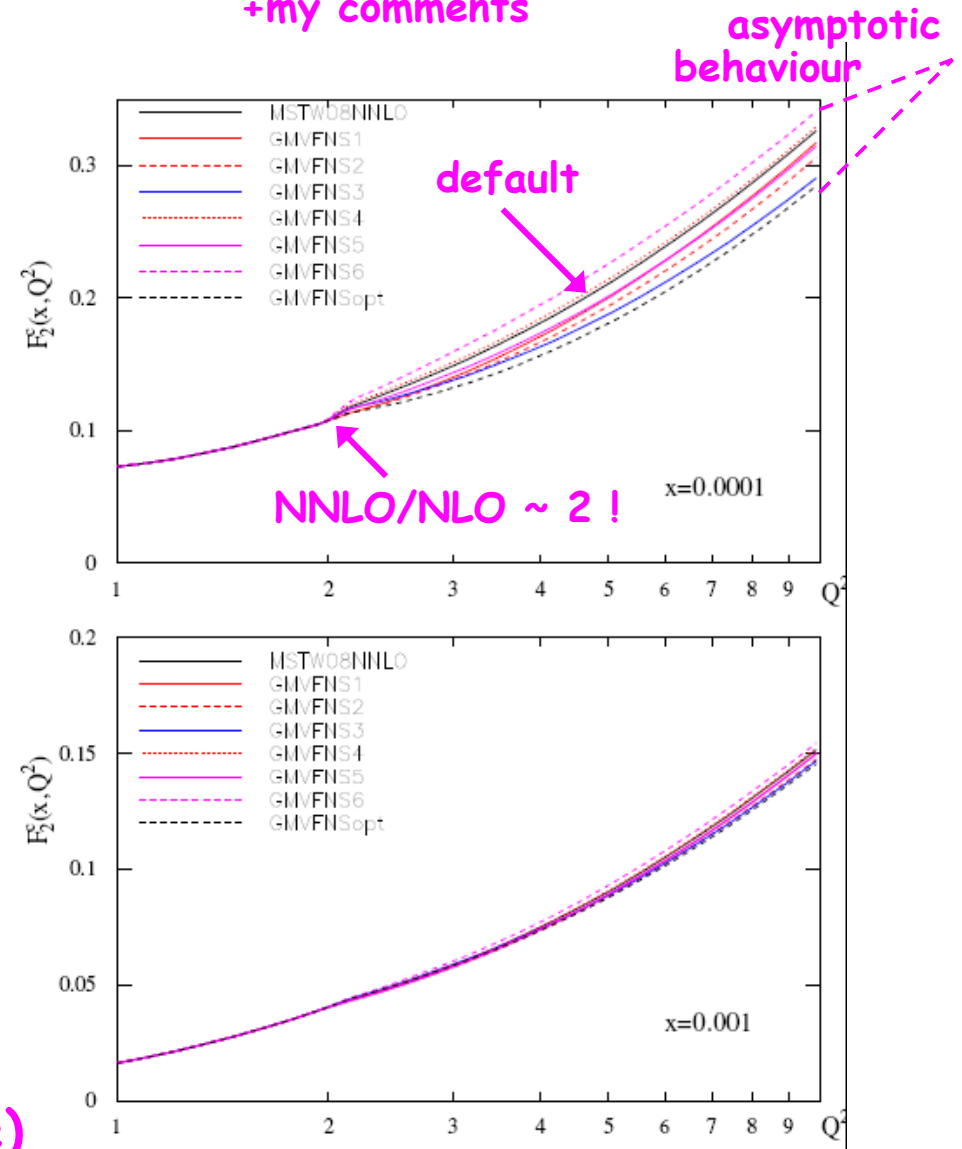
Variations in $F_2^c(x, Q^2)$ near the transition point due to different choices of GM-VFNS at NNLO.

Very much reduced, almost zero variation until very small x .

Shows that NNLO evolution effects most important in this regime.

still noticeable effect at low x

GMVNS1 smooth both at NLO and NNLO (see curves), better χ^2 for global MRST fit (Thorne)





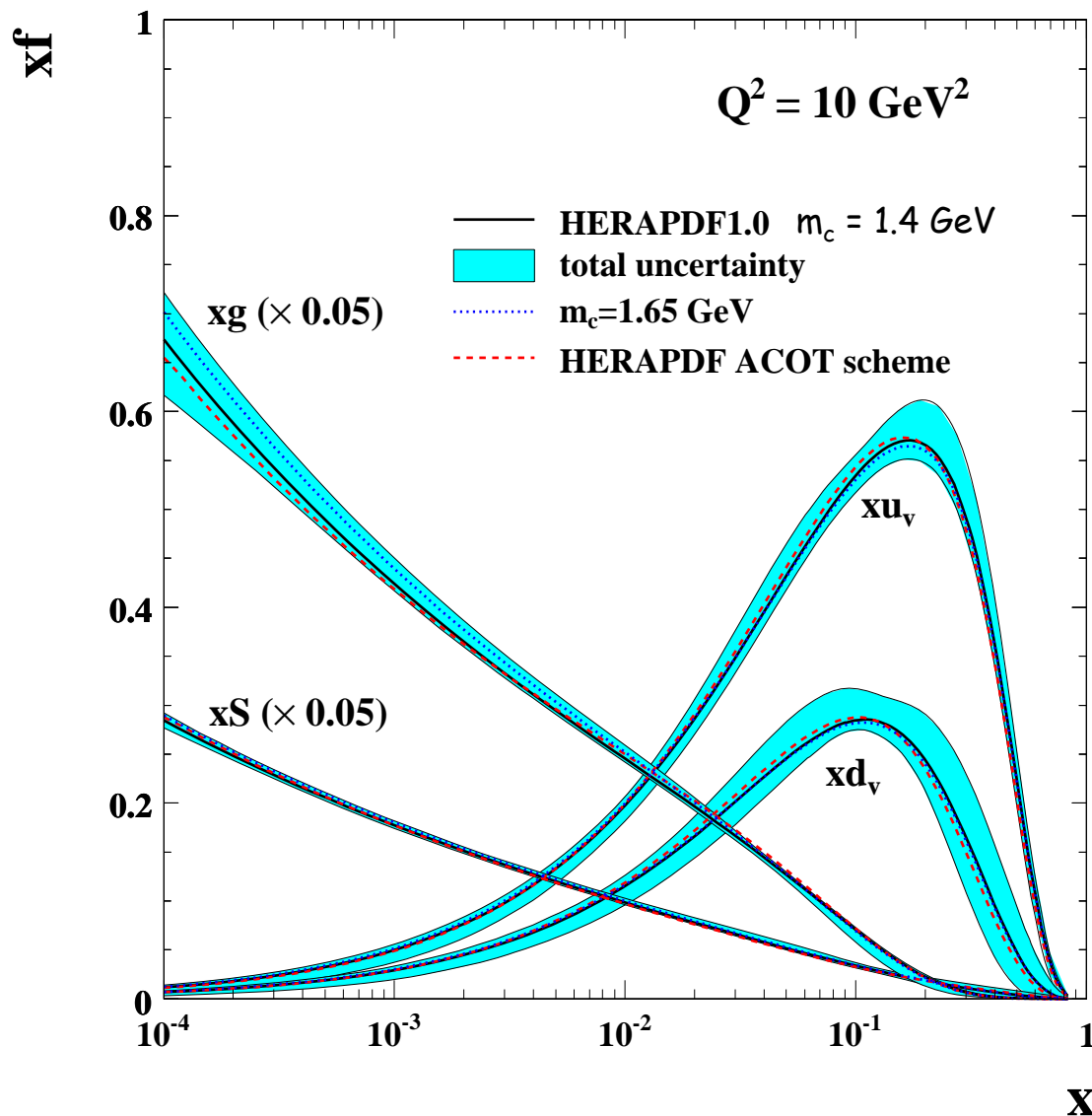
■ Can we reduce the heavy quark mass uncertainty on LHC standard candle processes by fitting m_c ?

-> for W/Z probably yes, see previous talk (R. Placakyte)
 m_c does NOT need to be physical parameter for this purpose
-> can use "any" scheme

Why does it work?

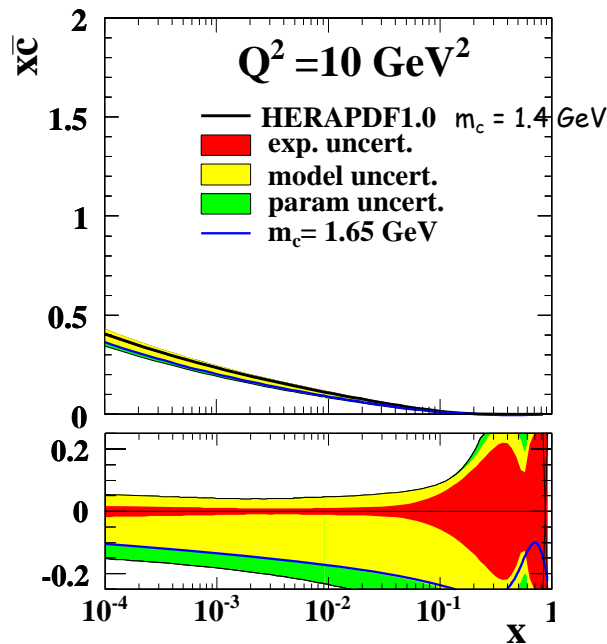
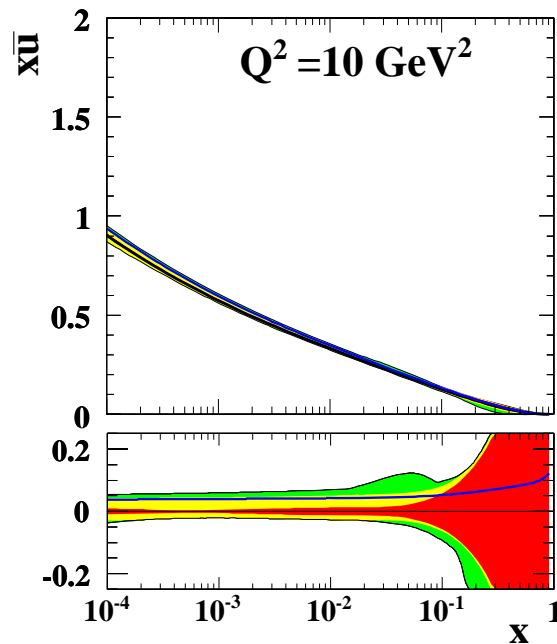
inspect e.g. HERAPDF

Cooper-Sarkar
+my comments



larger m_c
-> more gluon at low x

H1 and ZEUS

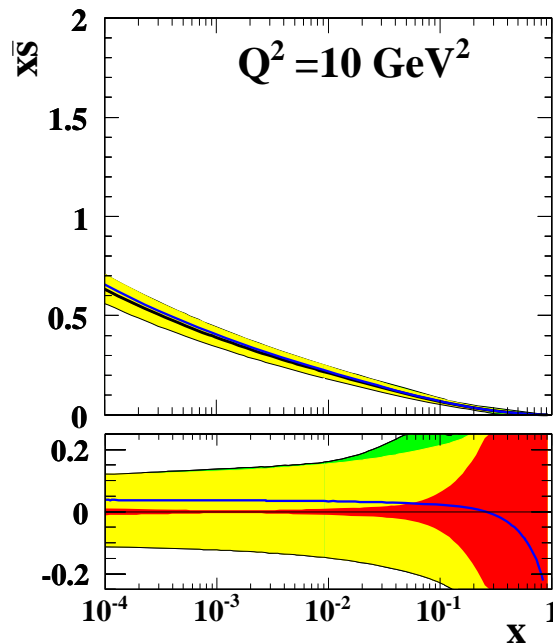
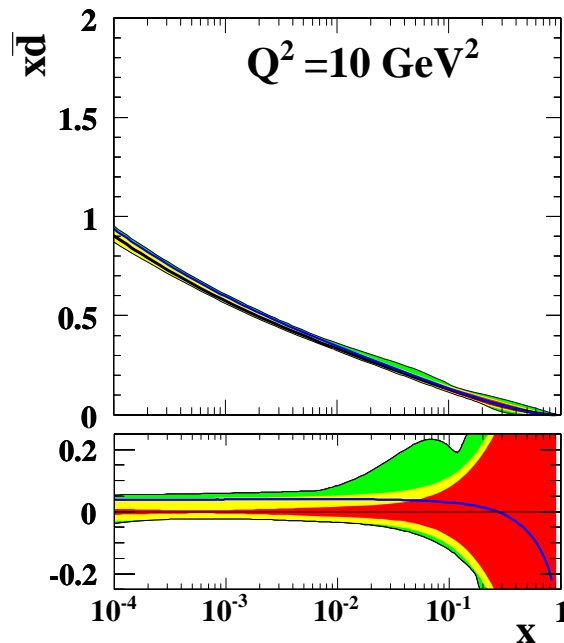


Cooper-Sarkar
+my comments

larger m_c


-> more gluon, less c

-> more light quarks



fitting m_c to describe F_2^c seems to yield almost the same sea quark flavour mixture in relevant x range for "any" scheme

-> reduced uncertainty on W/Z predictions at LHC

- 
- Will this also work for gluon-initiated processes (Higgs, top, ...)?
 - will it work for other kinematic regions?

not clear (under investigation)

until clarified: propose to keep enveloping mass

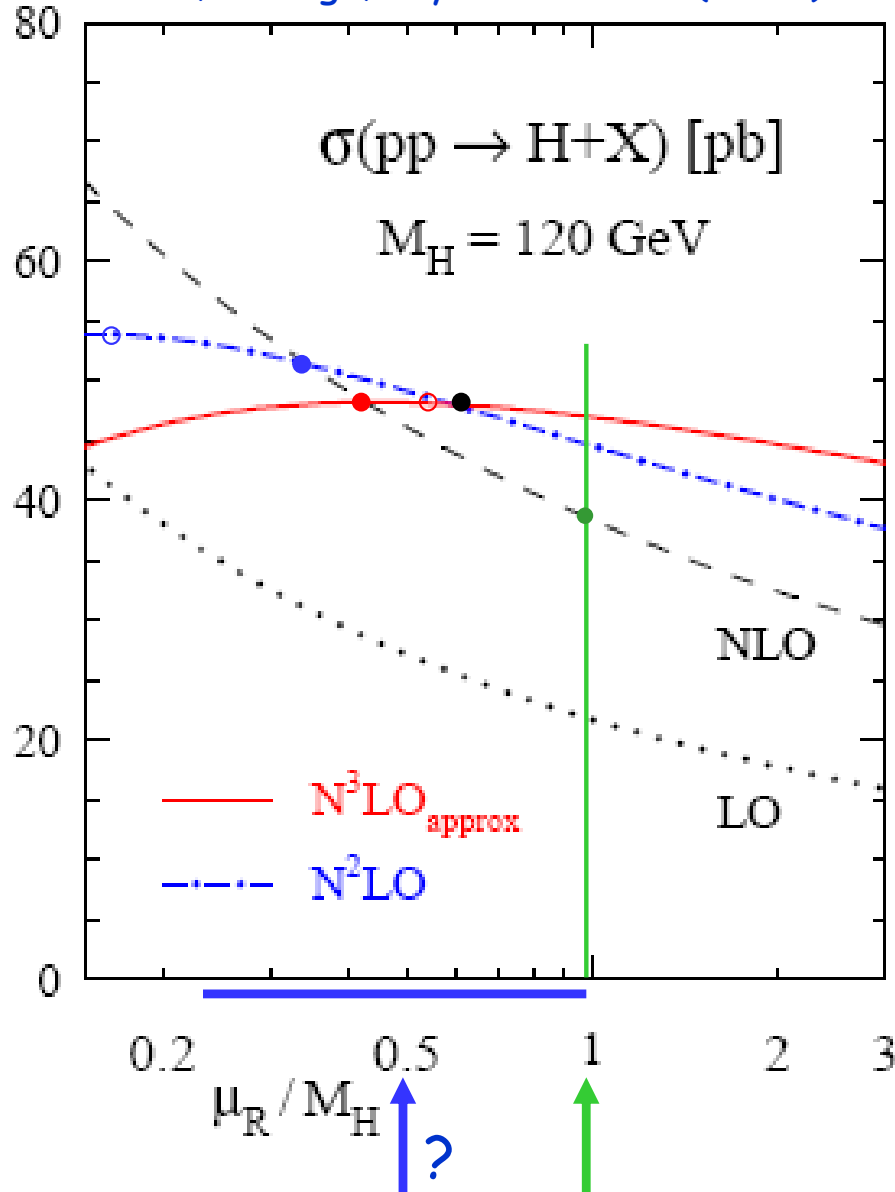
variation as PDF uncertainty, e.g. $m_c = 1.50 \pm 0.15 \text{ GeV}$

Variants of FFNS scheme

- Variants of GM-VFNS discussed repeatedly
- Less well known: variants of FFNS scheme exist, too !
All FFNS schemes have fixed number of real flavours in PDFs + heavy flavours generated dynamically, but
 - fixed flavour α_s evolution, e.g. MRST04FF, CTEQ5F3, CTEQ5F4, Riemersma et al., HVQDIS (Harris & Smith)
-> heavy flavour loops consistently removed from theory (?)
or
 - variable flavour α_s evolution, e.g. ABKM, GJR
-> heavy flavour loops treated explicitly (and partially resummed?)
- Schemes differ by $\alpha_s \log(\mu^2/m^2)$ terms in α_s and by corresponding heavy flavour loop terms in (both light and heavy) matrix elements/coefficient functions to avoid double counting of loops
- consider highest HERA energies: $\log(10000/1.5^2) \sim 8$,
for PDFs and matrix elements evaluated at LHC energies even larger
-> variable flavour α_s evolution scheme seems preferable for global fits and/or precise LHC predictions

NLO scale choice? example: Higgs at LHC

S. Moch, A. Vogt, Phys.Lett. B631 (2005) 48



in principle arbitrary, but

NNLO stability:

- NNLO = NLO
- $d\sigma_{NNLO}/d\mu = 0$

N³LO stability:

- $N^3LO = NLO$
- $N^3LO = NNLO$
- $d\sigma_{NLO+NLL}/d\mu = 0$

— "natural" scale

NNLO/N³LO calculations, where available, often suggest ren./fact. scale ~ half "natural" scale for NLO

more details: [arXiv:0711.1983](https://arxiv.org/abs/0711.1983) [hep-ex]

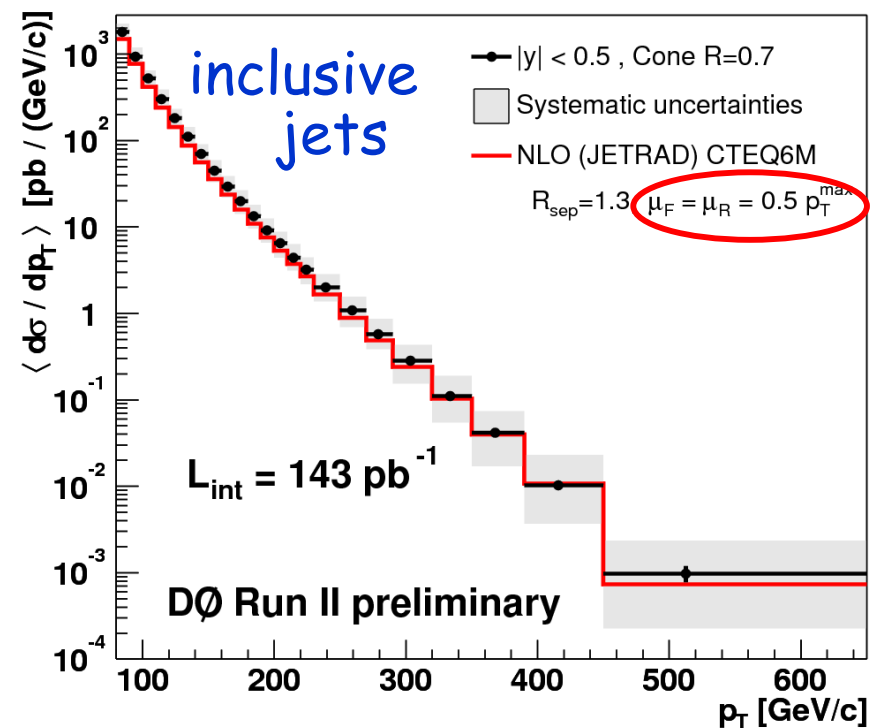
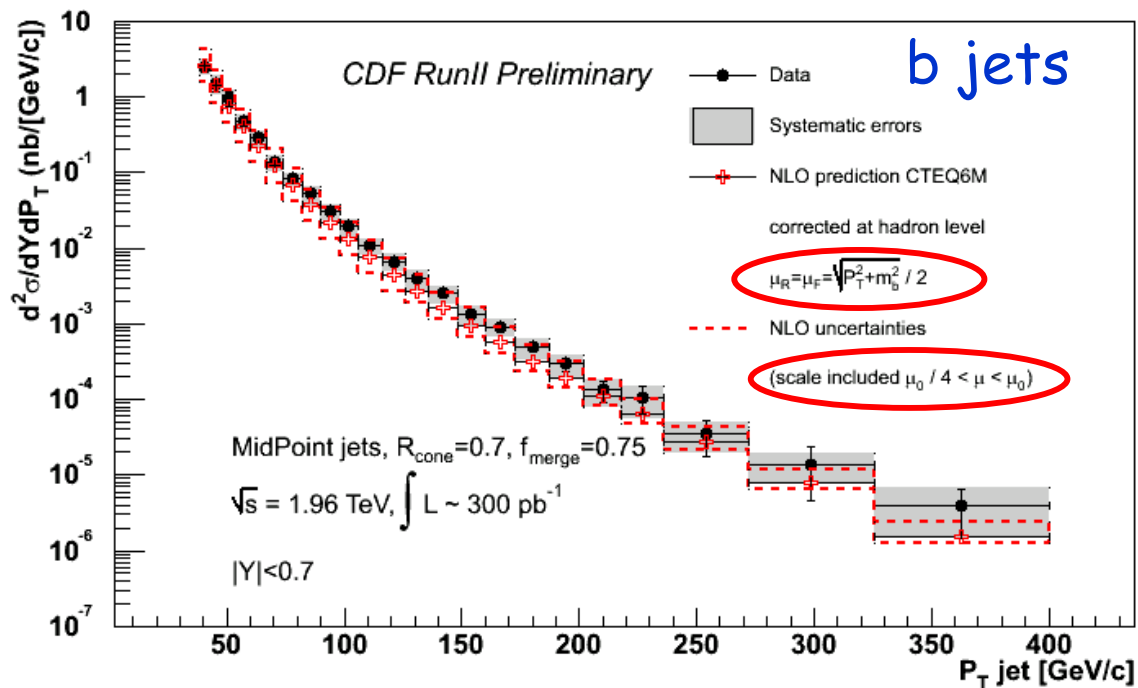
personal remark:

either dedicated scale study, or

- consider to use default QCD scale $\mu_0/2$ for your favourite NLO cross section predictions, including LHC, in particular before claiming discrepancies

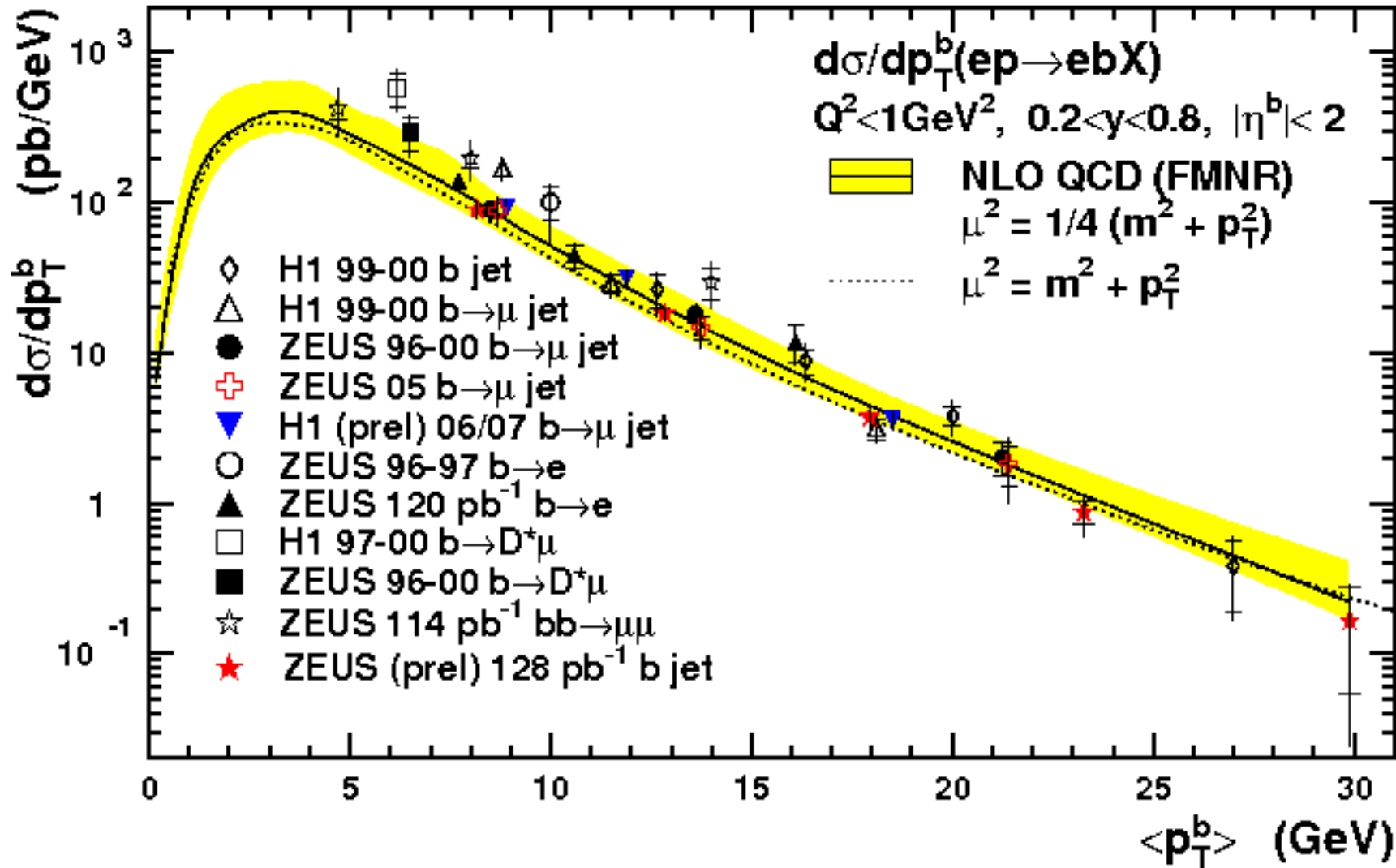
some people are doing this already:
(also see talks J. Huston and M. Grazzini)

more details: [arXiv:0711.1983](https://arxiv.org/abs/0711.1983) [hep-ex]



beauty in photoproduction

HERA



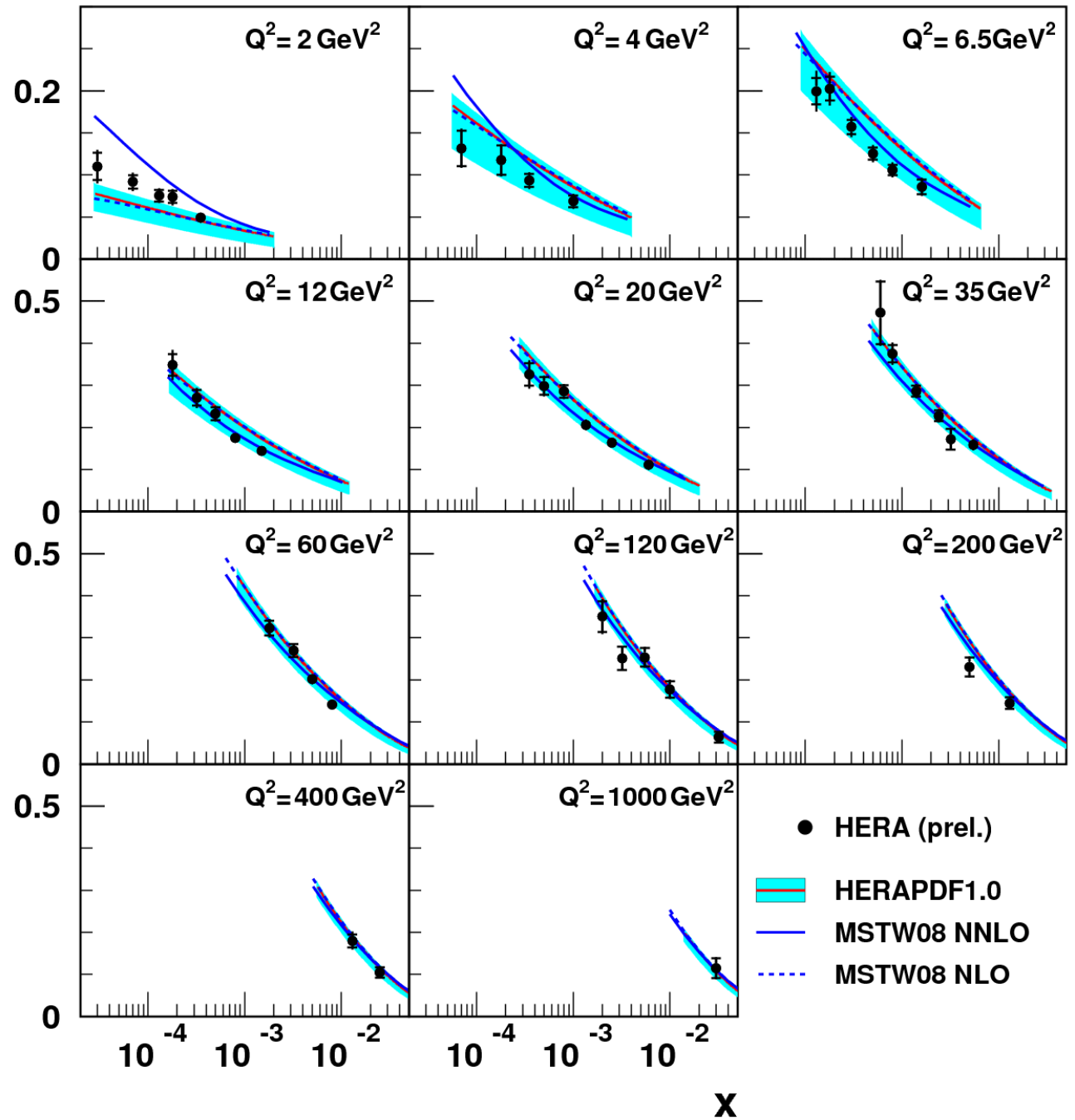
reasonably
described
by
NLO QCD

Summary and conclusions

- Use FFNS with \overline{MS} mass renormalization scheme to measure physical b and c masses? need NNLO?
- Uncertainty on W/Z standard candle cross sections can probably be reduced through empirical m_c fit to F_2^c in each scheme. Need further studies to check whether this is also true for other cross sections (might not).
- Different FFNS schemes exist. Variable flavour α_s evolution scheme preferable for global fits/LHC predictions?
- Consider to use \sim half natural scale for NLO cross section predictions (photo- and hadroproduction)

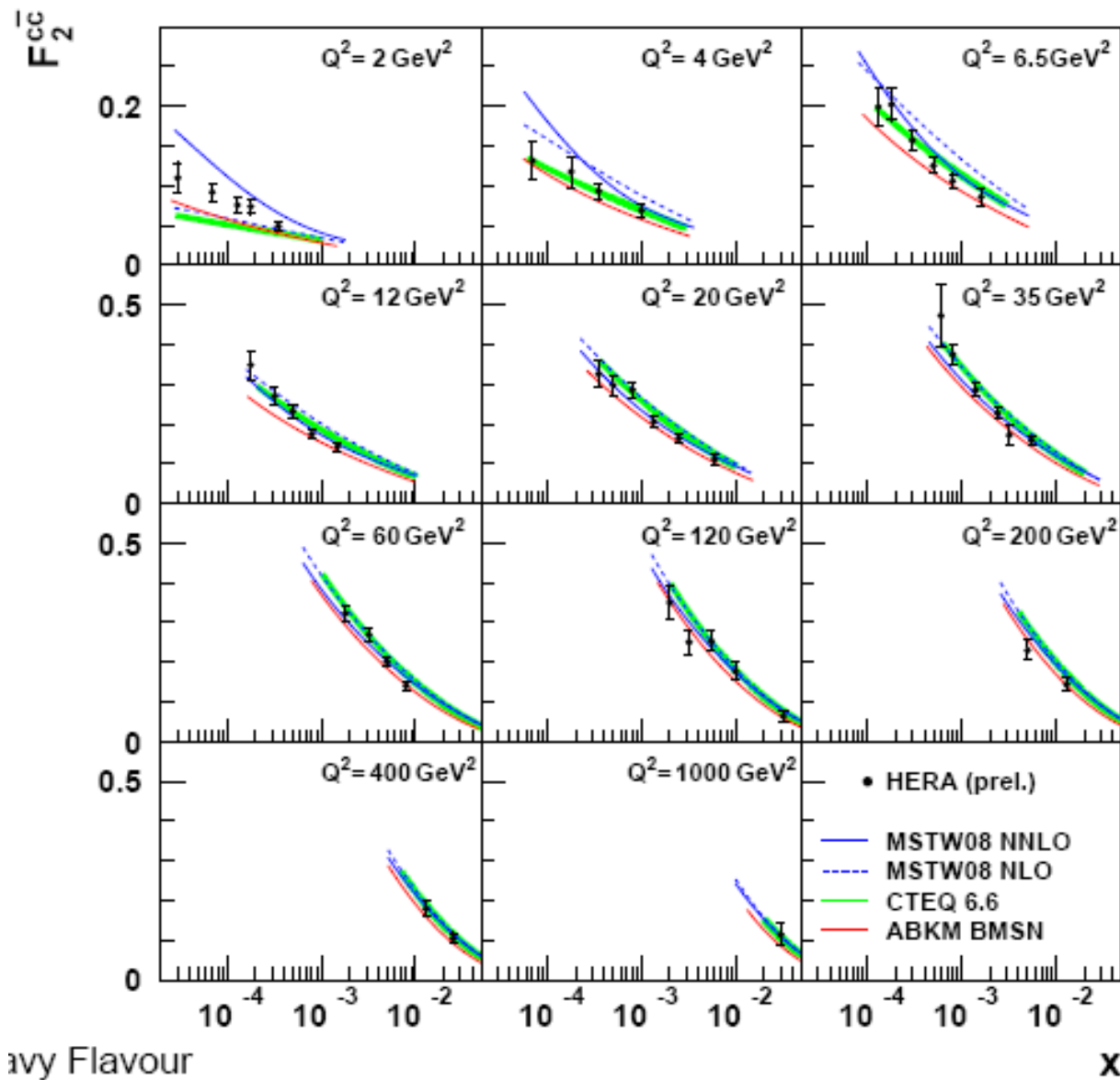
Backup slides



 $\overline{F_2^{cc}}$ 

October 2009

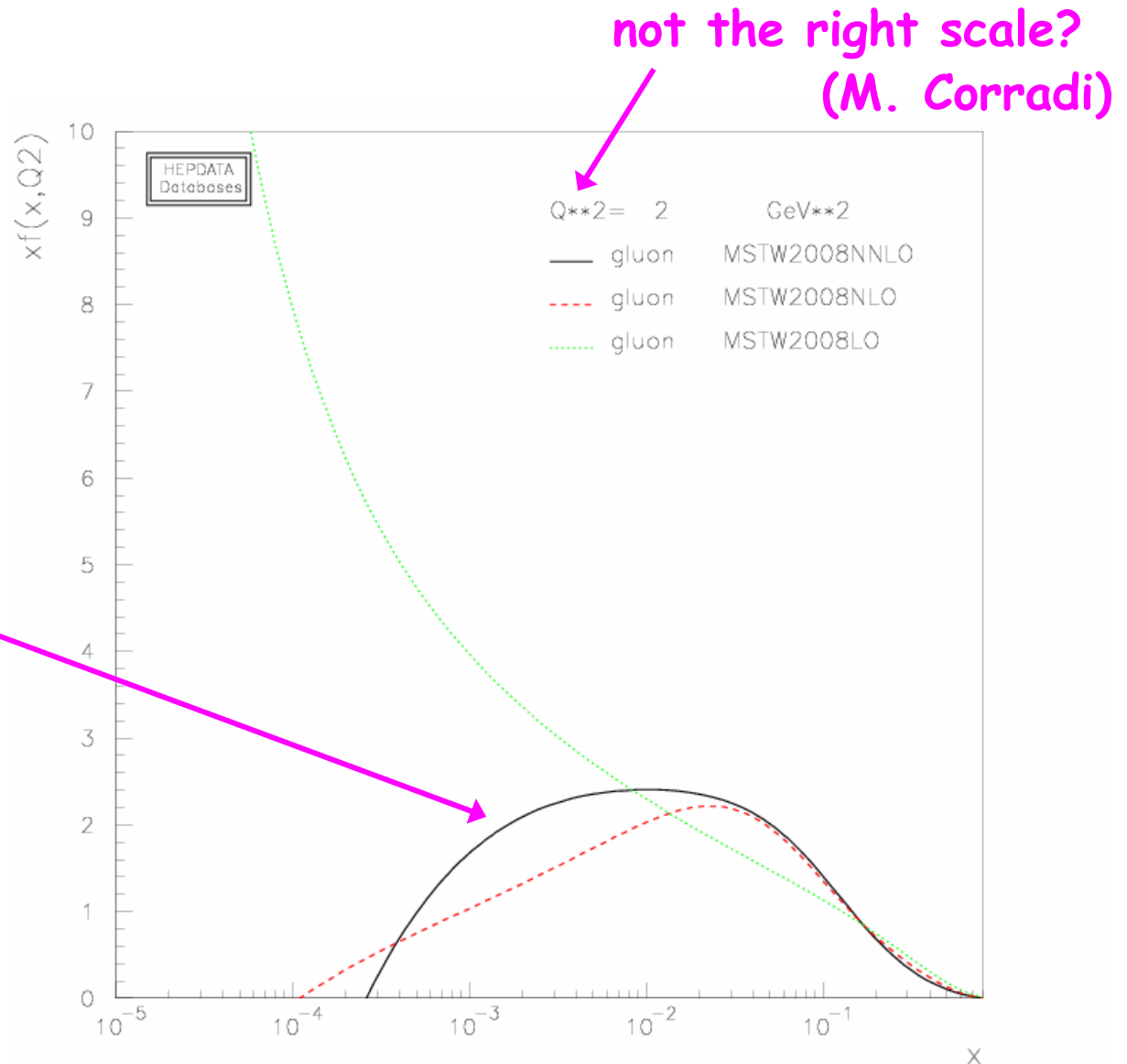
HERA Heavy Flavour Working Group



avy Flavour

Why so large ratio NNLO/NLO for MRST F2c ?

- NNLO only approximate, is it reliable?
- large contribution from change of gluon (courtesy O. Behnke / PDFplotter)



why does gluon change so much?