

# Resolved $\gamma^*$ & NLO QCD @ HERA

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on behalf of H1 and ZEUS

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physics at low/medium  $Q^2$ :

- what is this talk about?
- jets, dijets, forward jets
- forward  $\pi^0$
- charm jets
- charm meson production



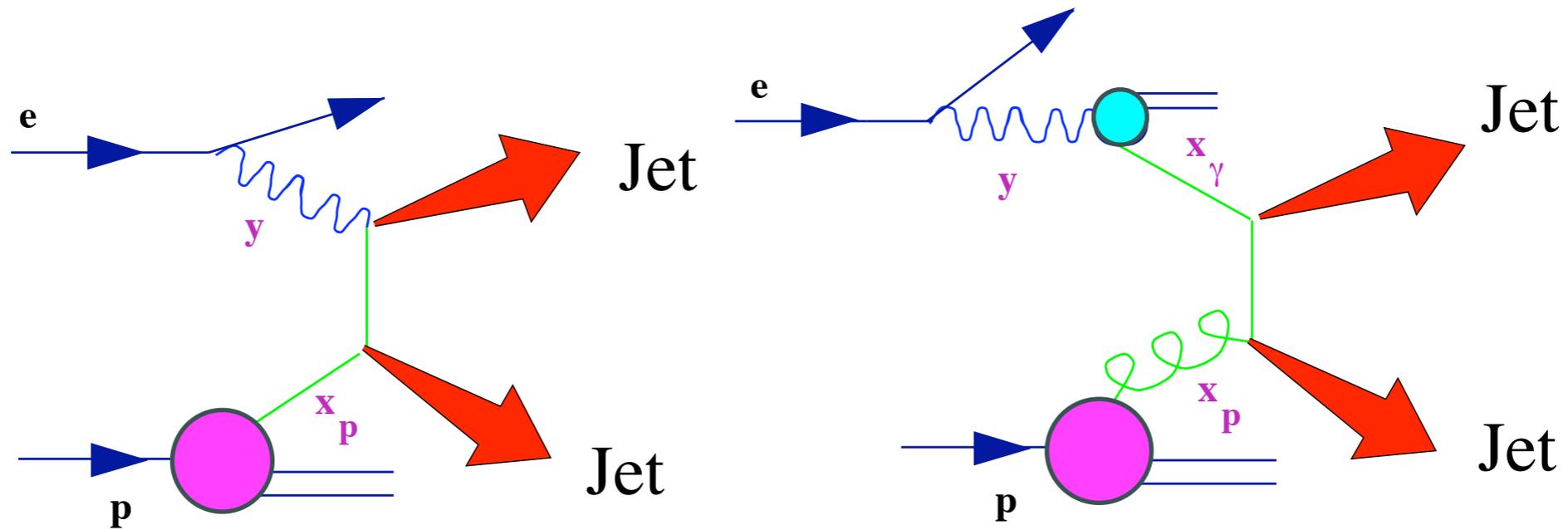
# Introduction

- **What this talk is not about:**
  - **Measurements of hard processes in photoproduction or at high  $Q^2$  at HERA and their comparison with NLO QCD calculations were already covered, for example in talks on jets by Bussey, Jimenez, Brownson, and on heavy flavours by Loizides and Boenig, ...**
    - **Conclusions have been:**
      - **NLO calculations provide reasonable description of the data**
      - **improvements may come from:**
        - **fitting PDFs of the photon using the high jet  $E_T$  data**
        - **also better understanding of multi-parton interactions at low jet  $E_T$**

# Introduction

- What this talk is (should be) about:
  - do NLO calculations of DIS require resolved  $\gamma^*$  contributions in NLO, i.e. beyond what is already included in LO via perturbative photon splitting to describe DIS data?
    - look at the transition region from DIS to photoproduction, i.e.  $\text{few GeV}^2 < Q^2 < 100 \text{ GeV}^2$ , where  $Q^2$  is still a hard scale
    - to probe the photon structure we require another hard scale, e.g. jet or particle  $E_T^2 \gg Q^2$
    - at HERA we have data with  $E_T^2 \approx Q^2$ , but also with  $E_T^2 \gg Q^2$
  - unfortunately, low  $Q^2$  at HERA implies more or less also low  $x$  (or large  $\ln 1/x$ ) and therefore sensitivity to other possible effects:
    - BFKL/CCFM, colour dipole model, cgc, modifying in one way or another the gluon evolution at low  $x$  and leading to effects not described by NLO calcs. in standard collinear factorisation with DGLAP evolution
- 👉 If data are not described by NLO QCD (without photon PDFs), there are many options of what might be missing besides the virtual photon structure

# Introduction



- $\sigma_{jj} \propto f_{\gamma^*/e}(y, Q^2) \otimes f_{i/\gamma^*}(x_\gamma, \mu_R, \mu_{F\gamma^*}) \otimes f_{j/p}(x_p, \mu_{Fp}) \otimes \sigma_{ij}(x_\gamma, x_p, \mu_R)$ 
  - $f_{i/\gamma^*}(x_\gamma, \mu_R, \mu_{F\gamma^*}) = f_{i/\gamma^*}^{\text{pert}} + f_{i/\gamma^*}^{\text{non-pert}}$
  - simplicity lets us typically set all scales equal to  $\mu^2 = E_T^2$  or  $E_T^2 + Q^2$
  - $x_\gamma = (E_{T,1} e^{-\eta_1} + E_{T,2} e^{-\eta_2}) / 2yE_e$  or  $x_\gamma = \sum_{1,2} (E - P_z) / \sum_{\text{had}} (E - P_z)$ ;  $\Delta\Phi = \Phi_1 - \Phi_2$

# MC tools for “DIS”

- **MEPJET<sup>1</sup>, DISENT<sup>2</sup>, DISASTER++<sup>3</sup>, NLOJET++<sup>4</sup> for dijet production at NLO**
- **NLOJET++ for trijet production at NLO**
  - 1 Mirkes, Zeppenfeld; 2 Seymour, Catani; 3 Graudenz;  
4 Nagy , Trocsany; 5 Pötter, Klasen, Kramer
- **JETVIP<sup>5</sup> for dijet production at NLO; it is based on the phase space slicing method. This MC simulates direct processes for dijets in NLO (like e.g. DISENT) and it allows to include resolved processes at NLO, using PDFs of the virtual photon, BUT**
  - the main author Björn Pötter has left HEP some years ago
  - for direct processes: disagreement with other programs (in some regions of phase space for distributions differential in jet variables)
  - for resolved processes: problem with the stability of NLO resolved (slicing parameter  $y_c$ )
  - 🖱️ difficult to draw quantitative conclusion
- **Program by Fontannaz et al. for forward hadron ( $\pi^0$ ) production, including direct and resolved processes at NLO**
- **all other NLO calculations contain a perturbative resolved contrib. in LO only**

# Inclusive jets: $5 < Q^2 < 100 \text{ GeV}^2$ , $0.2 < y < 0.7$

- H1 preliminary result

- $43.5 \text{ pb}^{-1}$
- jets: incl. kt algo in Breit frame
- $E_T > 5 \text{ GeV}$
- $-1 < \eta_{\text{lab}} < 2.5$

- NLOJET++

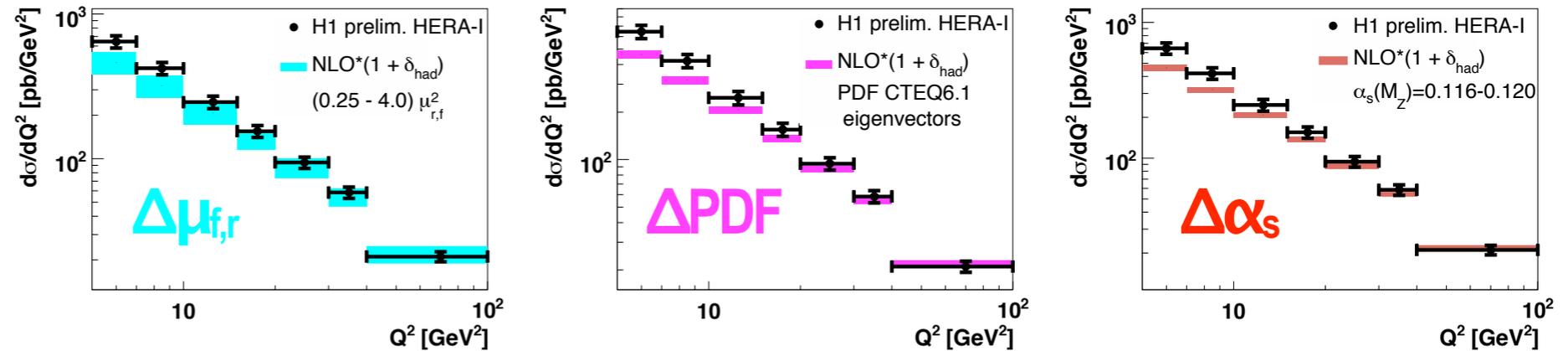
- CTEQ6.1M
- $\mu_R^2 = E_T^2$ ,  $\mu_F^2 = Q^2$
- corr. for had. effects

■ NLO QCD describes data for  $Q^2 > 10$  to  $20 \text{ GeV}^2$  & jet  $E_T > 10 \text{ GeV}$

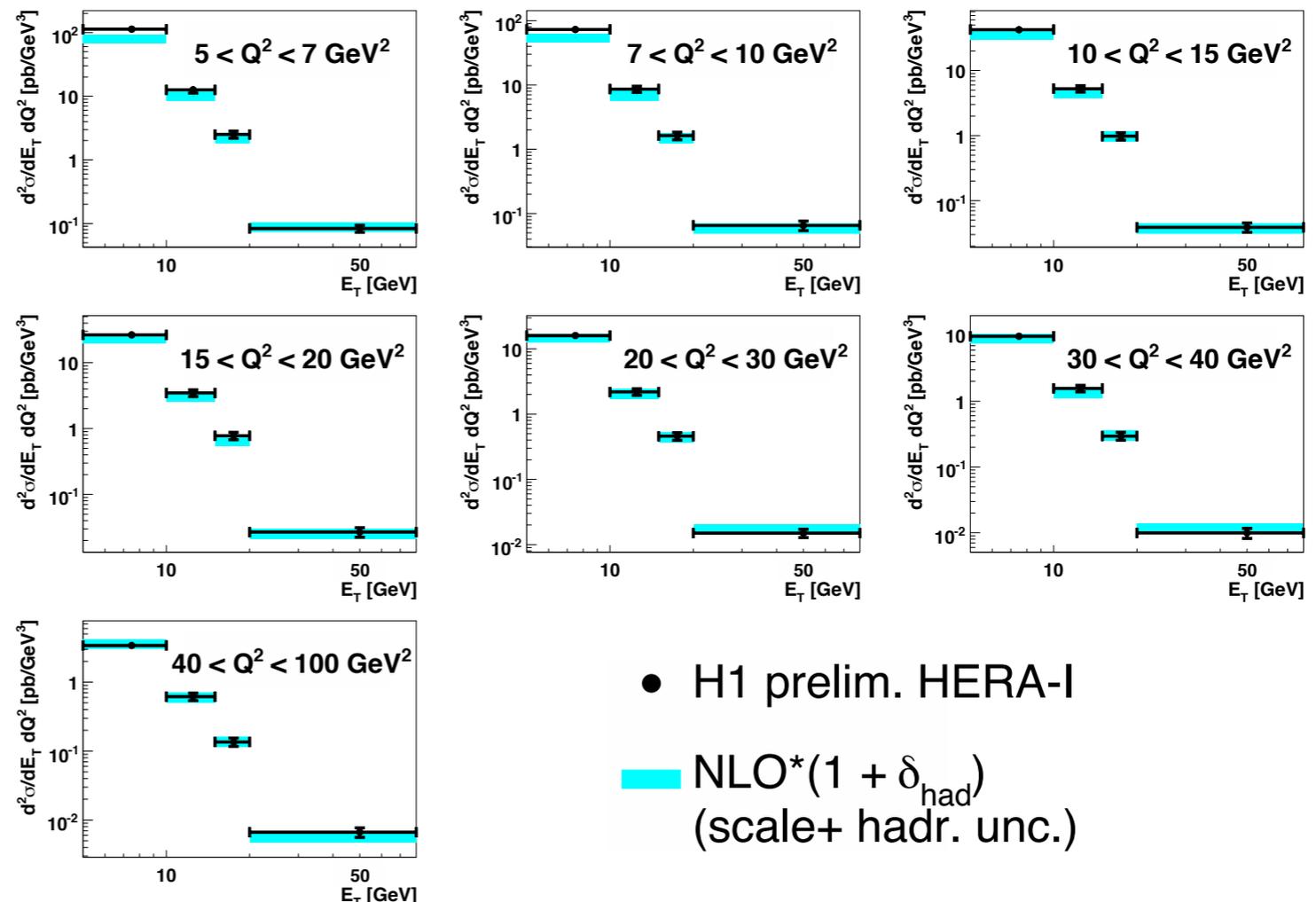
■ large scale uncertainties, perhaps underestimated?

■ room for higher orders or resolved contrib. at low  $Q^2$

## H1 Inclusive Jet Cross Sections $\frac{d\sigma}{dQ^2}$



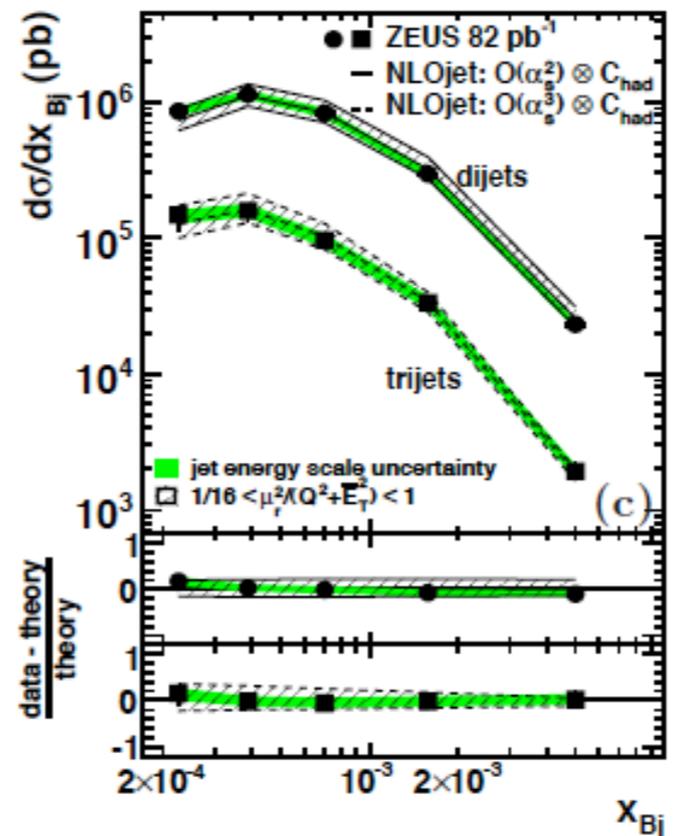
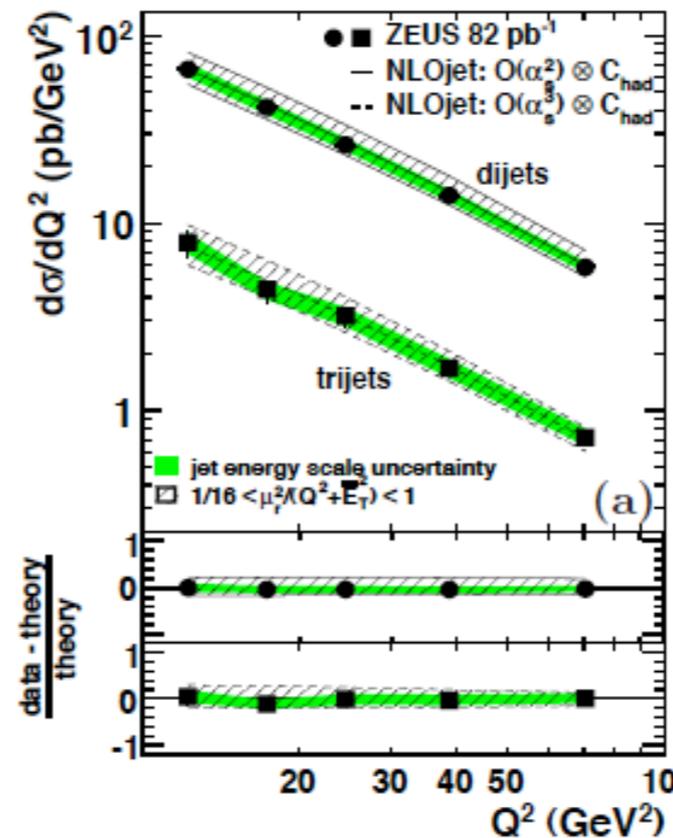
## H1 Inclusive Jet Cross Sections $\frac{d^2\sigma}{dQ^2 dE_T}$



- H1 prelim. HERA-I
- NLO\*(1 +  $\delta_{\text{had}}$ ) (scale+ hadr. unc.)

# Dijets & trijets: $10 < Q^2 < 100 \text{ GeV}^2$ , $0.1 < y < 0.6$

- ZEUS, DESY-07-062, arXiv:0705.1931
  - 82 pb<sup>-1</sup>
  - jets: incl. kt algo in HCM frame
  - $E_{T1} (E_{T2,3}) > 7 (5) \text{ GeV}$
  - $-1 < \eta_{\text{lab}} < 2.5$
- NLOJET++:  $O(\alpha_s^2)$  &  $O(\alpha_s^3)$ 
  - CTEQ6M
  - $\mu_r^2 = \mu_f^2 = (Q^2 + \langle E_T \rangle^2) / 4$
  - correct for had. effects



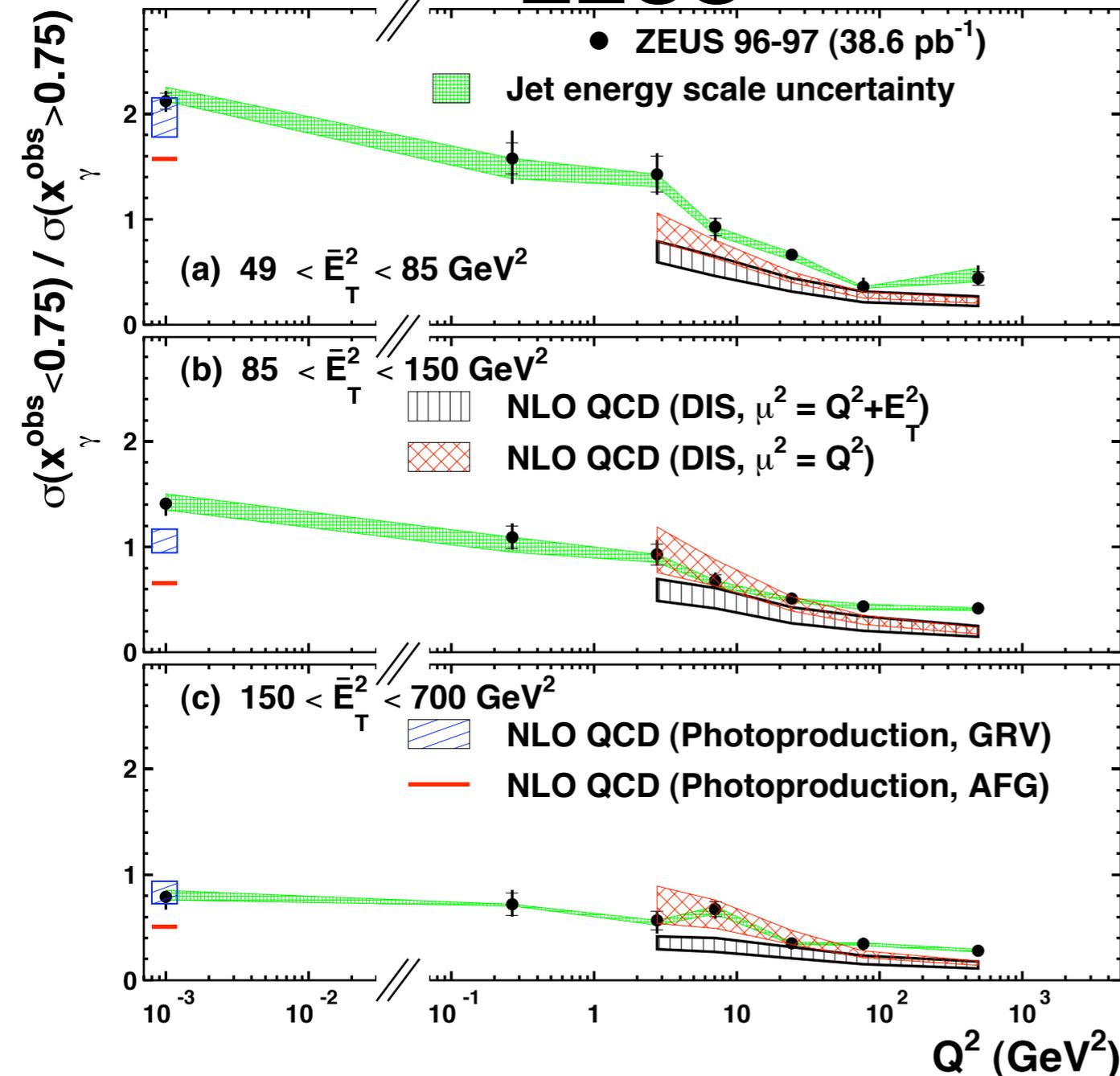
☞  $d\sigma/dQ^2$  &  $d\sigma/dx_{Bj}$  for dijets and trijets well described by NLO QCD for  $Q^2 \geq 10 \text{ GeV}^2$ ,  $x_{Bj} \geq 1.7 \cdot 10^{-4}$

☞ no need for resolved contributions beyond pert. LO

# Dijets: $0 \leq Q^2 < 2000 \text{ GeV}^2$ , $0.2 < y < 0.55$

$\sigma(\text{"resolved"}) / \sigma(\text{"direct"})$

## ZEUS

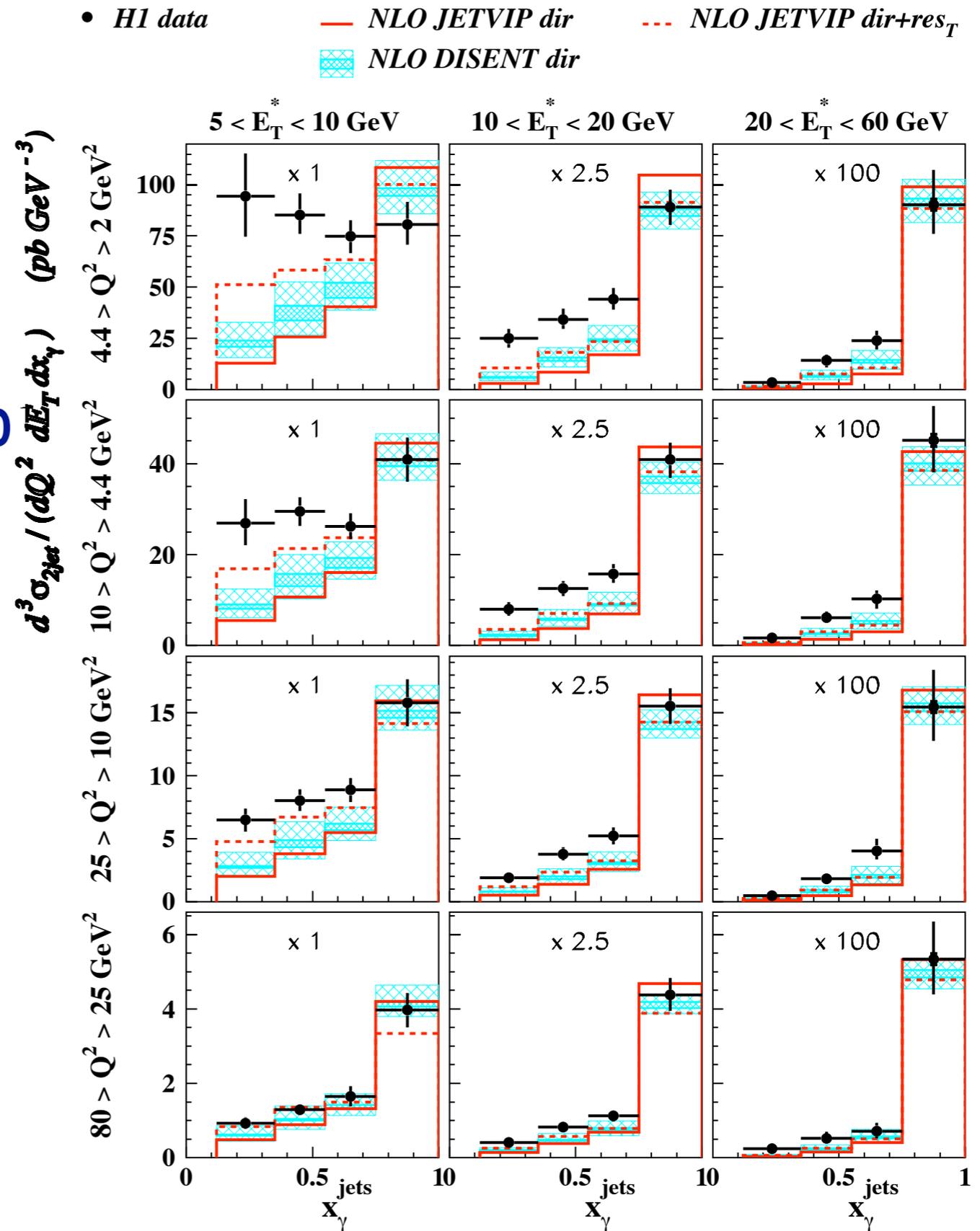


- ZEUS EPJC 35, 487, 2004 ( $38.6 \text{ pb}^{-1}$ )
- jets: incl.  $k_t$  algo in the  $\gamma^*p$ -frame
  - $E_{T1} (E_{T2}) > 7.5 (6.5) \text{ GeV}$
  - $-3 < \eta_{1,2}^* < 0$
  - $\sigma(x_\gamma < 0.75) / \sigma(x_\gamma > 0.75)$  vs.  $Q^2$  and for three regions in mean jet  $E_T^2$
- NLO calculations:
  - FMNR,  $\mu^2 = E_T^2$ , CTEQ5M1 & GRV or AFG
  - DISASTER++,  $\mu^2 = Q^2 + E_T^2$ , CTEQ5M1
- photoprod. reasonably well described
- DIS: NLO QCD below the data
- with  $\mu^2 = Q^2$  as scale in DIS, the theory uncertainty increases significantly
- previous new result, with smaller scale  $\mu^2$ , agrees with data !?!

# Dijets: $2 \leq Q^2 < 80 \text{ GeV}^2$ , $0.1 < y < 0.85$

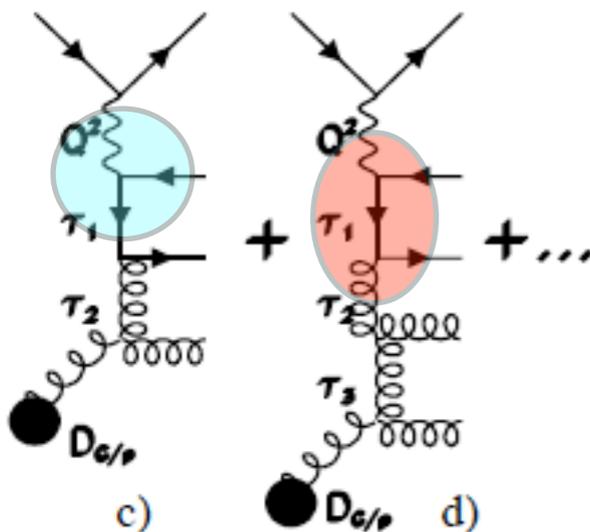
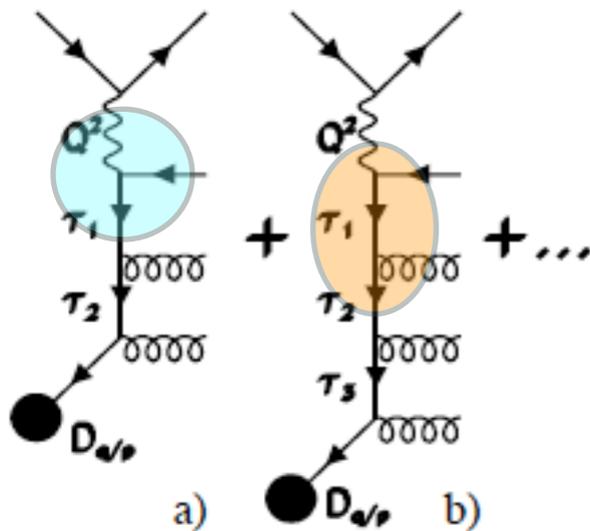
- H1 EPJ C37, 141, 2004 ( $57\text{pb}^{-1}$ )
- jets: incl.  $k_t$  algo in the  $\gamma^*p$ -frame
  - $E_{T1} (E_{T2}) > 7 (5) \text{ GeV}$
  - $-2.5 < \eta_{1,2} < 0$
- DISSENT:  $\mu_r^2 = E_{T1}^2$ ,  $\mu_f = 9 \text{ GeV}$ , CTEQ6M
- JETVIP:  $\mu_r^2 = \mu_f^2 = E_{T1}^2$ , CTEQ6M, SAS1D
- for DISSENT hadronisation and scale uncertainty are shown
- $x_\gamma > 0.75$  reasonably well described by direct NLO calculations
- $x_\gamma < 0.75$  data not well described, particularly at small  $x_\gamma$ , low  $Q^2$  and jet  $E_T$
- JETVIP dir  $\neq$  DISSENT; both are LO calc. for  $x_\gamma < 1$
- JETVIP dir+res still below the data for small  $x_\gamma$

mean jet  $E_T$

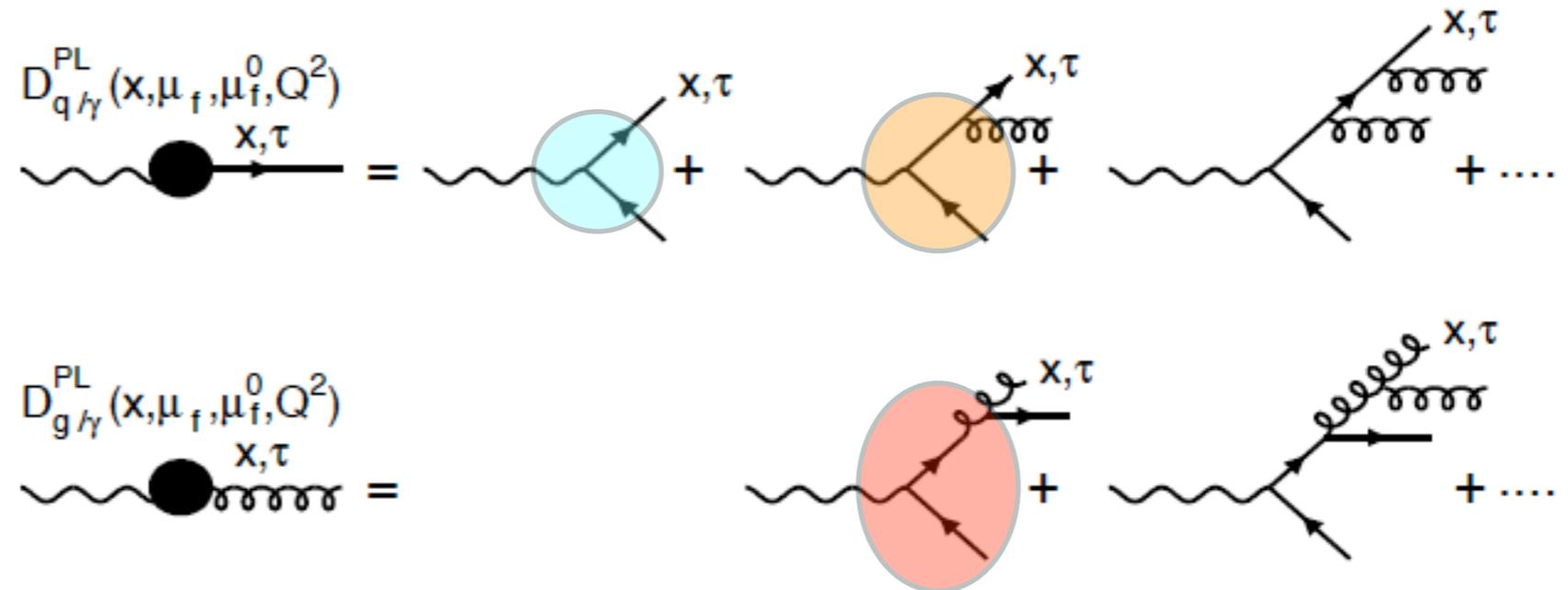


# NLOJET++: dijets at NLO & trijets at NLO

3-jets at LO    3-jets at NLO



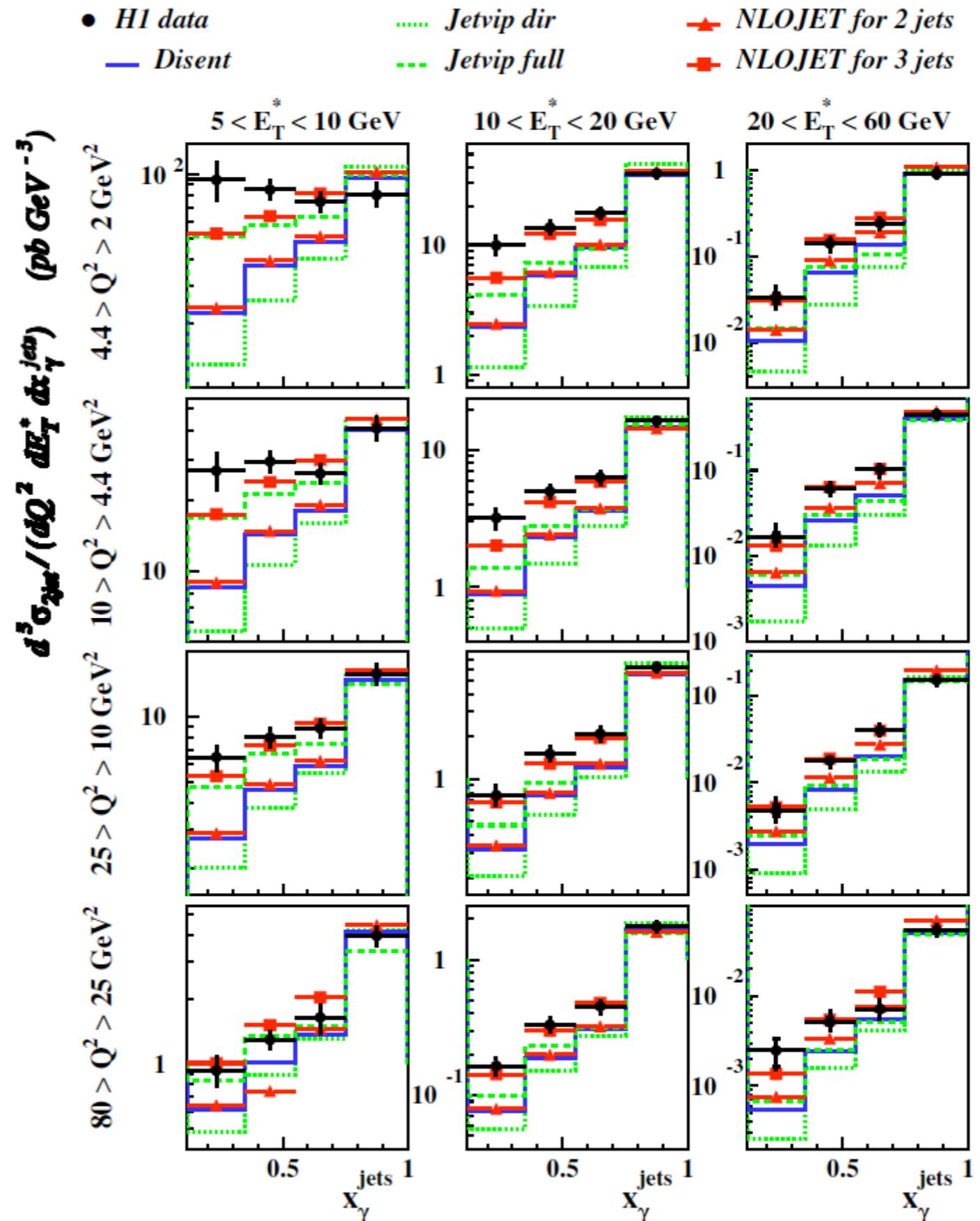
point-like part of the photon structure



- point-like part of the photon structure provides an approximate method for including higher order direct contributions
- $x_\gamma < 1$  is populated by events with at least three jets (at parton level), for programs like DISSENT, JETVIP dir this is LO only; with NLOJET++ can calculate at NLO

# Dijets: $2 \leq Q^2 < 80 \text{ GeV}^2$ , $0.1 < y < 0.85$

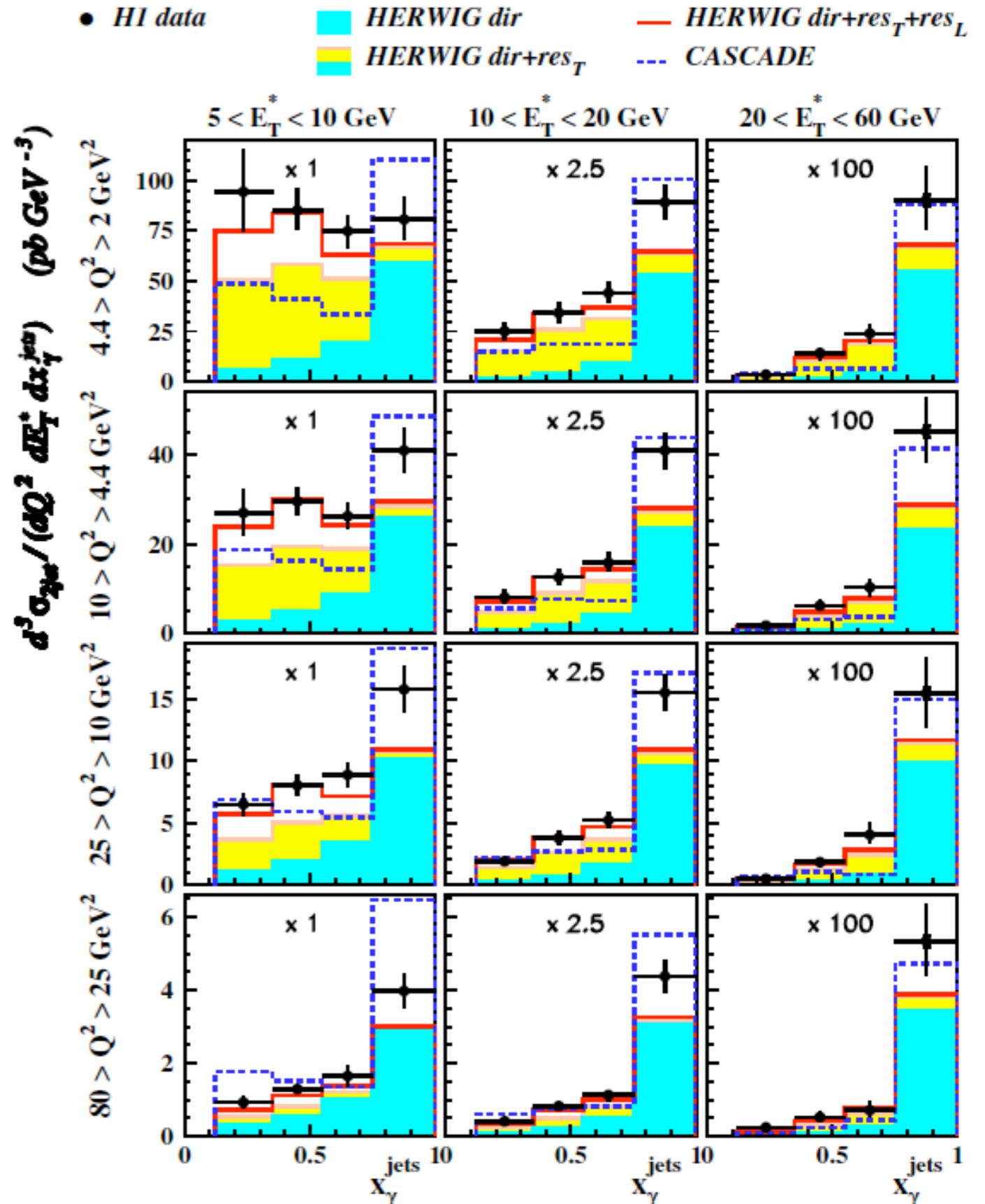
- H1 data as on previous page & J.Chyla et al., EPJ C40, 469, 2005 (NLOJET++ predictions)
- good agreement for 2-jets at NLO, i.e. **DISENT = NLOJET++**
- best description of data for 3-jets at NLO using **NLOJET++** (k-factor  $\approx 2$  at lowest  $x_\gamma$  and  $Q^2$  bin !)
- remaining difference, mainly at low  $x_\gamma$  and  $Q^2$ , indicate need for still higher orders or resummation as in case of photon structure
- what are the scale uncertainties for **NLOJET++** (3-jets) ?
- need reliable full NLO calculation of resolved contribution



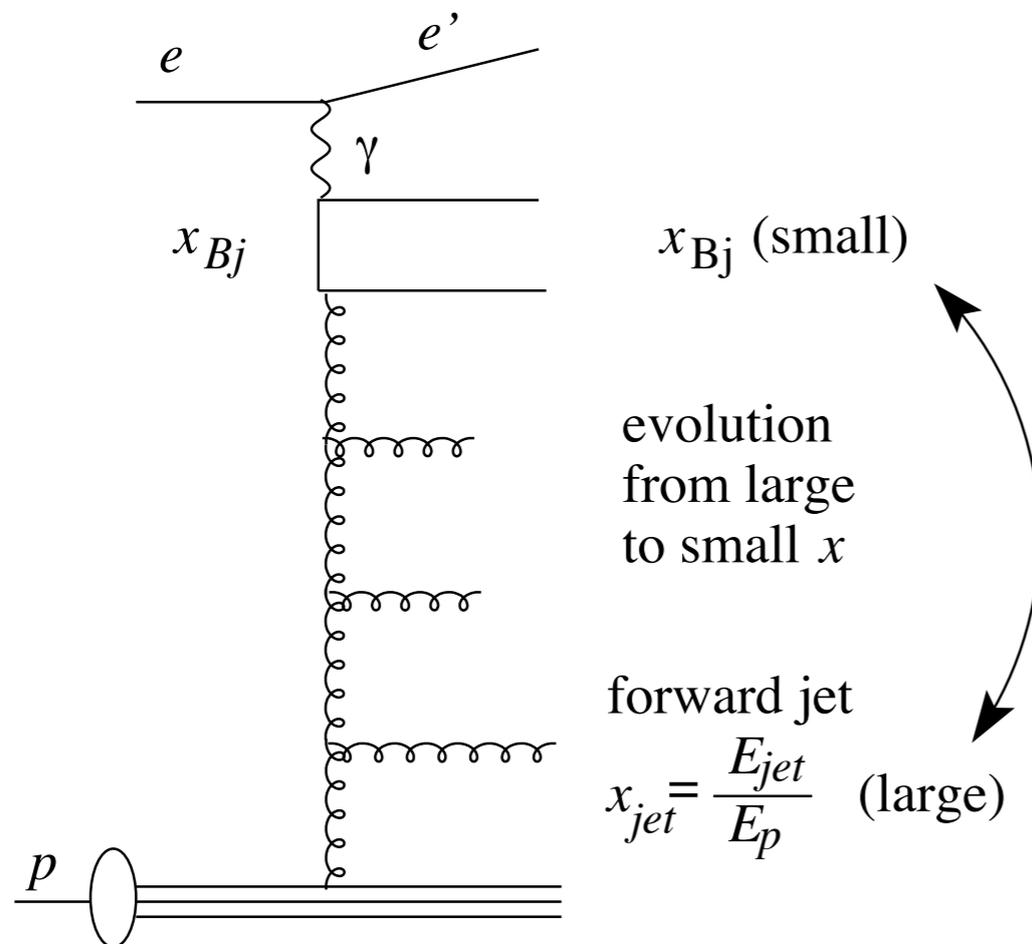
# Dijets: $2 \leq Q^2 < 80 \text{ GeV}^2$ , $0.1 < y < 0.85$

a small aside on LO+PS MC:  
a comparison with HERWIG

- same dijet data as on previous page
- best description of these data by HERWIG, LO+PS and with a virtual photon structure for transverse and longitudinal photons (calc. by J.Chyla)



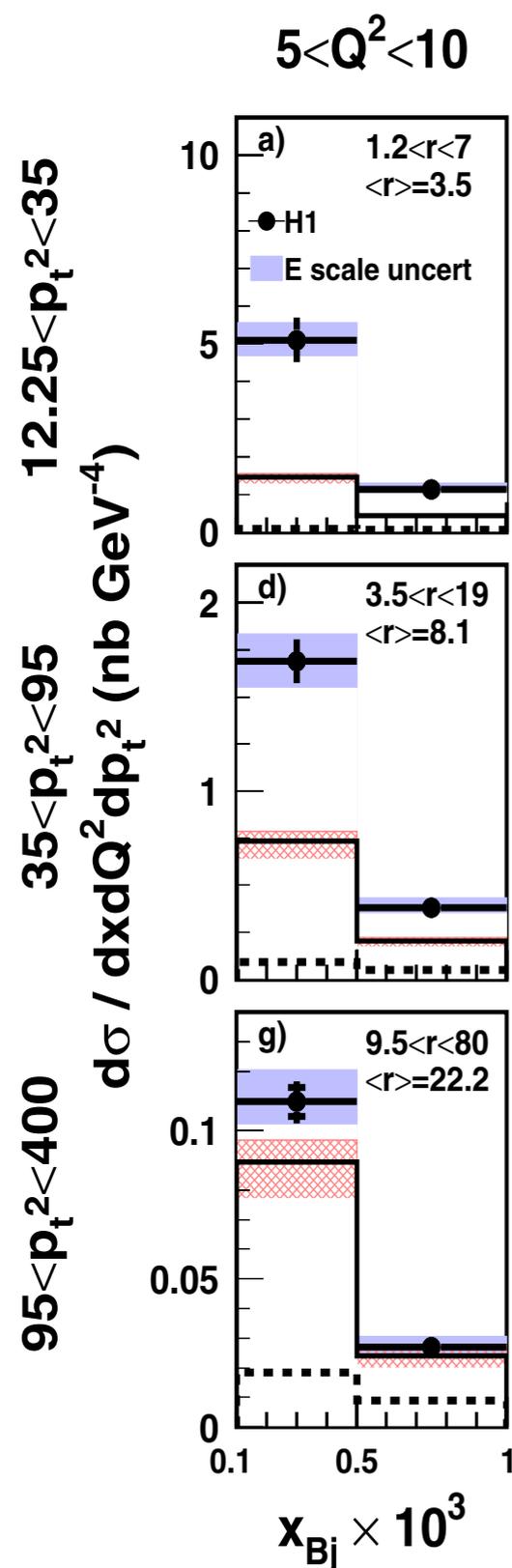
# Forward jets in DIS: $5 < Q^2 < 85 \text{ GeV}^2$ , $0.1 < y < 0.7$



H1 Collab., Eur. Phys. J. C **46** (2006) 27  
[\[arXiv:hep-ex/0508055\]](https://arxiv.org/abs/hep-ex/0508055)

- in DGLAP the strong ordering in virtuality gives softest  $p_t$  gluon closest to proton
- suppress DGLAP:  $p_{T,jet}^2 \sim Q^2$
- in BFKL the gluon  $p_T$  close to the proton can be hard; strong ordering occurs in  $x$
- enhance BFKL:  $x_{jet} \gg x_{Bj}$
- fwd jet (incl.  $k_T$  in Breit frame) & cuts in HERA frame
  - $7^\circ (2.79) < \theta_{jet} (\eta_{jet}) < 20^\circ (1.74)$
  - $p_{T,jet} > 3.5 \text{ GeV}$
  - $x_{jet} = E_{jet}/E_p > 0.035$
- enhance resolved contributions  $\rightarrow$  large  $r = p_{T,jet}^2/Q^2$

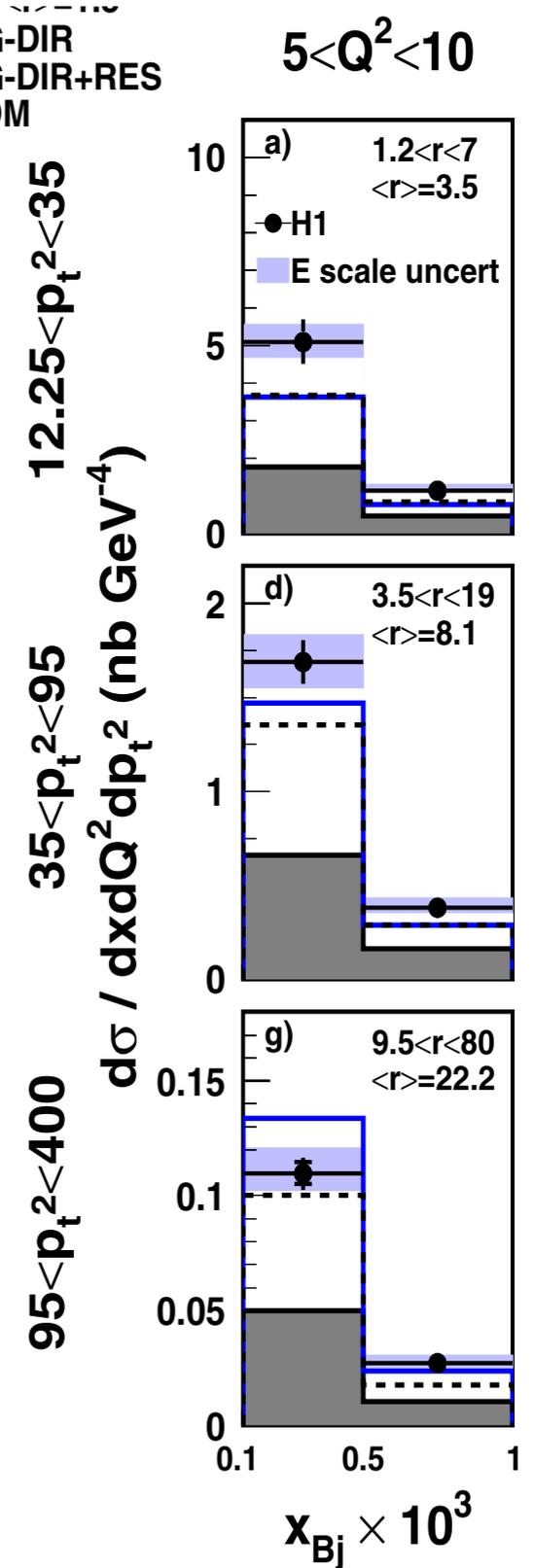
# $d\sigma/dx_{Bj}$ for events with large $r = p_{T,jet}^2/Q^2$



— NLO DISENT  $1+\delta_{HAD}$   
 0.5 $\mu_{r,f}<\mu_{r,f}<2\mu_{r,f}$   
 PDF uncert.  
 - - - LO DISENT  
 $1+\delta_{HAD}$

- NLO significantly below data for low  $x_{Bj}$
- LO  $\ll$  NLO, fwd-jets in LO suppressed by kinematics  $\rightarrow$  NLO  $\approx$  LO for fwdjets
- bands includes scale ( $1/4$  &  $4 \mu^2 = p_{T,jet}^2$ ) and PDF uncertainty; using  $Q^2$  as scale the prediction increases by 35% at low  $x_{Bj}$  and  $Q^2$  and the uncertainty increases
- RAPGAP direct fails, direct+resolved and CDM give good description of data except at lowest  $x_{Bj}$

■ RG-DIR  
 - - - RG-DIR+RES  
 — CDM

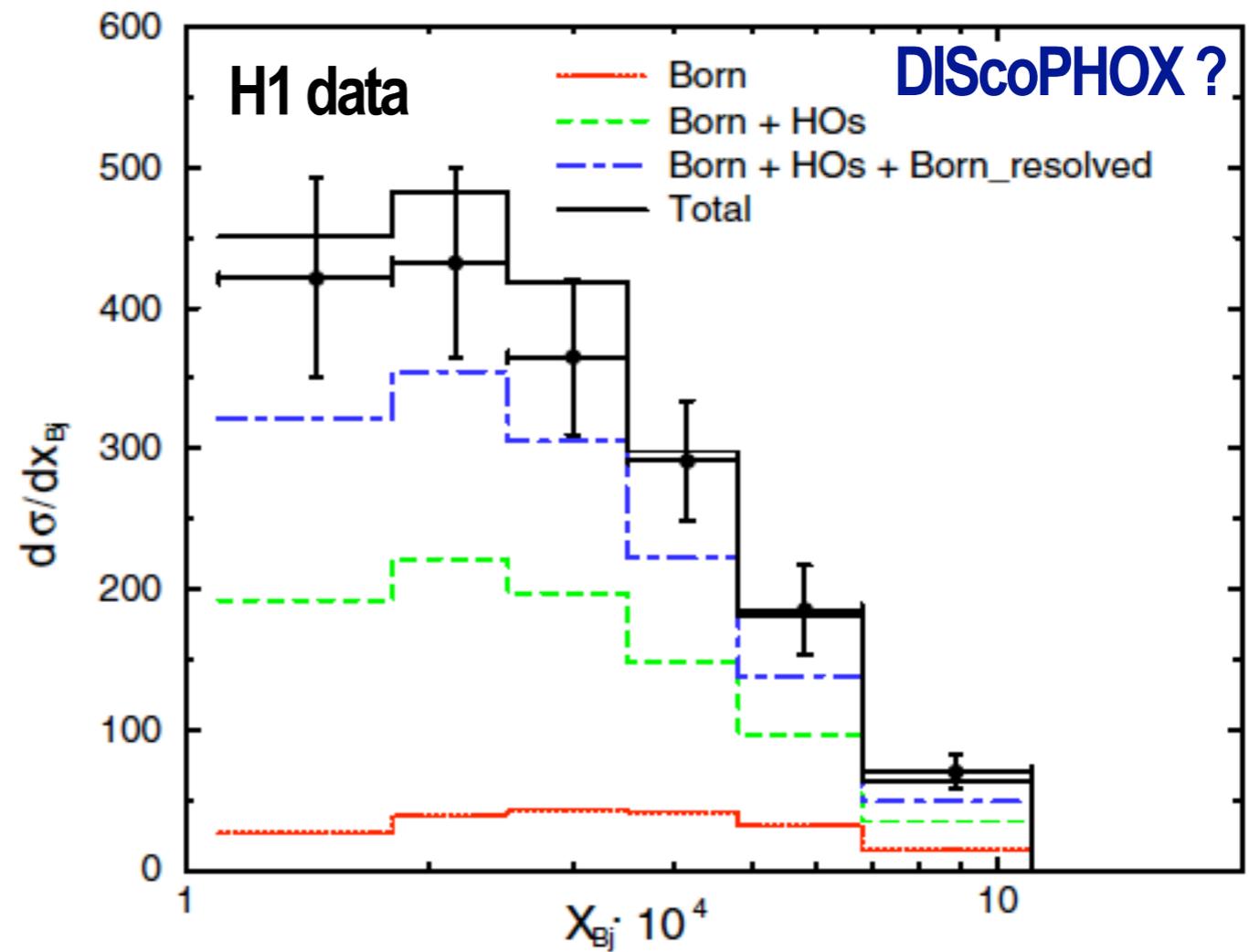


see paper for more

- triple diff. x-sections ( $x_{Bj}$ ,  $Q^2$ ,  $p_{T,jet}$ )
- x-sections for events with additional dijets

# $d\sigma/dx_{Bj}$ for events with a forward $\pi^0$

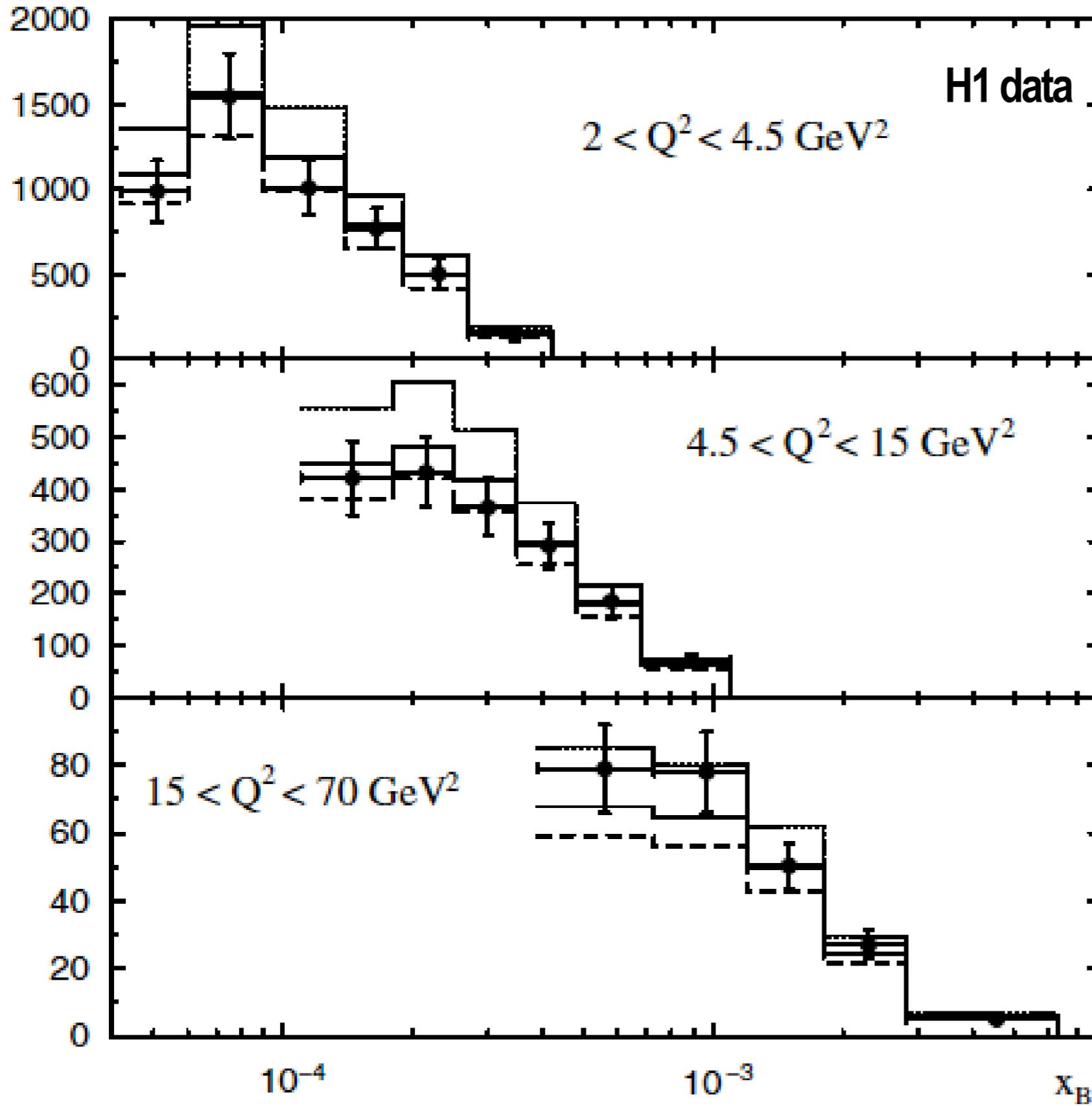
- H1: EPJ C 36, 441 (2004); 21pb-1
  - $4.5 (2) < Q^2 < 15 (70) \text{ GeV}^2$
  - $0.1 < y < 0.6$
  - $5^\circ < \theta_\pi < 25^\circ$
  - $x_\pi > 0.1$
  - $E_{T,\pi}^* > 2.5 \text{ GeV}$
- NLO calc. by Fontannaz
  - includes virtual photon struct. in NLO
  - CTEQ6M,  $\gamma^*$  PDF also by Fontannaz
  - all scales =  $\mu^2 = E_{T,\pi}^{*2} + Q^2$
  - Kniehl, Kramer, Pötter frag. function



NLO from Aurenche et al., EPJ C 42, 43 (2005)

- ☞ good description of the data
- ☞ all corrections LO dir to NLO dir, LO resolved to NLO resolved are large (at least for the chosen scale)

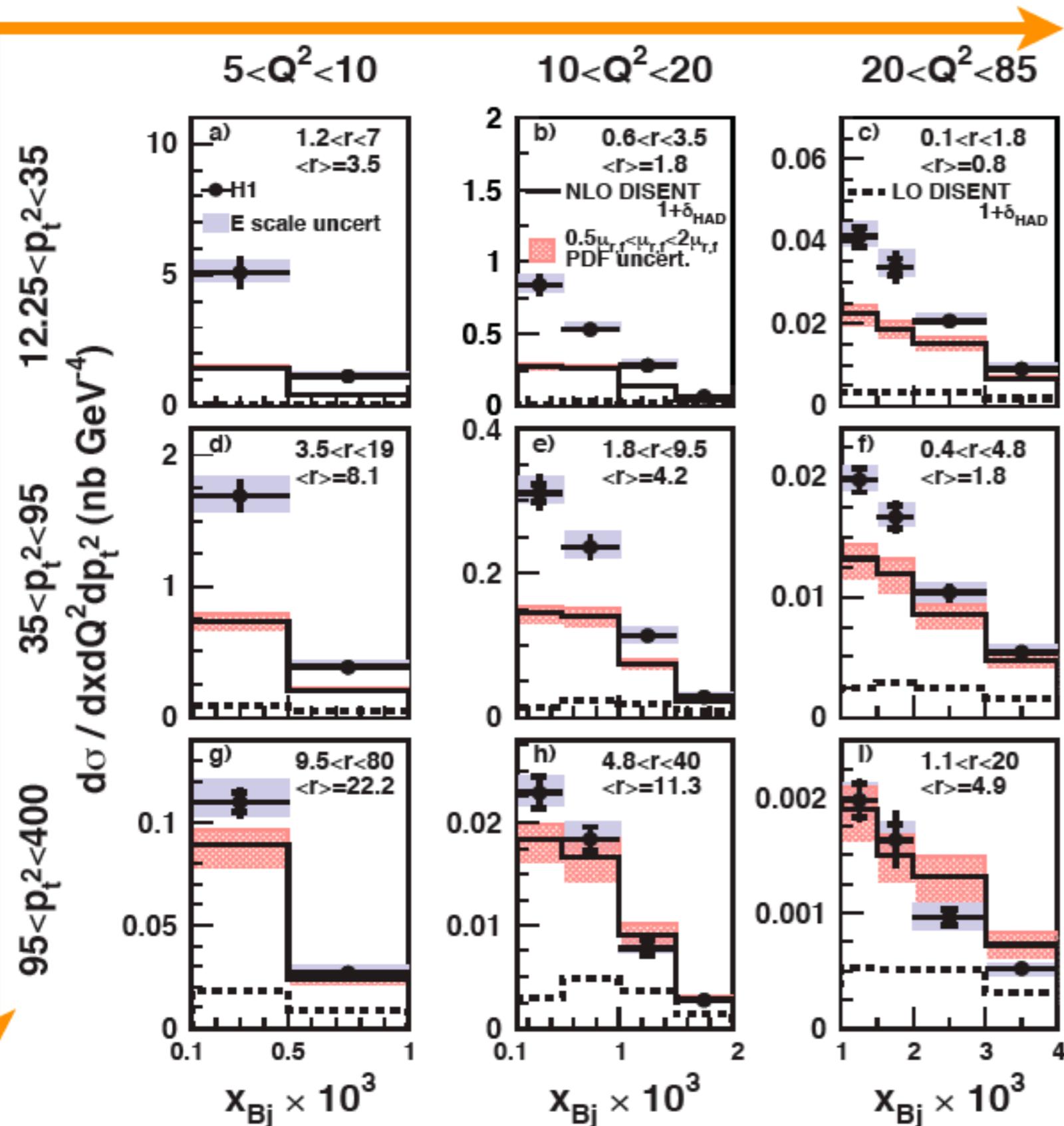
# $d\sigma/dx_{Bj}$ for events with a fwd $\pi^0$ : scale dep.



NLO from Aurenche et al.,  
EPJ C 42, 43 (2005)

- $\mu^2 = 0.5 (E_{T,\pi}^* + Q^2)$
- $\mu^2 = E_{T,\pi}^* + Q^2$
- $\mu^2 = 2 (E_{T,\pi}^* + Q^2)$
-  large scale dependence; see detailed study in theory paper

# Fwd-jet: more differential distributions



- would be interesting to have these data & data on fwd-jet + dijets compared to the NLO calc. by Fontannaz !
- remember  $r = p_{T,\text{jet}}^2 / Q^2$ 
  - large  $r \rightarrow$  resolved region

# $D^{*\pm}$ , D meson & $D^{*\pm}$ + jet production

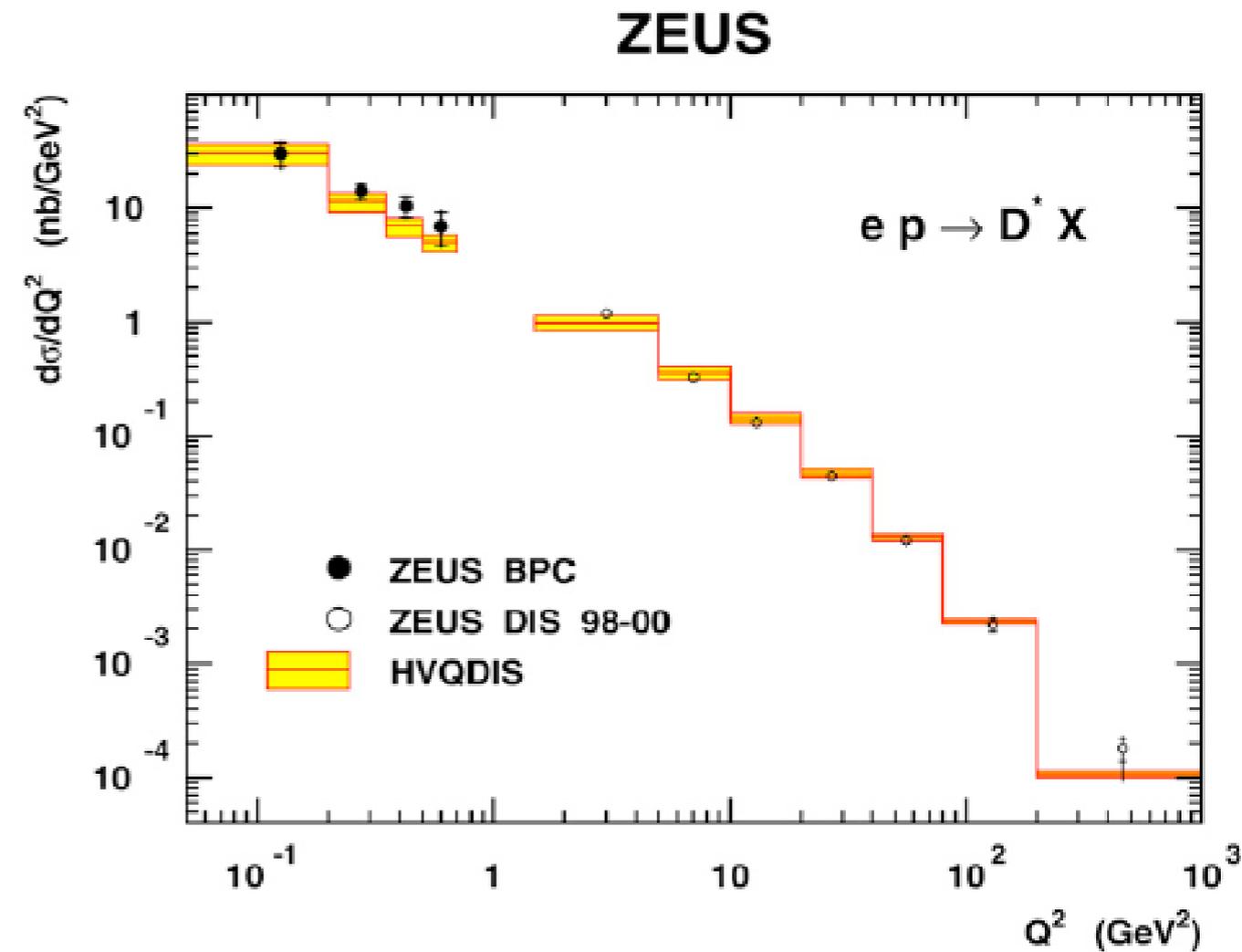
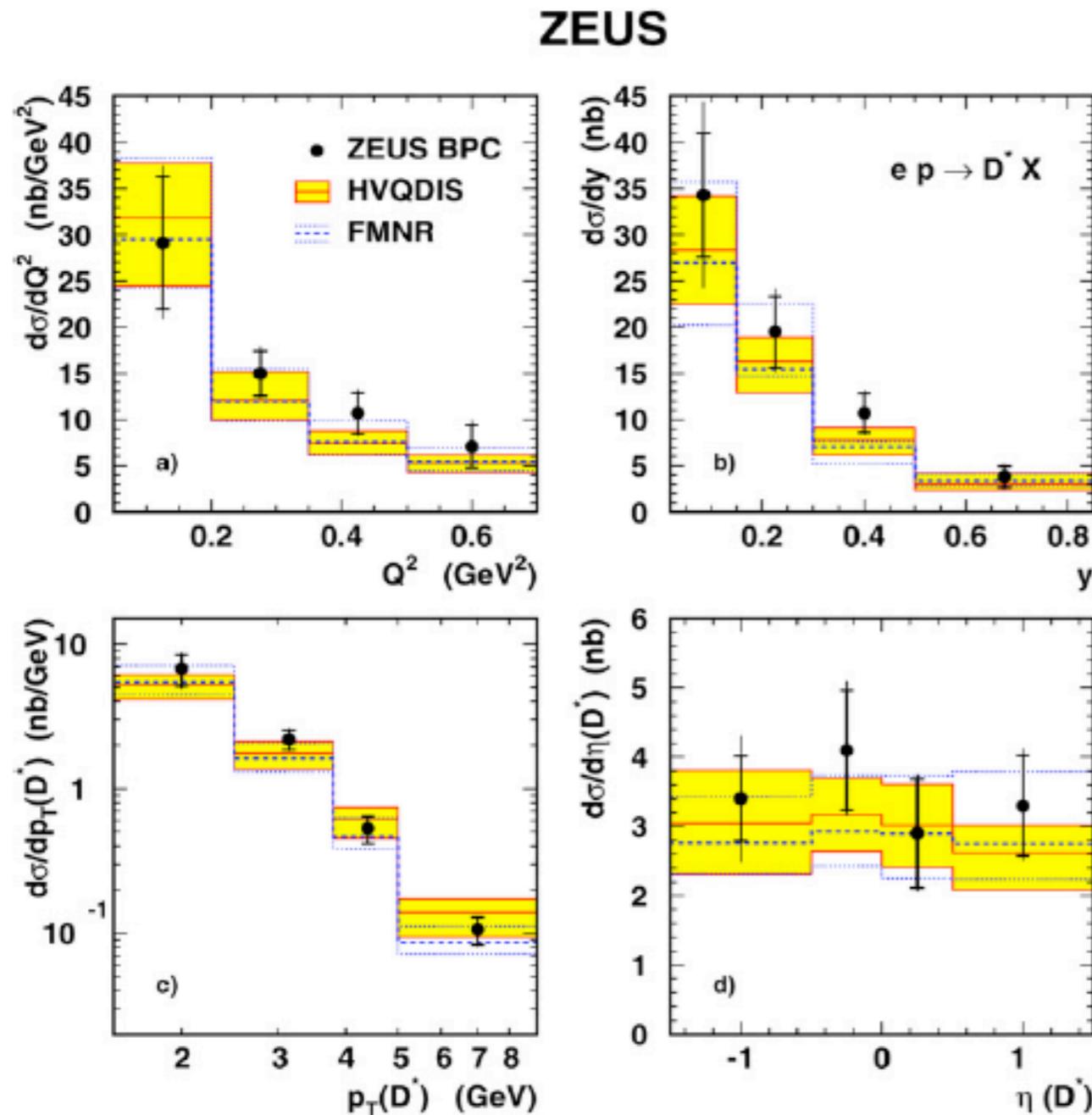
## ■ ZEUS:

- “Measurement of  $D^{*\pm}$  meson production in  $e^+p$  scattering at low  $Q^2$ ”, PLB 649 (2007) 111
  - $0.05 < Q^2 < 0.7 \text{ GeV}^2$ ,  $0.02 < y < 0.85$ ;  $1.5 < p_T(D^*) < 9 \text{ GeV}$ ,  $|\eta(D^*)| < 1.5$ ;  $82 \text{ pb}^{-1}$
  - compare data to FMNR (+WW) & HVQDIS using ZEUS PDF-fit in FFNS,  $m_c = 1.35 \text{ GeV}$ ,  $\mu_r^2 = \mu_f^2 = p_T^2 + m_c^2$  (FMNR),  $\mu_r^2 = \mu_f^2 = Q^2 + 4m_c^2$  (HVQDIS)
  - also compare to data from PRD 69 012004 (2004) with  $1.5 < Q^2 < 1000 \text{ GeV}^2$
- “Measurement of D mesons prod. in DIS at HERA”, DESY-07-052, arXiv:0704.3562, JHEP

## ■ H1:

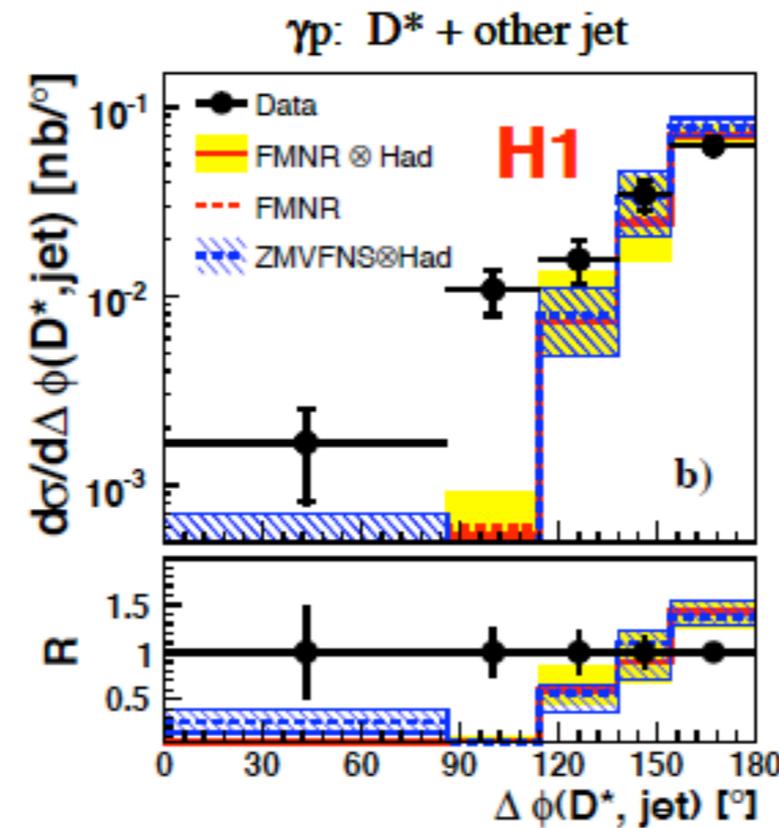
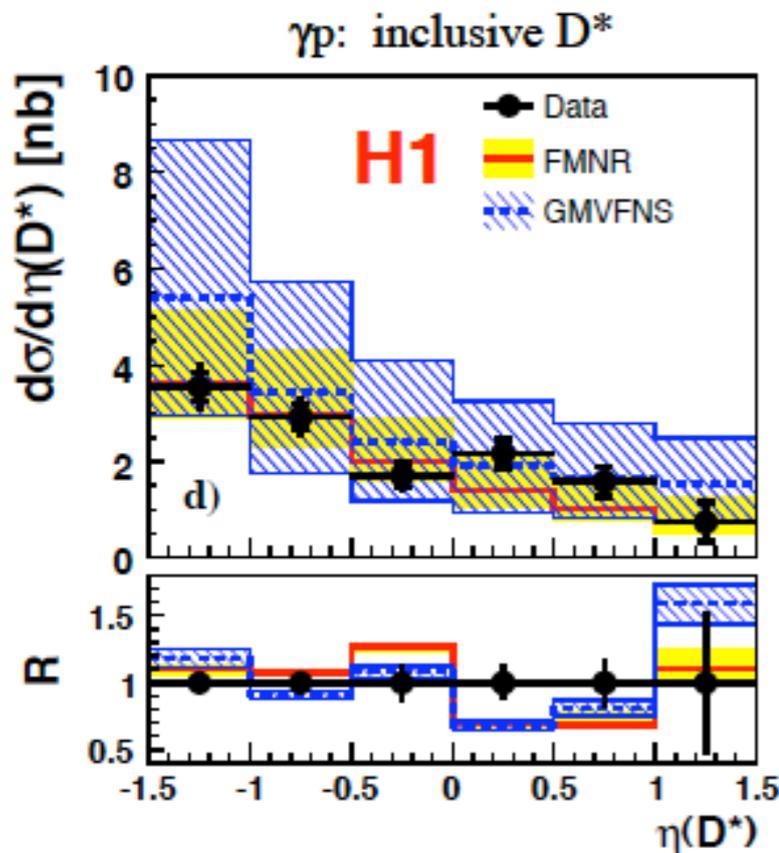
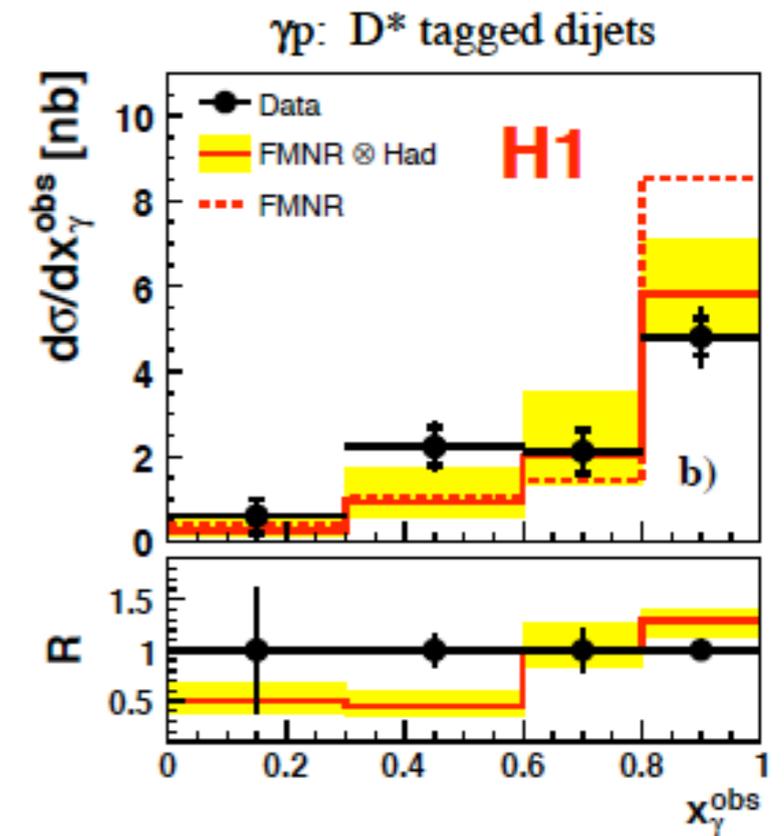
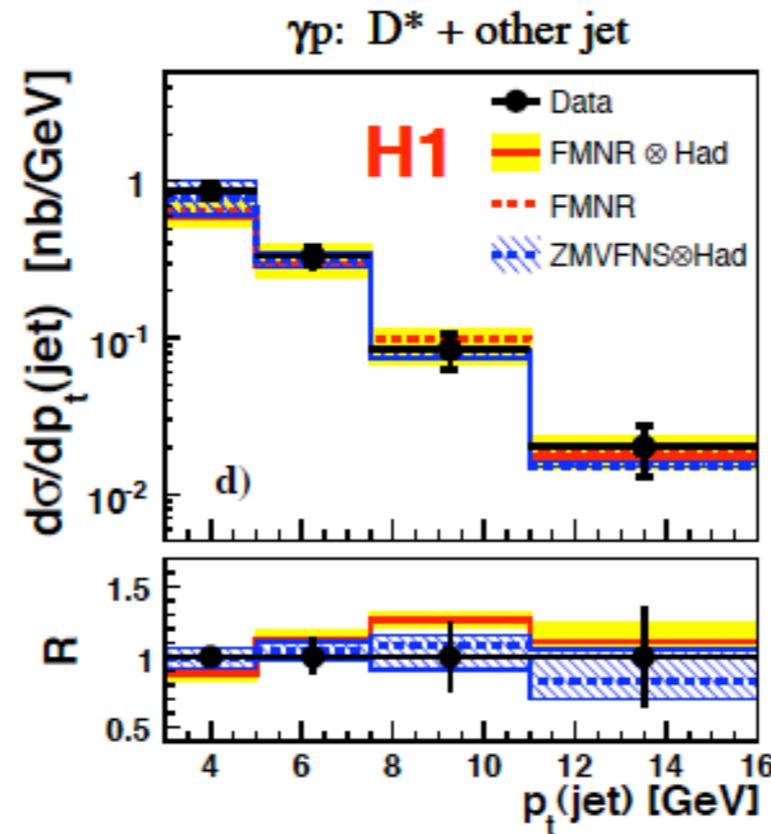
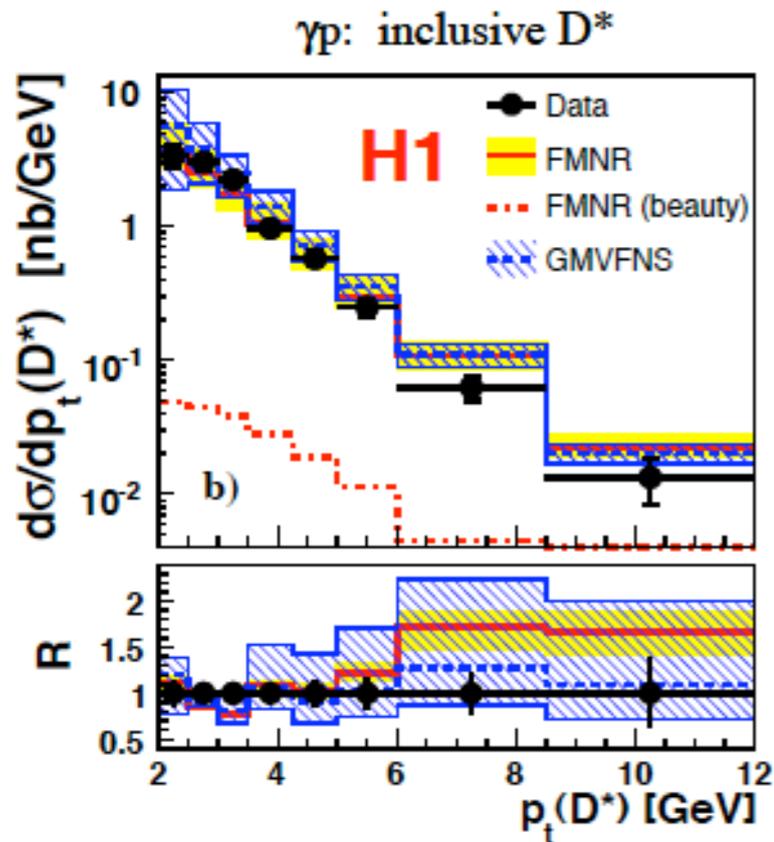
- “Incl.  $D^{*\pm}$  meson cross sects. and  $D^{*\pm}$  - jet correlations in photoprod. ...”, EPJ C50 (2007) 251
  - $Q^2 < 0.01 \text{ GeV}^2$ ,  $0.29 < y < 0.65$ ; a)  $p_T(D^*) > 2 \text{ GeV}$ ,  $|\eta(D^*)| < 1.5$ ;  $51 \text{ pb}^{-1}$
  - b)  $D^{*\pm}$  + other jet:  $p_{T,\text{jet}} > 3 \text{ GeV}$ ; c)  $D^{*\pm}$  tagged dijet:  $p_{T,\text{jet1}} (p_{T,\text{jet2}}) > 4 (3) \text{ GeV}$
  - compare data to FMNR using CTEQ6M, GRV-G HO,  $m_c = 1.5 \text{ GeV}$ ,  $\mu_r^2 = m_c^2 + (p_{T,c}^2 + p_{T,c}^2)/2$ ,  $\mu_f^2 = 2\mu_r^2$
- “Production of  $D^{*\pm}$  mesons with dijets in DIS at HERA”, EPJ C51 (2007) 271
  - $2 < Q^2 < 100 \text{ GeV}^2$ ,  $0.05 < y < 0.7$ ;  $1.5 < p_T(D^*) < 15 \text{ GeV}$ ,  $|\eta(D^*)| < 1.5$ ;  $47 \text{ pb}^{-1}$
  - compare data to HVQDIS using CTEQ5F3,  $m_c = 1.5 \text{ GeV}$ ,  $\mu_r^2 = \mu_f^2 = Q^2 + 4m_c^2$

# D\*± meson production in “DIS”



- $d\sigma/dQ^2$  at very low  $Q^2$  described by HVQDIS and FMNR (within large exp. and theory uncertainties)
- HVQDIS describes the whole range in  $Q^2$
- no resolved contr. needed beyond LO in HVQDIS

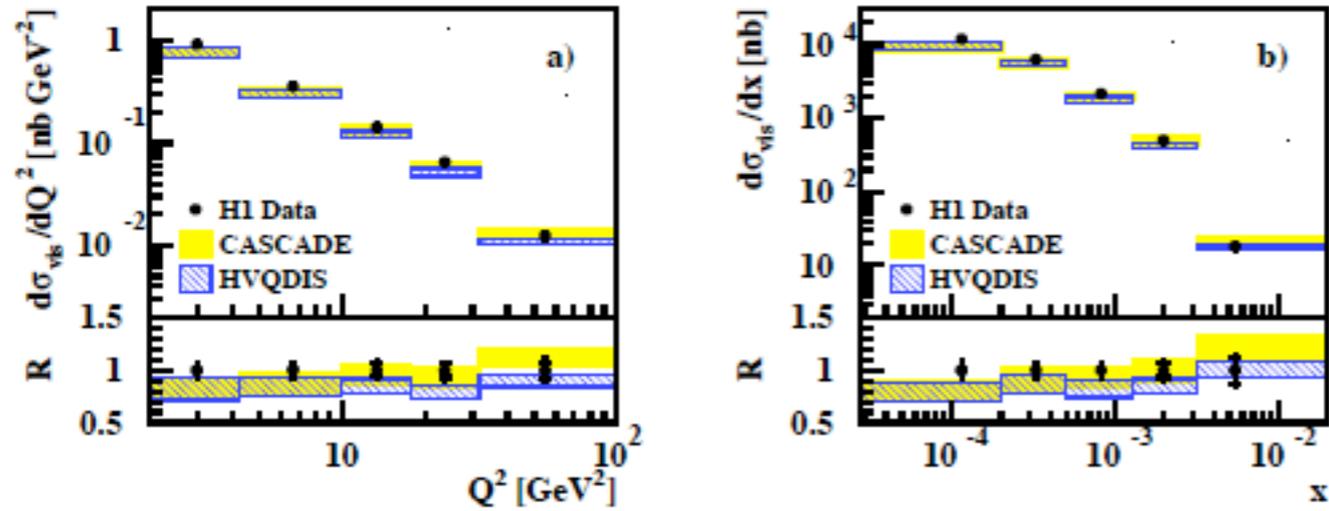
# $D^{*\pm}$ & $D^{*\pm} + \text{jets}$ production in $\gamma p$



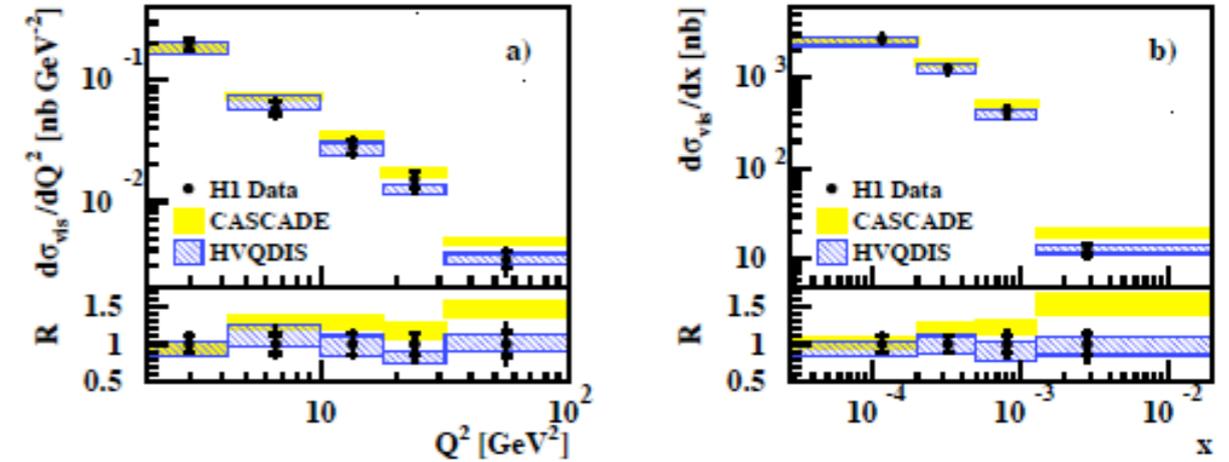
- $R = (1/\sigma \frac{d\sigma}{dv})_{\text{data}} / (1/\sigma \frac{d\sigma}{dv})_{\text{theo}}$  reduces exp. & theo uncertainties
- overall reasonable description of data by NLO
- except for  $\Delta\Phi$ ; data are more de-correlated than NLO (also seen by ZEUS)
- uncertainty in theory often larger than in data

# $D^{*\pm}$ & $D^{*\pm} + \text{jets}$ production in DIS

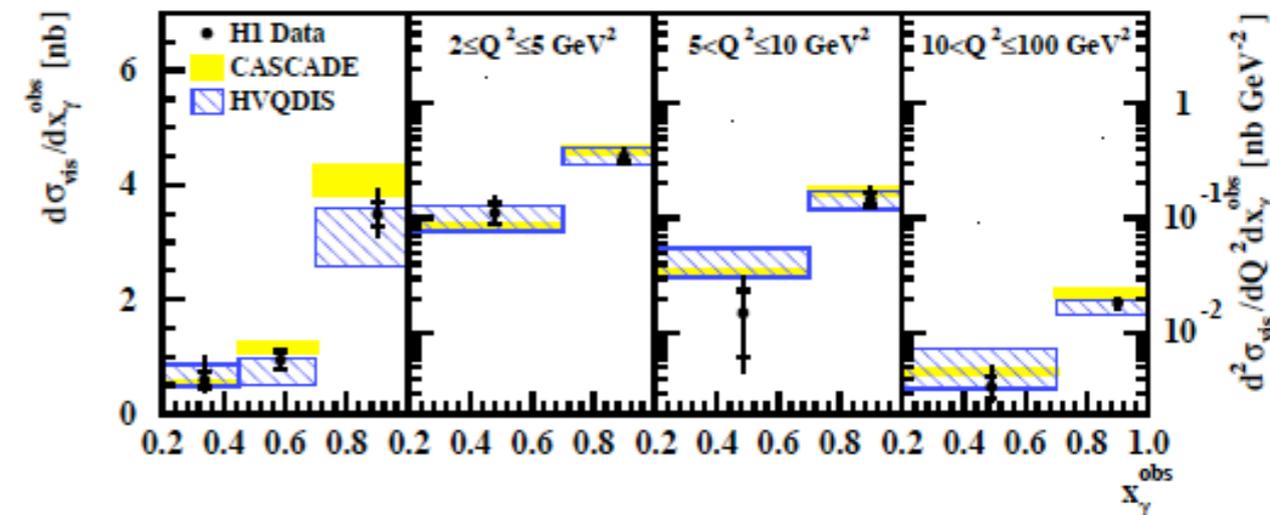
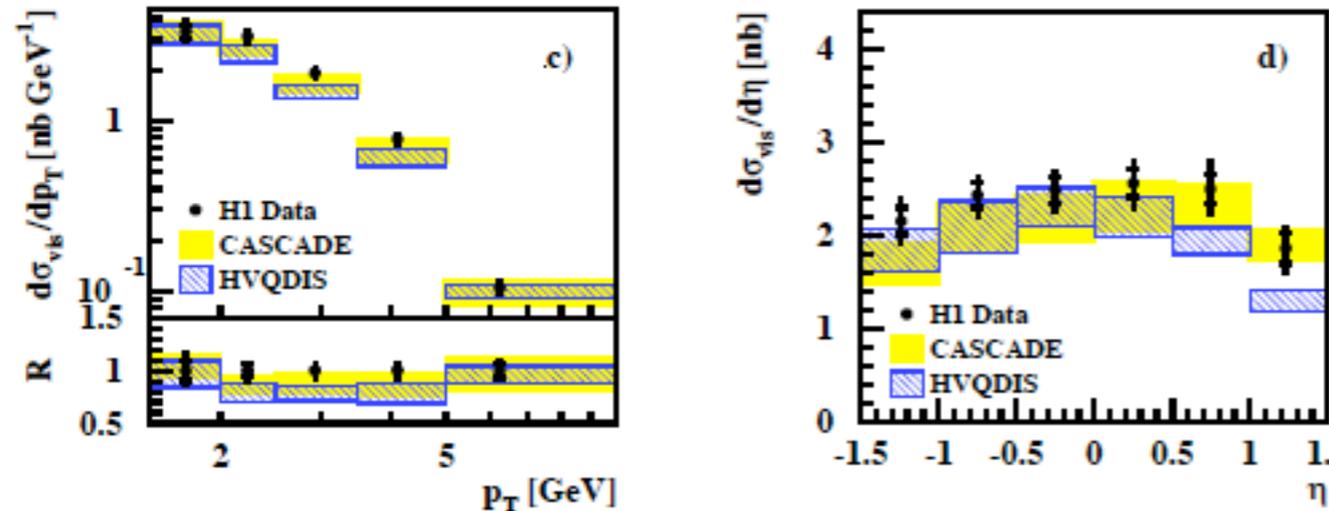
H1  $ep \rightarrow eD^{*\pm}X$



H1  $ep \rightarrow eD^{*\pm}jjX$



H1  $ep \rightarrow eD^{*\pm}jjX$



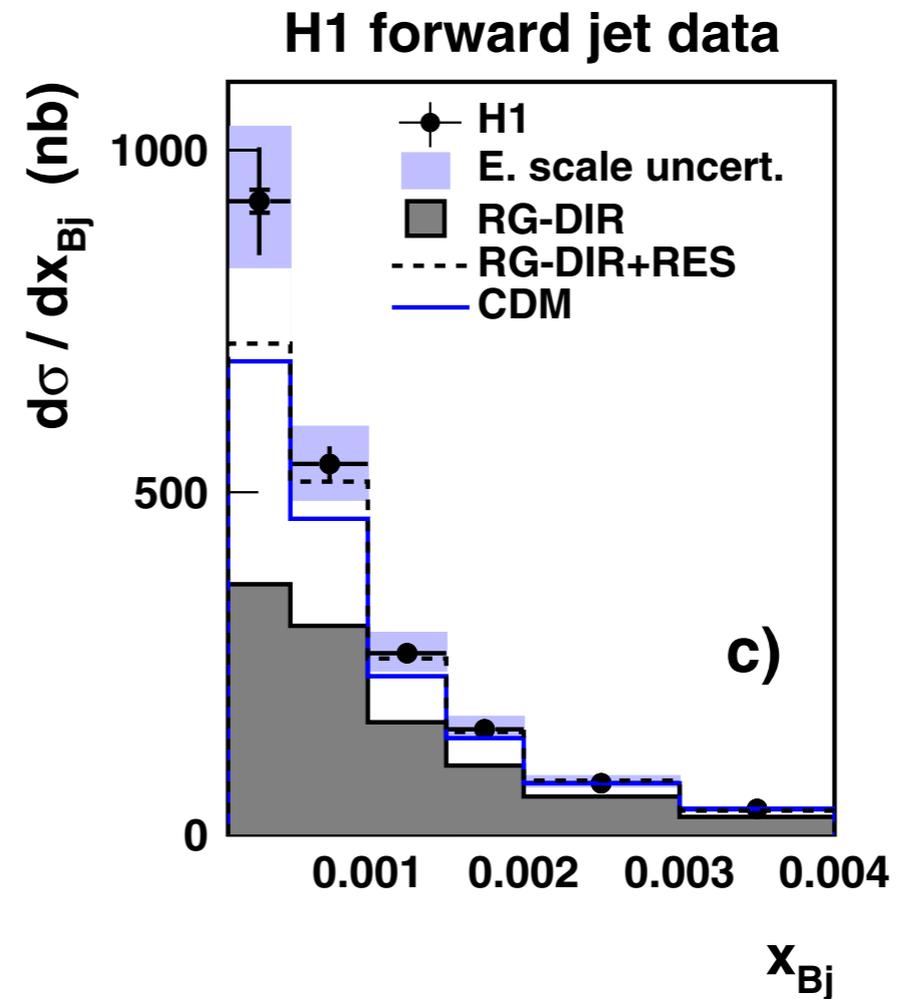
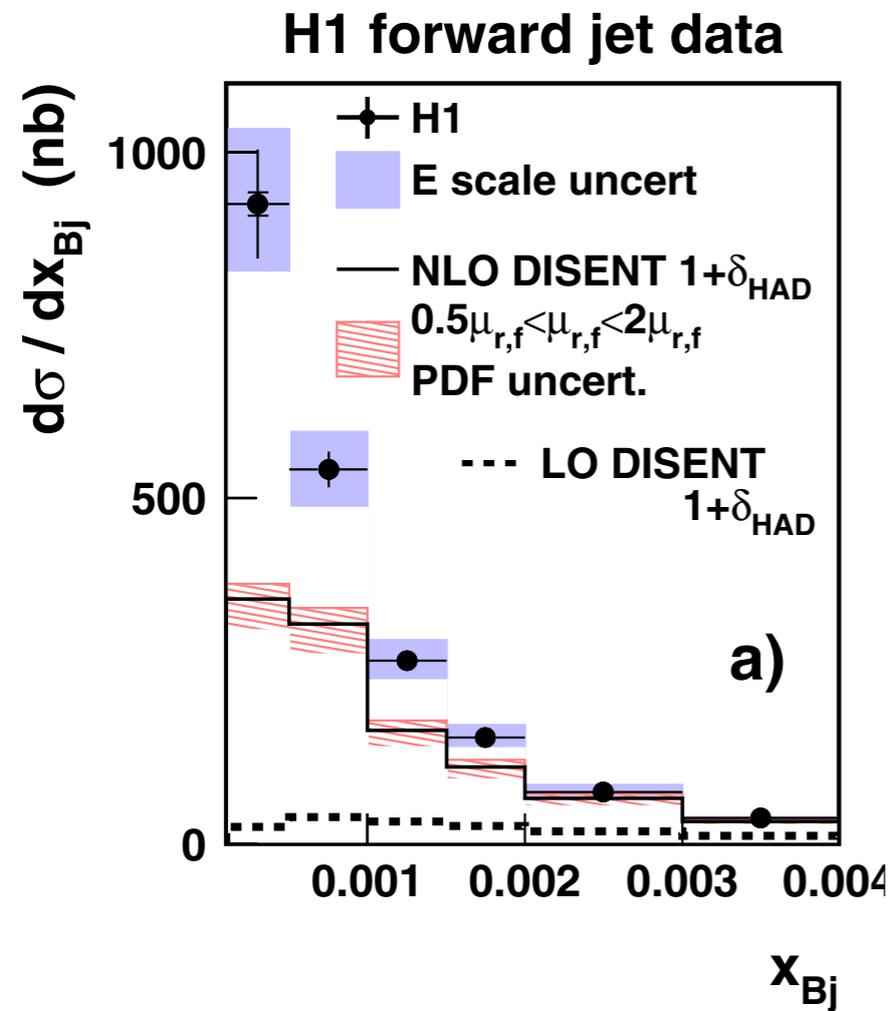
- $R$  is the ratio of the normalized x-sections of data to NLO

- ☞ overall reasonable description of data (better for jets) by NLO (HVQDIS)
- ☞  $x_\gamma$  described by NLO without need for additional resolved contribution
- ☞ scale uncertainty  $\geq$  data uncertainty

# Summary/Conclusions

- NLO calc. for DIS appear to work remarkably well even down to very low  $Q^2$  (e.g. D meson prod.)
- NLO at low  $Q^2$  is however not very predictive, the scale uncertainties are very large, “clever” scale choices have to be made, e.g.  $\mu^2 = (E_T^2 + Q^2) / 4$  or  $E_T^2$  or  $Q^2$
- NLO has problems - not surprisingly -
  - the more differential the distributions, e.g.
    - for  $d\sigma_{\text{dijet}} / dQ^2 dE_T^2 d x_Y$ , particularly at low  $Q^2$ , low  $x_Y$  and low jet  $E_T$
    - for observables like  $\Delta\Phi < \pi$  and  $x_Y < 1$
  - when only processes at tree level contribute
- Charm data are overall described by NLO, within the large scale uncertainties no need for resolved contributions, beyond what is included at LO direct
  - in charm prod. there is always a hard scale, i.e. the charm mass, could this be the reason?
- Deficiencies & typically large scale uncertainties indicate the need for higher orders, direct and/or resolved, or for other approaches like  $k_T$ -factorization and unintegrated parton densities, etc.
- To make progress two NLO programs with resolved contrib. at NLO, like JETVIP and “DIScoPHOX, would be needed

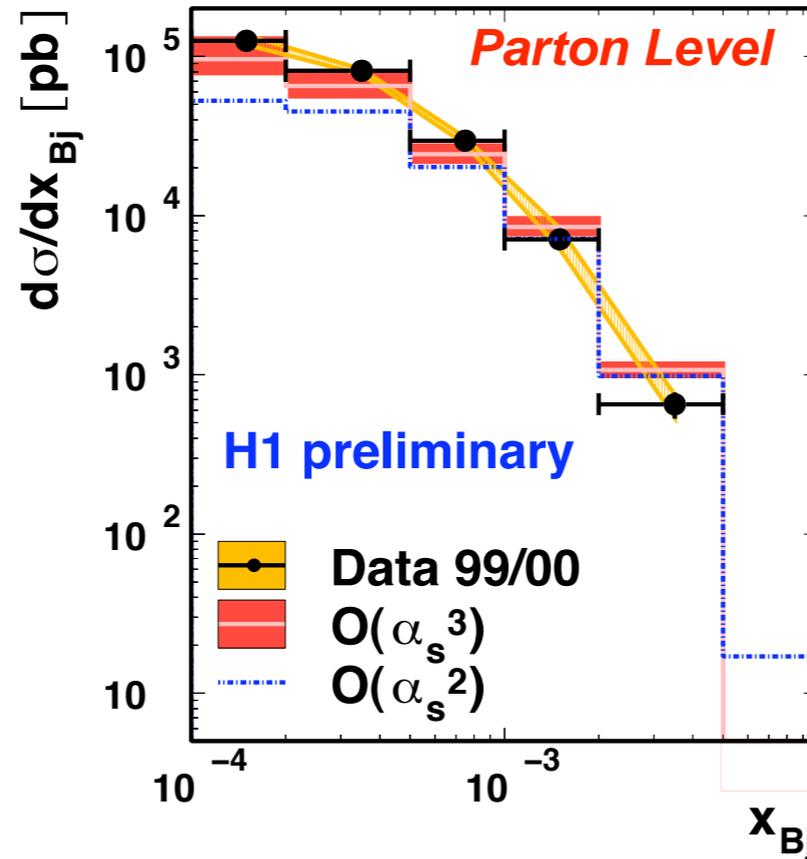
# extra plots



# extra plots

- $E_{T,jet1,2,3} > 4 \text{ GeV}$  &  
 $E_{T,jet1} + E_{T,jet2} > 9 \text{ GeV}$
- fwd-jet:  $\theta_{lab} < 20^\circ$ ,  
 $x_{jet} > 0.035$
- central jets:  
 $-1 < \eta_{lab} < 1.3$
- fwd-jets are mainly  
gluon jets; at lowest  
 $x_{Bj}$  gluon emissions  
unordered in  $p_T$

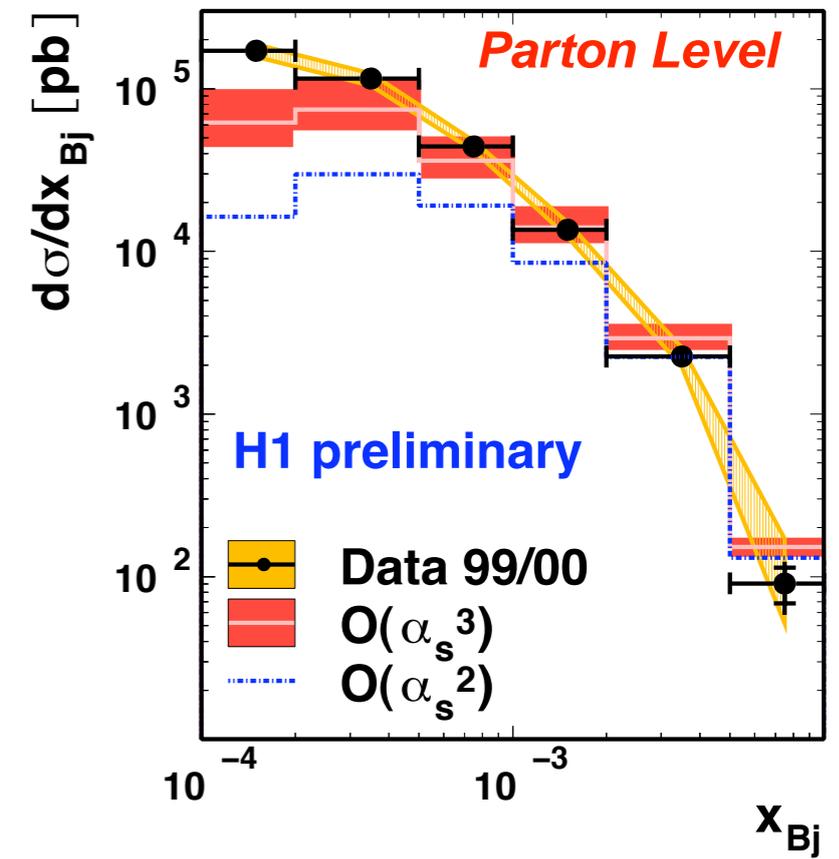
1 fwd-jet & 2 central jets



NLOJET++

- o NLO describes data well

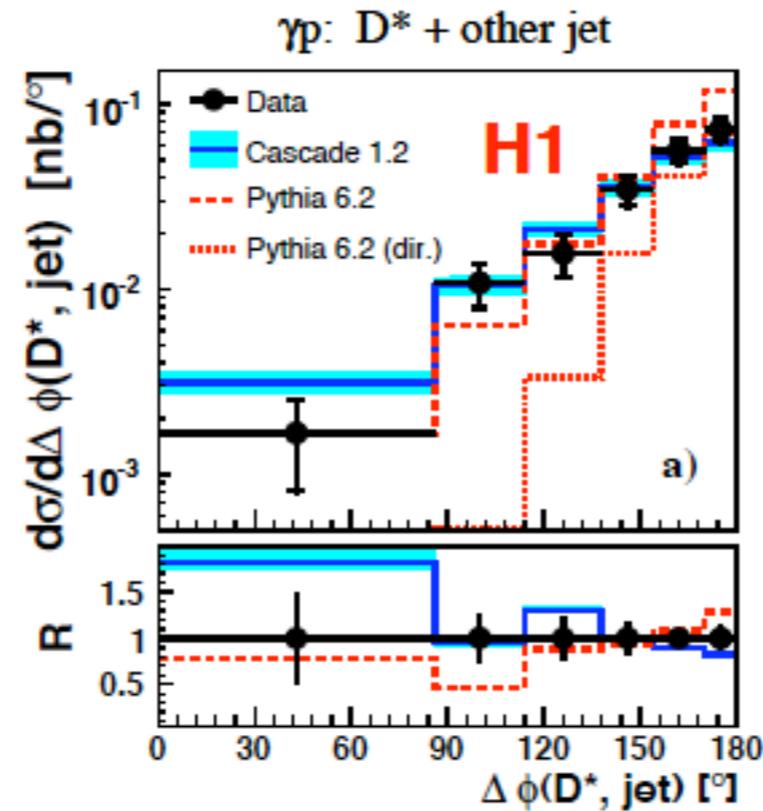
2 fwd-jets & 1 central jet



NLOJET++

- o NLO describes data well  
for  $x_{Bj} > 5 \cdot 10^{-4}$
- o NLO fails at lowest  $x_{Bj}$

# extra plots



H1  $ep \rightarrow eD^{*\pm}jjX$

