



**Electroweak Physics
at High Q^2
at HERA**

Iris Abt (MPI München)
on behalf of

A scenic view of a riverbank in Darmstadt, Germany. In the foreground, a white boat with red trim is docked at a wooden pier. The background features a dense cluster of historic buildings with red roofs, including a prominent castle-like structure with towers. The sky is clear and blue.

Outline

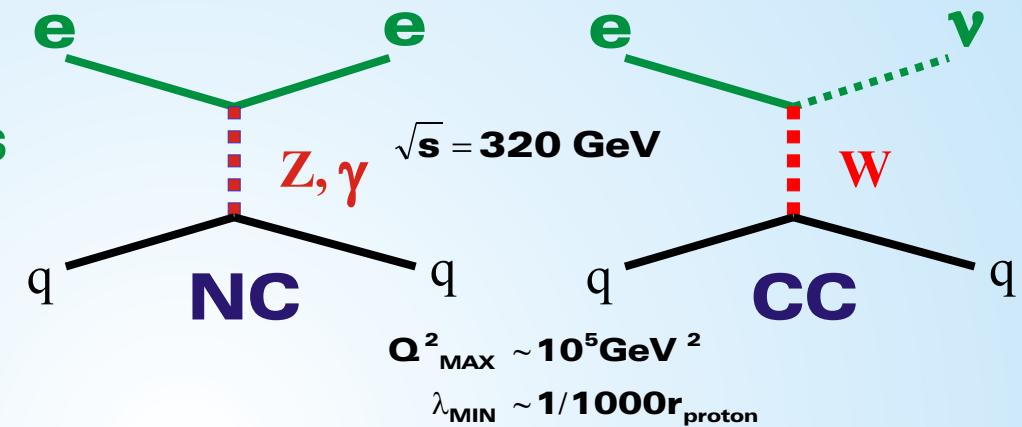
- **Introduction**
- **HERA**
- **Deep Inelastic Scattering -- ep**
- **Kinematic Range and Experiments**
- **Polarization and DIS**
- **High Q^2 Results for CC and NC**
- **Elektroweak Plus QCD Fit Results**
- **Single W Production [Isolated Leptons]**
- **Health Bulletin of the SM, i.e. Summary**

No attempt was made to be complete

ep Scattering

Only at HERA

27.5 GeV Electrons
probing
920 GeV Protons



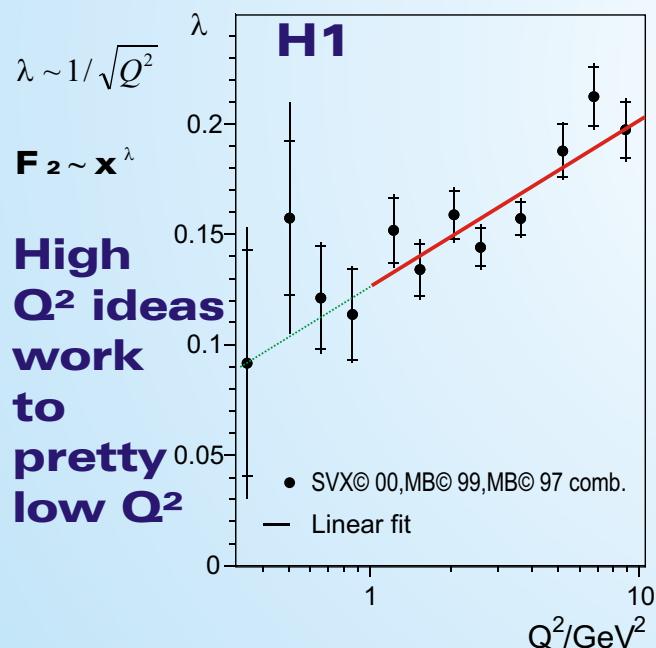
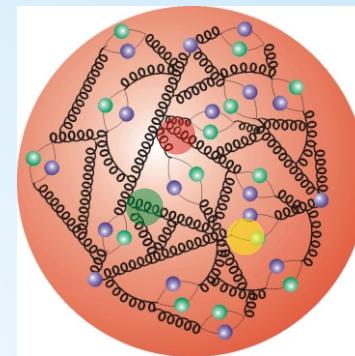
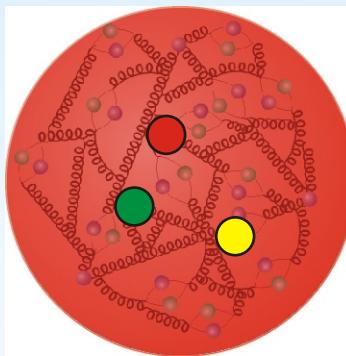
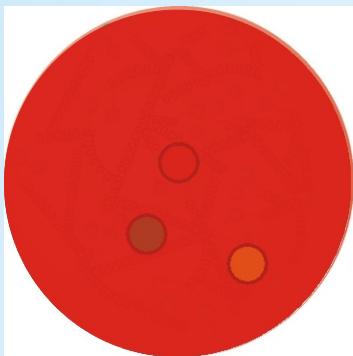
DIS at $Q^2 \approx \text{EW-scale}$

Note: DIS was not invented here, but at $Q^2 \approx 0$. The Z, i.e. NC, was first seen in νN interactions.



Would need a 50 TeV beam in a fixed target experiment

What Is High Q^2

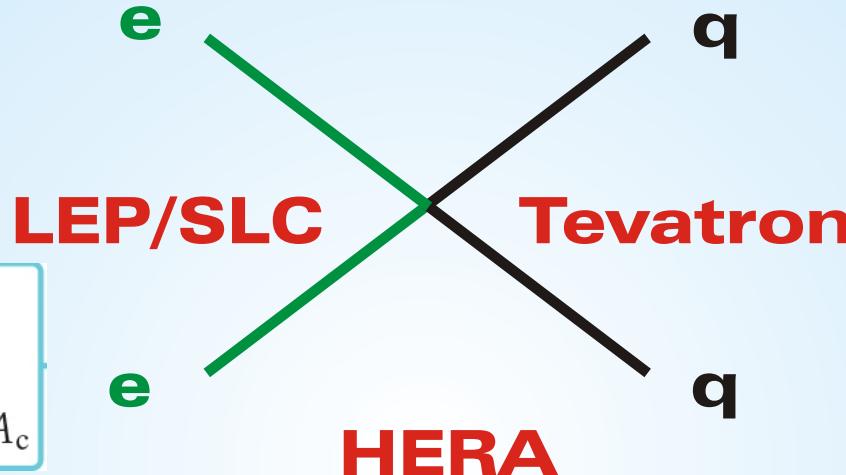


Anything that reveals the internal structure of the proton. Increasing Q^2 first reveals valence quarks and then sea quarks and glue.

Who Works at the EW Scale?

LEP:
 $m_Z, \Gamma_Z, \sigma_h^0, R_l^0, A_{FB}^{0,l}$
 $P_\tau \rightarrow A_l$
 $Q_{FB} \rightarrow \sin^2 \theta_{eff}^{lept}$

SLD: A_l
LEP+SLD:
 $R_b^0, R_c^0, A_{FB}^{0,b}, A_{FB}^{0,c}, A_b, A_c$



$p\bar{p}$: m_t
LEP+ $p\bar{p}$: m_W, Γ_W

HERA:

t-channel exchange of gauge bosons

- γZ interference
- propagator masses

Parton Distribution Functions needed

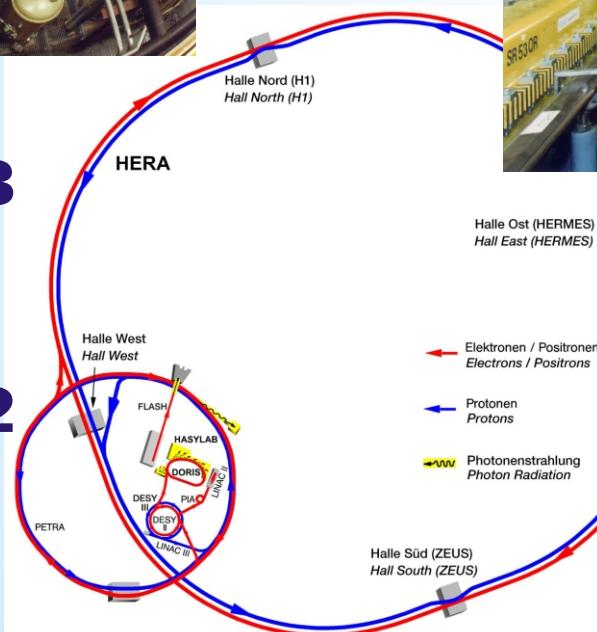
HERA in Memoriam



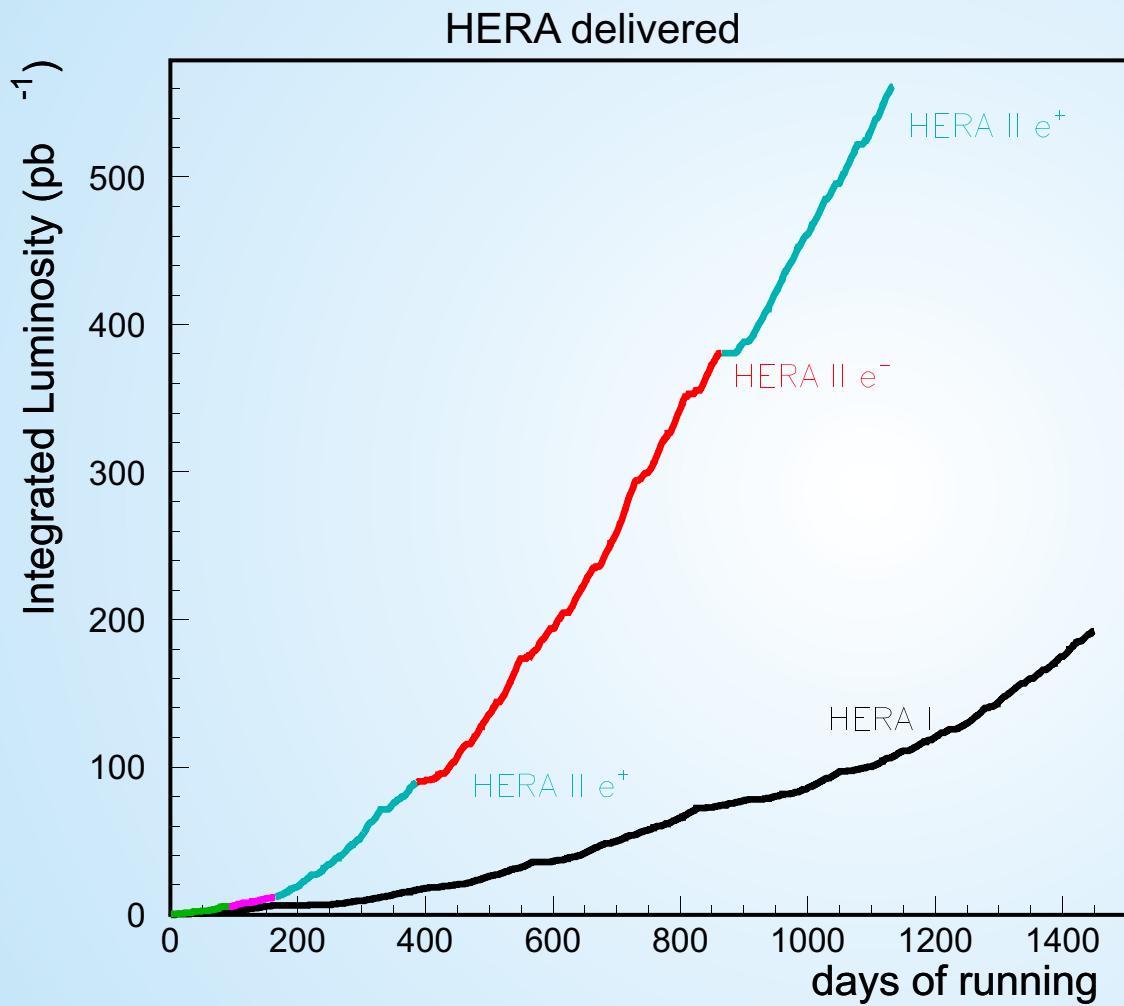
*
24.5.1993

**
1.11.2002

**last beam
30.6.2007**



HERA Luminosity



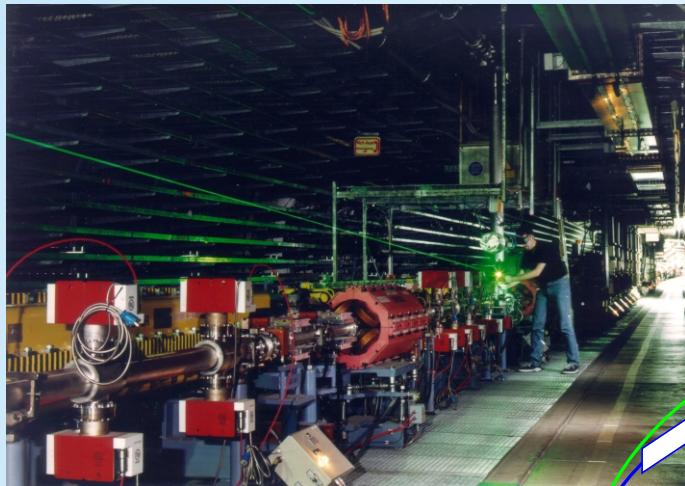
562 pb^{-1}

**Nov. 2002
Mar. 2007**

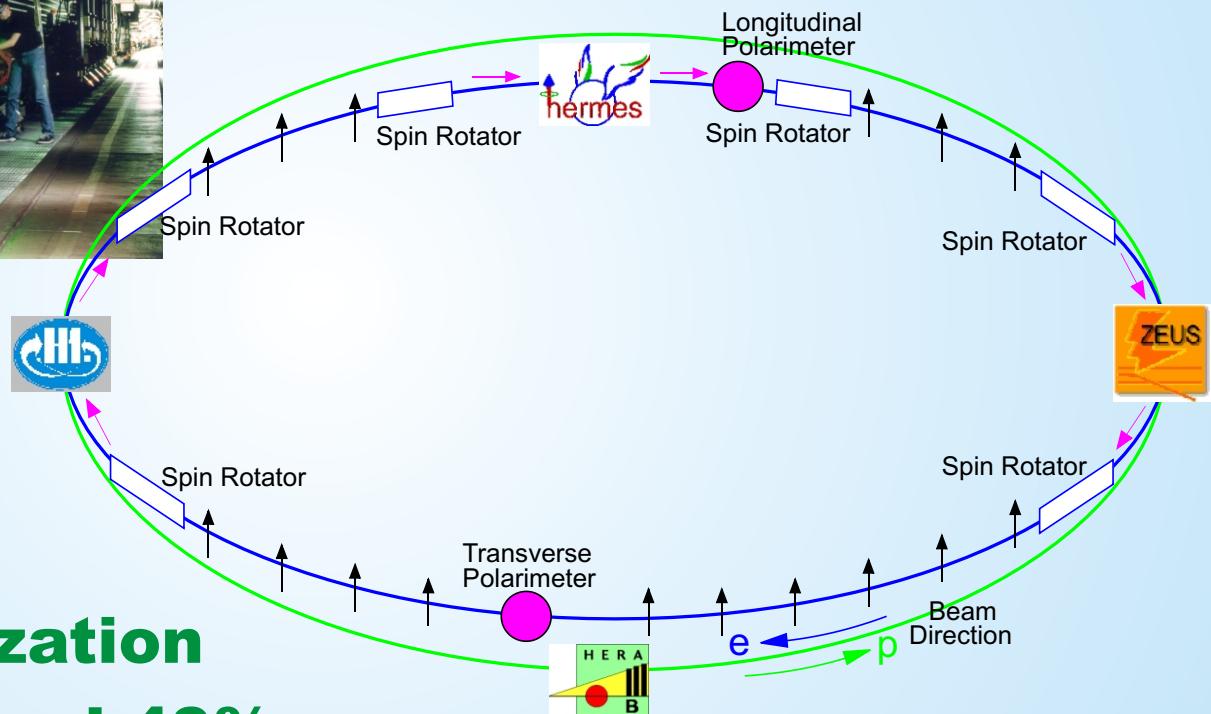
193 pb^{-1}

**May 1993
Aug. 2000**

Polarization at HERA



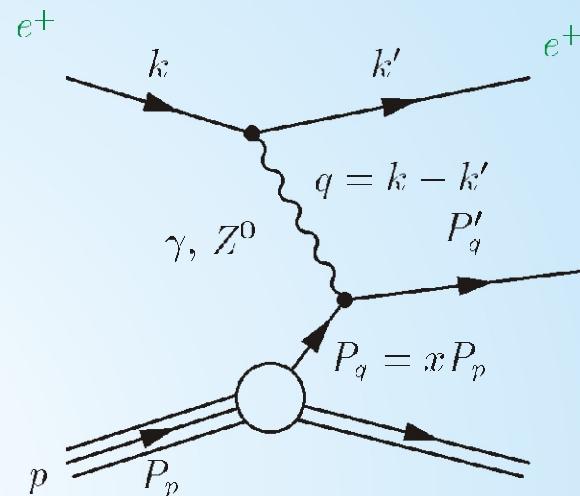
Longitudinally polarized electrons and positrons.



Average polarization
between 30% and 40%

Kinematics

- Virtuality: $Q^2 = -(k - k')^2$
→ Spatial resolution of probe $\lambda \sim 1/\sqrt{Q^2}$
- Bjorken scaling variable: $x = Q^2 / 2pq$
→ Momentum fraction of struck parton
- Inelasticity: $y = pk / pq$
→ Energy transfer to proton (in p rest frame)



$$Q^2 = xys$$

**Experiment measures Cross-sections:
Structure Functions (SFs) are deduced.**

$$\frac{d^2\sigma}{dxdQ^2} = \frac{2\pi\alpha^2}{Q^4} Y_+ F_2$$

$$Y_+ = 1 + (1 - y)^2$$

If proton is point like $\rightarrow \frac{d^2\sigma}{dxdQ^2} = \frac{2\pi\alpha^2}{Q^4} Y_+$

(The longitudinal SF, F_L , is neglected)

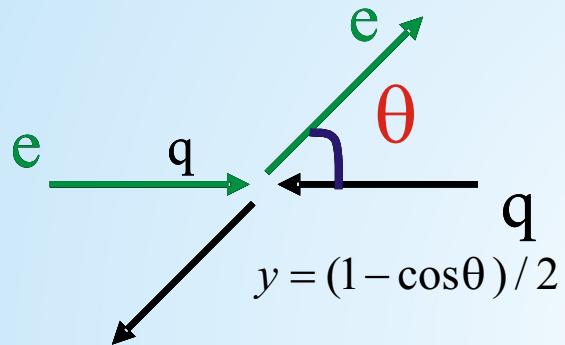
Deviation from pointlike behavior

Quark Parton Model -- QPM

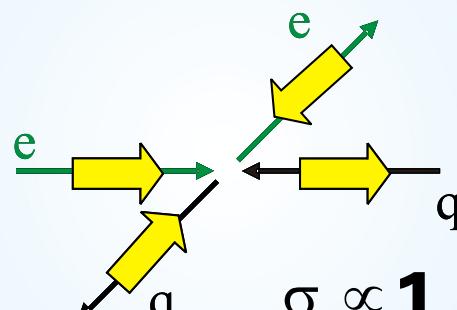
Kinematic variable y corresponds to angle θ between e and quark.

eq c.m.s.

$$\frac{d^2\sigma}{dx dQ^2} = \frac{2\pi\alpha^2}{Q^4} Y_+ F_2$$



$$V: \quad Y_+ = 1 + (1 - y)^2 \quad A: \quad Y_- = 1 - (1 - y)^2$$



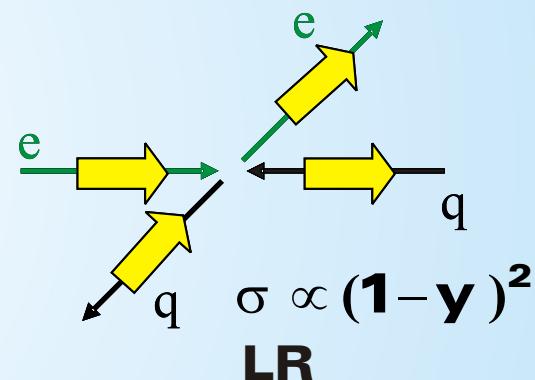
**Low Q^2 , i.e.
photon exchange:**

- vector component only
- F_2 is based on based on quark charges

$$F_2 = x \sum e_q^2 (q + \bar{q})$$



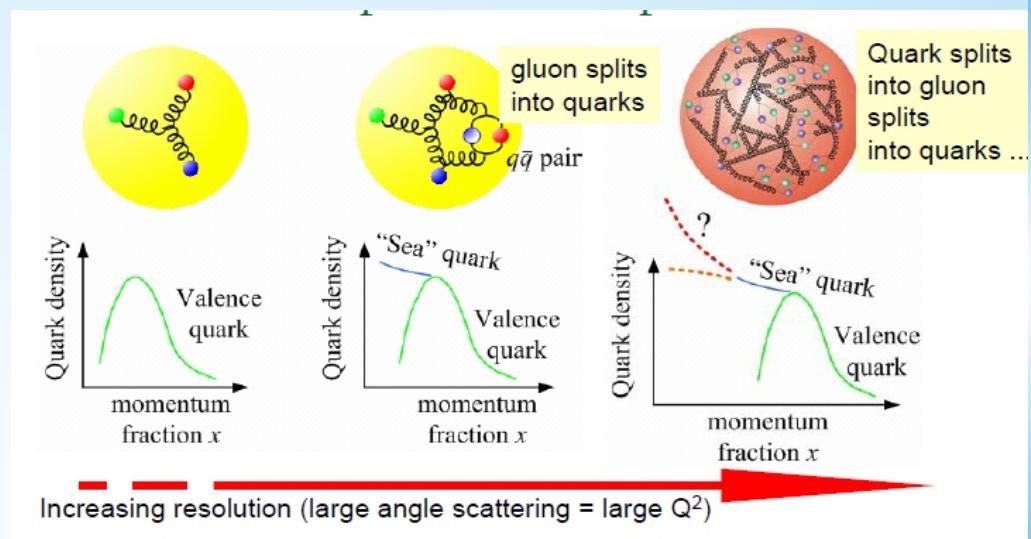
$\sigma =$ **coupling • propagator • kinematic factor • charges² • PDFs**



QCD Evolution and Gluons

beyond QPM

- PDFs are not static
→ “evolution” as Q^2 grows.
- Structure depends on the power to see.
- pQCD can describe this evolution: “DGLAP eq.”



$$\frac{\partial}{\partial \ln Q^2} \left(\frac{\Sigma}{xg} \right) = \alpha_s \begin{pmatrix} P_{qq} & P_{gq} \\ P_{gq} & P_{gg} \end{pmatrix} \otimes \left(\frac{\Sigma}{xg} \right)$$

$$\frac{\partial}{\partial \ln Q^2} q_{NS} = \sigma_s P_{qq} \otimes q_{NS}$$

$$\text{At low-}x: \frac{\partial F_2}{\partial \ln Q^2} \propto \alpha_s x g$$

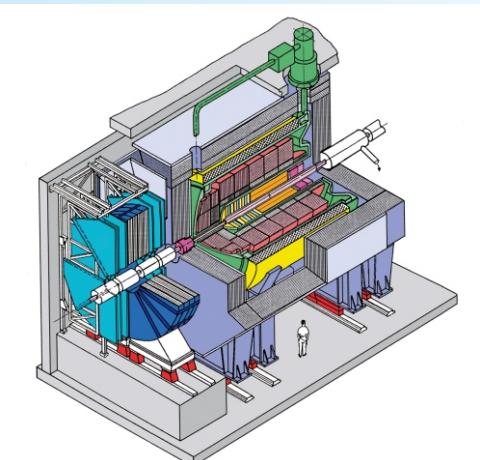
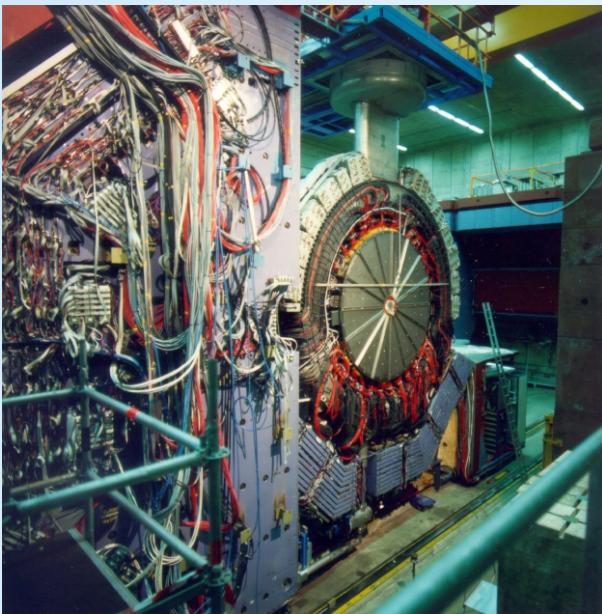
- F_2 is sum of $q / q\bar{q}$ PDFs
→ Gluons not directly in F_2 (in LO)
- Gluons cause “slope” of F_2 in $\log Q^2$ evolution

pQCD cannot predict x - dependence of PDFs a priori
PDFs are determined by a global fit to experimental data

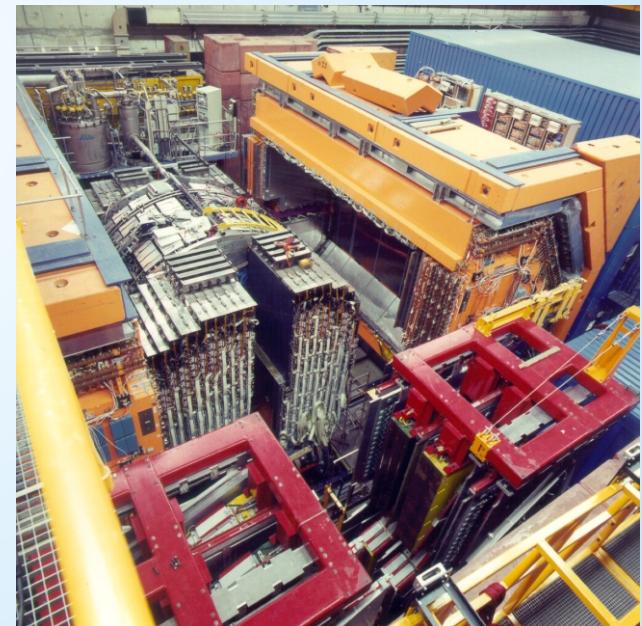
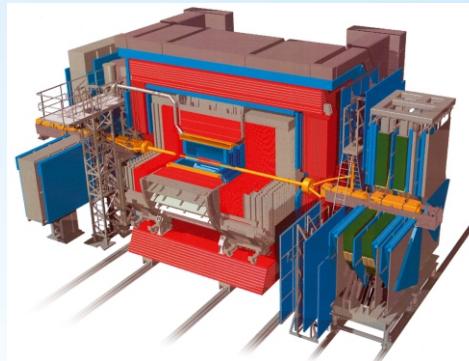
Experiments

H1

went
for
LAr

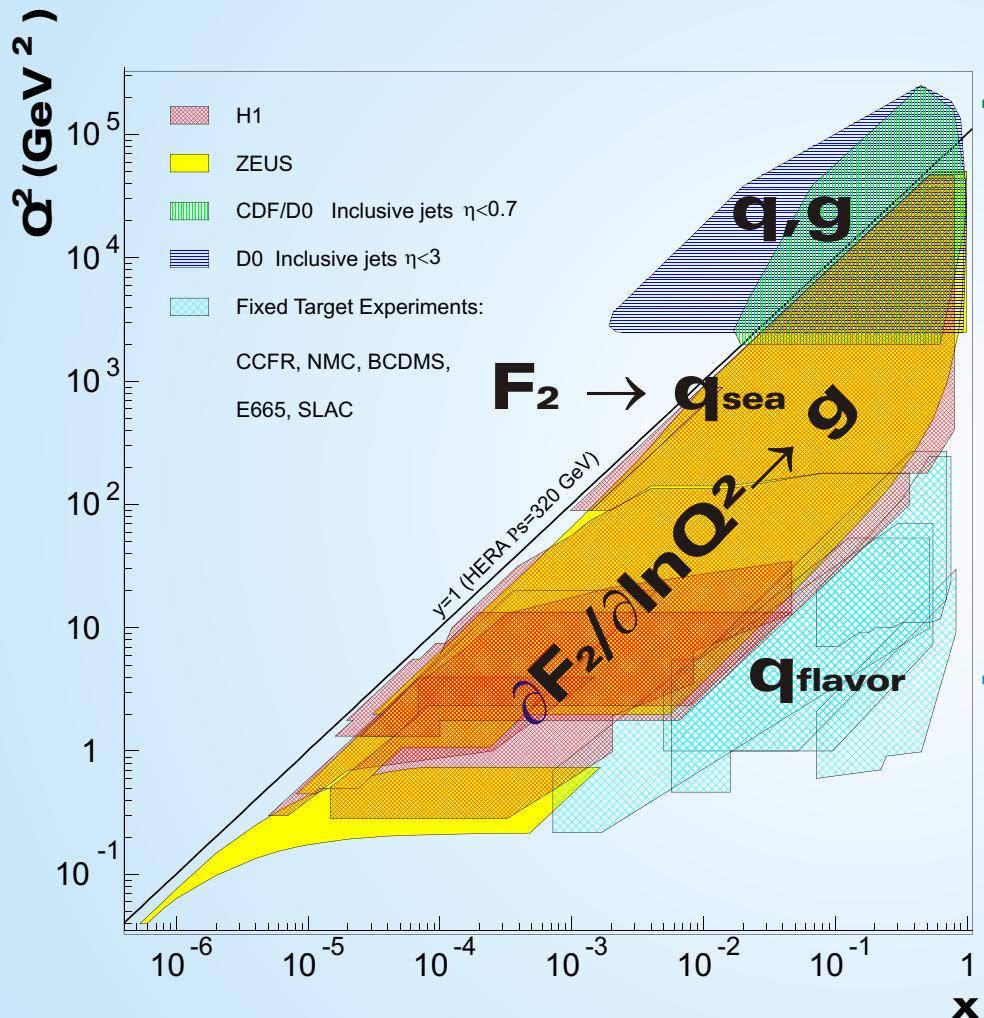


ZEUS
went
for
compen-
sation



Let's have
a final look
while they
are still
there....

Kinematic Range



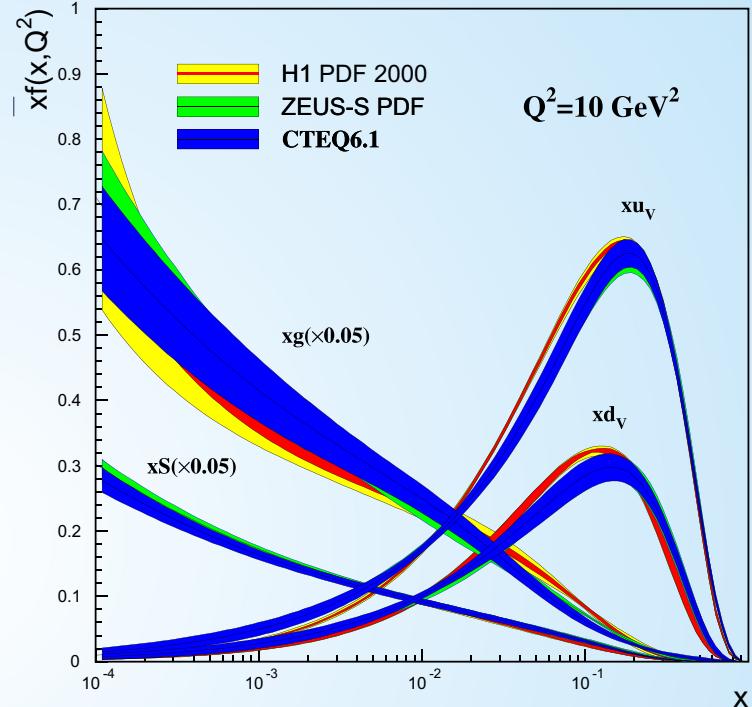
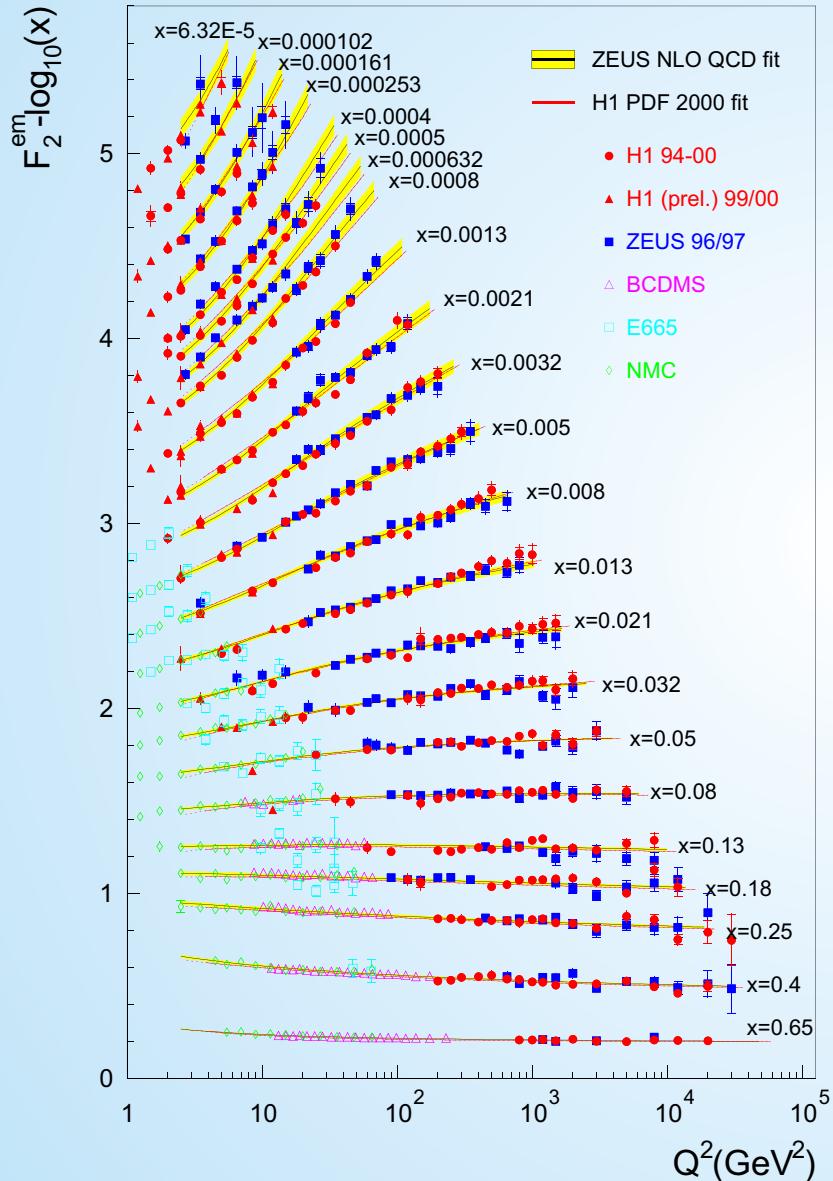
Tevatron Jet Distributions

HERA: specialised in sea quarks determined via F_2 .

Fixed Target

Note, that kinematic range is not static.

Structure Function F_2 and PDFs



NLO pQCD describes F_2 over

- 4 orders in Q^2
- 3 orders in x

Scaling Violations reveal the DIS invisible gluon.

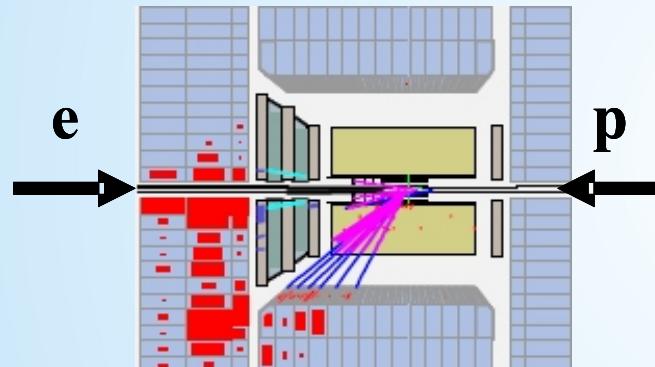
**PDFs known \Rightarrow
EW can be studied**

CC at High Q²

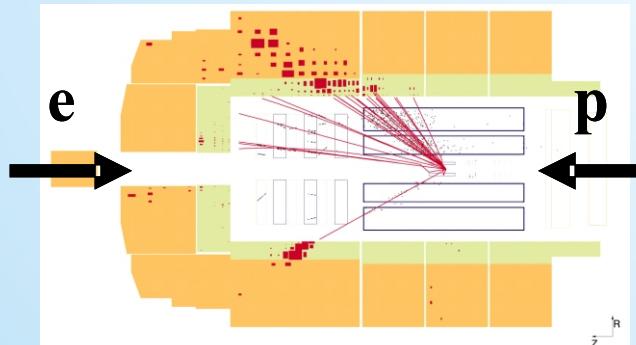
CC ep → vX : Pure Weak (only L)

$$\frac{d^2\sigma(e^+ p)}{dx dQ^2} = \frac{G_F^2}{2\pi} \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2 \{ (\bar{u} + \bar{c}) + (1-y)^2 (d + s) \}$$

$$\frac{d^2\sigma(e^- p)}{dx dQ^2} = \frac{G_F^2}{2\pi} \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2 \{ (u + c) + (1-y)^2 (\bar{d} + \bar{s}) \}$$



... while NC event looks like:



e-p:

-- charge selecting nature:
only up-type q (downtype
anti-q)

-- anti-q receives (1-y)²
helicity suppression

- Selection: large missing transverse energy: Pt,miss
- Kinematics reconstructed using hadrons (only possibility)

- Selection: presence of high p_T scattered electron, scattered at large angle
- Kinematics well reconstructed using either electrons or hadrons (or both)

NC at High Q²

- NC ep → eX: Z effects at high Q²
 - F₂ receives additional terms
 - “Axial” SF, F₃, comes into

$$\frac{d^2\sigma(e^\pm P)}{dx dQ^2} = \frac{2\pi\alpha^2}{Q^4} \{ Y_+ \tilde{F}_2 \mp Y_- x \tilde{F}_3 \}$$

Nb.: xF₃ is written as F₃ in the equations below for simplicity

$$\begin{aligned}\tilde{F}_2 &= \Sigma A_q x(q + \bar{q}) = F_2^\gamma - v_e \chi_Z F_2^{\gamma Z} \\ \tilde{F}_3 &= \Sigma B_q x(q - \bar{q}) = -a_e \chi_Z F_3^{\gamma Z}\end{aligned}$$

$$\chi_Z = \frac{1}{\sin^2 2\theta_w} \frac{Q^2}{M_Z^2 + Q^2}$$

Propagator term

For axial: sign flips between particles and anti-particles

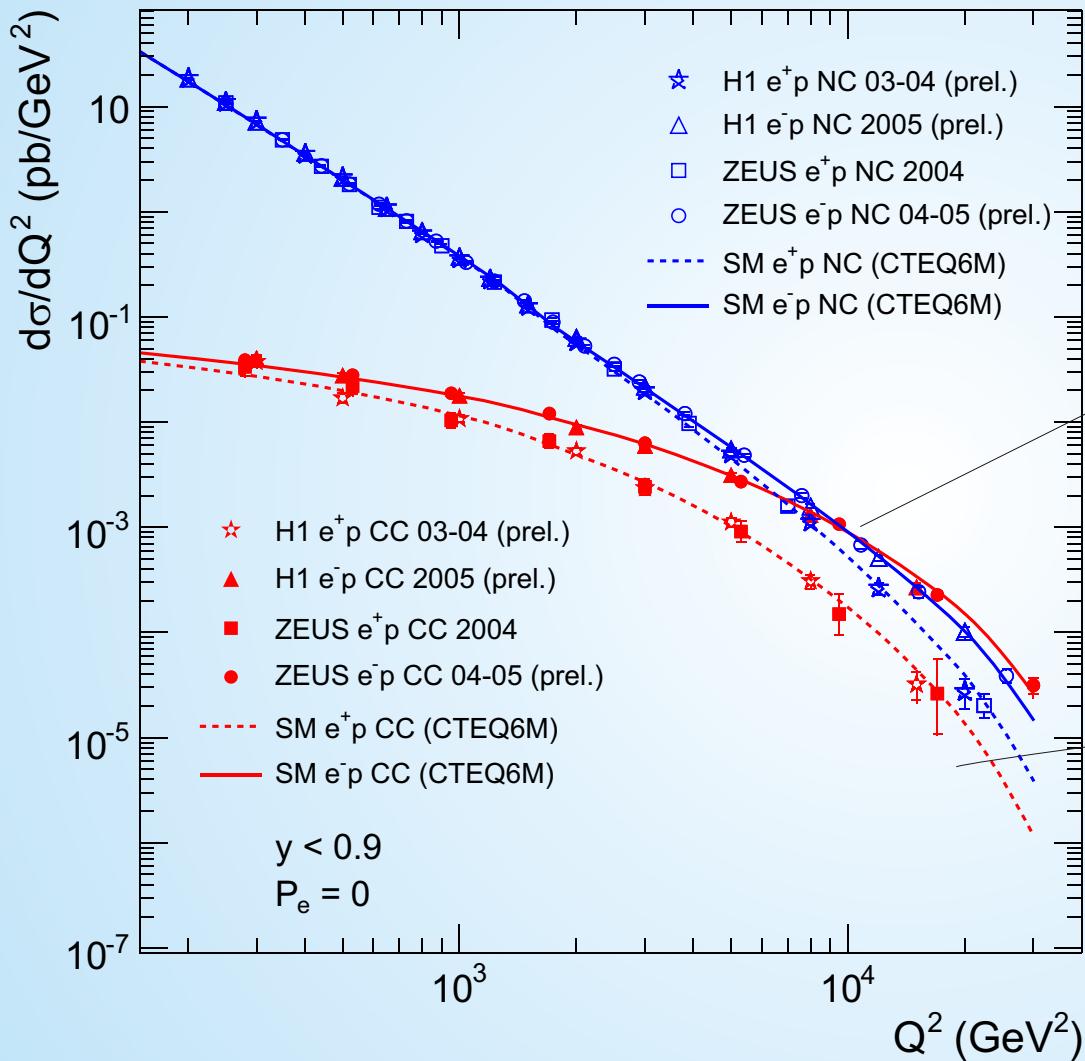
- Sign flips between e+/e-
- q/qbar contributes to xF₃ with different sign

→ xF₃ is proportional to valence q

1 st -order V	2 nd -order V
$v_e \chi_Z F_2^{\gamma Z}$	$(v_e^2 + a_e^2) \chi_Z^2 F_2^Z$
$a_e \chi_Z F_3^{\gamma Z}$	$2v_e a_e \chi_Z^2 F_3^Z$
$\begin{matrix} 1^{\text{st}}-\text{order A} \\ \gamma Z \text{ interference} \end{matrix}$	
$\begin{matrix} 2^{\text{nd}}-\text{order A} \\ \text{Pure Z} \end{matrix}$	

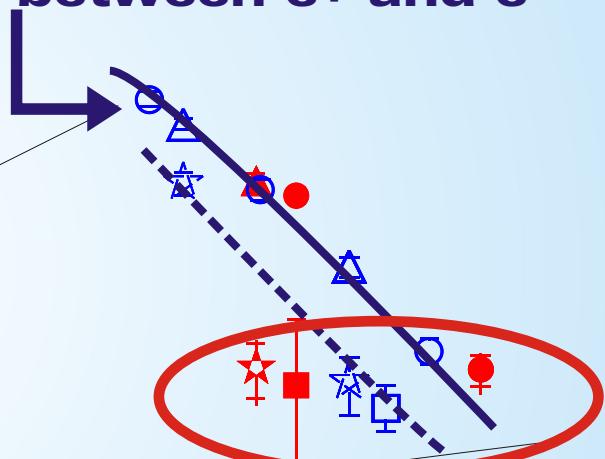
- F₂: 2nd order only $\sim a_e^2 \chi_z^2 F_2^Z$
- F₃: 1st order γZ interf. $\sim a_e \chi_Z F_3^{\gamma Z}$

EW “unification”



HERA II

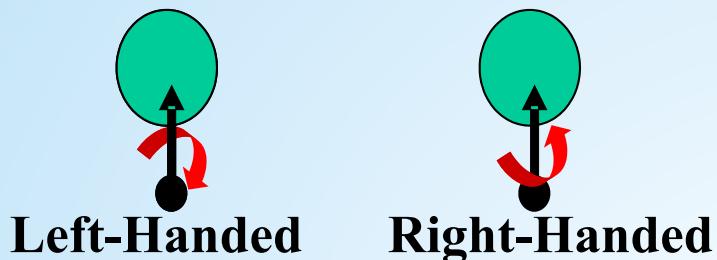
In NC axial component
 xF_3 is seen as difference
between e^+ and e^-



NC and CC cross-
sections become
similar.

Remaining differences
mainly due to PDFs

EW Physics With Polarized Leptons



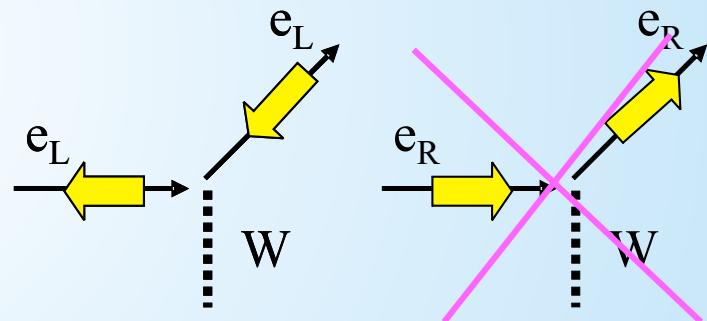
Charged-current DIS

- “Pure” Weak
 - Chiral structure of weak int. is directly visible as a function of Polarization
- Weak = “100% parity violated” (no RH)
 - Zero cross section @ Pol=1 (-1 for e+)
 - $\sigma(\text{Pol}) = (1+\text{Pol}) \sigma(\text{Unpol})$

Neutral-current DIS

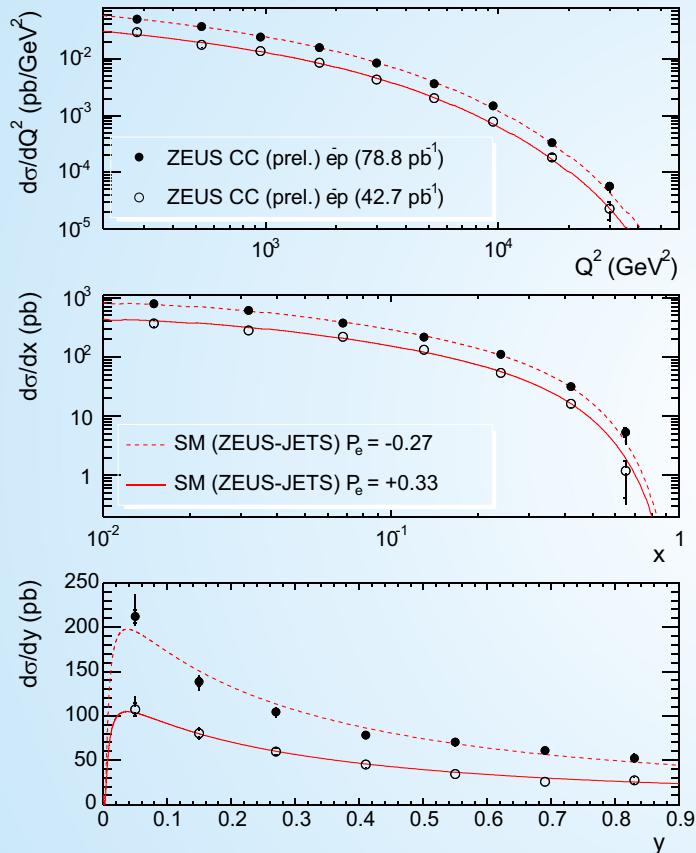
- Weak parity violating effect through γZ interference and pure Z
 - visible only at large Q^2
- Such γZ and Z terms contain EW parameters,
i.e. quark couplings to Z , $\sin\theta_W$, M_Z

- Polarization = Asymmetry of Helicity states:
 $P = (N_R - N_L) / (N_R + N_L)$
- Helicity = Chirality (if mass is neglected)
 - By means of Pol, chiral structure can be tested.
- $RH \neq LH \Leftrightarrow$ parity violation

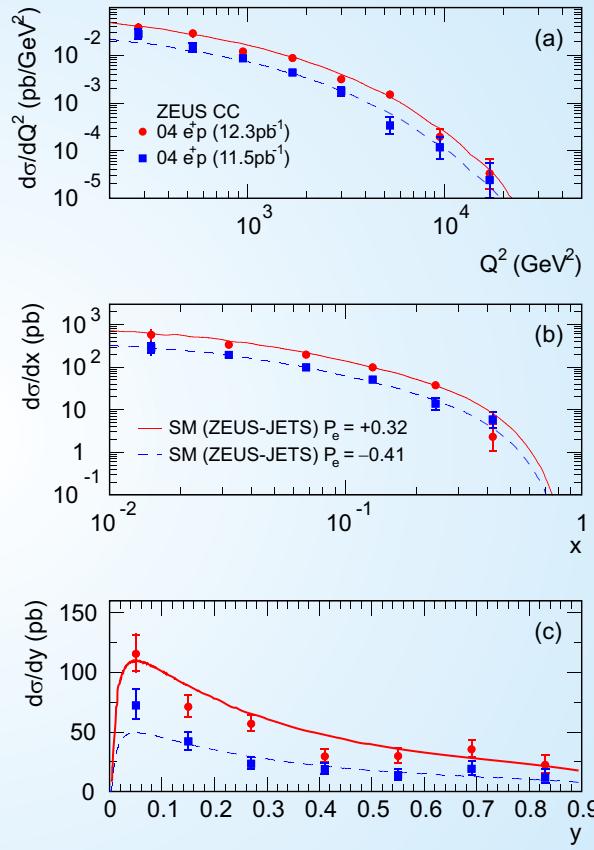


Charged Current Cross-Section

Electron



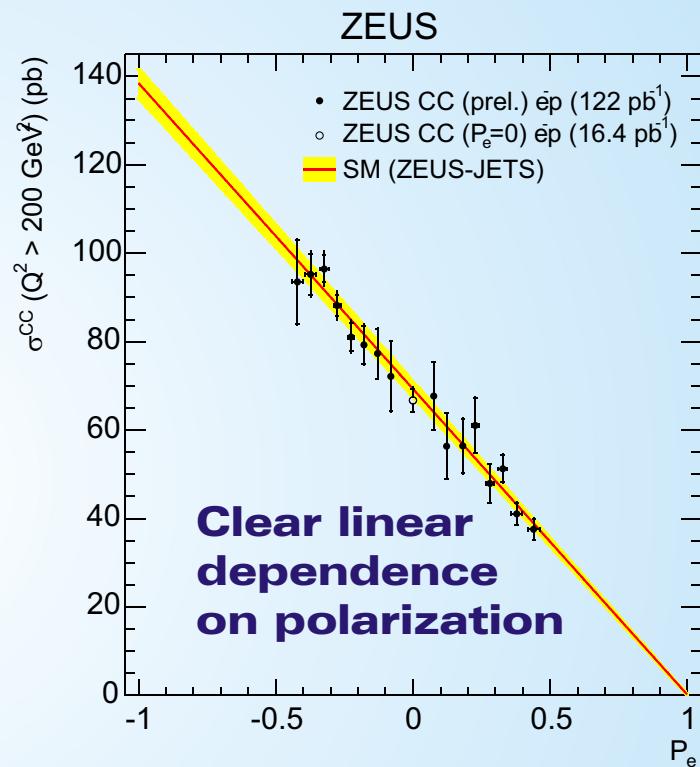
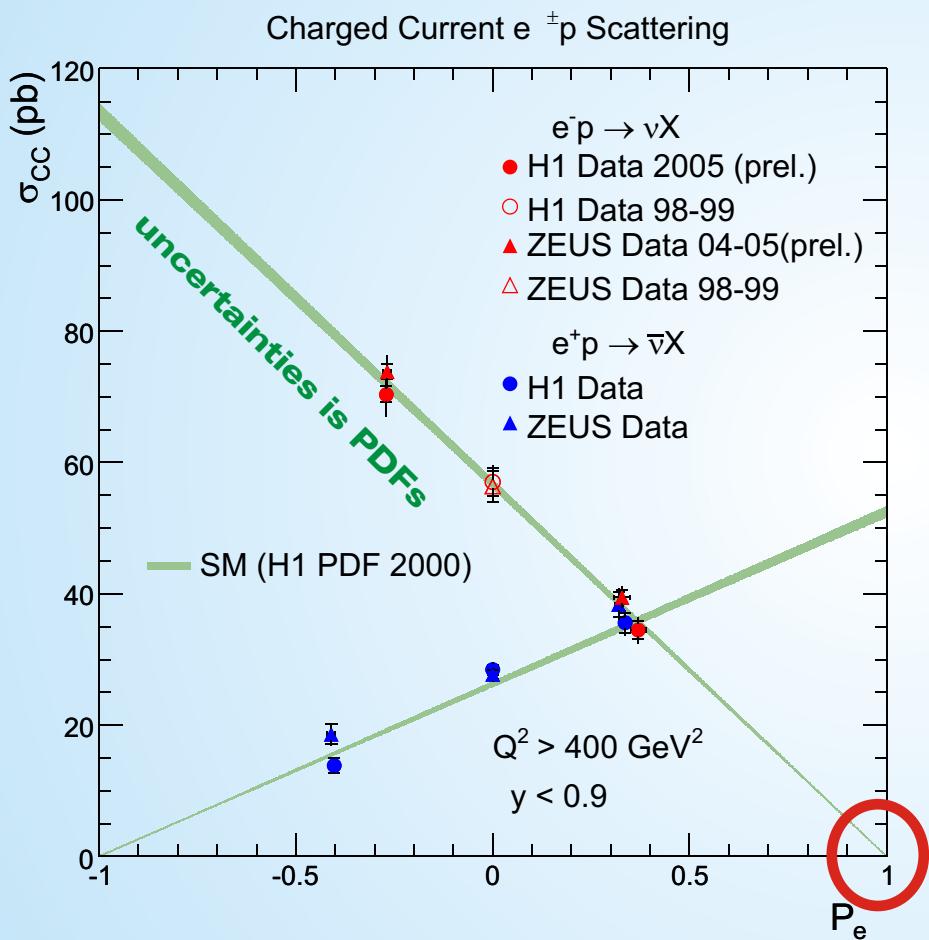
Positron



ZEUS

**Clear difference in cross-sections
between negative and positive polarization**

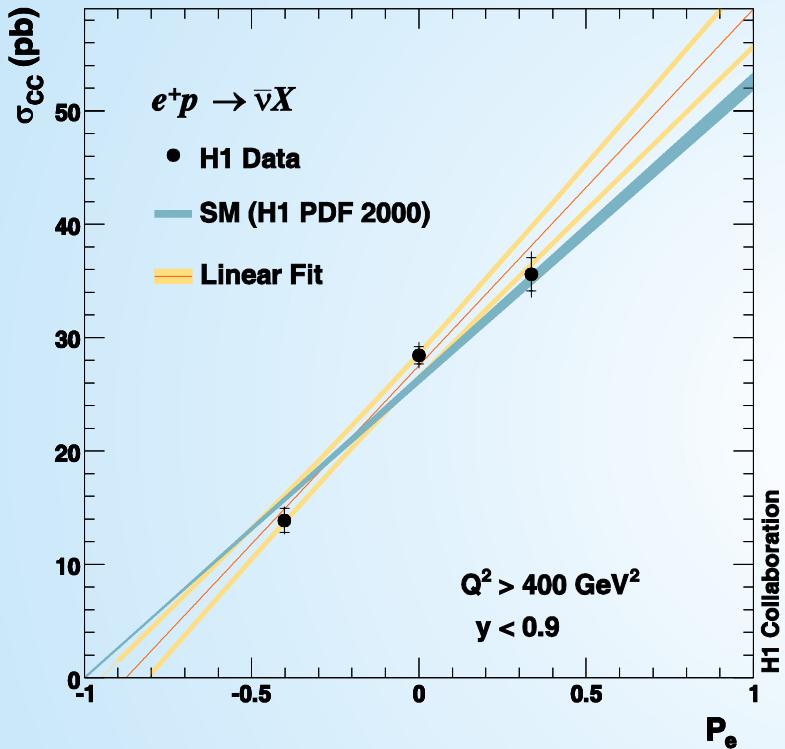
Polarization Dependence of CC Cross-Section



Consistent with SM prediction
 $\sigma(CC_{RH}) = 0$

W_R Mass Limit

H1 positron

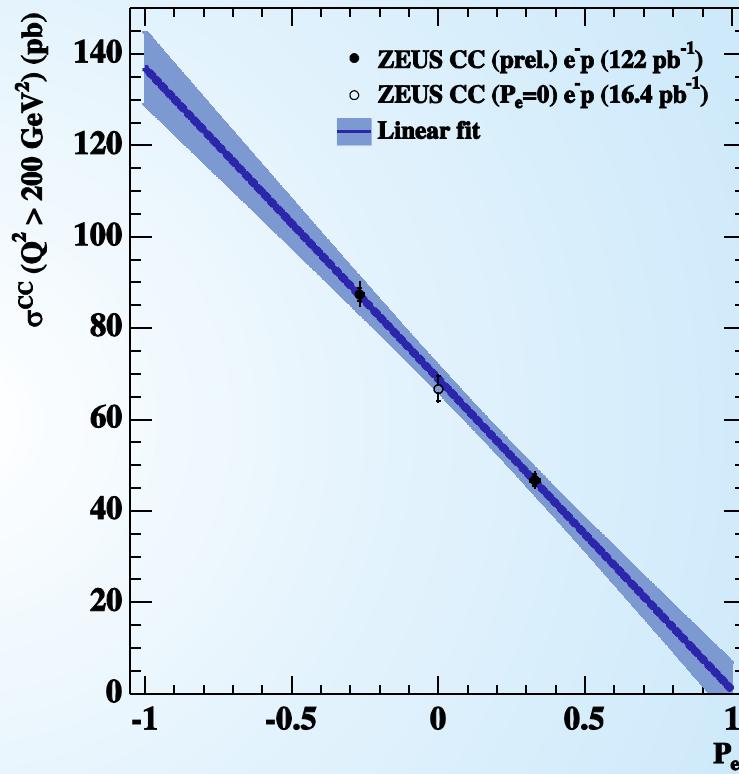


Assume $g_L = g_R$ and ν_R is light:

H1 e- $\rightarrow W_R > 186 \text{ GeV}$

H1 e+ $\rightarrow W_R > 208 \text{ GeV}$

ZEUS electron



ZEUS e- $\rightarrow W_R > 180 \text{ GeV}$

> 310 GeV from polarized ^{12}N β decay

> 786 GeV from $W' \rightarrow e\nu, \mu\nu$

Polarization Effects in NC

$$\begin{aligned}\tilde{F}_2 &= F_2^\gamma - (\textcolor{blue}{v}_e \pm P_e \textcolor{red}{a}_e) \chi_Z F_2^{\gamma Z} + ((v_e^2 + a_e^2) \pm P_e 2\textcolor{blue}{v}_e \textcolor{red}{a}_e) \chi_Z^2 F_2^Z \\ \tilde{F}_3 &= -(\textcolor{red}{a}_e \pm P_e \textcolor{blue}{v}_e) \chi_Z F_3^{\gamma Z} + ((2\textcolor{red}{v}_e \textcolor{blue}{a}_e \pm P_e (v_e^2 + a_e^2)) \chi_Z^2 F_3^Z\end{aligned}$$

Nb.: xF_3 is written as F_3 for simplicity

- Polarization modifies γZ and Z terms:

- Axial in F_2 , vector in F_3
- dependent on size of P_e

$$\textcolor{blue}{a}_e \approx 0$$

- F_2 : 1st order, $\sim \pm P_e a_e \chi_Z F_2^{\gamma Z}$
- F_3 : 2nd order only, $\sim \pm P_e a_e^2 \chi_Z^2 F_3^Z$

Unpol:

$$\sigma(e^+) - \sigma(e^-) \rightarrow F_3^{\gamma Z}$$

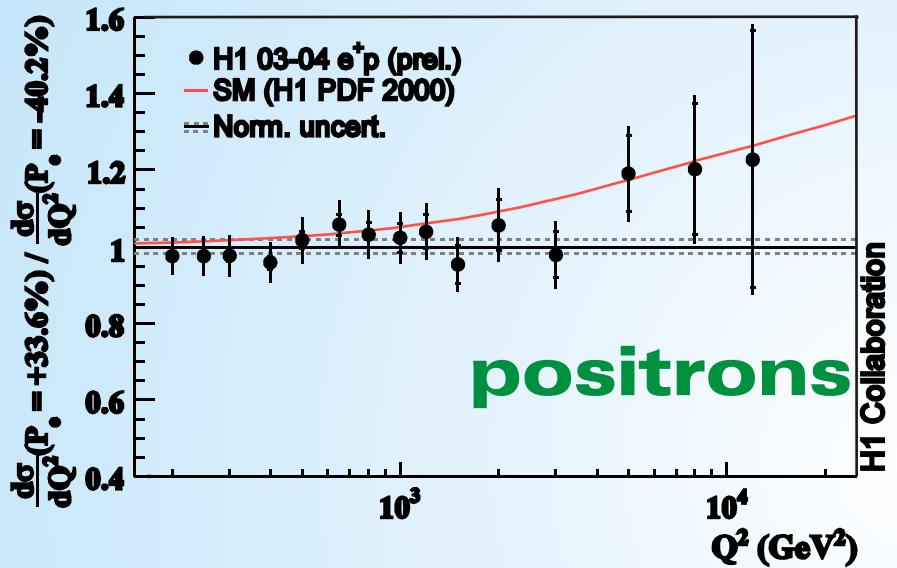
Pol :

$$\sigma(P_e \rightarrow) - \sigma(P_e \leftarrow) \rightarrow F_2^{\gamma Z}$$

- Polarization effects expected only at EW scale, i.e large Q^2

NC Cross-Section vs. Polarization

H1

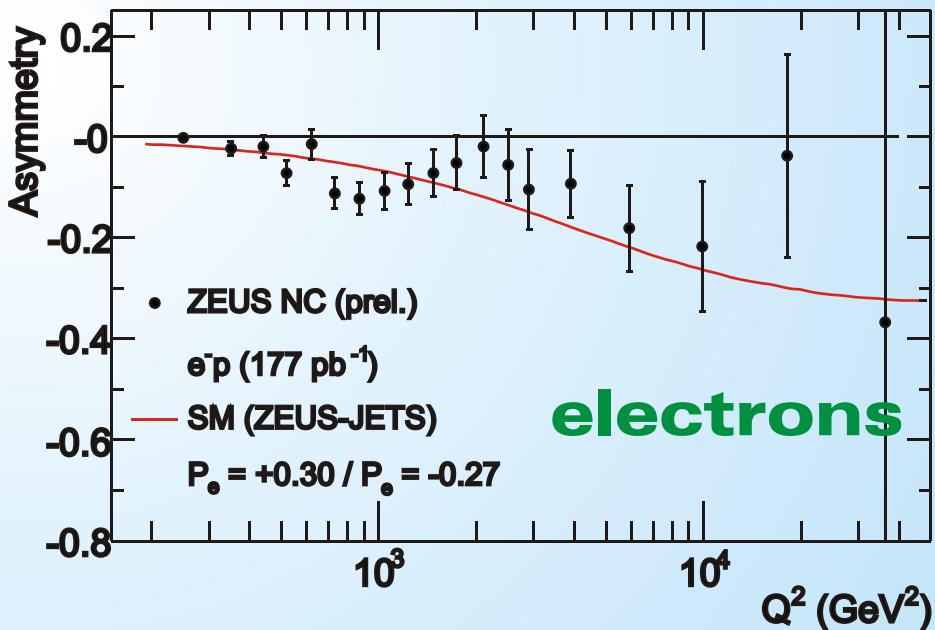


**Observation
of Parity Violation
at EW scale**

**$d\sigma/dx$ and $d\sigma/dy$
do not strongly
depend on P_e**

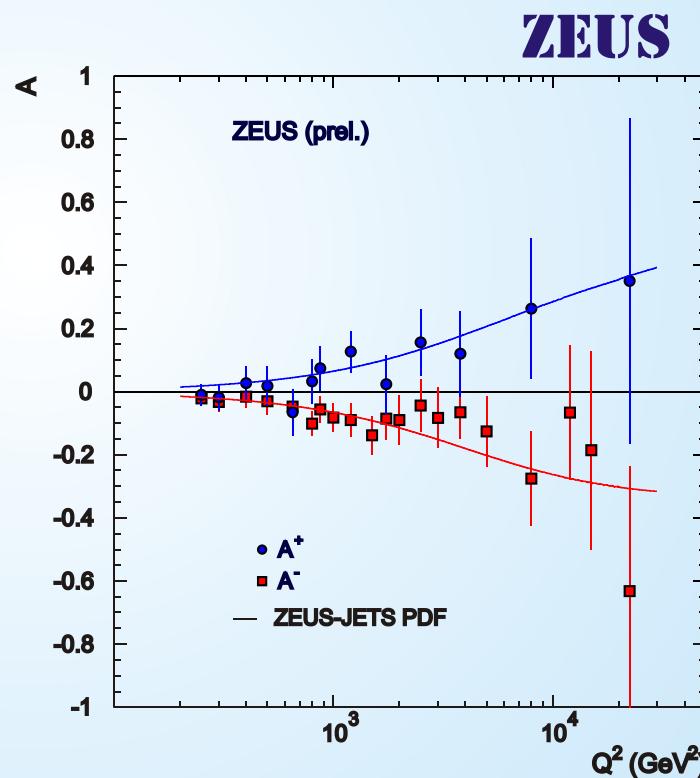
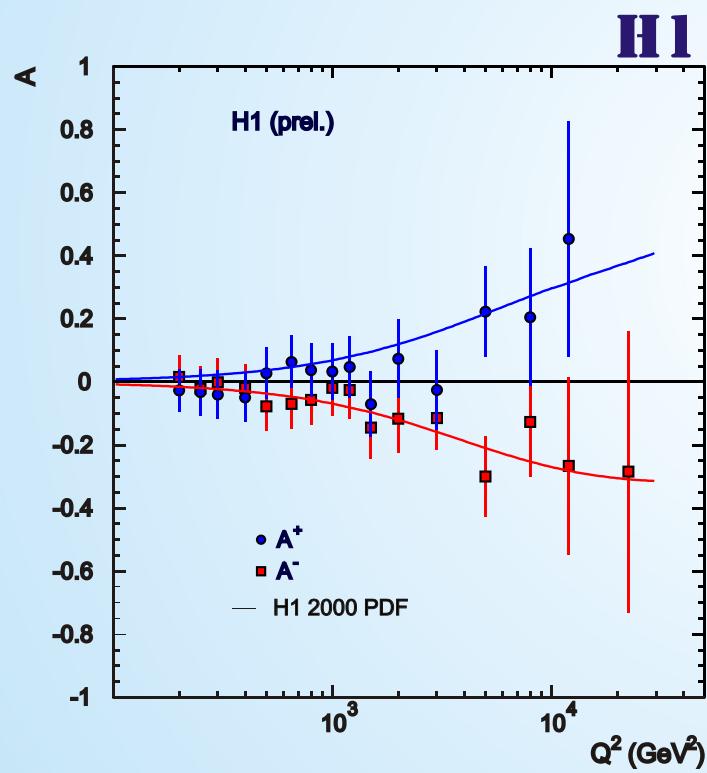
$d\sigma/dQ^2$ does

ZEUS

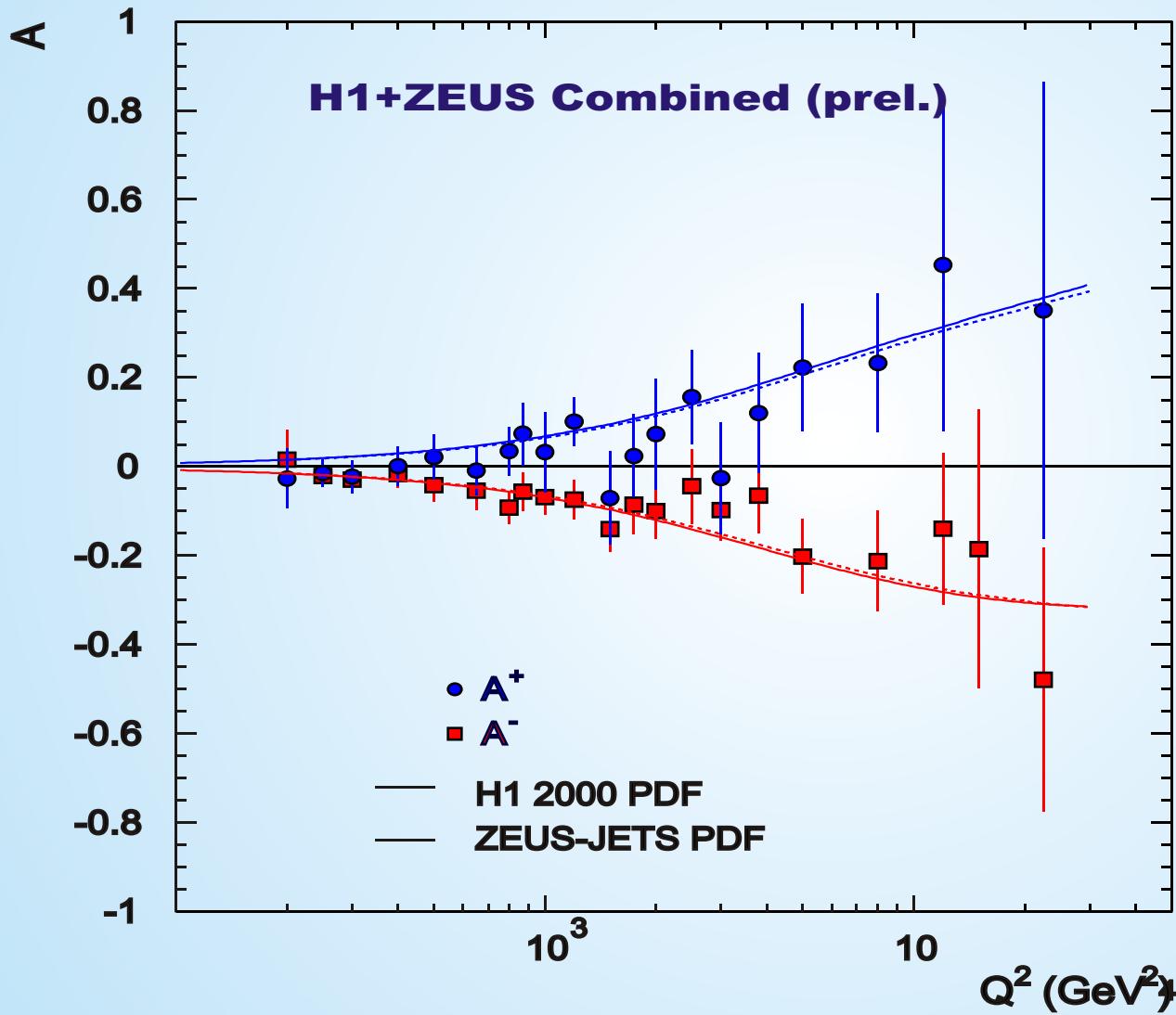


NC Cross-Section Asymmetries

$$A_{\pm} = \frac{2}{P_R - P_L} \frac{\sigma_{\pm}(P_R) - \sigma_{\pm}(P_L)}{\sigma_{\pm}(P_R) + \sigma_{\pm}(P_L)}$$



NC Cross-Section Asymmetries



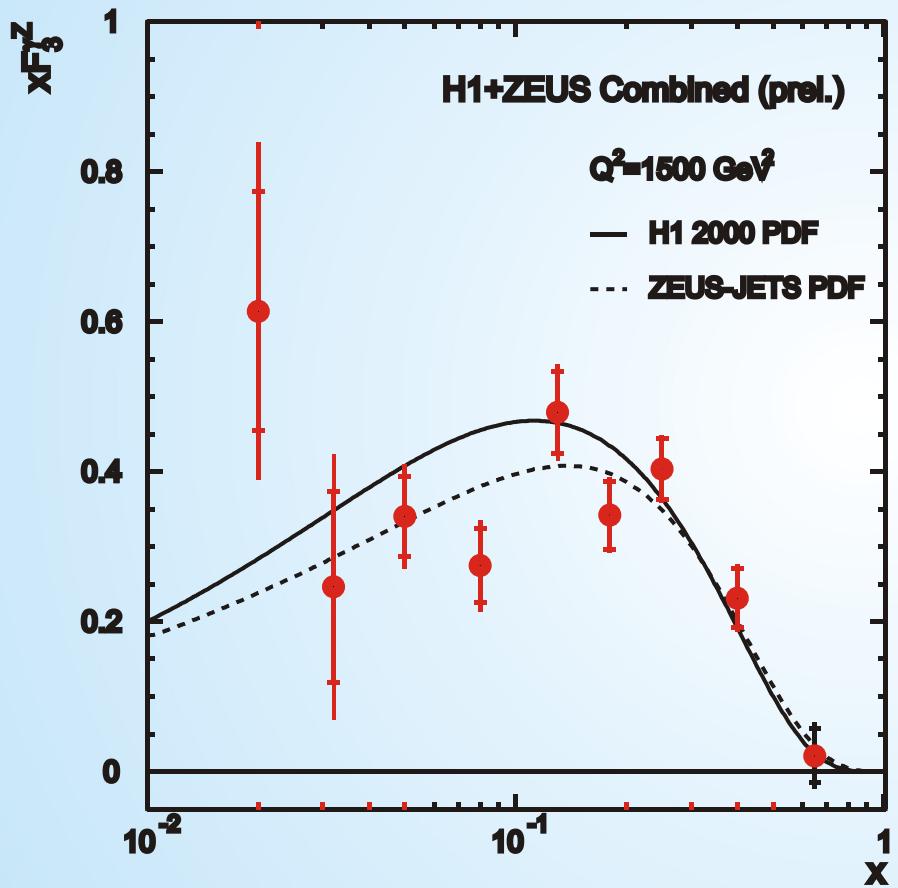
Parity Violation
due to γZ
interference

At high x ,
assuming
SM couplings

$$A_{\pm} \sim \frac{u_v + d_v}{4u_v + d_v}$$

constrain d/u

H1 + ZEUS



Combine all data
and correct to
zero polarization.

Assuming
SM couplings

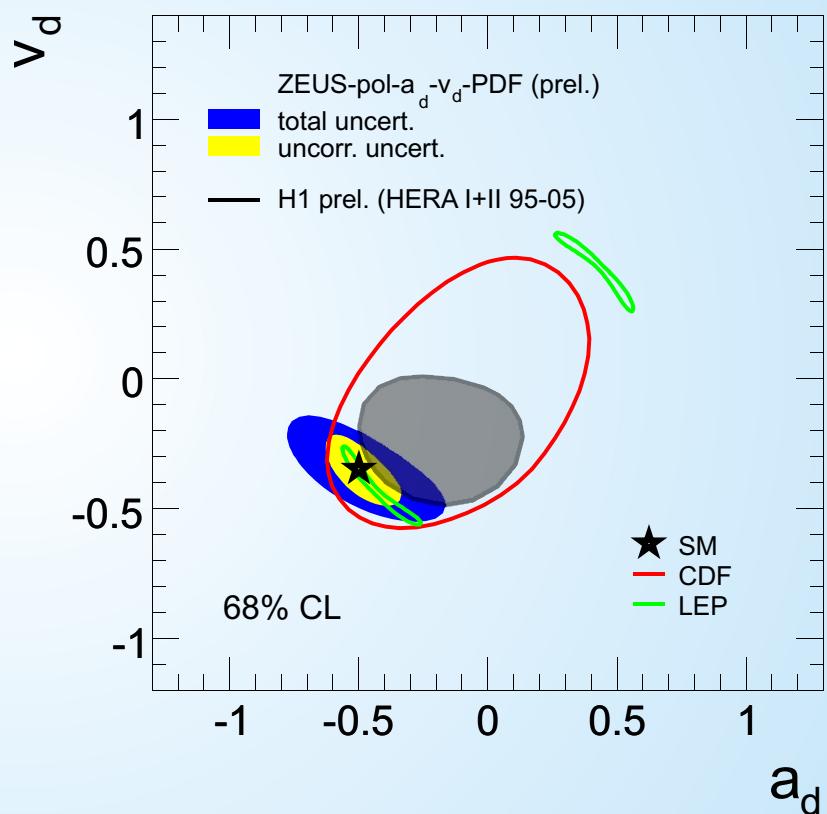
$$F_3 \sim 2/3 u_v + 1/3 d_v$$

Constrain
valence quarks

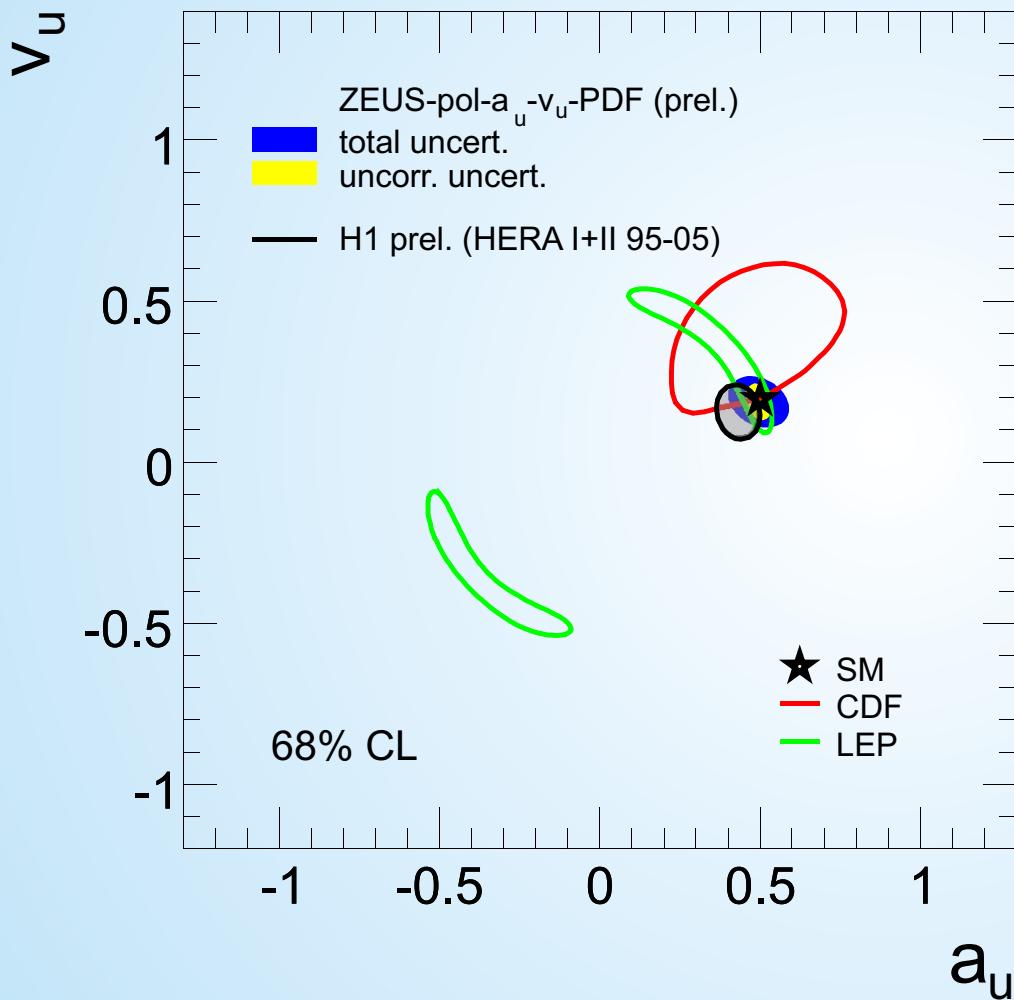
EW plus QCD Fit

**Fit both PDFs and
EW parameters
simultaneously**

**Experiments fit
seperately to
simplify handling
of systematic
errors**



EW plus QCD Fits



Precise determination of light quark couplings to Z

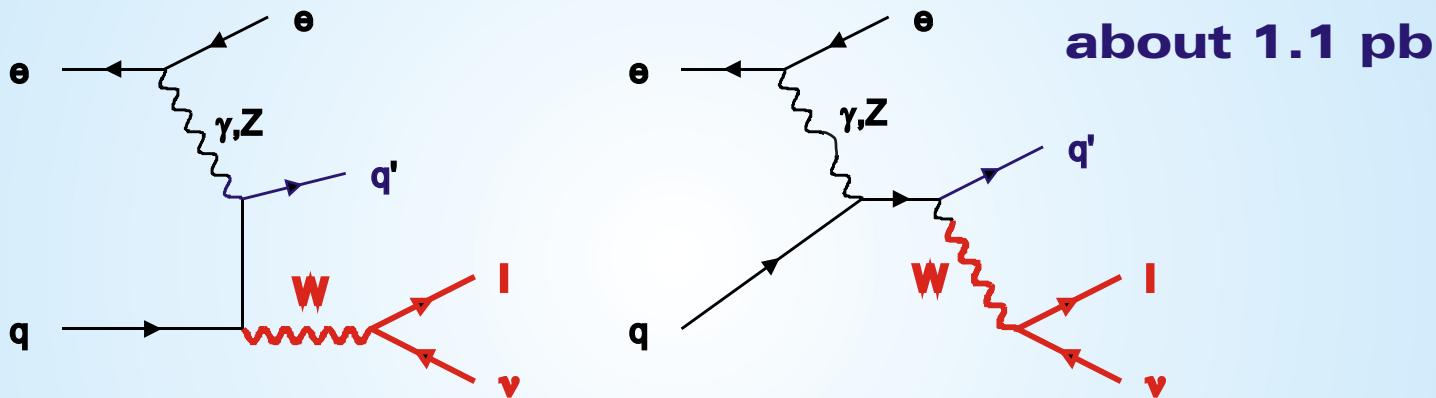
Luminosity helps with a

Polarization helps with v

Smack on SM prediction...

Isolated leptons

In the SM isolated leptons are only produced through single W production:

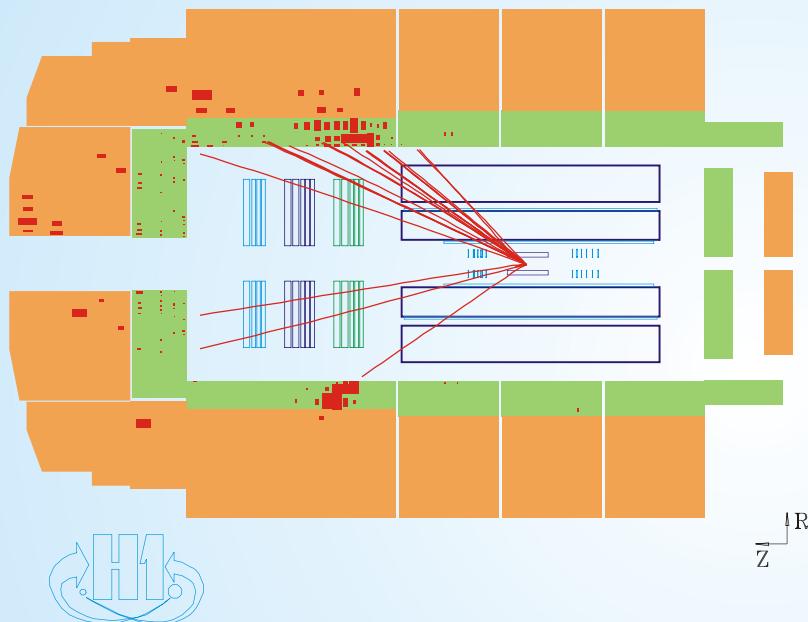


Select

quark jet with large transverse momentum p_{Tq}
only small fraction of SM events/expected for single t

isolated lepton with large transverse momentum p_{Tl}
large missing transverse momentum $p_{T\text{miss}}$

Event in H1

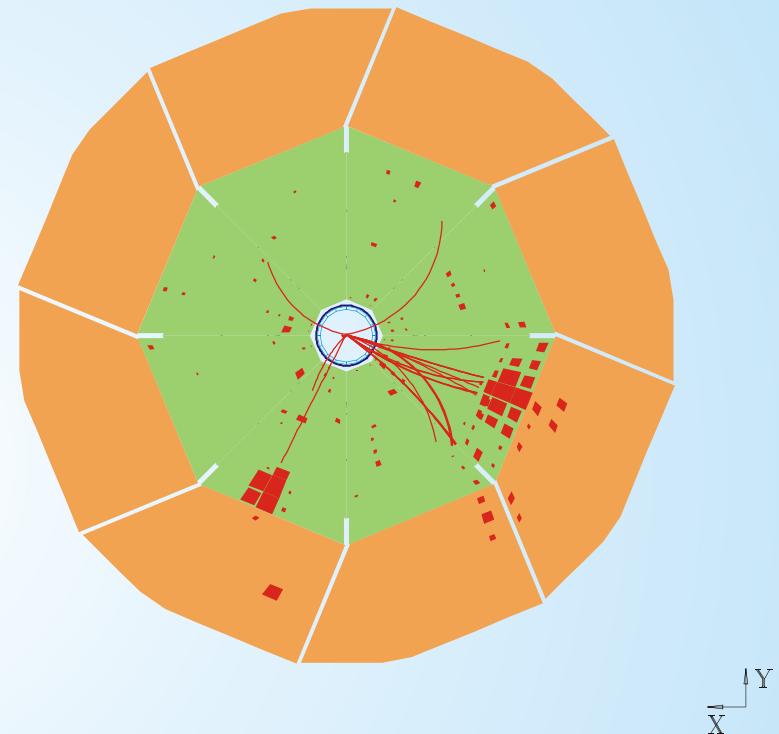


$p_{T_x} > 29 \text{ GeV}$

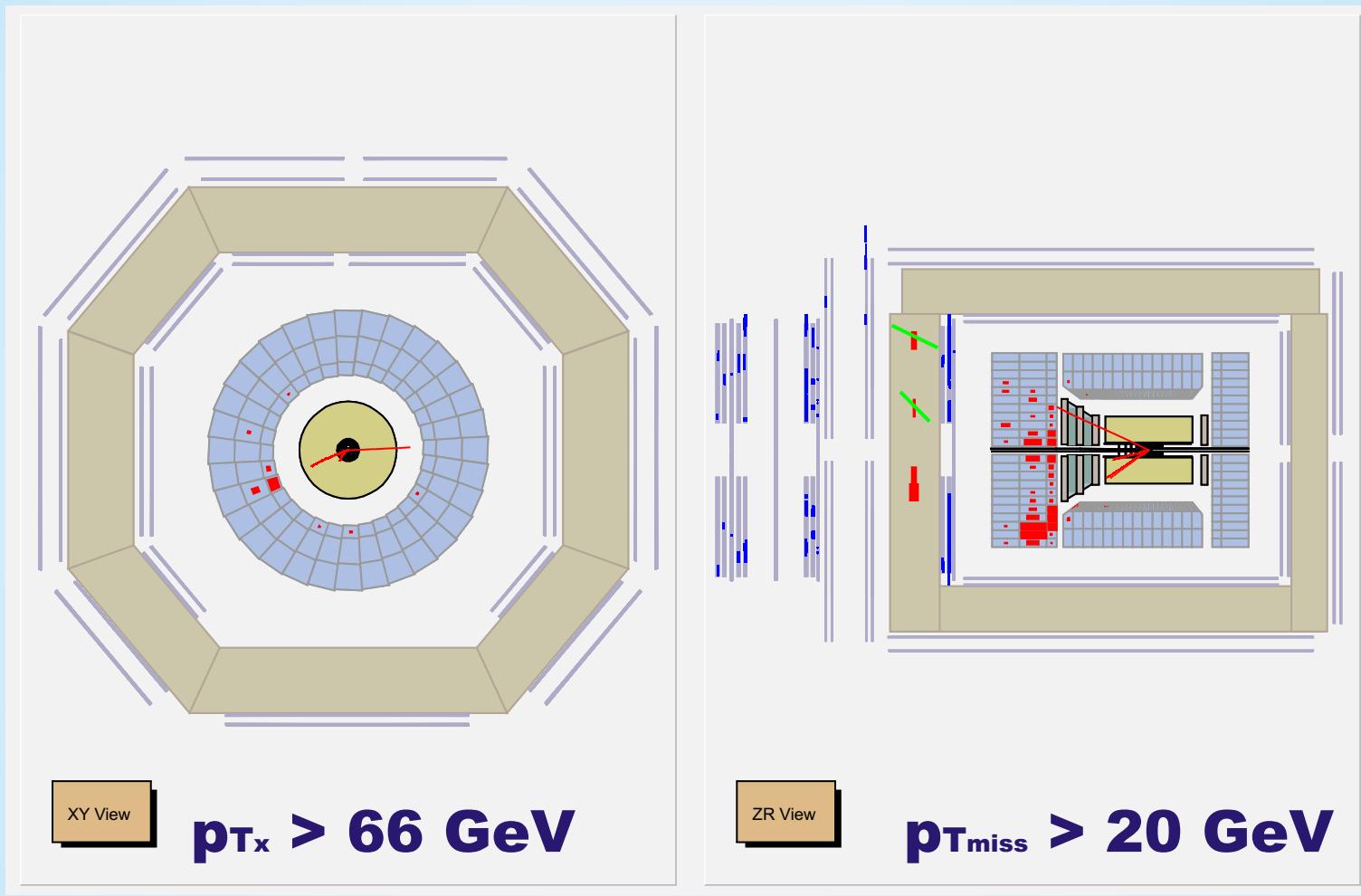
$p_{T_{\text{miss}}} > 44 \text{ GeV}$

$p_{T_l} > 37 \text{ GeV}$

large acoplanarity



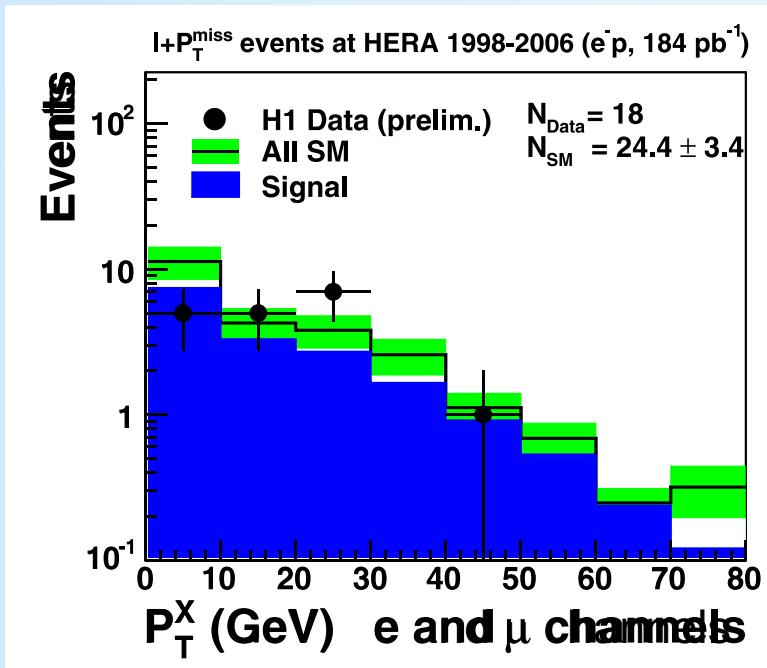
Event in ZEUS



large acoplanarity

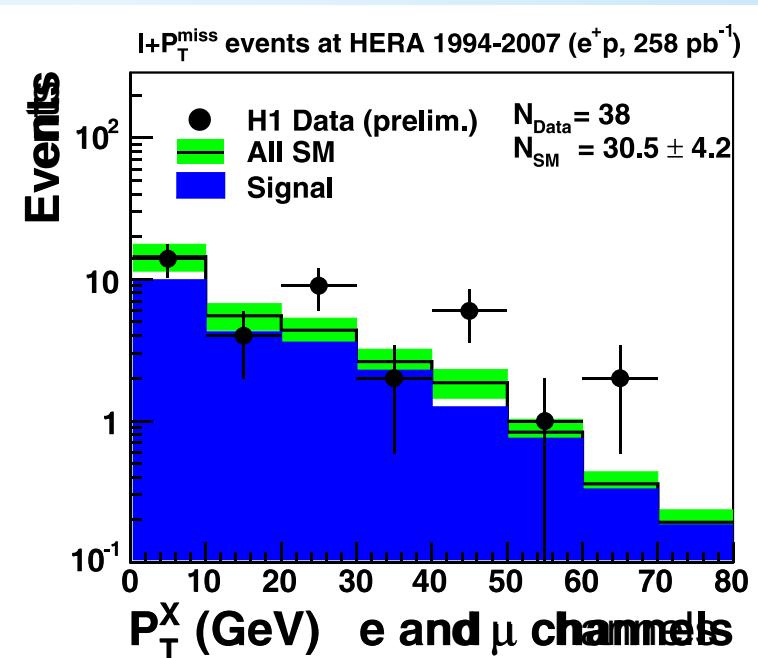
Transverse Momentum Distributions

H1



electrons

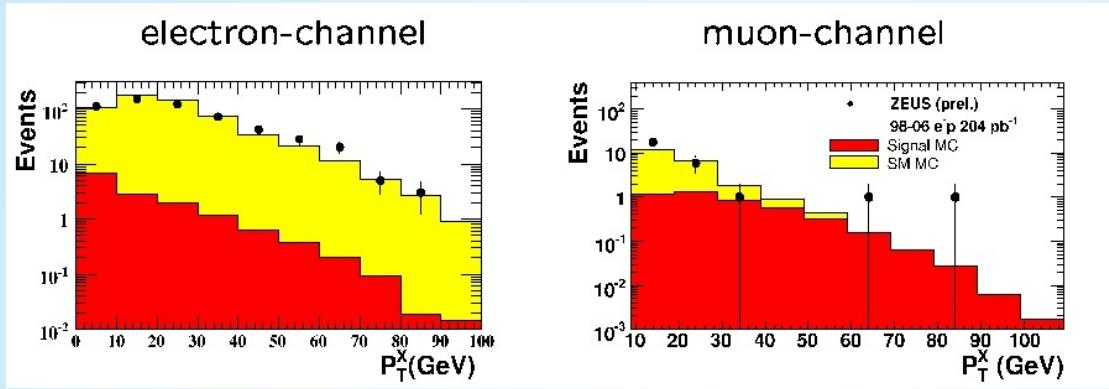
Clean single W signal



positrons

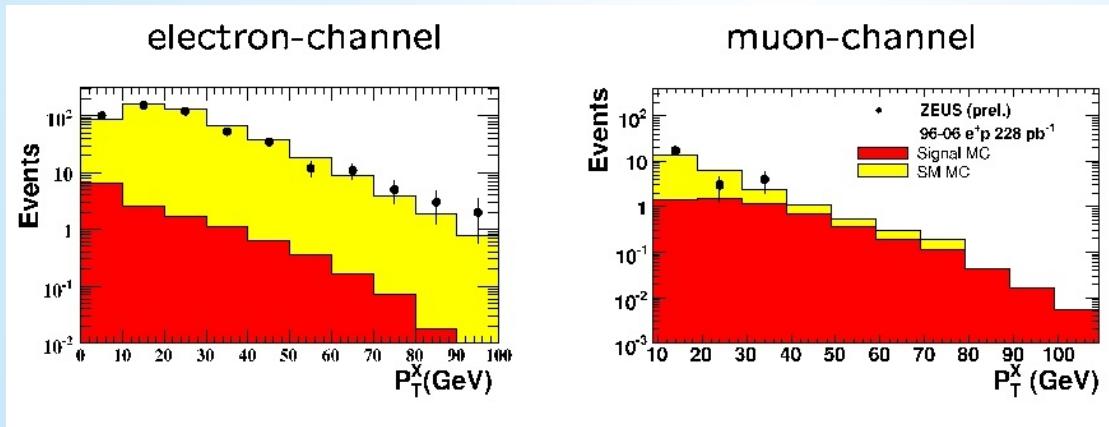
Transverse Momentum Distributions

ZEUS



positrons

misidentified
NC is main
background
in electron
channel



electrons

muon
channel
is a lot
cleaner

Isolated Leptons

$P_T > 25 \text{ GeV}$		electrons Data/SM	muons Data/SM
e⁺p	H1	294 pb^{-1}	$11/4.7 \pm 0.9$
	ZEUS	228 pb^{-1}	$1/3.2 \pm 0.4$
e⁻p	H1	184 pb^{-1}	$3/3.8 \pm 0.6$
	ZEUS	204 pb^{-1}	$5/3.8 \pm 0.6$



**Cristinel Diaconu
DIS 2007**

**Experiments are checking on things
like acceptance:**

**H1 has larger acceptance, but the extra events
are in common region.**

like different cuts:

Sample purities are, however, similar.

Statistics struck again?

Leptoquarks



Compositeness

Extra Dimensions

**Sorry,
NOTHING**

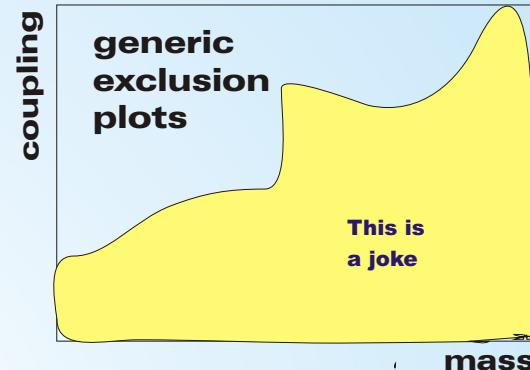
H1 → **see extra talk**
ZEUS → **some examples**

Quark Form Factors

Resonance Production

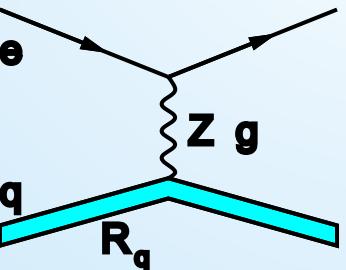
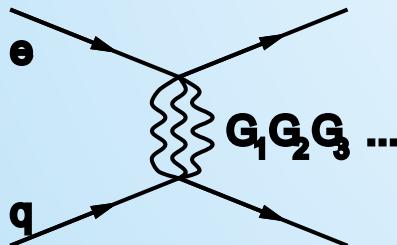
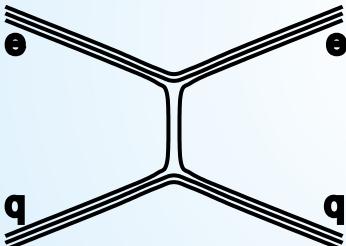
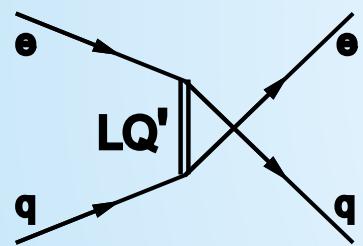
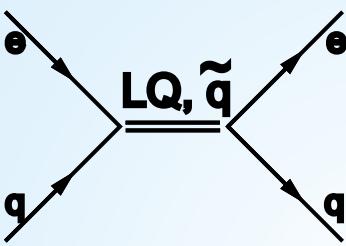
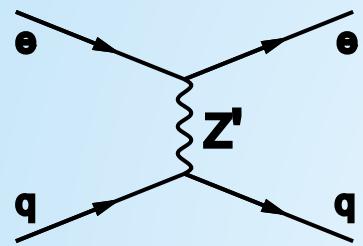
Anomalous Whatever

Supersymmetry

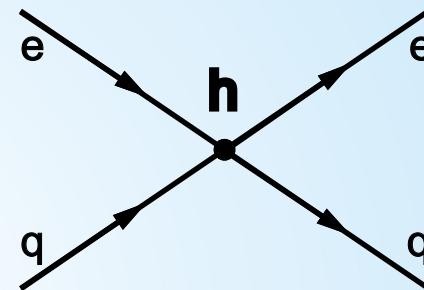


Contact Interactions

Possible “new physics” in NC DIS:



$\sqrt{s} \ll \text{process scale } \Lambda$
 \Rightarrow effective parametrization



eeqq contact interactions (CI)

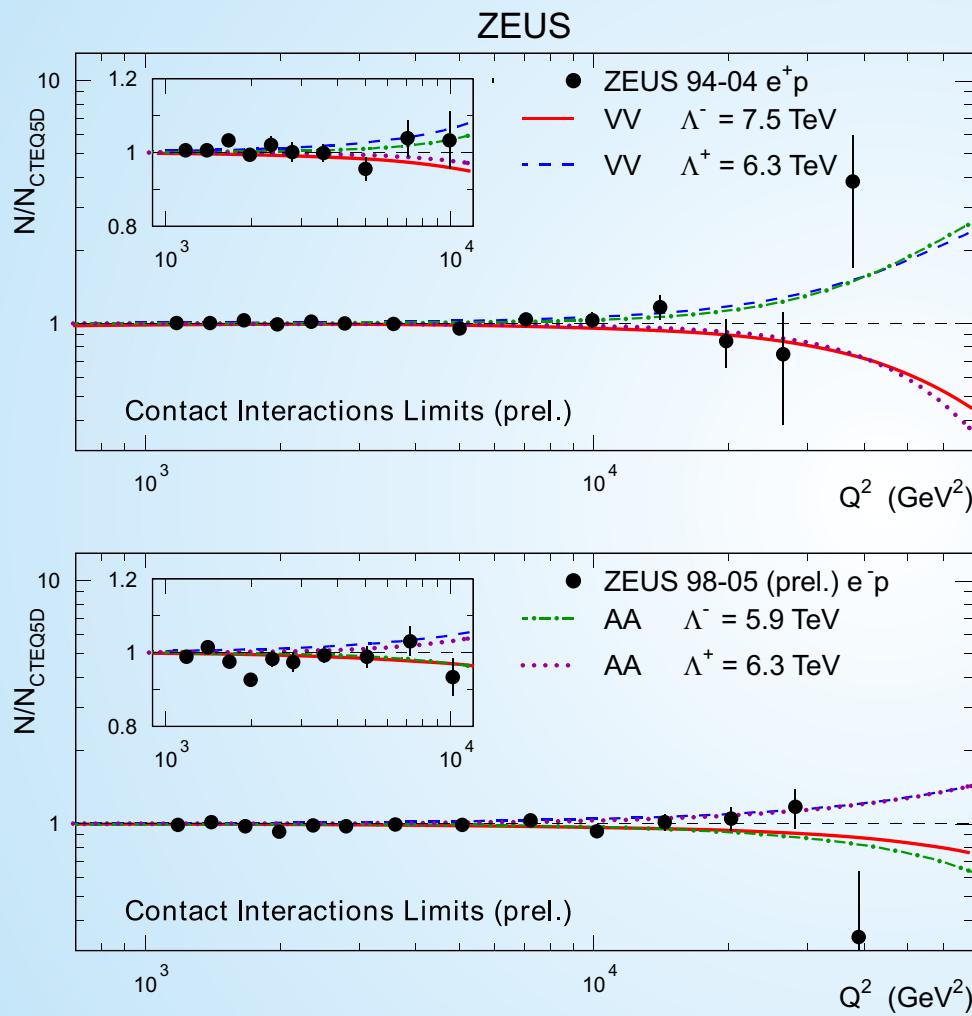
Contribution to the NC cross-section:

$$\frac{d\sigma}{dx dQ^2}(\eta) =$$

$$\frac{d\sigma^{SM}}{dx dQ^2} \cdot [1 + A(x, Q^2) \eta + B(x, Q^2) \eta^2]$$

General formula for all CI type models

Compositeness Models



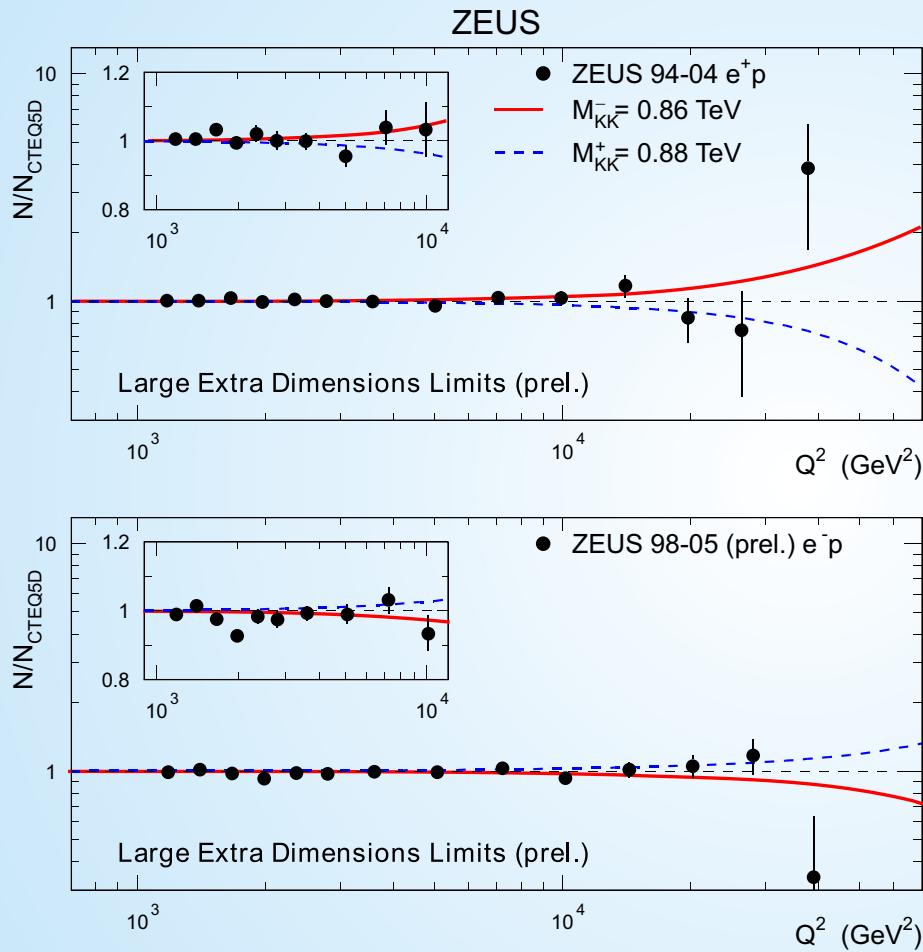
Contact interaction coupling given by:

$$\eta = \pm 4\pi / \Lambda^2$$

Λ : **compositeness scale**

Limits range from 2.0 to 7.5 TeV

Large Extra Dimensions



Kaluza-Klein
Graviton
Exchange

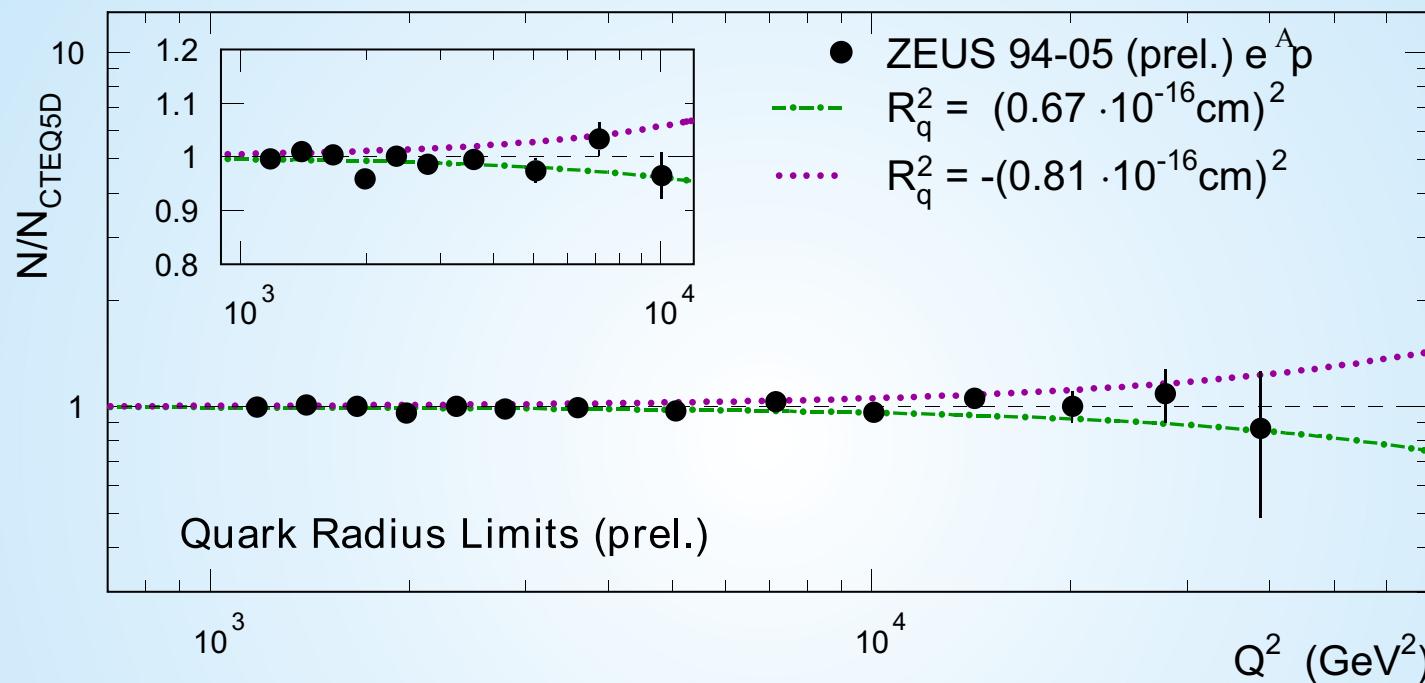
Effective Scale

$$\eta \sim s / M_{KK}^4$$

$M_{KK}^+ > 0.88 \text{ TeV}$

Quark Form Factor

ZEUS



Semi-classical model:
 R_q – quark radius

$$\frac{d\sigma}{dQ^2} = \frac{d\sigma^{\text{SM}}}{dQ^2} \left(1 - \frac{R_q^2}{6} Q^2 \right)$$

$$R_q < 0.67 \cdot 10^{-16} \text{ cm}$$

Summary

**HERA II has provided a large amount of high Q^2 data
We are in the middle of an exciting period of analysis.**

The PDFs will be the legacy and stand for a long time.

The Standard Model has survived HERA without a scratch.

Unfortunately [?] nothing unexpected has happened at HERA.

Still doing the low energy run for gluons.

Shut down June 30, i.e. this Saturday.

