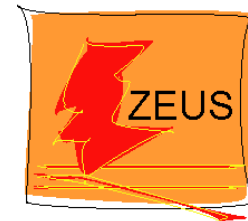


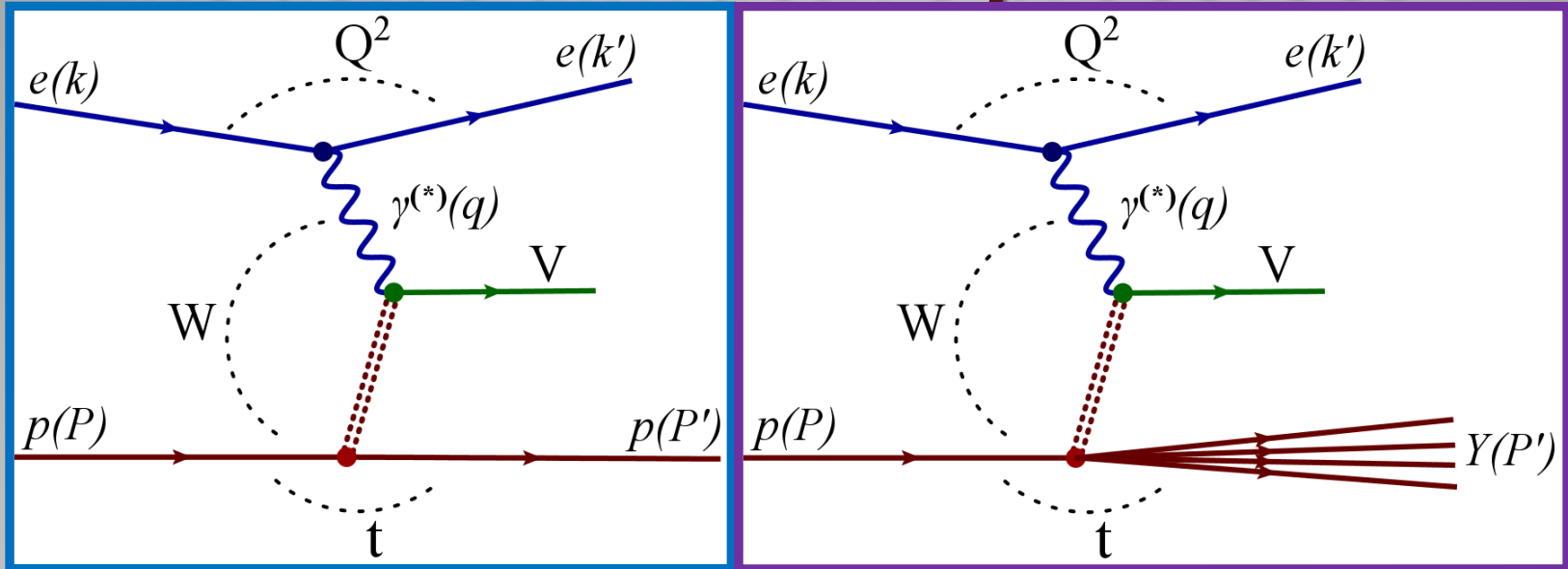
Vector Meson production and DVCS at HERA

On behalf of the H1 and ZEUS Collaborations



Janusz Tomasz Malka
University of Łódź

Exclusive diffractive processes



Diffraction - no quantum numbers are exchanged in the interaction btw γ and $p \rightarrow$ no colour flux \rightarrow large rapidity gap

V - Vector Meson ($\rho, \omega, \phi, J/\psi, \psi', \Upsilon$) photon (γ)

Q^2 - photon virtuality $Q^2 = -q^2 = -(k-k')^2$

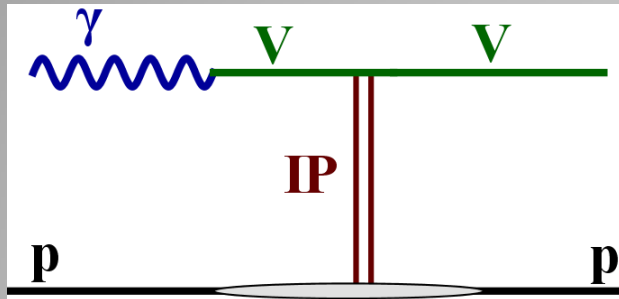
$(Q^2 \approx 0$ - photoproduction, $Q^2 > 0$ - electroproduction)

W - c.m. energy of γp system $W^2 = (q+p)^2$

t - (4-mom. transfer)² at p-vertex $t = (P - P')^2$

The proton can stay intact $p(P')$ or dissociate $Y(P')$

VDM & Regge theory (Soft physics)



Vector Dominance Model (VDM):

- The photon fluctuates into a vector meson, V , which carries the same quantum numbers as the photon ($\gamma p \rightarrow Vp$)
- The vector meson scatters elastically off the incoming proton ($Vp \rightarrow Vp$)

Predictions:

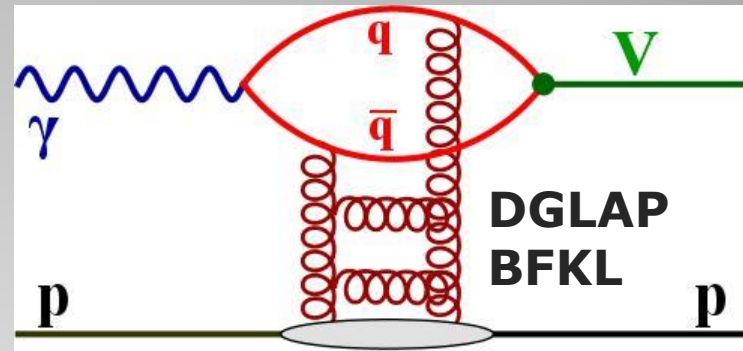
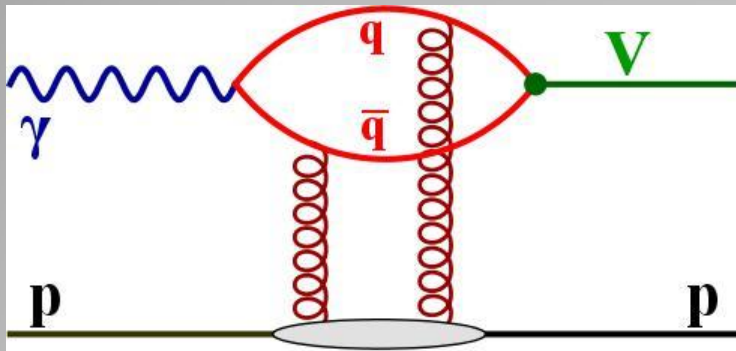
$$\frac{d\sigma(\gamma p \rightarrow Vp)}{dt} \propto e^{-b_0 t} \left(\frac{W^2}{W_0^2} \right)^{2(\alpha(t)-1)}$$

Experimental observations:

- $\alpha(t) = \alpha(0) + \alpha' t$, $\alpha(0) = 1.08$, $\alpha' = 0.25$; (**Donnachie-Landschof**)
- **Shrinkage of the diffractive peak** $b(W) = b_0 + 4\alpha' \ln \left(\frac{W}{W_0} \right)$
- **Weak energy dependence of cross section**

$$\sigma \propto W^\delta, \quad \delta \approx 4 \left[\alpha(0) - \frac{\alpha'}{b_0} - 1 \right], \quad \delta \sim 0.2;$$

pQCD models (Hard physics)



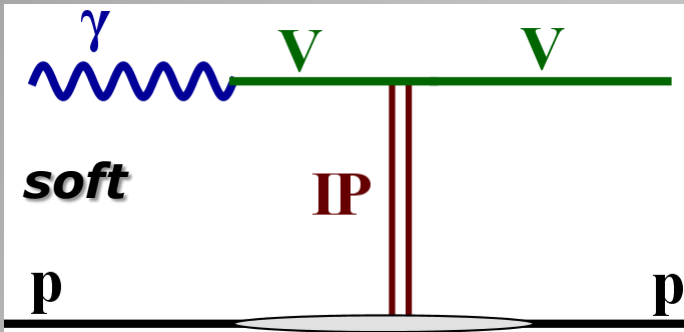
- the photon fluctuates into a $q\bar{q}$ state
- the $q\bar{q}$ pair scatters off the proton target
- the scattered $q\bar{q}$ pair turns into a vector meson

Predictions:

- $\sigma \propto [xg(x,\mu^2)]^2$ where $\mu^2 = f(Q^2, M_V^2, t)$
- **Increasing W corresponds to going to small x**
- **fast increase of the $\gamma^*p \rightarrow Vp$ cross section with energy W**
- **universal exponential t dependence**
- $b \sim 4 - 5 \text{ GeV}^{-2} \Rightarrow \alpha' \approx 0$

$$x \approx \frac{Q^2}{W^2} \text{ at low } x$$

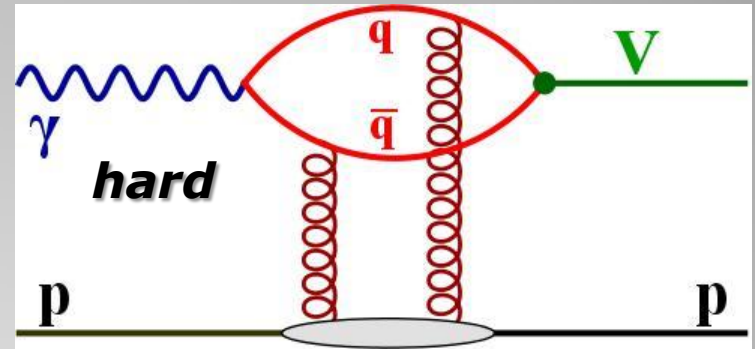
Transition from soft to hard physics



$$\delta \approx 0.2$$

$$b \approx 10 \text{GeV}^{-2}$$

$$\alpha' \approx 0.25 \text{GeV}^{-2}$$



$$\sigma \propto W^\delta$$

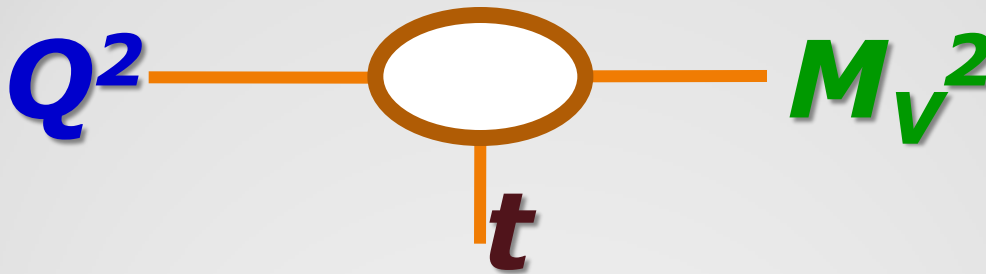
$$\frac{d\sigma}{d|t|} \propto e^{-b|t|}$$

$$\alpha(t), b(W)$$

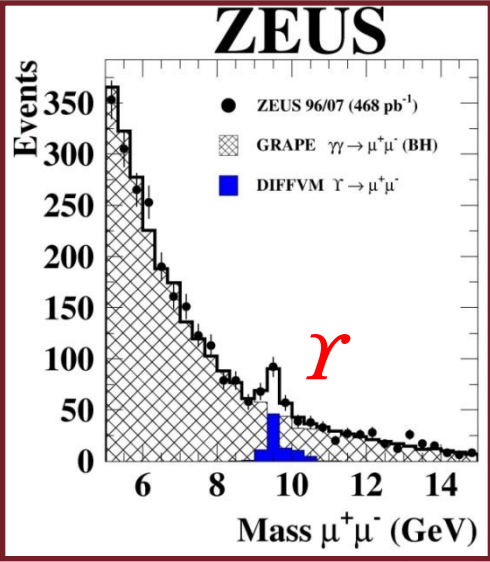
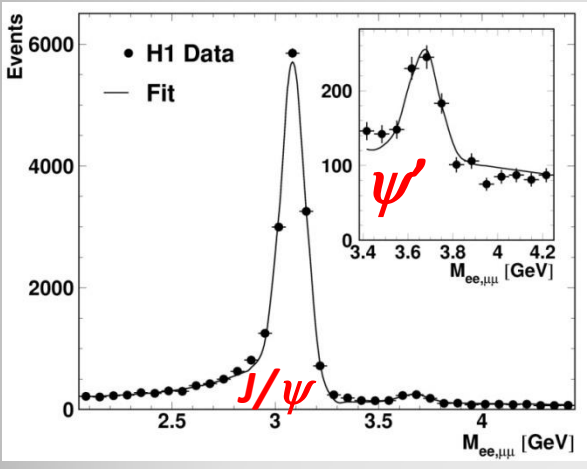
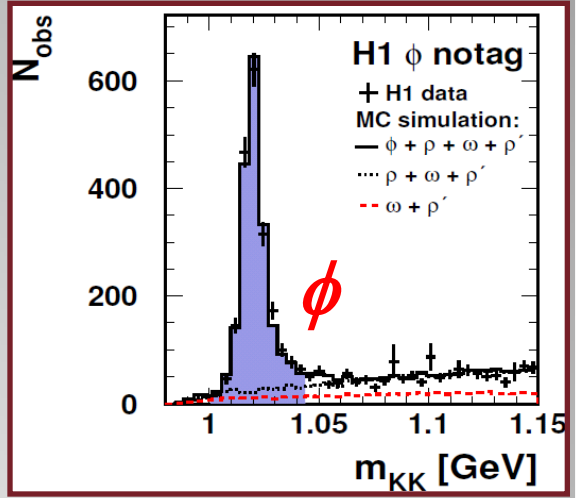
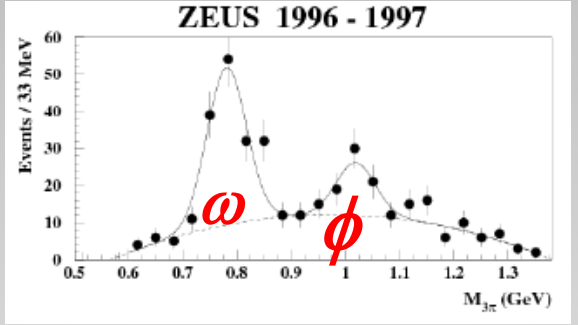
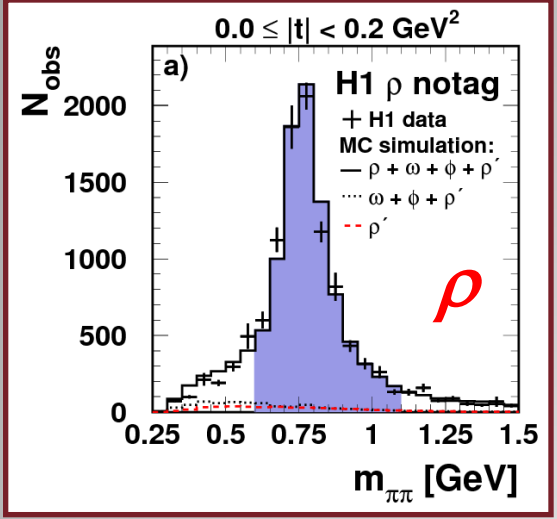
$$\delta \leq 0.7$$

$$b \approx 4-5 \text{GeV}^{-2}$$

$$\alpha' \approx 0$$

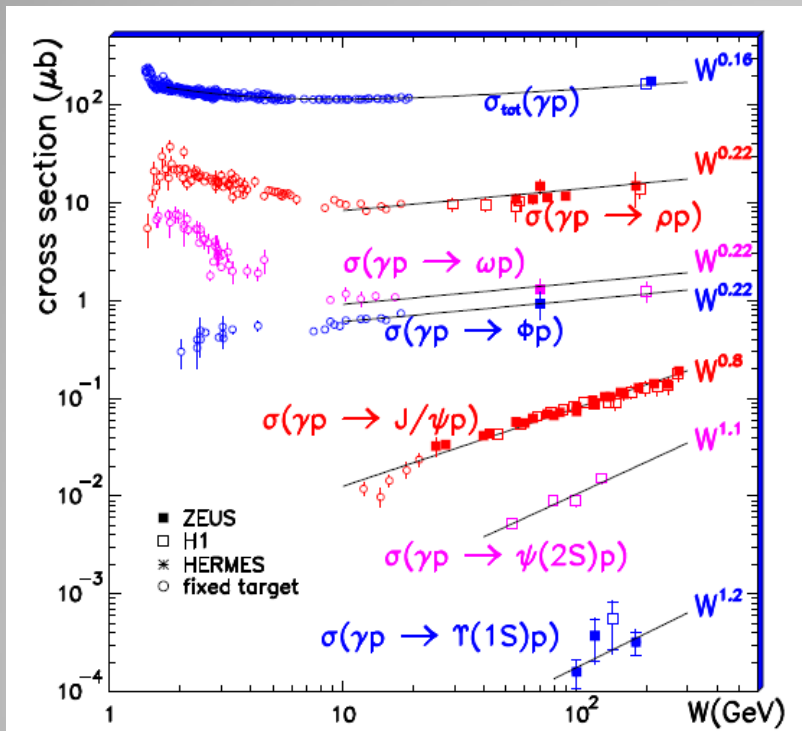


Vector mesons mass distributions



W dependence of the photoproduction of VM

The parameterisation of the cross section: $\sigma \propto W^\delta$



**Process becomes hard
(steeper W dependence)
as mass of VM becomes larger**

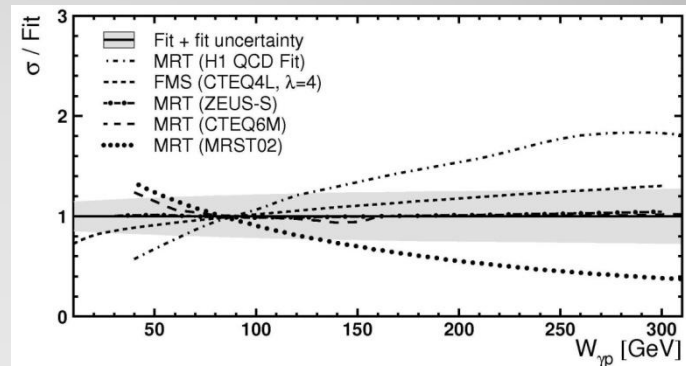
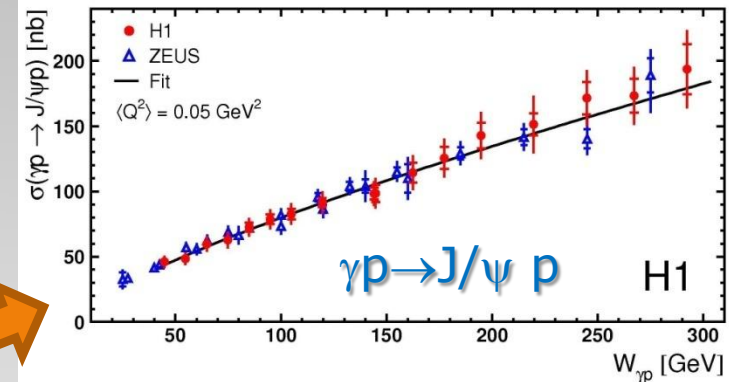
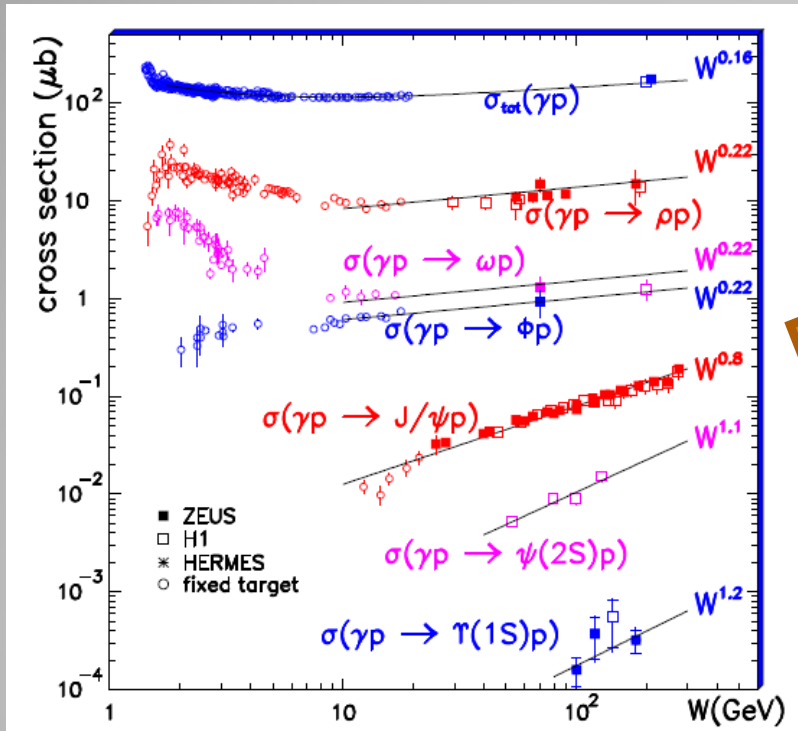


Apply pQCD models

W dependence of the photoproduction of VM

The parameterisation of the cross section:

$$\sigma \propto W^\delta$$

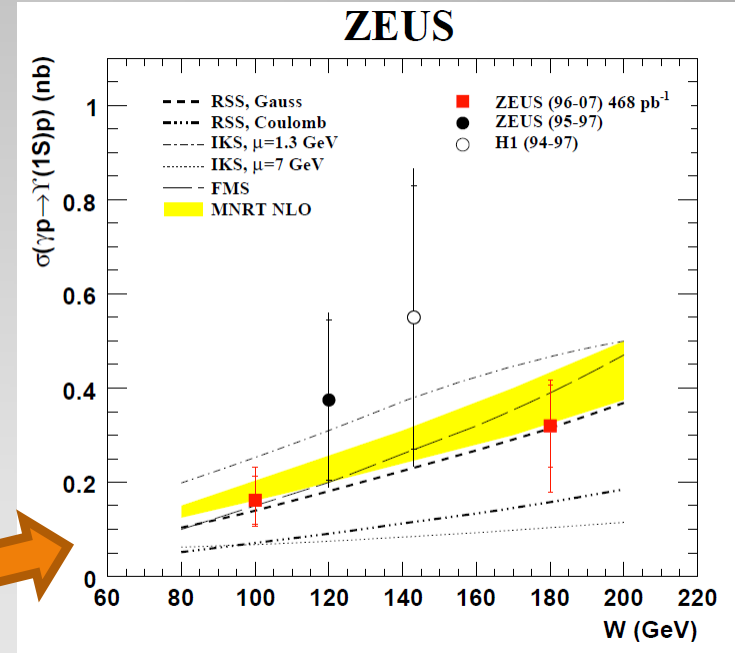
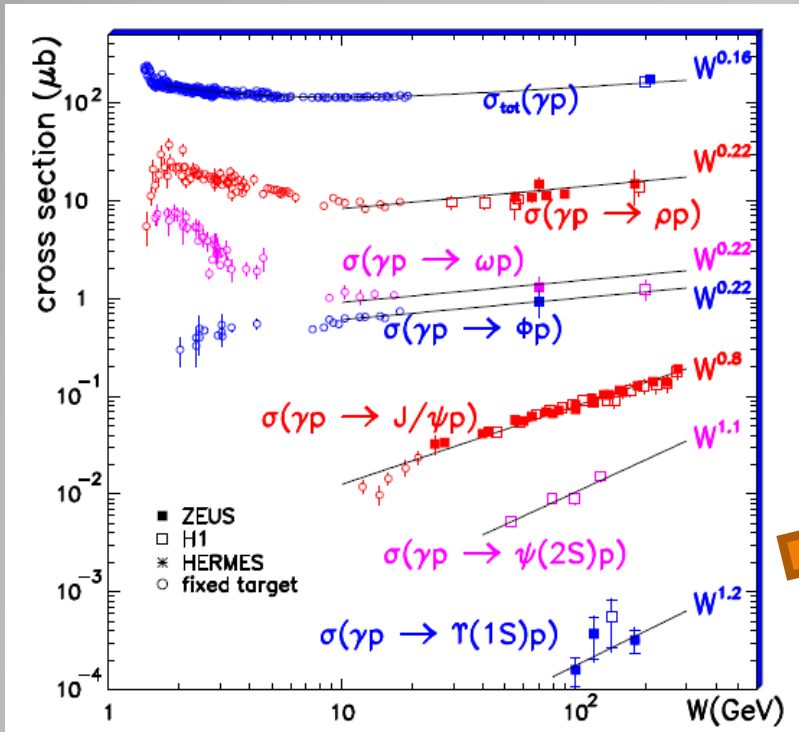


FMS – L. Frankfurt, M. McDermott and M. Strikman, JHEP **0103** (2001)
MRT – A. D. Martin, M. G. Ryskin and T. Teubner, Phys. Rev. D **62** (2000) 014022

Sensitivity to gluon density

W dependence of the photoproduction of VM

The parametrisation of the cross section: $\sigma \propto W^\delta$



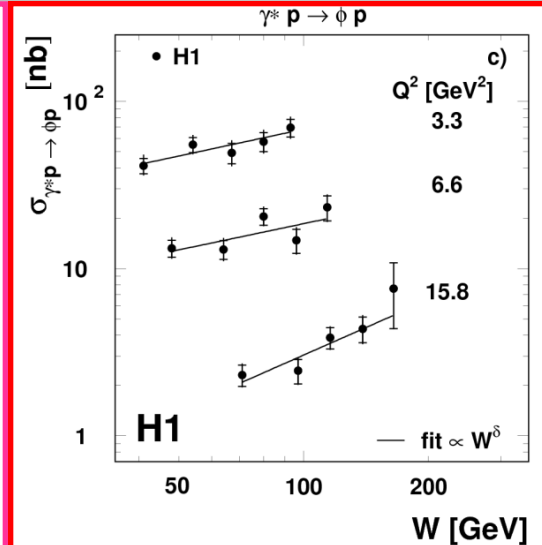
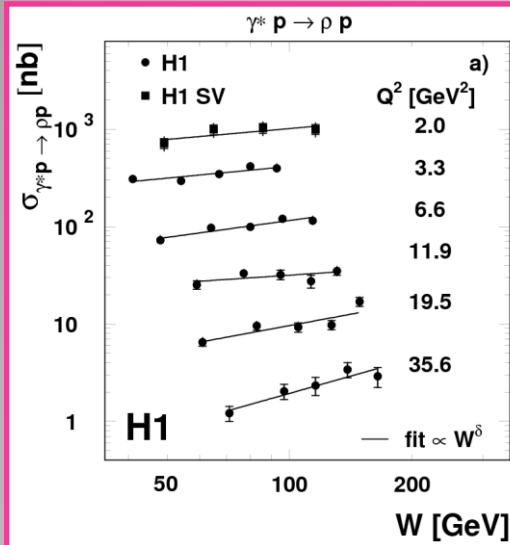
RSS (k_T) – A. Rybarska, W. Schäfer, A. Szczurek, *Phys. Lett. B* 668 (2008), p. 126.

IKS(NLO) – D.Yu. Ivanov, G. Krasnikov, L. Szymanowski, *Nucl. Phys. B (Proc. Suppl.)* 146 (2005), p. 134.

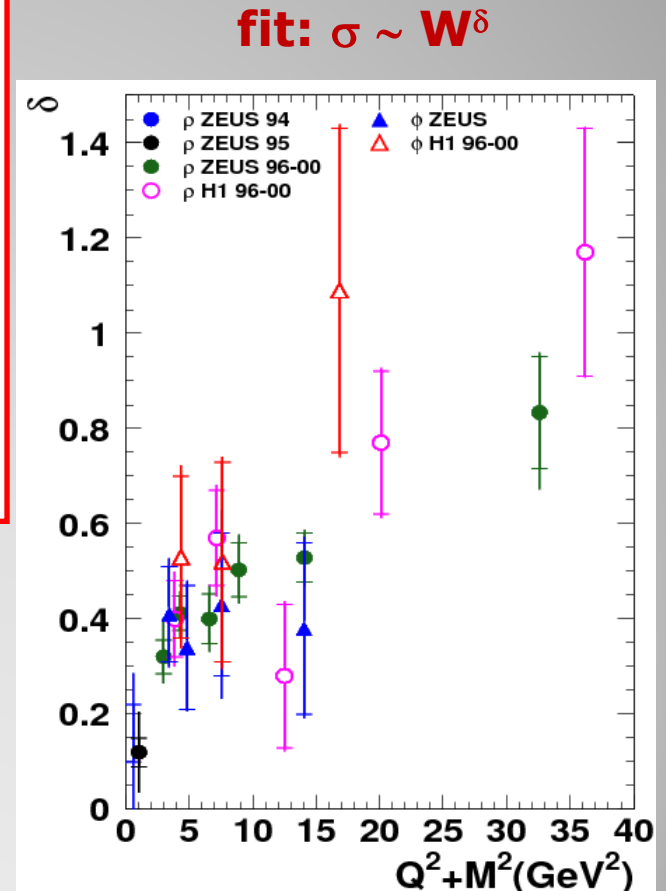
FMS(CTEQ4L) – L.L. Frankfurt, M.F. McDermott, M. Strikman, *JHEP* 9902 (1999), p. 002.

MNRT(HERA J/ψ) – A.D. Martin, C. Nockles, M. Ryskin, T. Teubner, *Phys. Lett. B* 662 (2008), p. 252.

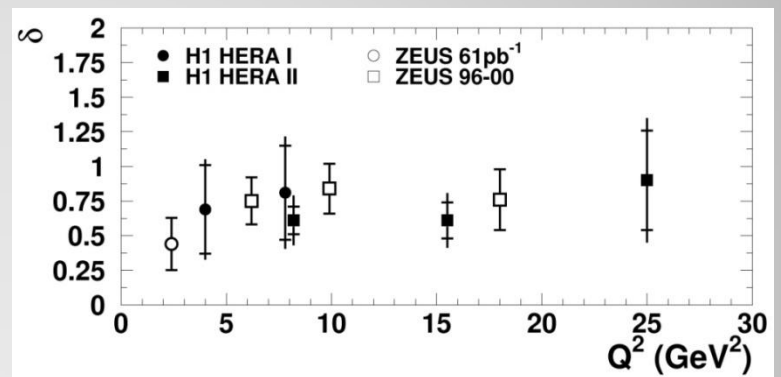
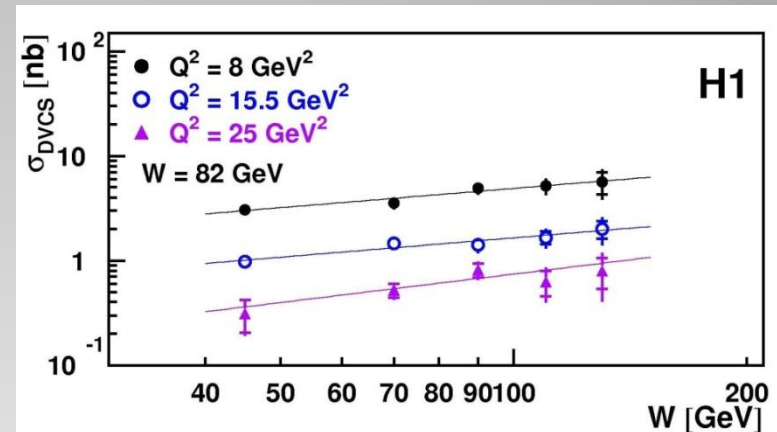
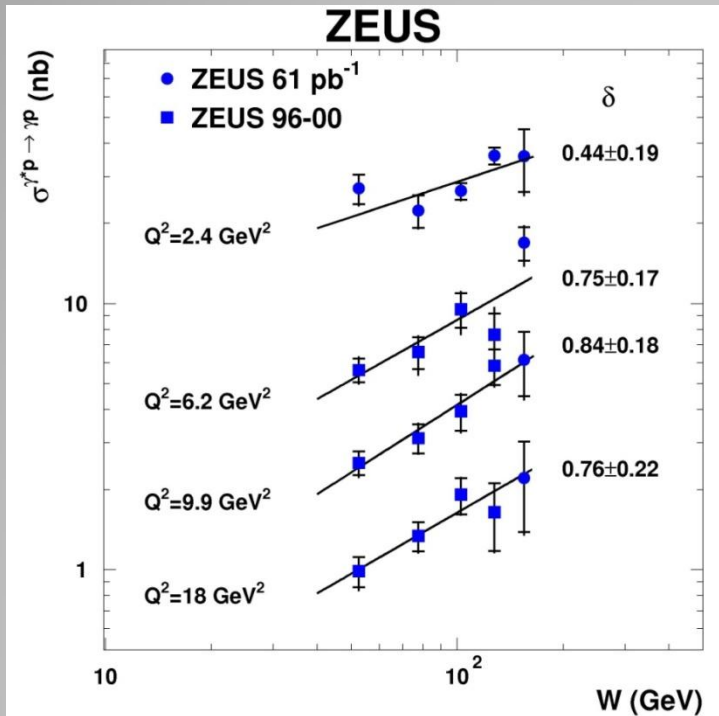
VM electroproduction (W dependence)



cross section W dependence becomes steeper at high Q^2 for ρ, ϕ from soft to hard regime

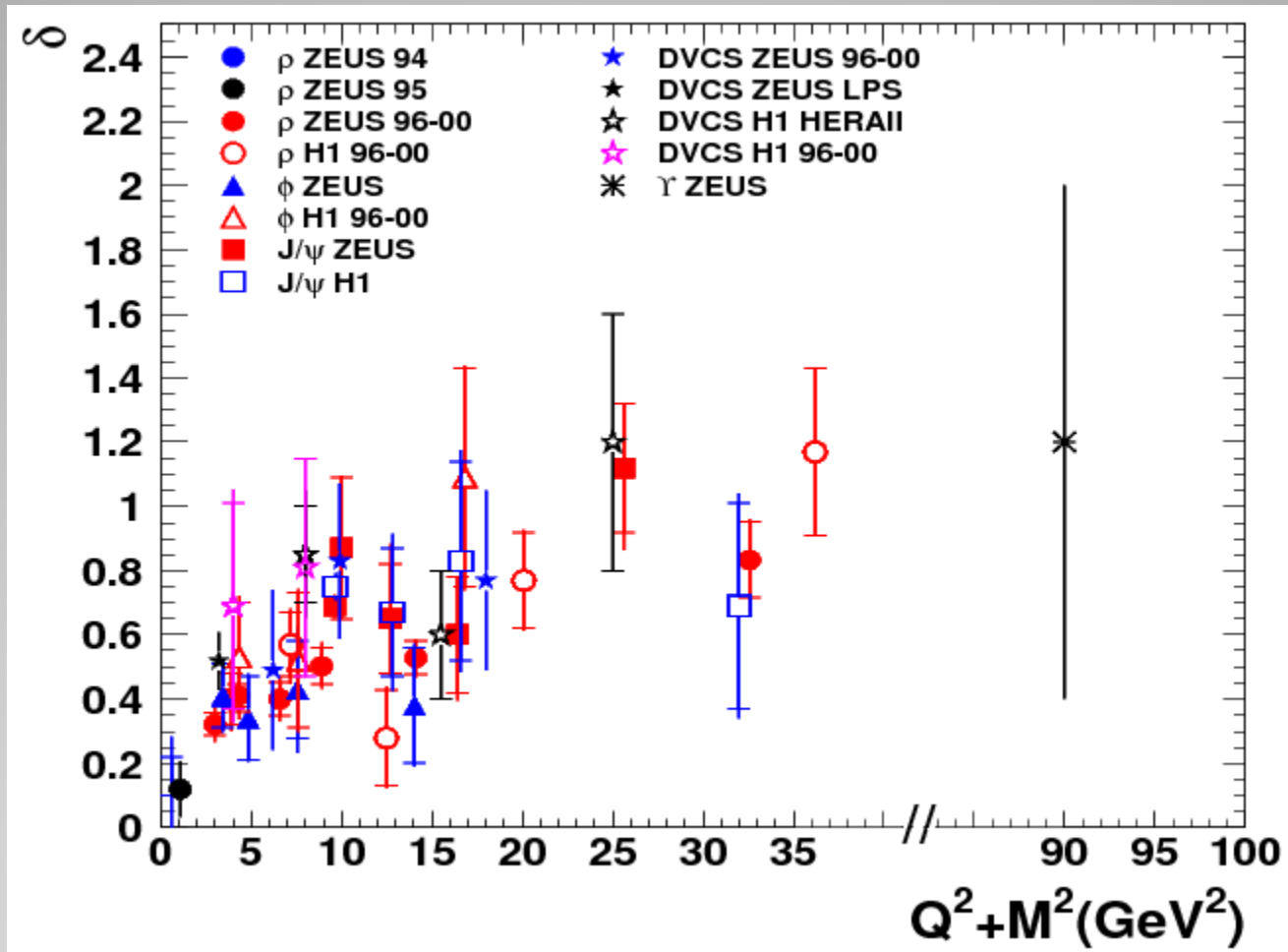


Deeply Virtual Compton Scattering (W dependence)



- fit: $\sigma \sim W^\delta$
- No δ dependence on Q^2 is observed
- Hard regime

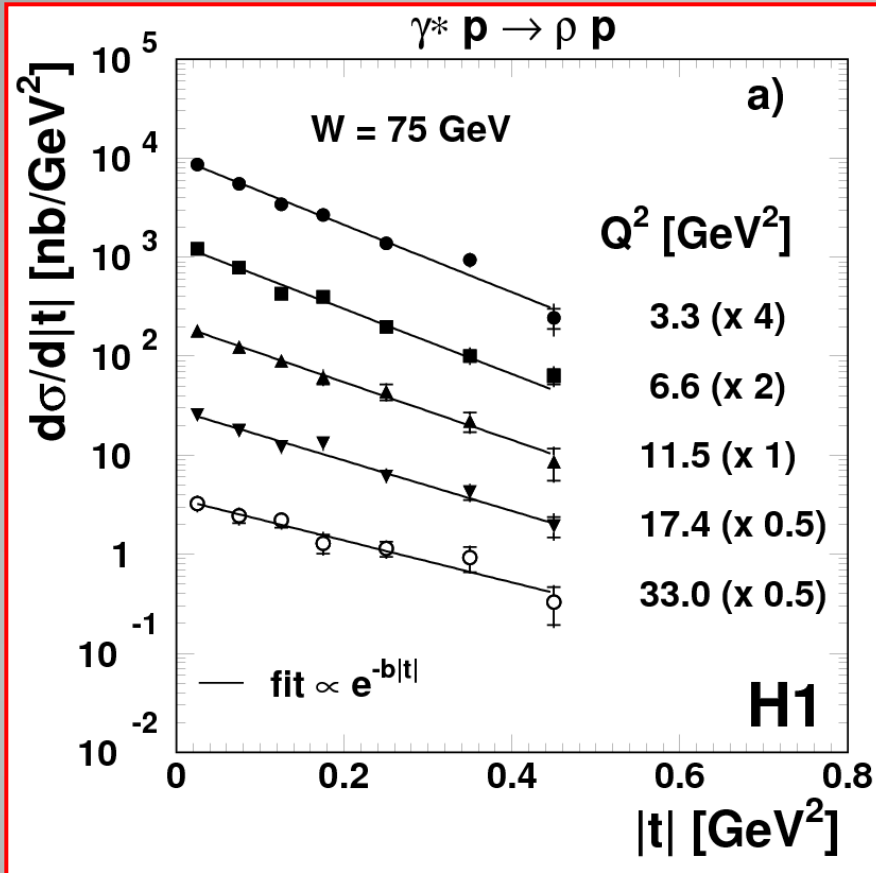
Summary (W dependence)



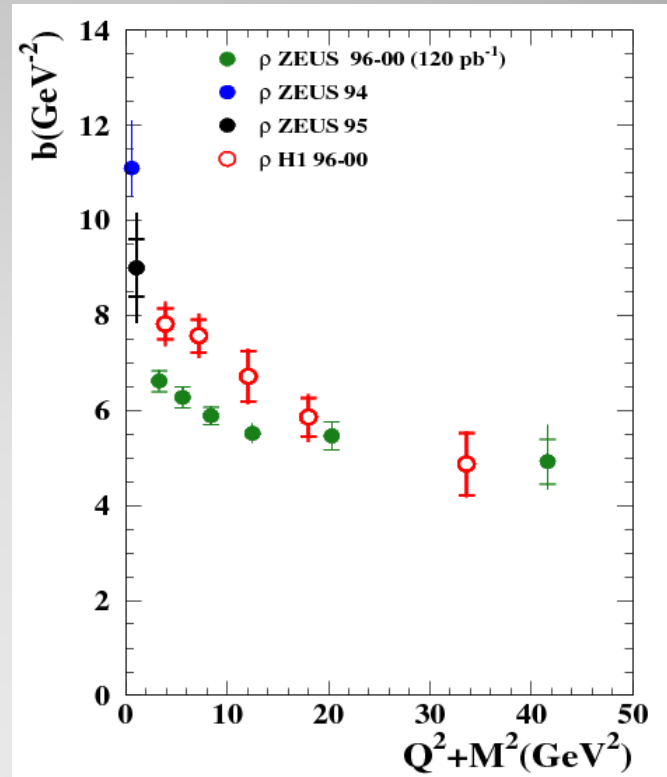
• process becomes hard as scale ($Q^2 + M^2$) becomes larger

$|t|$ dependence (ρ)

$d\sigma/d|t| \sim \exp(-b|t|)$ in bins of Q^2



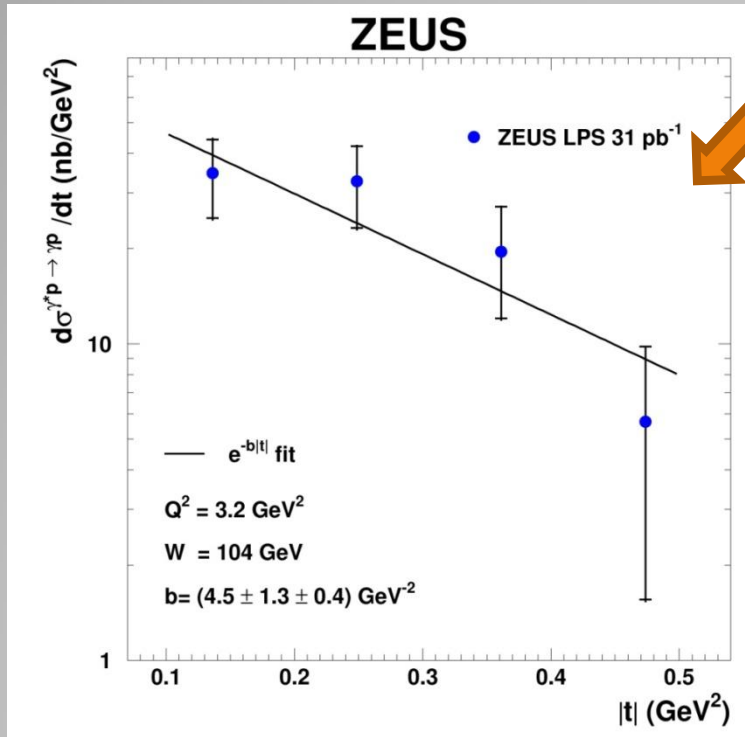
b describes the transverse size of the interaction region $b \propto R^2$



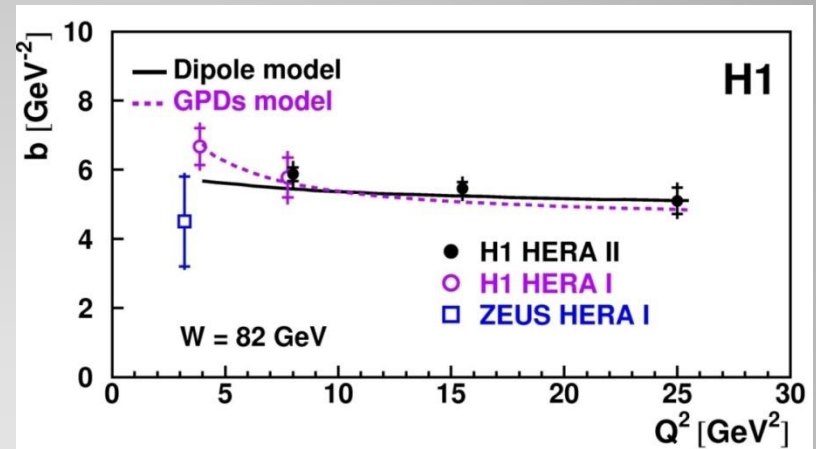
Transverse size of dipole decreases with Q^2

b decreases from soft values to pQCD expected values ($\sim 4-5$ GeV²)

$|t|$ dependence (DVCS)



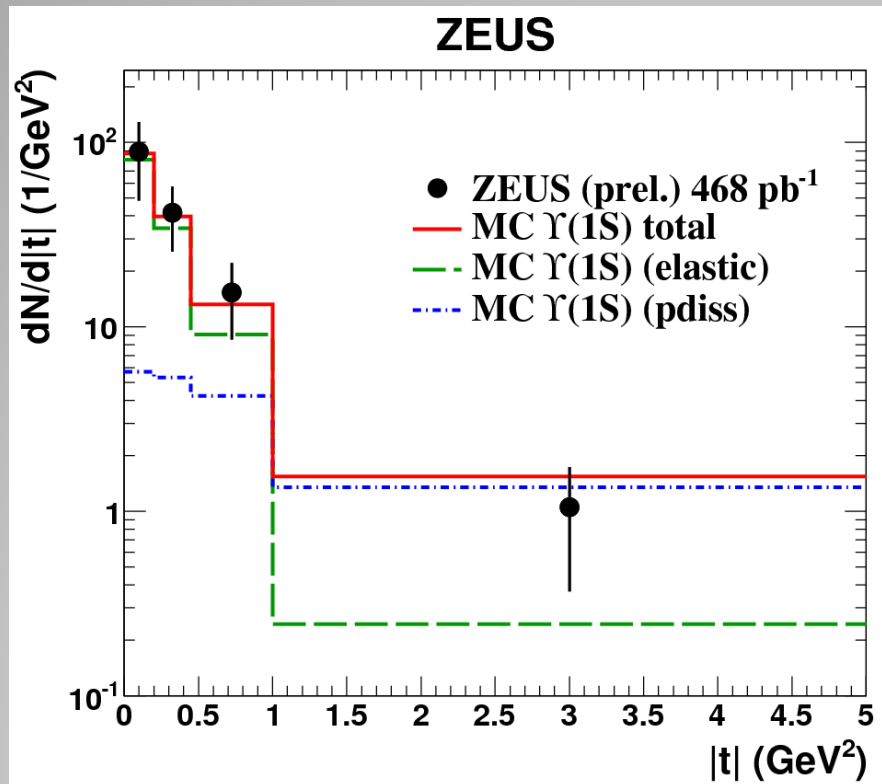
direct t measurement using Leading Proton Spectrometer



H1: $b = 5.41 \pm 0.14 \pm 0.31$ GeV⁻²
ZEUS: $b = 4.5 \pm 1.3 \pm 0.4$ GeV⁻²

- DVCS point don't show Q^2 dependence

$|t|$ dependence ($Upsilon$) photoproduction



2 times far from present results
at $Q^2 + M^2 = 89.5 \text{ GeV}^2$

Elastic: $b = 4.3^{+1.7}_{-1.1} \quad +0.5_{-0.5} [\text{GeV}^{-2}]$

$|t|$ dependence

$$\frac{d\sigma}{dt} \sim e^{-b|t|}$$

transverse size:

$$b = b_V + b_p$$

Vector Meson
transverse size

$$b_V \sim \frac{1}{Q^2 + M_V^2}$$

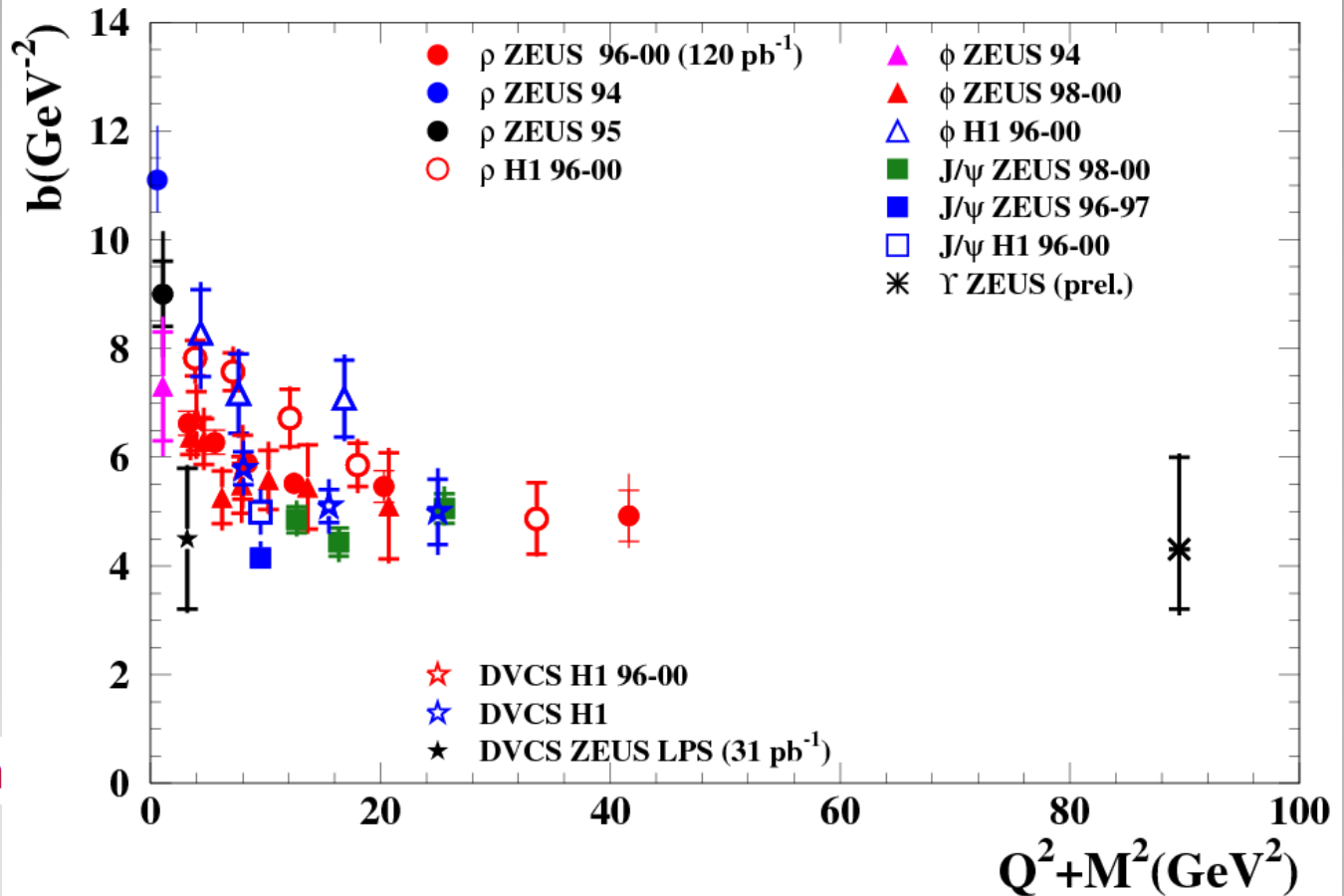
Target size:

$$b_p \approx 5 \text{ GeV}^{-2}$$

corresponds to

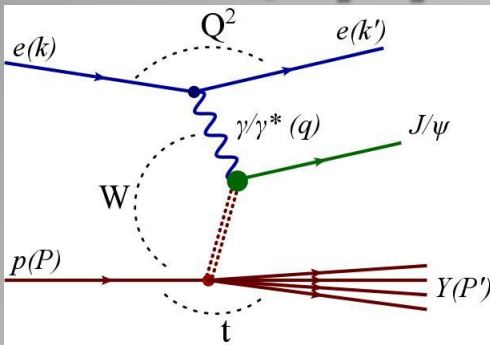
$$r_{\text{gluons}} \approx 0.5 \text{ fm}$$

$$r_{\text{em}} \approx 0.8 \text{ fm}$$



**b decreases with from $\sim 10 \text{ GeV}^{-2}$ (soft) to $\sim 5 \text{ GeV}^{-2}$ (hard)
size of scattered VM getting smaller with scale**

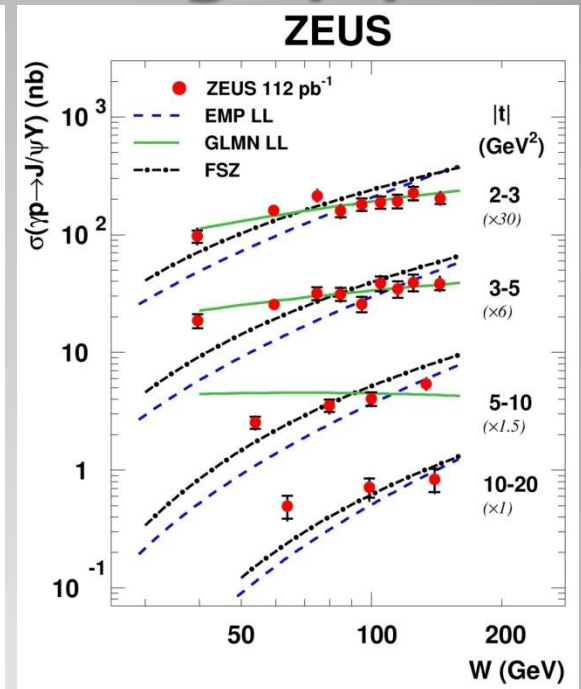
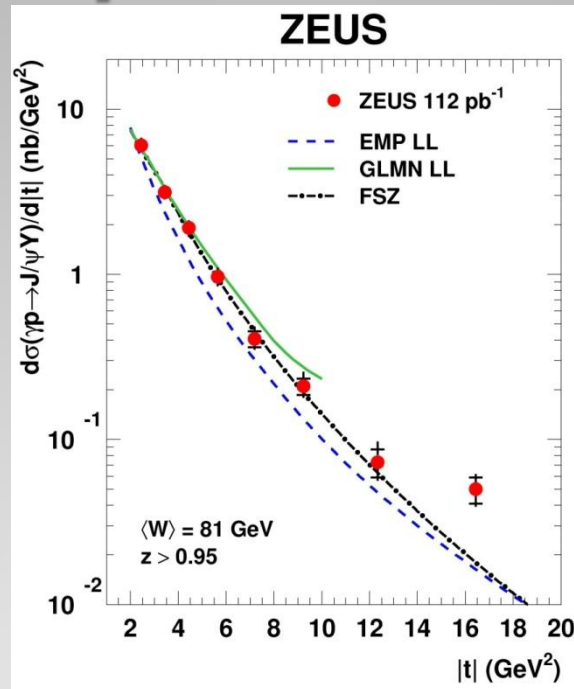
J/ψ photoproduction at high $|t|$



Hard scale provided by $|t|$ and mass
 t dependence no longer exponential

$$\frac{d\sigma}{d|t|} \propto t^n$$

$n = -1.9 \pm 0.1$, $2 < |t| < 4 \text{ GeV}^2$
 $n = -3.0 \pm 0.1$, $4 < |t| < 16 \text{ GeV}^2$



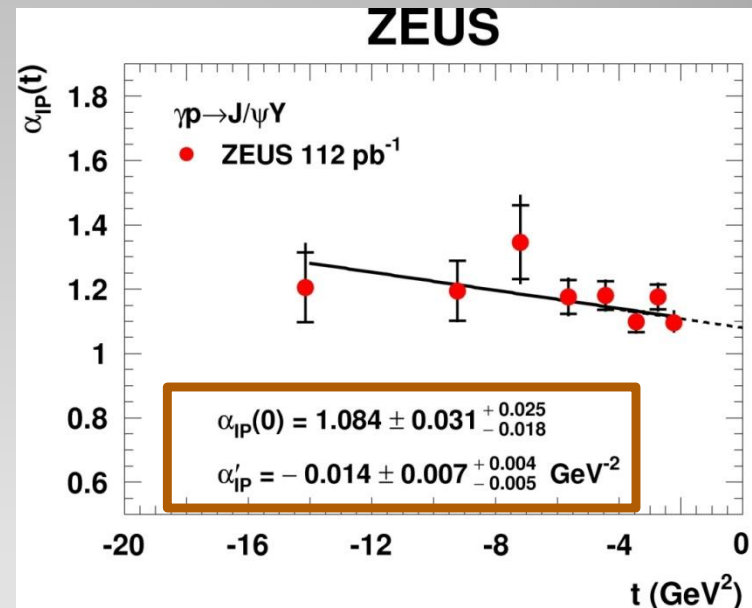
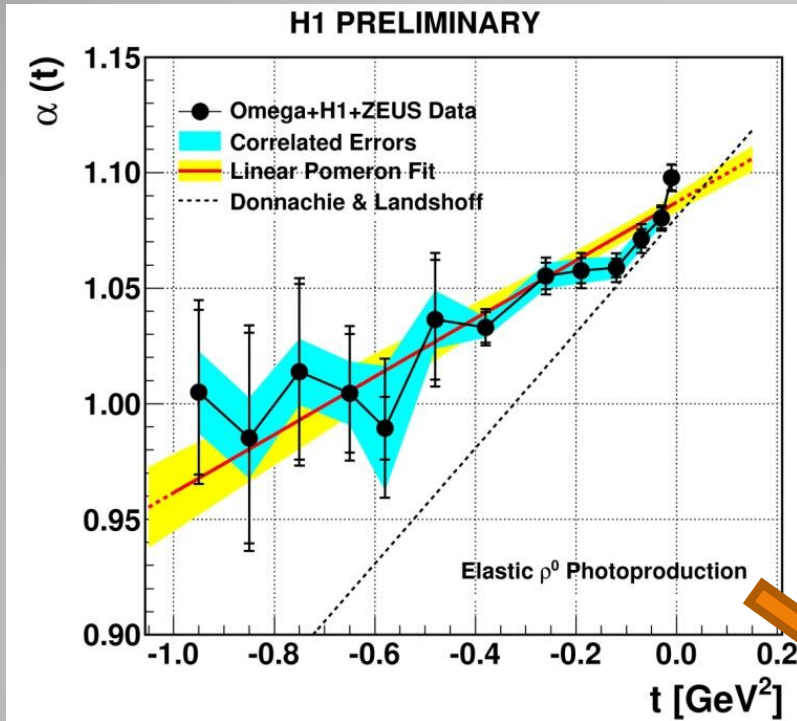
σ vs W in t ranges: data rise with W for all t
EMP (BFKL) below data
GLMN (DGLAP) fails at $|t| > 5 \text{ GeV}^2$
FSZ gives good description up $|t| = 12 \text{ GeV}^2$
BFKL LL and FSZ are too steep (W dependence)

None of the models describes the data over the full phase space

effective Pomeron trajectory

$$\frac{d\sigma}{d|t|} \propto \left(\frac{W}{W_0}\right)^{4(\alpha(t)-1)}$$

$$\alpha(t) = \alpha(0) + \alpha't$$



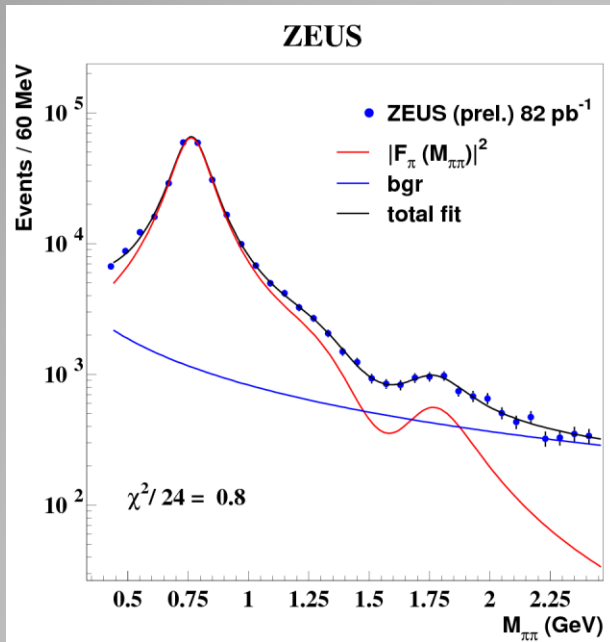
$$DL: \alpha(t) = 1.08 + 0.25t$$

$$\alpha(0) = 1.0871 \pm 0.0026 \pm 0.0030$$

$$\alpha' = 0.126 \pm 0.013 \pm 0.012 \text{ GeV}^{-2}$$

$\alpha(0)$ consistent with 1.08 from soft pp scattering
 $\alpha' \sim$ twice smaller than 0.25 GeV^{-2}

Two pion diffractive electroproduction



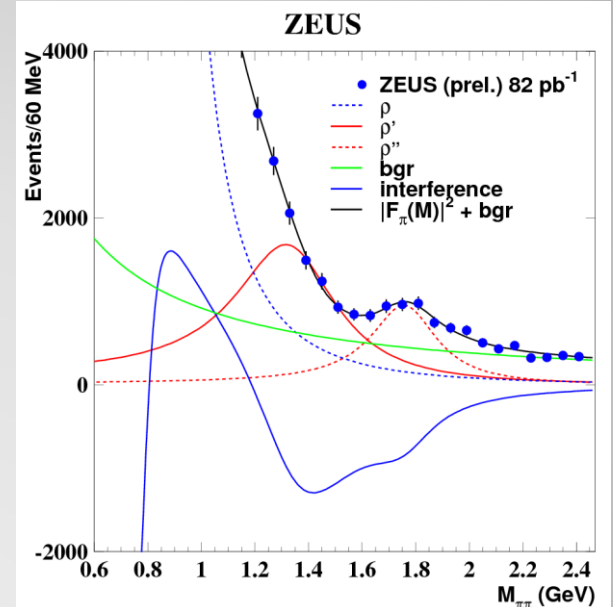
Measure two pion mass distribution
 $0.4 < M_{\pi\pi} < 2.4$ GeV in $2 < Q^2 < 80$ GeV²

$$\frac{dN(M_{\pi\pi})}{dM_{\pi\pi}} = N \left[|F_{\pi}(M_{\pi\pi})|^2 + \frac{B}{M_{\pi\pi}^n} \right]$$

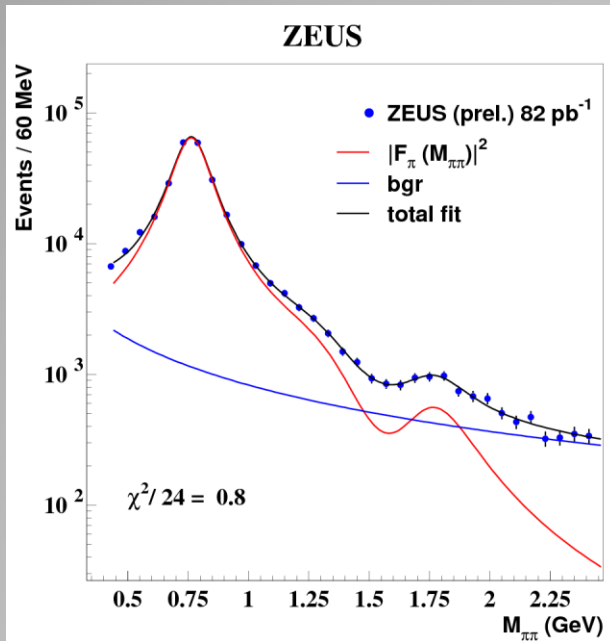
$$F_{\pi}(M_{\pi\pi}) = \frac{BW(\rho) + \beta BW(\rho') + \gamma BW(\rho'')}{1 + \beta + \gamma}$$

11 parameters fit:

$N, B, n, M(\rho), \Gamma(\rho), M(\rho'), \Gamma(\rho'), M(\rho''), \Gamma(\rho''), \beta, \gamma$,
 where N - total normalization factor, B, n - described
 background, M, Γ - masses and widths of vector mesons
 and β, γ - relative amplitudes



Two pion diffractive electroproduction



Measure two pion mass distribution
 $0.4 < M_{\pi\pi} < 2.4$ GeV in $2 < Q^2 < 80$ GeV²

$$\frac{dN(M_{\pi\pi})}{dM_{\pi\pi}} = N \left[|F_\pi(M_{\pi\pi})|^2 + \frac{B}{M_{\pi\pi}^n} \right]$$

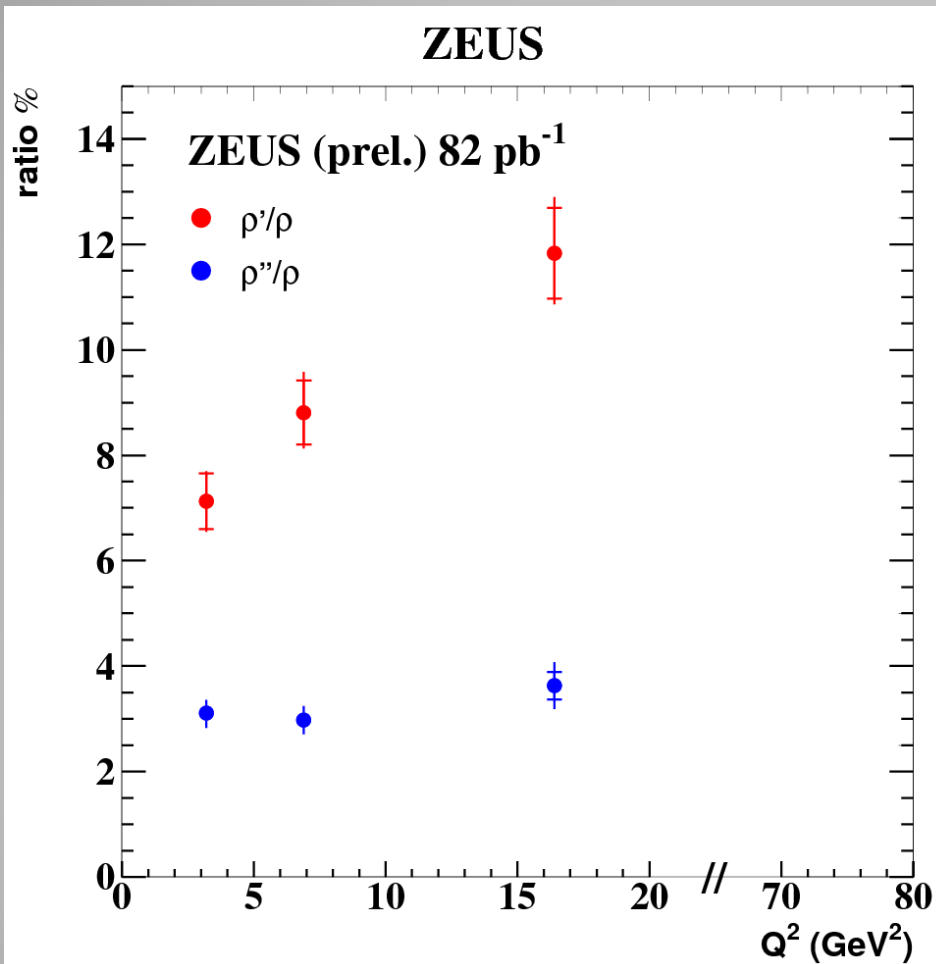
Par.	ZEUS (prel.)	PDG
M_ρ	$772 \pm 2^{+2}_{-1}$	775.49 ± 0.34
Γ_ρ	$155 \pm 5 \pm 2$	149.4 ± 1
$M_{\rho'}$	$1360 \pm 20^{+20}_{-30}$	1465 ± 25
$\Gamma_{\rho'}$	$460 \pm 30^{+40}_{-45}$	400 ± 60
β	$-0.27 \pm 0.02 \pm 0.02$	
$M_{\rho''}$	$1770 \pm 20^{+15}_{-20}$	1720 ± 20
$\Gamma_{\rho''}$	$310 \pm 30^{+25}_{-35}$	250 ± 100
γ	$0.10 \pm 0.02^{+0.02}_{-0.01}$	

$$F_\pi(M_{\pi\pi}) = \frac{BW(\rho) + \beta BW(\rho') + \gamma BW(\rho'')}{1 + \beta + \gamma}$$

11 parameters fit:

N , B , n , $M(\rho)$, $\Gamma(\rho)$, $M(\rho')$, $\Gamma(\rho')$, $M(\rho'')$, $\Gamma(\rho'')$, β , γ ,
 where N - total normalization factor, B , n - described
 background, M , Γ - masses and widths of vector mesons
 and β , γ - relative amplitudes

ρ, ρ', ρ'' electroproduction (cross section ratios)



ρ'/ρ —increases with Q^2

the **anomalous behaviour of ρ'/ρ production** ratio with Q^2 was predicted in works of J. Nemchik, B. Kopeliovich, N. Nikolaev, B. Zakharov, (see hep-ph/9605208)

ρ''/ρ —constant with Q^2

Summary

New measurements of ρ , ϕ and Υ at HERA.

The measurements allow us to study the transition from the soft to hard regime.

Effective Pomeron trajectory has smaller slope than that extracted from soft hadron-hadron scattering.

DVCS asymmetries and $|t|$ slope measurements provide access to GPDs.

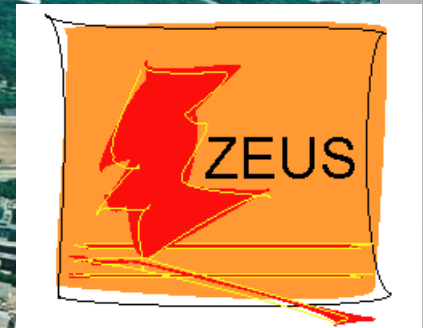
Measurement of ρ , ρ' , ρ'' production shows anomalous behaviour of ρ'/ρ cross section ratio

pQCD expectations are in general compatible with the data, but still a lot to be understood.

Backup

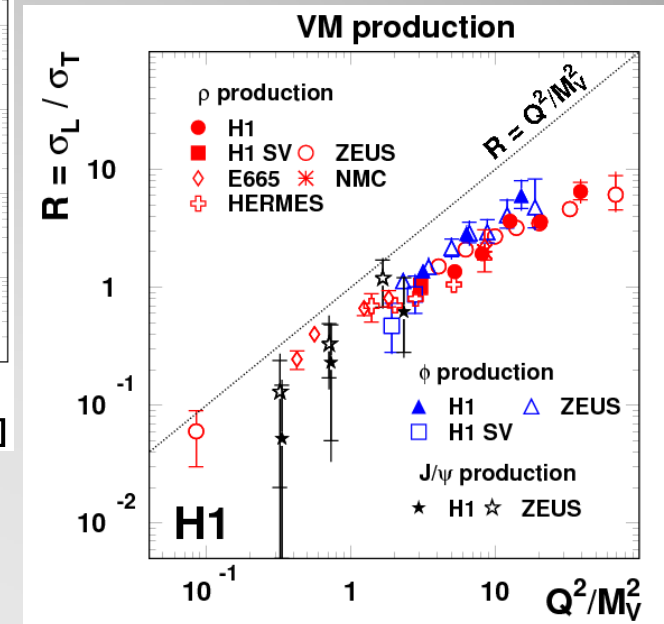
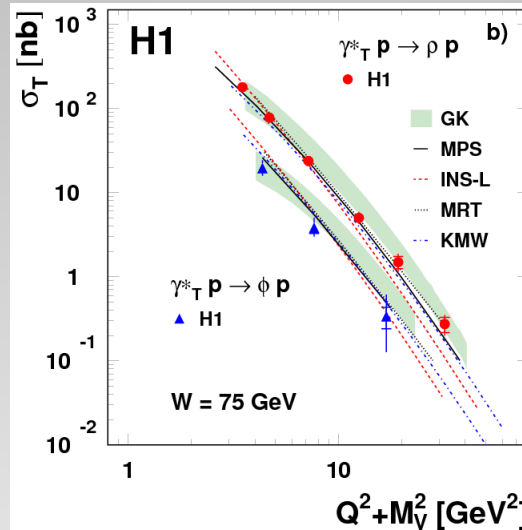
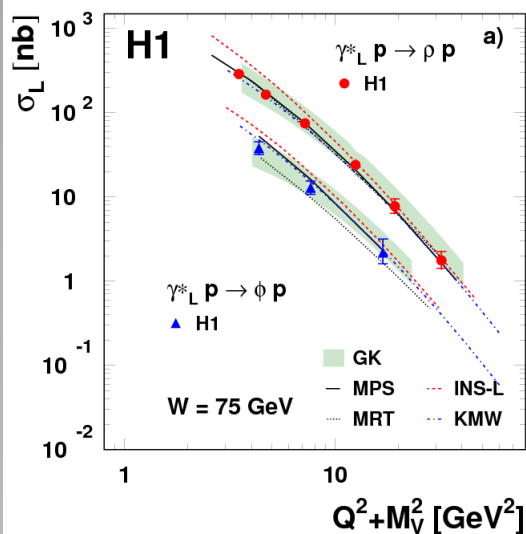
HERA experiments

27.5 GeV electrons/positrons with 920 GeV protons beam
CMS energy = 318 GeV



σ_L, σ_T cross sections

- Unique opportunity to extract $R = \sigma_L / \sigma_T$ $\sigma = \sigma_T + \epsilon \sigma_L$ $\epsilon \approx 0.996$
- R measured from angular distributions $f(\cos \theta_H, r^{04}_{00})$, in SCHC approximation



expectation: $R = \sigma_L / \sigma_T \sim Q^2 / M^2$
as the scale gets harder σ_L dominates

σ_L and σ_T have different $Q^2 + M^2$ dependence

Models based on pQCD describe well σ_L , but not σ_T

INS - I.P. Ivanov, N.N. Nikolaev and A.A. Savin, Phys.Part.Nucl. 37 (2006) 1

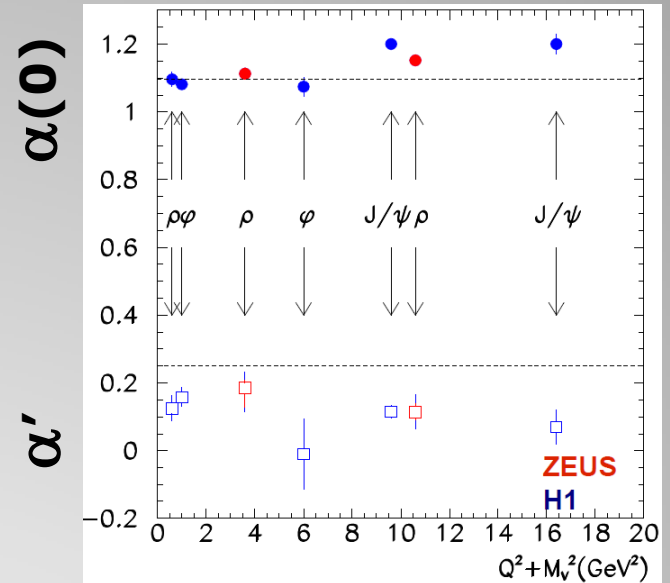
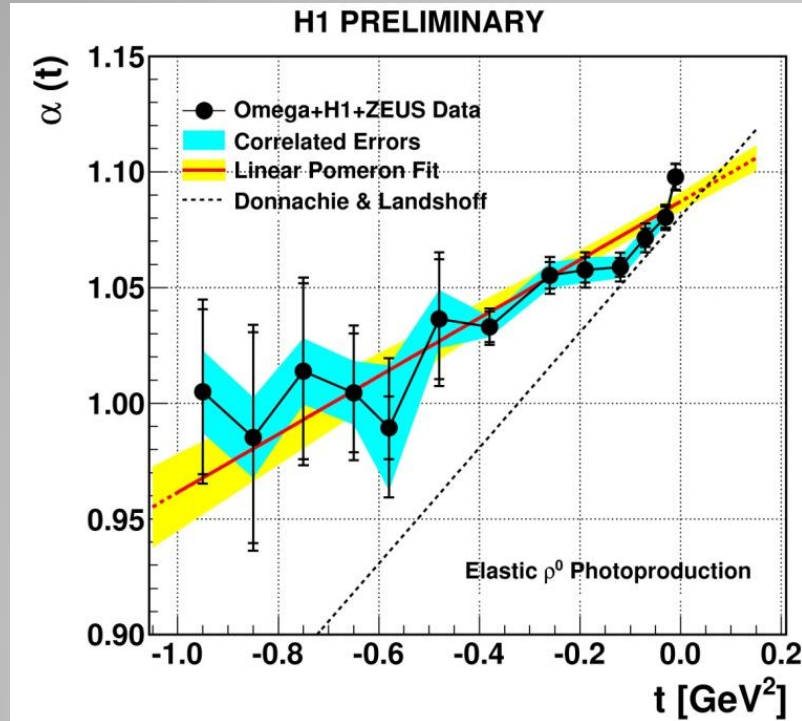
GK - S.V. Goloskokov and P. Kroll, arXiv:hep-ph/0708.3569 (2007)

MPS - C. Marquet, R. Peschanski and G. Soyez, Phys. Rev. D76 (2007) 034011

effective Pomeron trajectory

$$\frac{d\sigma}{d|t|} \propto \left(\frac{W}{W_0}\right)^{4(\alpha(t)-1)}$$

$$\alpha(t) = \alpha(0) + \alpha't$$



$$DL: \alpha(t) = 1.08 + 0.25t$$

$$\alpha(0) = 1.0871 \pm 0.0026 \pm 0.0030$$

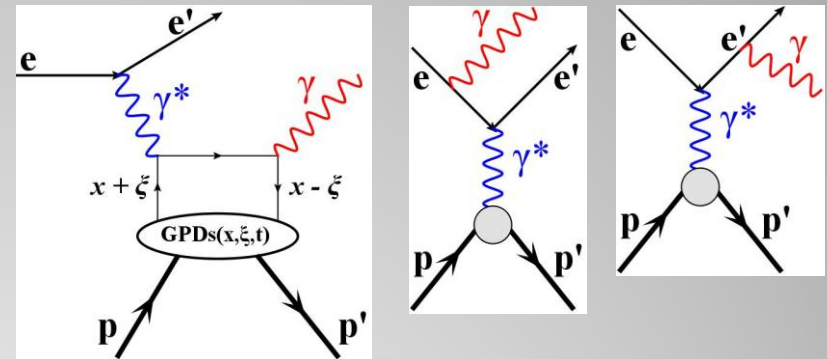
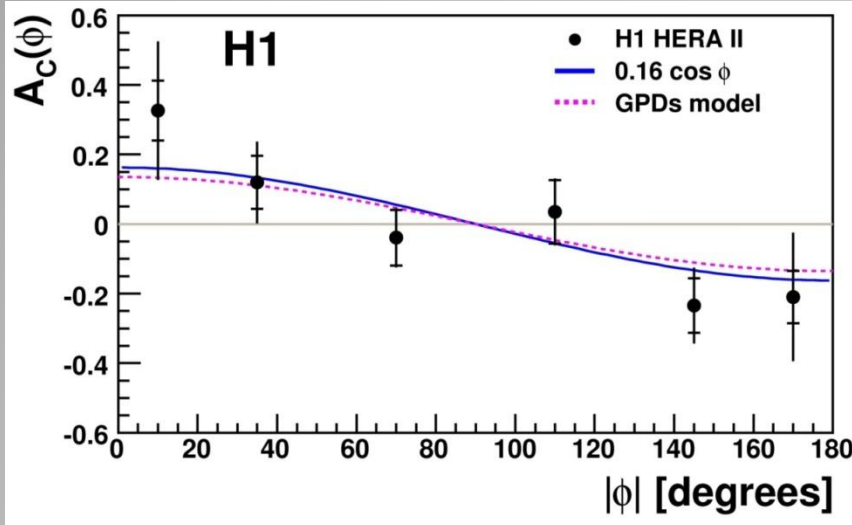
$$\alpha' = 0.126 \pm 0.013 \pm 0.012 \text{ GeV}^{-2}$$

$\alpha(0)$ consistent with 1.08 from soft pp scattering
 $\alpha' \sim$ twice smaller than 0.25 GeV^{-2}

DVCS – Beam Charge Asymmetry

DVCS and Bethe-Heitler have the same initial and final states and they are indistinguishable. The QCD-QED interference term is sensitive to the real part of the QCD amplitude. It changes sign with lepton beam charge:

$$\sigma = \sigma_{DVCS} + \sigma_{BH} \pm \sigma_{interf.}$$



$x_{1,2} = x \pm \xi$ – parton longitudinal momentum fraction, ξ – fraction of the momentum transfer

$$A_C = \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-} = p_1 \cos(\phi) + \dots$$

$$p_1 = 0.16 \pm 0.04 \pm 0.06$$

ϕ is the angle between two planes defined by incoming and outgoing electron and γ^* and outgoing proton

DVCS gives access to Generalized Parton Distributions (GPD), which describe the correlations between two partons (x_1, x_2) which differ by longitudinal ($x_1 \neq x_2$) and transverse (t) momentum at given Q^2

GPD's based model compatible with data