

## EXCLUSIVE DIFFRACTIVE FINAL STATES IN ELECTRON-PROTON COLLISIONS

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The exclusive diffractive production of vector mesons and real photons in  $ep$  collisions has been studied at HERA in a wide kinematic range. The most recent experimental results are presented.

### 1. Introduction

The diffractive scattering is a process where the colliding particles scatter at very small angles and without any color flux in the final state. This involves a propagator carrying the vacuum quantum numbers, called Pomeron and described, in the soft regime, within the Regge theory. The discovery of a big amount of diffractive events in DIS regime provided a hard scale which can be varied over a wide range and therefore it is an ideal testing for QCD models.

In particular, the diffractive production of Vector Mesons (VMs) and real photons in  $ep$  collisions allows to study the transition from the soft to the hard regime in strong interactions. The hard regime (high energy and low Bjorken- $x$ ) is dominated by the exchange of a hard Pomeron well described by perturbative QCD (pQCD), while at low- $x$  the interaction is well described within the Regge phenomenology. Indicating with  $Q^2$  the virtuality of the exchanged photon and with  $M^2$  the square mass of the produced VM, HERA data suggested a universal hard scale,  $Q^2 + M^2$ , for the diffractive exclusive production of VM and real photons, which indicates the transition from the soft to the hard regime.

### 2. $Q^2$ and $W$ dependence of the cross section

A new precision measurement of the reaction  $\gamma^*p \rightarrow \rho^0p$  was published by ZEUS <sup>1</sup>. It was found that the cross section falls steeply with the increas-

ing of  $Q^2$  but, unlike it was observed for the  $J/\psi$  electroproduction<sup>2,3</sup>, it cannot be described by a simple propagator term like  $\sigma \propto (Q^2 + M^2)^{-n}$ , in particular an  $n$  value increasing with  $Q^2$  appears to be favored. Figure 1 reports the cross section for the  $\rho^0$  electroproduction versus  $Q^2$  compared with several theoretical predictions: the KWM model<sup>4</sup> based on the saturation model, the FSS model<sup>5</sup> with and without saturation and the DF model<sup>6</sup>. None of the available models gives a good description of the data over the full kinematic range of the measurement.

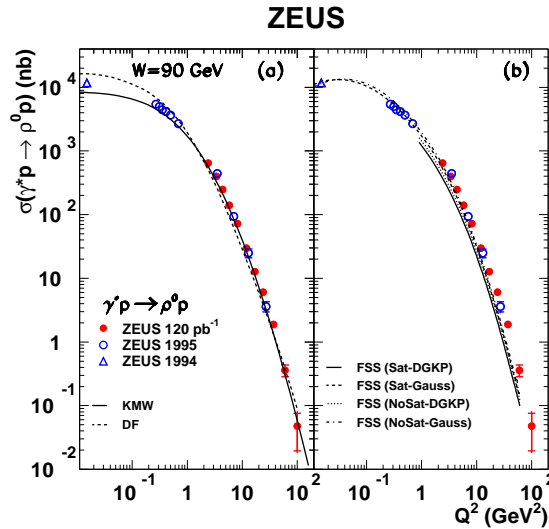


Figure 1. The  $\gamma^*p \rightarrow \rho^0 p$  cross section as a function of  $Q^2$  measured at  $W = 90 \text{ GeV}^2$  and compared in (a) and (b) with different models as described in the text.

The soft to hard transition can be observed looking at the dependence of the VM photoproduction ( $Q^2 = 0$ ) cross section from the  $\gamma^*p$  centre of mass energy,  $W$ , where the scale is provided by  $M^2$ . Figure 2 collects the  $\sigma(\gamma^*p \rightarrow Vp)$  as a function of  $W$  from the lightest vector meson,  $\rho^0$ , to the heaviest,  $\Upsilon$ , compared to the total cross section. The cross section rises with the energy as  $W^\delta$ , where the  $\delta$  exponent increases with the hard scale  $M^2$  as expected for a transition from the soft to the hard regime. New results on the  $\Upsilon$  photoproduction<sup>7</sup>, recently published by ZEUS, confirmed the steeper rise of  $\sigma(W)$  for higher vector meson masses.

The transition from the soft to the hard regime can also be studied

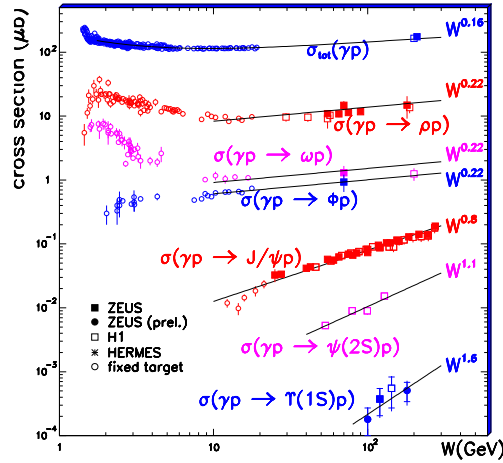


Figure 2. The  $W$  dependence of the cross section for exclusive VM photoproduction together with the total photoproduction cross section. Lines are the result of a  $W^\delta$  fit to the data at high  $W$ -energy values.

varying  $Q^2$ . Recent results were achieved by H1<sup>8</sup> and ZEUS<sup>9</sup> for the exclusive production of a real photon, the Deeply Virtual Compton Scattering (DVCS), where the hard scale is provided only by the photon virtuality,  $Q^2$ . Figure 3 shows the H1 (left) and the ZEUS (right) results. The steep rise with  $W$  of the cross section even at low- $Q^2$ , seems to suggest that the most sensitive part to the soft scale comes from the wave function of the produced VM. A similar result was obtained for the  $J/\psi$  electroproduction<sup>2,3</sup>.

The electroproduction of a large variety of VMs was studied at different  $Q^2$  values and the corresponding slope  $\delta$  is reported in Fig. 4 (left) versus the scale  $Q^2 + M^2$ , including the DVCS measurements. Data show a logarithmic shape  $\sigma \propto \ln(Q^2 + M^2)$  and the behaviour seems to be universal with  $\delta$  increasing from 0.2 at low scale, as expected from a soft Pomeron exchange<sup>10</sup> to  $\sim 0.8$  at large scale values.

### 3. $t$ dependence of the cross section

The differential cross section as a function of  $t$  can be parametrised by a fit  $\frac{d\sigma}{dt} \propto e^{b|t|}$ . Figure 4 (right) reports the collection of the  $b$  values versus the scale  $Q^2 + M^2$  for the electroproduction of VMs and DVCS, with  $b$  decreasing from  $\sim 11 \text{ GeV}^{-2}$  to  $\sim 5 \text{ GeV}^{-2}$  as expected in hard regime.

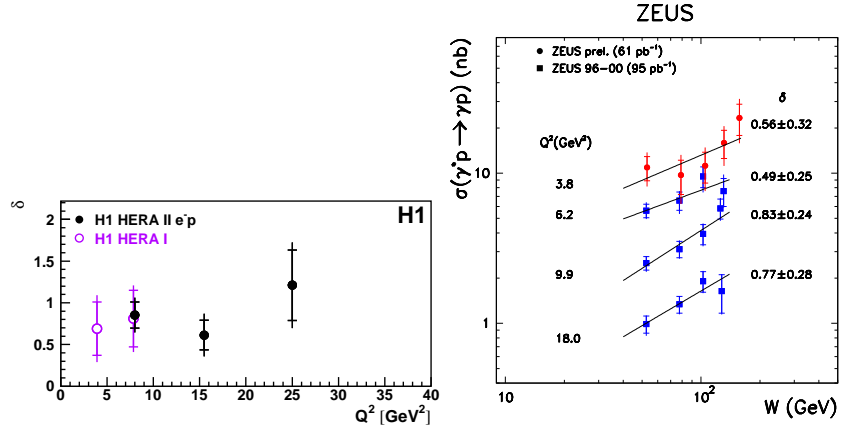


Figure 3. The  $W$  dependence of the cross section for a DVCS process. Lines come from a  $W^\delta$  fit to the data. Left: the H1 measurement of the  $\delta$  slope as a function of  $Q^2$ . Right: the new ZEUS preliminary measurement at low  $Q^2$  (dots) together with the published measurements (squares).

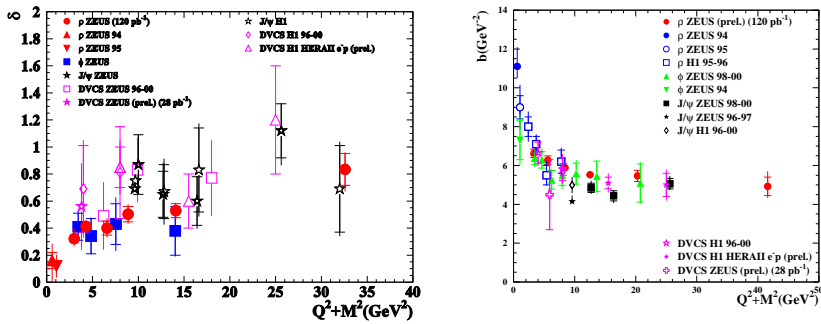


Figure 4. The dependence on the hard scale  $Q^2 + M^2$  of the value  $\delta$  (left) extracted from a fit  $W^\delta$  and of the slope  $b$  (right) extracted from a fit  $\frac{d\sigma}{dt} \propto e^{b|t|}$  for the exclusive VM electroproduction. DVCS is also included.

The measurement of  $d\sigma/d|t|$  for the DVCS process, recently published by the H1 Collab<sup>8</sup>, where  $t$  was obtained from the transverse momentum distribution of the photon, studied  $b$  versus  $Q^2$  and  $W$  as shown in Fig. 5.  $b$  seems to decrease with  $Q^2$  up to the value expected for a hard process but it doesn't depend on  $W$ . A new preliminary ZEUS measurement<sup>9</sup> of  $d\sigma/d|t|$  has been achieved from a direct measurement of the proton final

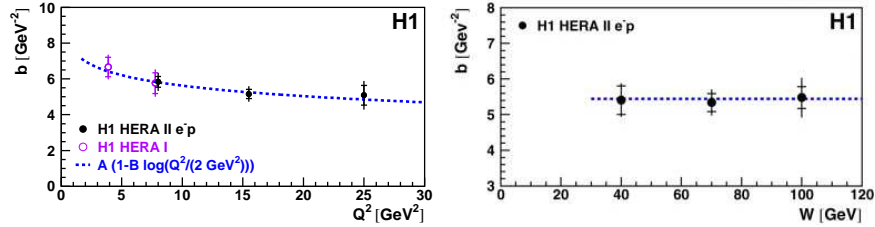


Figure 5. The  $t$  slope parameter  $b$  as a function of  $Q^2$  (left) and  $W$  (right).

state of using a spectrometer based on the roman pot technique. The result  $b = 4.4 \pm 1.3$  (*stat.*)  $\pm 0.4$  (*syst.*)  $\text{GeV}^{-2}$ , measured at  $Q^2 = 5.2 \text{ GeV}^2$  and  $W = 104 \text{ GeV}$ , is consistent, within the large uncertainties due to the low acceptance of the spectrometer, with the H1 result <sup>8</sup> of  $b = 5.45 \pm 0.19$  (*stat.*)  $\pm 0.34$  (*syst.*)  $\text{GeV}^{-2}$  at  $Q^2 = 8 \text{ GeV}^2$  and  $W = 82 \text{ GeV}$ .

Since  $b$  value can be related via a Fourier transform to the impact parameter and assuming that the exclusive process in the hard regime is dominated by gluons, the relation  $\langle r^2 \rangle = b(\hbar c)^2$  can be used to obtain the radius of the gluon confinement area in the proton.  $b \sim 5 \text{ GeV}^2$  corresponds to  $\langle r^2 \rangle \sim 0.6 \text{ fm}$  smaller than the proton radius ( $\sim 0.8 \text{ fm}$ ) indicating that the gluons are well contained within the charge-radius of the proton.

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