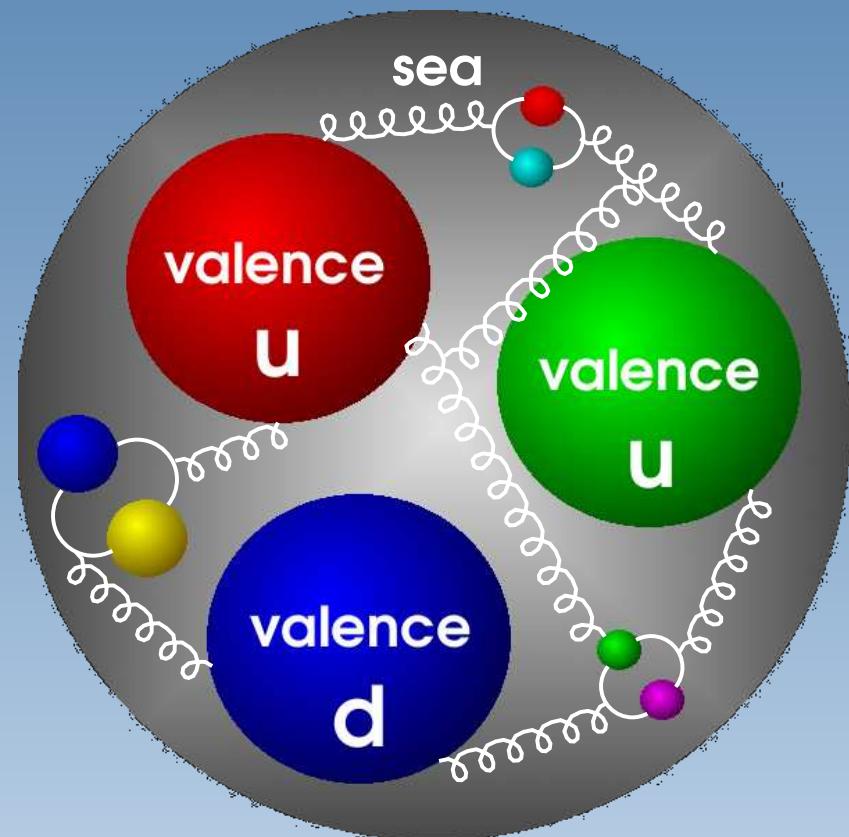


Proton Structure and Low x Physics

Outline

- Introduction
- Parton distributions
- Evolution at low x
- Non-linear dynamics



Victor Lendermann
Universität Heidelberg

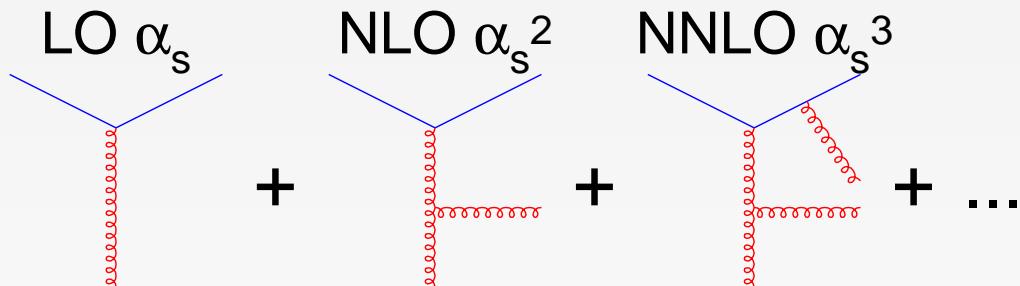
DPG-Frühjahrstagung
Heidelberg, 9.03.2007



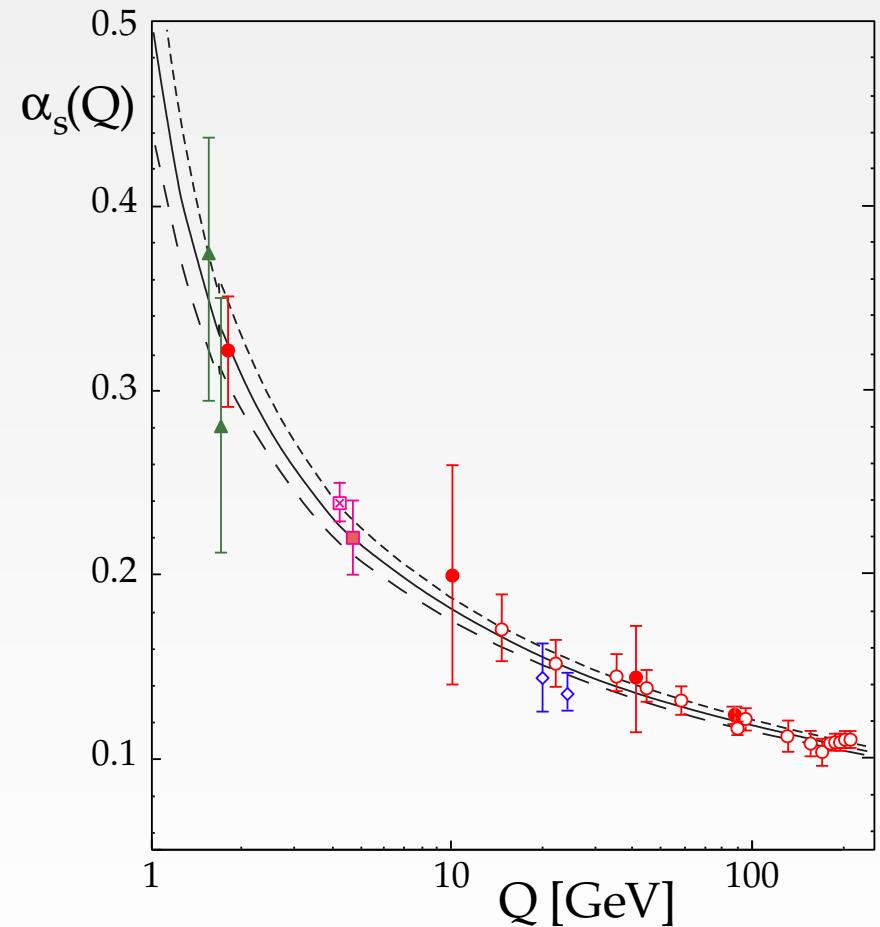
QCD and Hadron Scattering

QCD: perturbative or non-perturbative? That is the question.

- ◆ pQCD calculates interactions of quarks and gluons using pert. expansion in α_s



- ◆ Actually, hadrons are in initial and final state due to confinement
- ◆ The limit for pQCD is where α_s is large

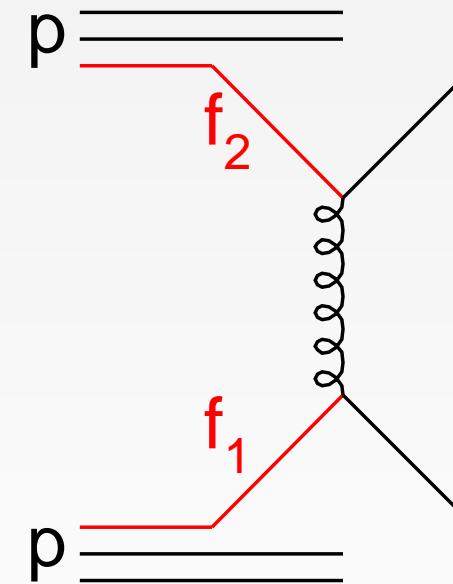
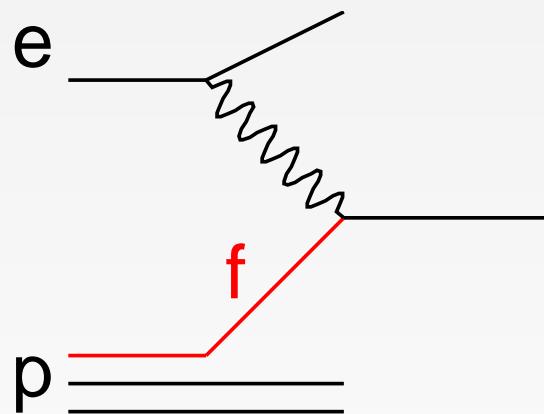


Factorisation Theorem

Non-perturbative input is given by universal parton distributions (PDF)

$$\sigma_{ep} = \sum_i f_i \otimes \hat{\sigma}_i$$

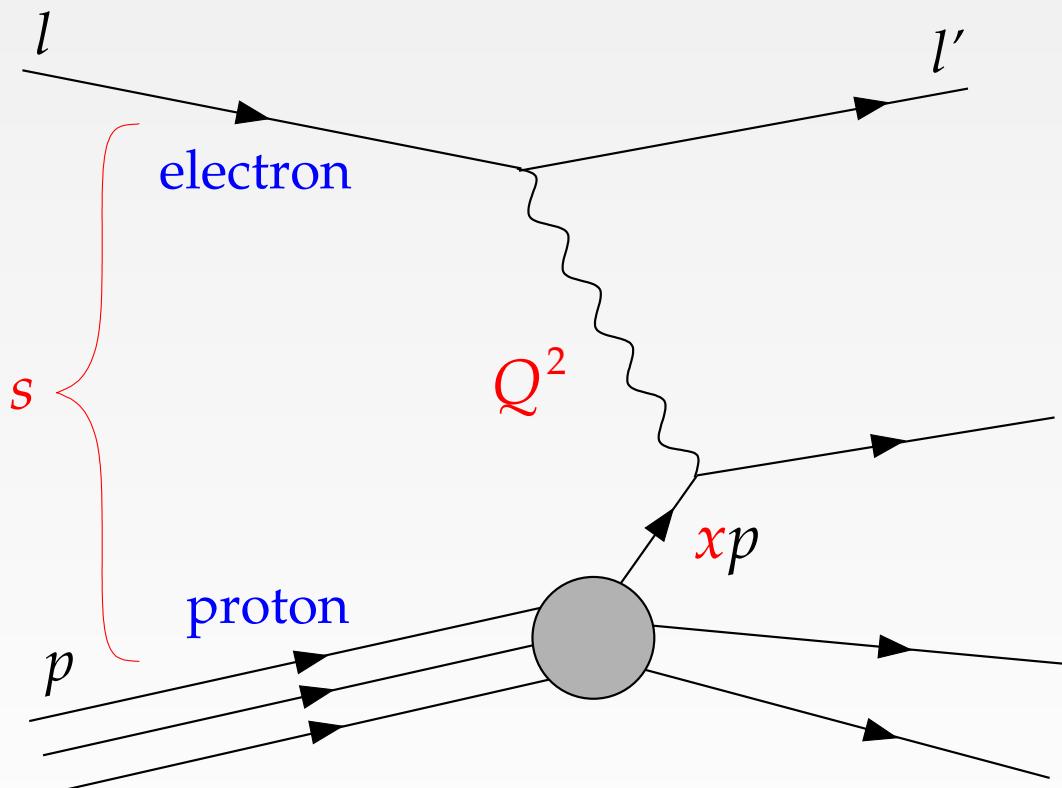
$$\sigma_{pp} = \sum_{ij} f_i \otimes f_j \otimes \hat{\sigma}_{ij}$$



PDFs must be determined experimentally
and can be used for other scattering processes

Inclusive Deep Inelastic ep Scattering

DIS – best method to determine the proton structure **with highest precision**

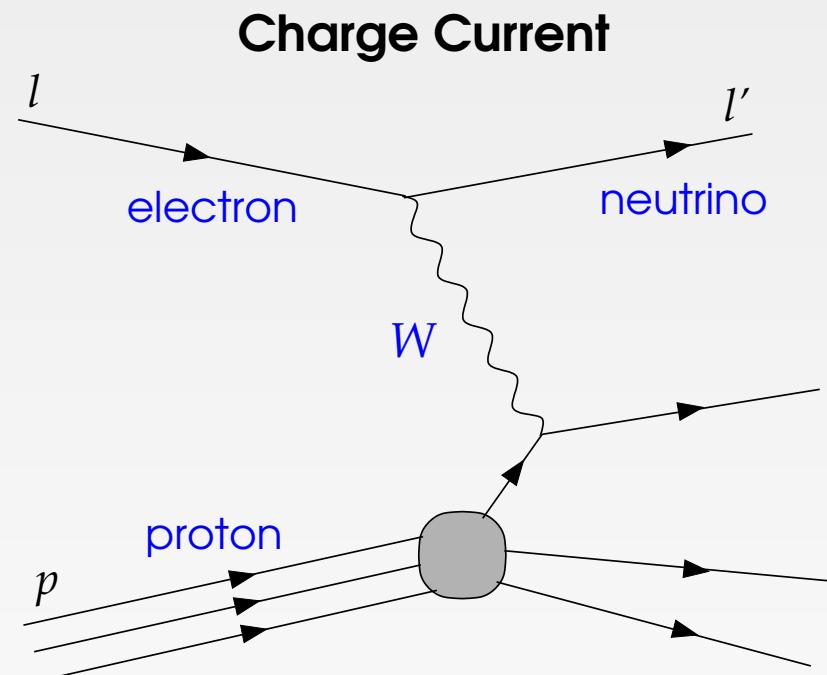
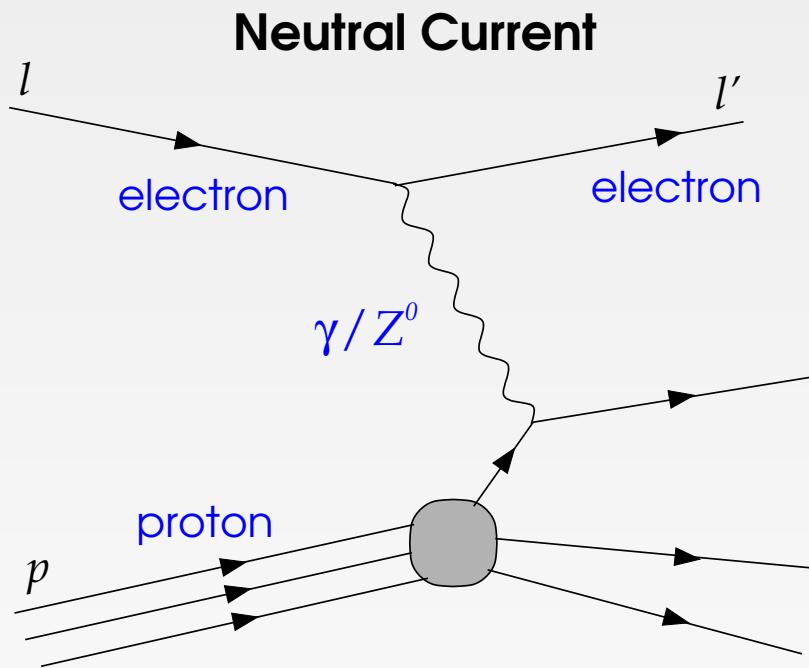


boson virtuality Q^2
= resolution scale

fractional momentum x
of struck quark

determine PDFs: $q_i(x, Q^2)$, $g(x, Q^2)$

Determination of Quark Distributions



Cross sections: $\sigma_{e^+p}^{NC}$, $\sigma_{e^-p}^{NC}$, $\sigma_{e^+p}^{CC}$, $\sigma_{e^-p}^{CC}$
have different sensitivity to different quark flavours

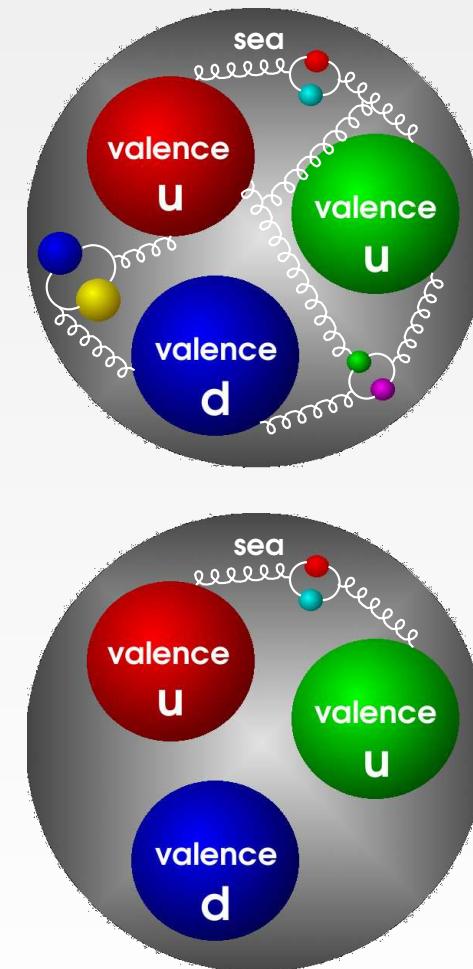
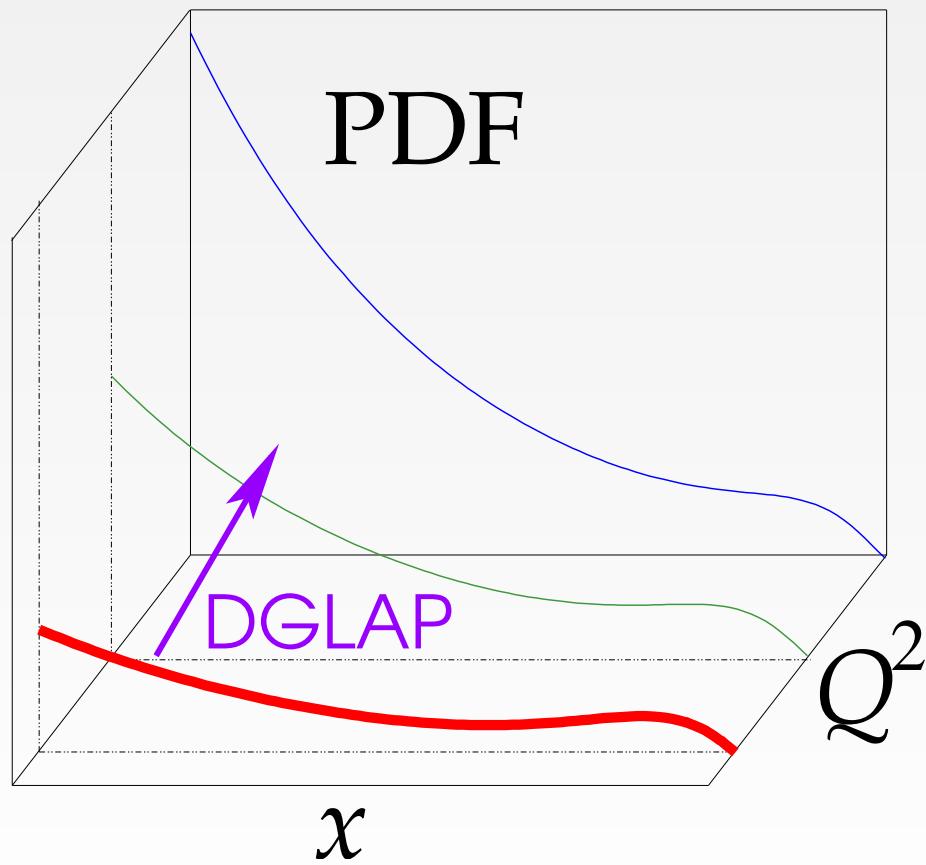
Almost full flavour decomposition is possible with HERA only

Highest precision by combining with fixed target DIS, Drell-Yan, ...

DGLAP Evolution of PDFs

PDFs – intrinsically non-perturbative but evolve perturbatively in Q^2

Dokshitzer, Gribov, Lipatov, Altarelli, Parisi (DGLAP)



Only need start-up distributions $q(x)$, $g(x)$ at starting scale Q_0^2

Inclusive DIS Data

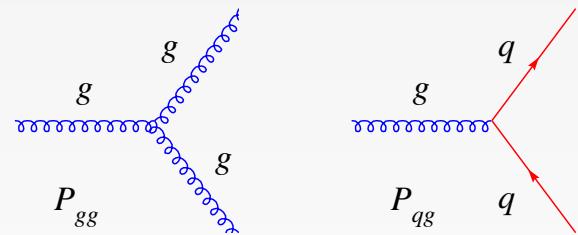
$$F_2^{\text{em}} = x \sum e_q^2 (q + \bar{q})$$

◆ 5 orders of magnitude in x and Q^2

Only HERA covers low x domain

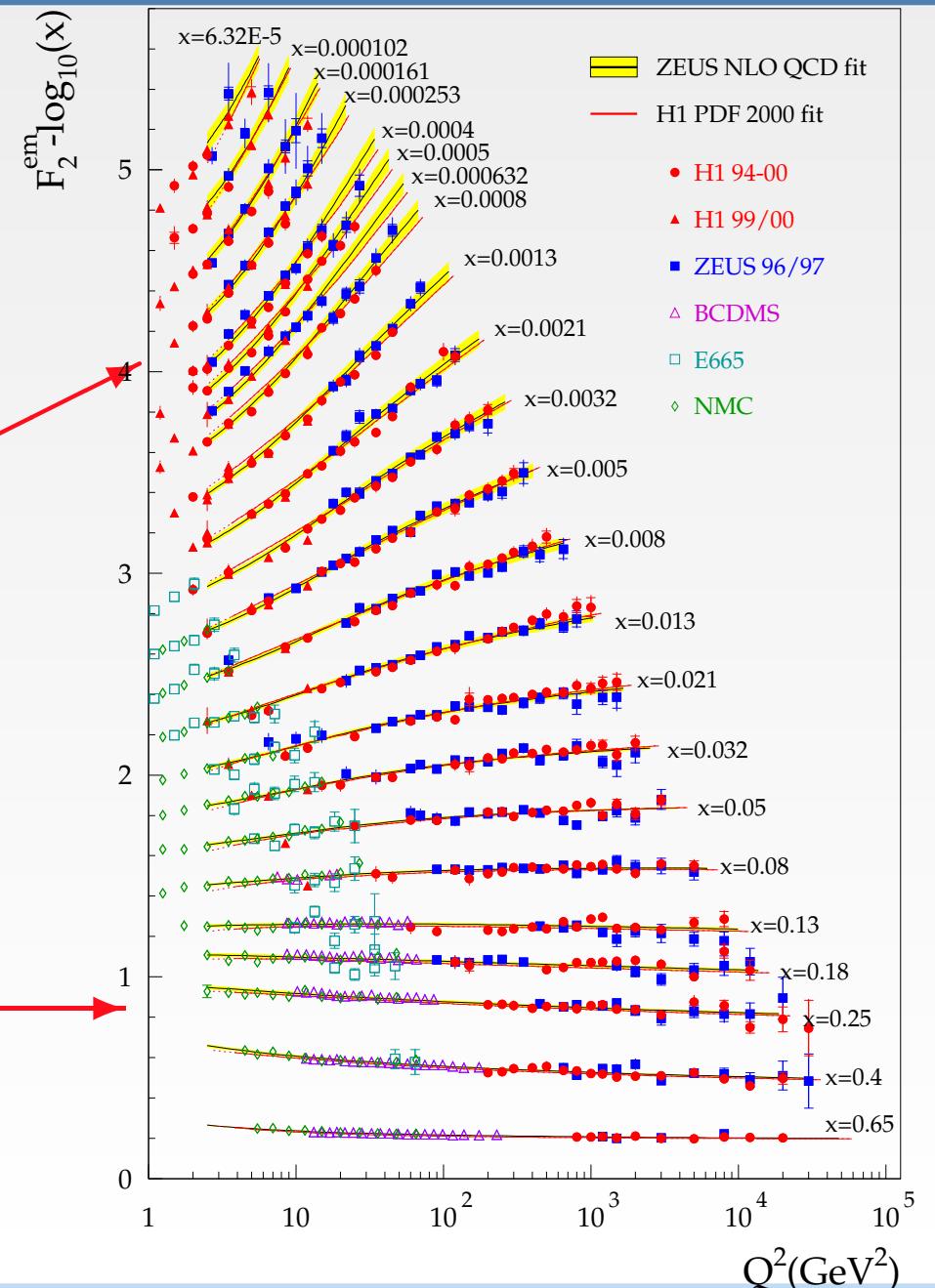
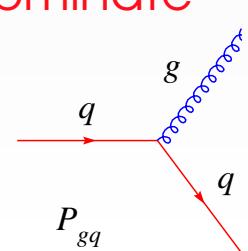
◆ At low x : rising F_2

gluons \rightarrow sea quarks dominate



◆ At high x : falling F_2

valence quarks dominate



Inclusive DIS Data

$$F_2^{\text{em}} = x \sum e_q^2 (q + \bar{q})$$

- ◆ 5 orders of magnitude in x and Q^2

NLO DGLAP fits well describe inclusive data

NLO = $\mathcal{O}(\alpha_s^2)$, first NNLO $\mathcal{O}(\alpha_s^3)$ fits available

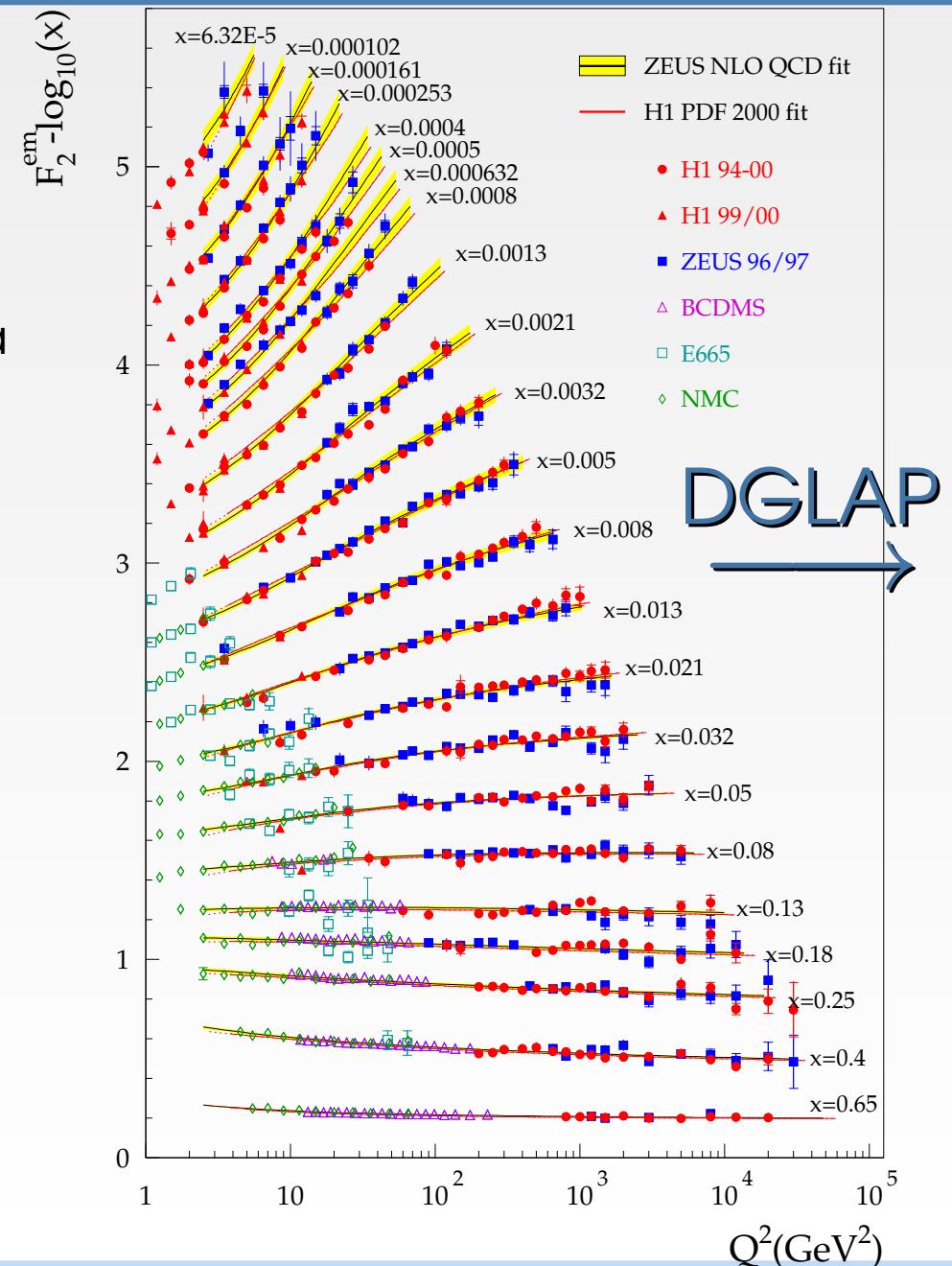
- ◆ Data precision 2 – 3% in bulk region

1 – 2% seem feasible

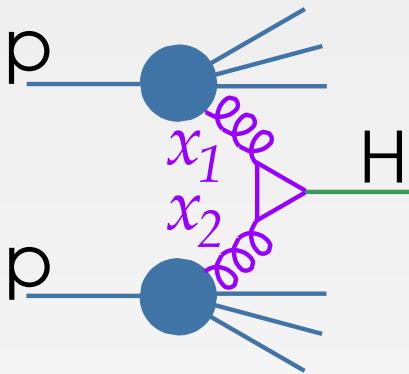
- ◆ High Q^2 dominated by stat. errors

Factor 2 reduction is aimed at with HERA II

So what is needed by LHC?



\mathcal{PDFs} for \mathcal{LHC}

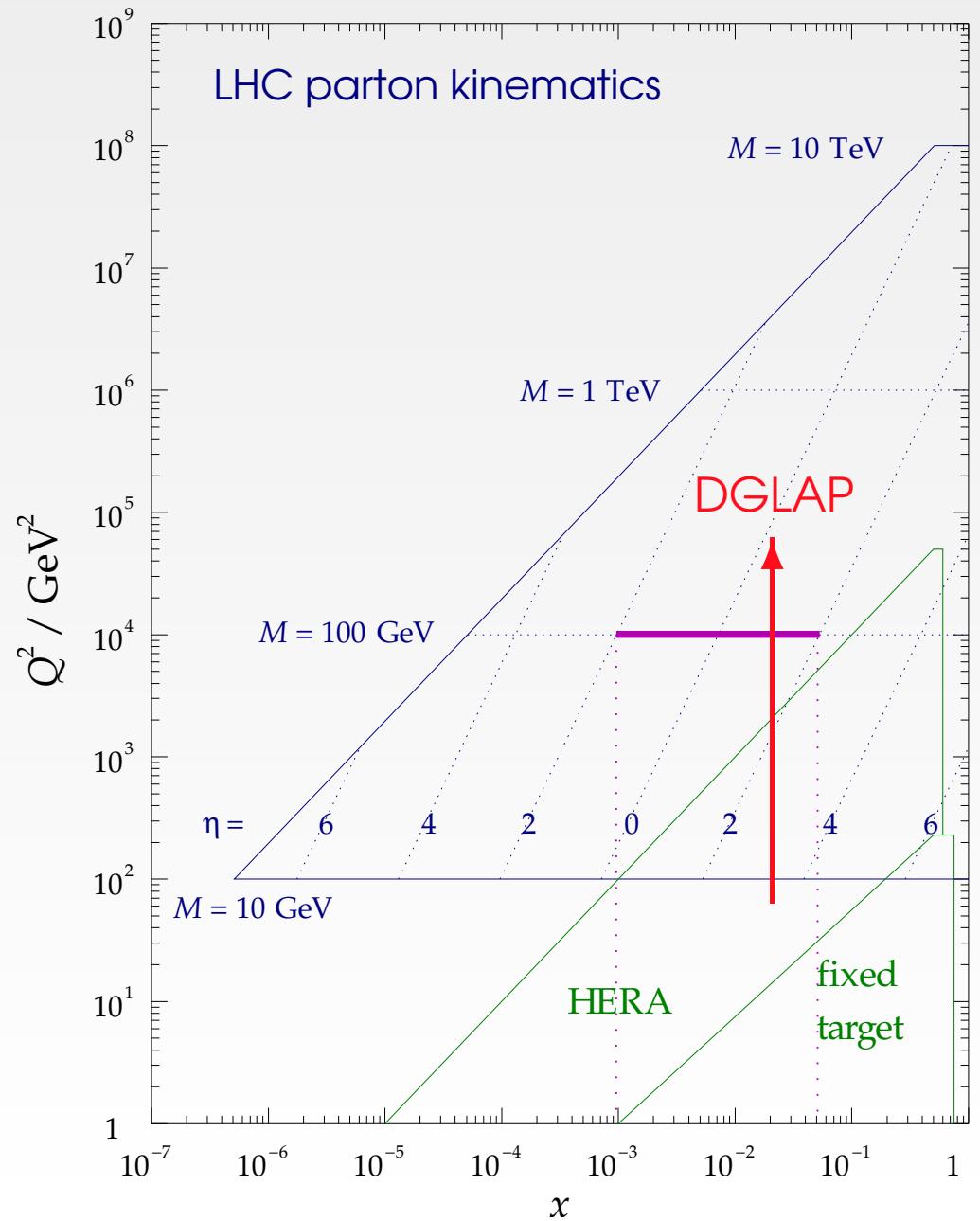


$$x_{1,2} = \frac{M}{\sqrt{S}} \exp(\pm \text{rapidity})$$

Precise quark and gluon densities
are required **in the whole x range**

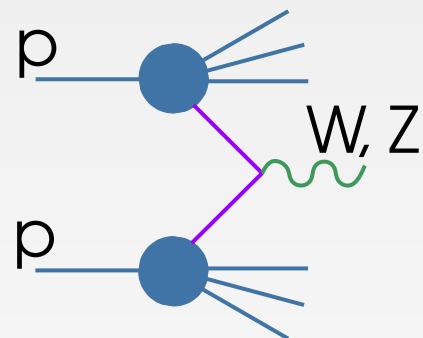
Solution:

Evolve from HERA Q^2 to LHC scales M^2



Luminosity Measurement at LHC with W, Z

$$\sigma = \frac{N}{L} \implies \text{Use reference 'standard candle' process to measure lumi: } L = \frac{N_{\text{ref}}}{\sigma_{\text{ref}}}$$

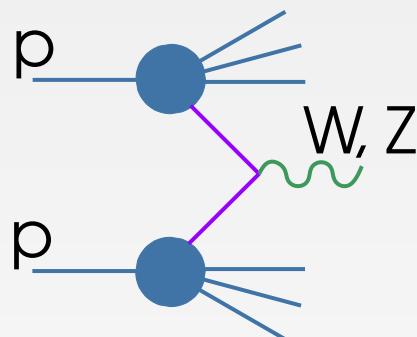


W and Z production is well suited

- ◆ High rate: few Hz at low lumi \rightarrow small stat. errors
- ◆ Well measurable: 1 – 2% syst. error for Z
- ◆ Well-known parton x-section: 1% at NNLO

Luminosity Measurement at LHC with W, Z

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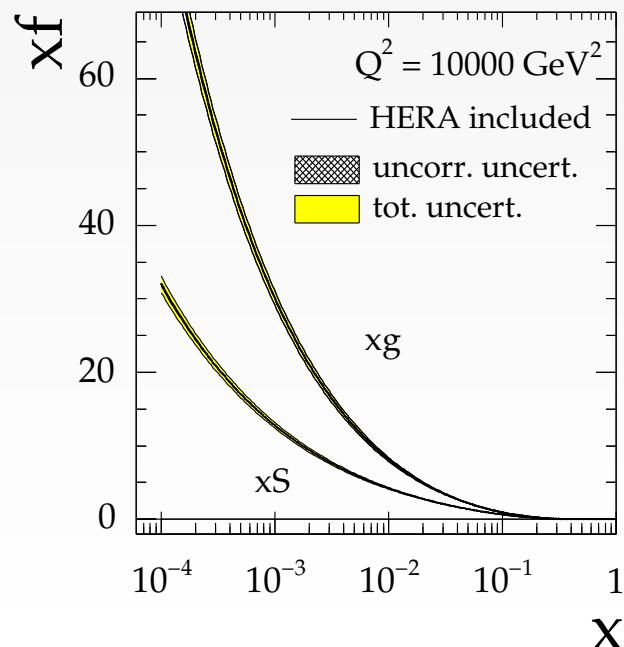


W and Z production is well suited

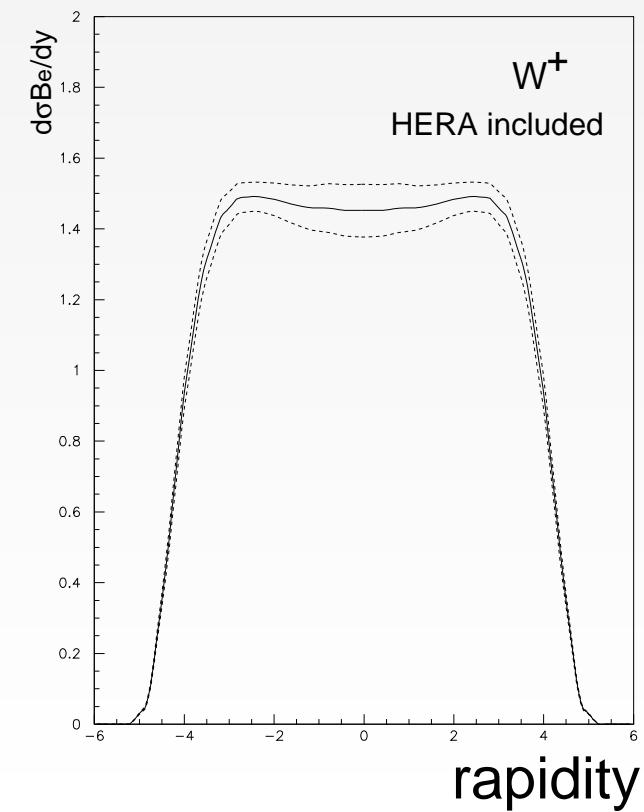
- ◆ High rate: few Hz at low lumi \rightarrow small stat. errors
- ◆ Well measurable: 1 – 2% syst. error for Z
- ◆ Well-known parton x-section: 1% at NNLO

Crucial issue – knowledge of PDFs:

$\sim 4\%$ for total σ , $\sim 8\%$ for central rap.

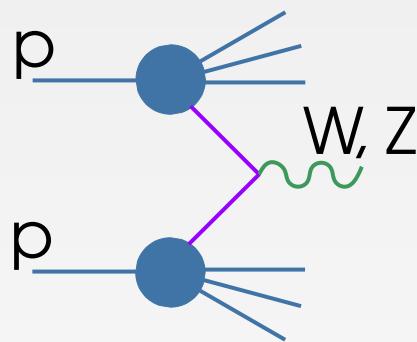


Tricoli, Cooper-Sarkar, Gwenlan



Luminosity Measurement at LHC with W, Z

$$\sigma = \frac{N}{L} \implies \text{Use reference 'standard candle' process to measure lumi: } L = \frac{N_{\text{ref}}}{\sigma_{\text{ref}}}$$

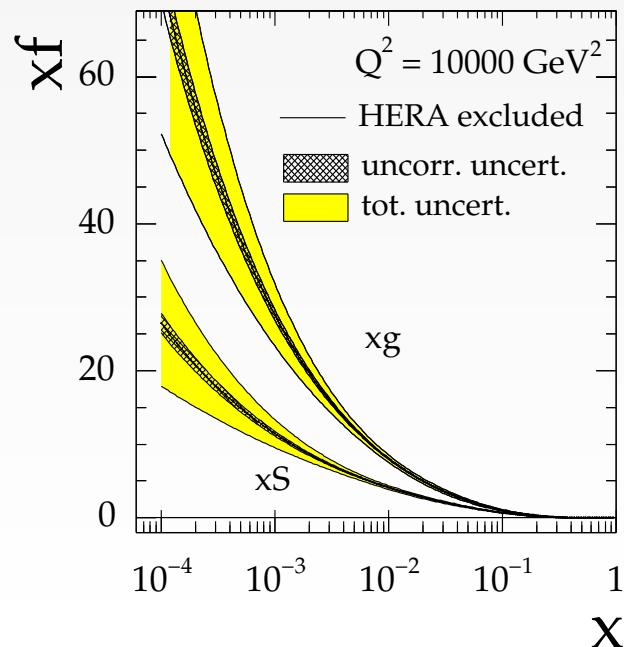


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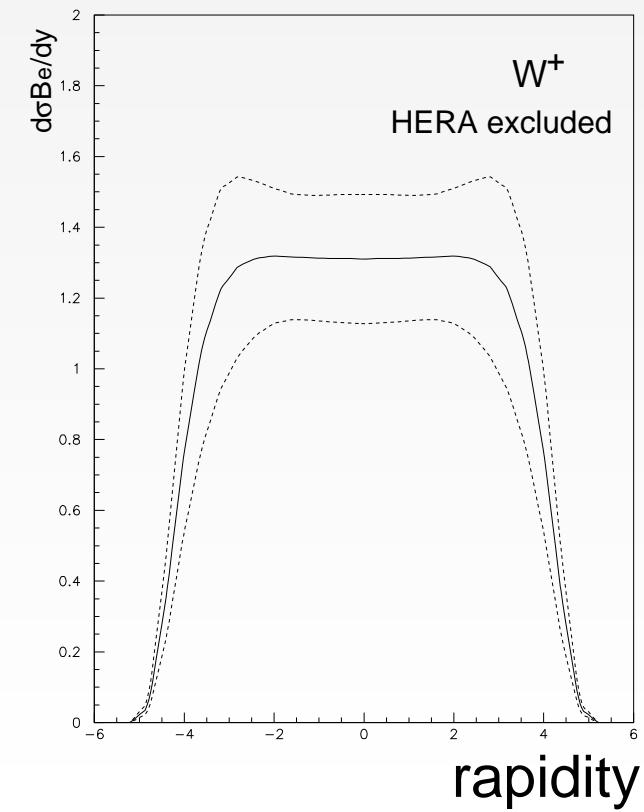
- ◆ High rate: few Hz at low lumi \rightarrow small stat. errors
- ◆ Well measurable: 1 – 2% syst. error for Z
- ◆ Well-known parton x-section: 1% at NNLO

Crucial issue – knowledge of PDFs:

Error would be $\sim 15\%$ without HERA



Tricoli, Cooper-Sarkar, Gwenlan



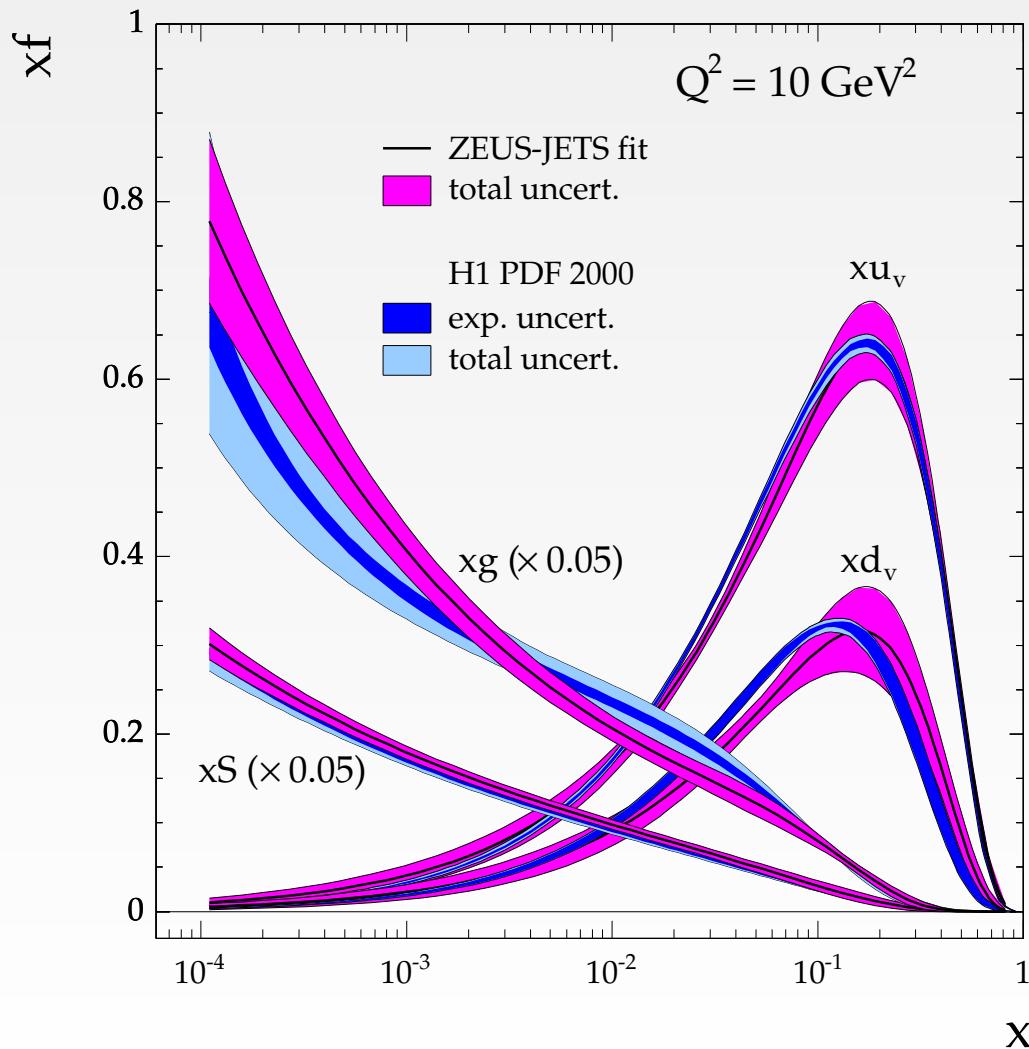
Questions

- ◆ How good is our knowledge of PDFs?
- ◆ Is DGLAP sufficient to extrapolate PDFs to LHC scales?
- ◆ Are non-linear QCD effects relevant?

Let's look...

\mathcal{PDF} Extraction

Typical uncertainties: **u density 2 – 5%** , **d density 5 – 8%** , **gluon density 10 – ... %**



Example

H1/ZEUS mostly agree but

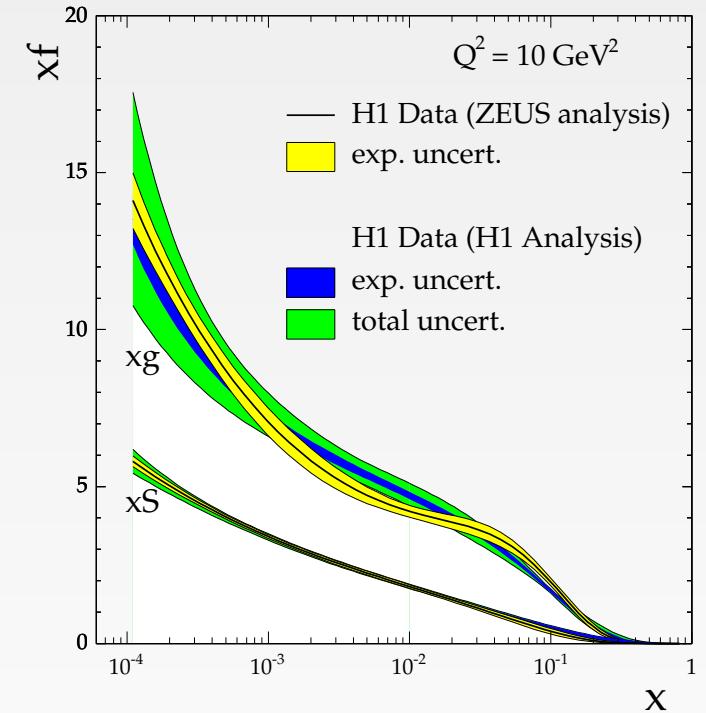
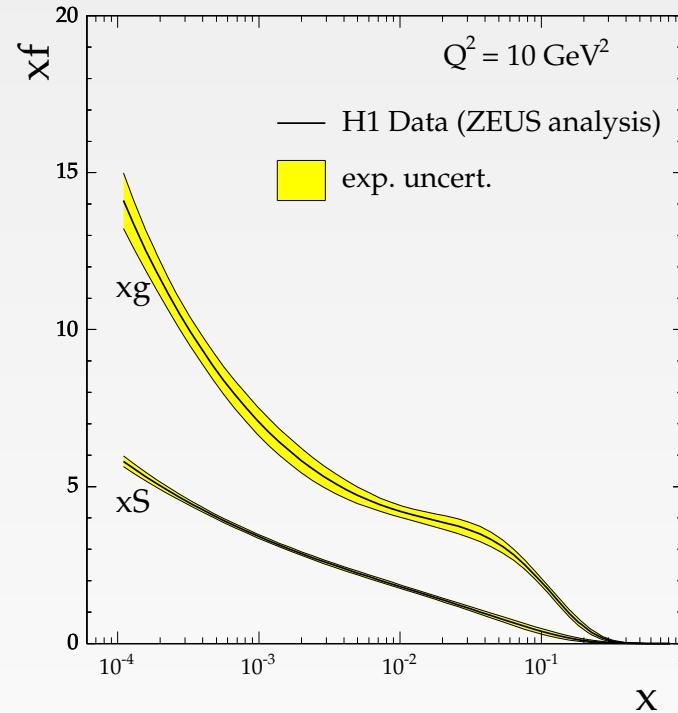
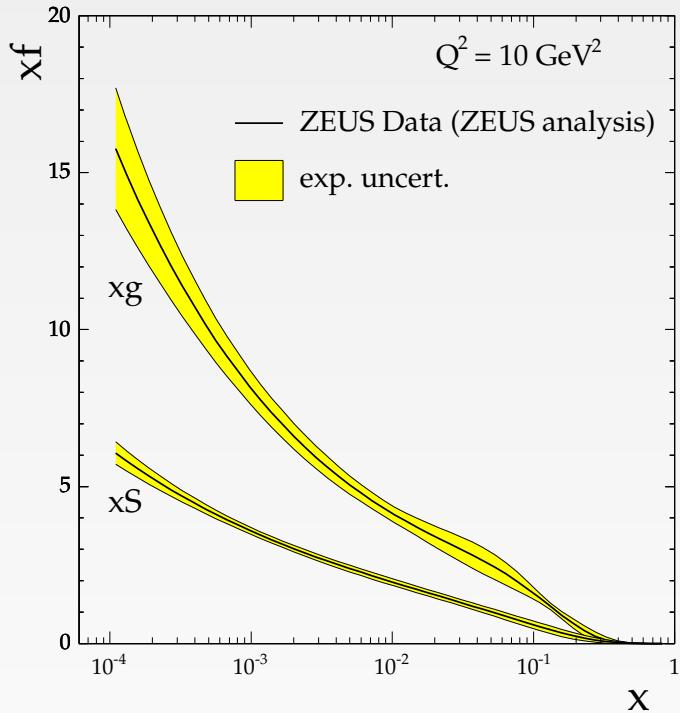
Differences

- ◆ Choice of data sets
- ◆ Treatment of systematic errors
- ◆ PDFs to extract
- ◆ Form of x distribution
- ◆ Number of parameters
- ◆ Constraints on parameters
- ◆ ...

Many effects understood!

Differences between H1 and ZEUS

Cooper-Sarkar and Gwenlan, HERA-LHC Workshop



Different data, same analysis

Same data, different analysis

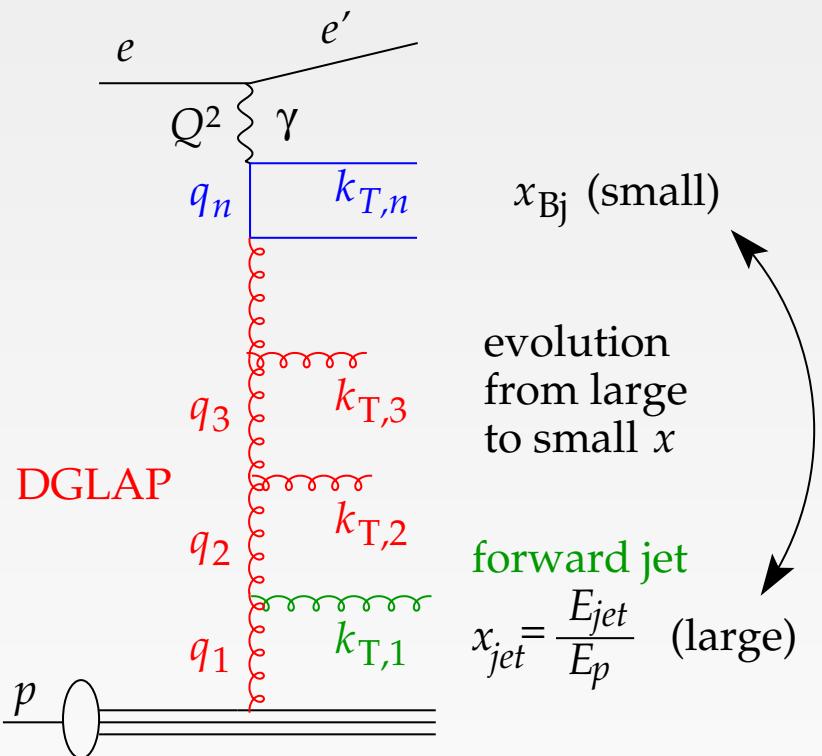
Differences already in the data

H1 and ZEUS intend to produce common data sets

Next Question

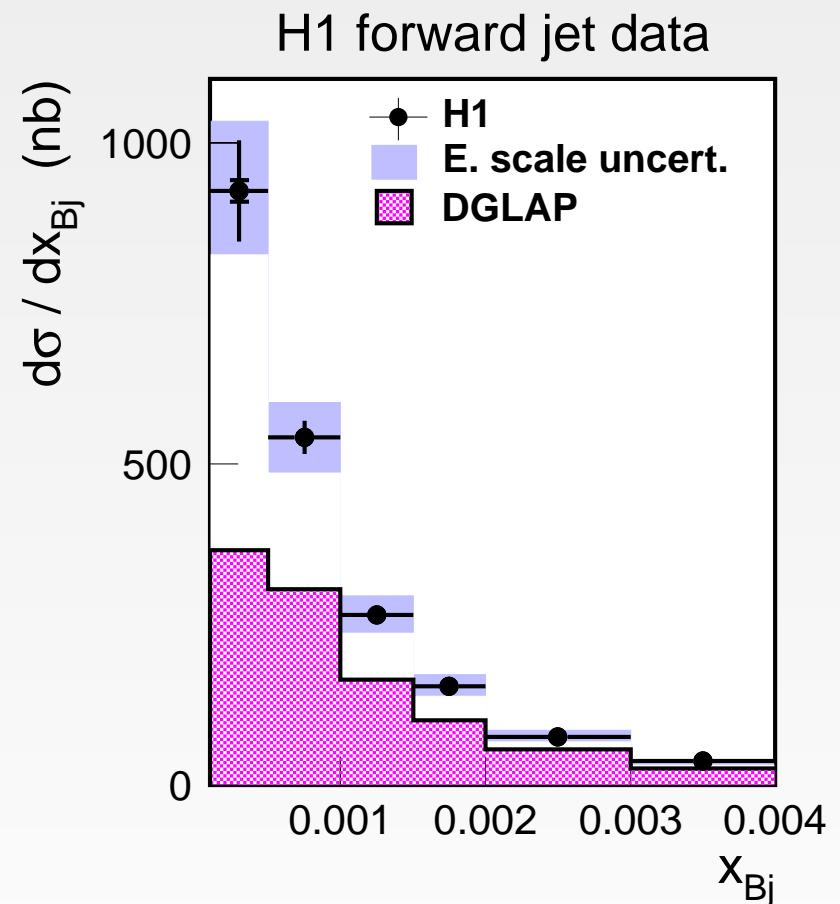
- ◆ How good is our knowledge of PDFs?
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Forward Jets at Low χ

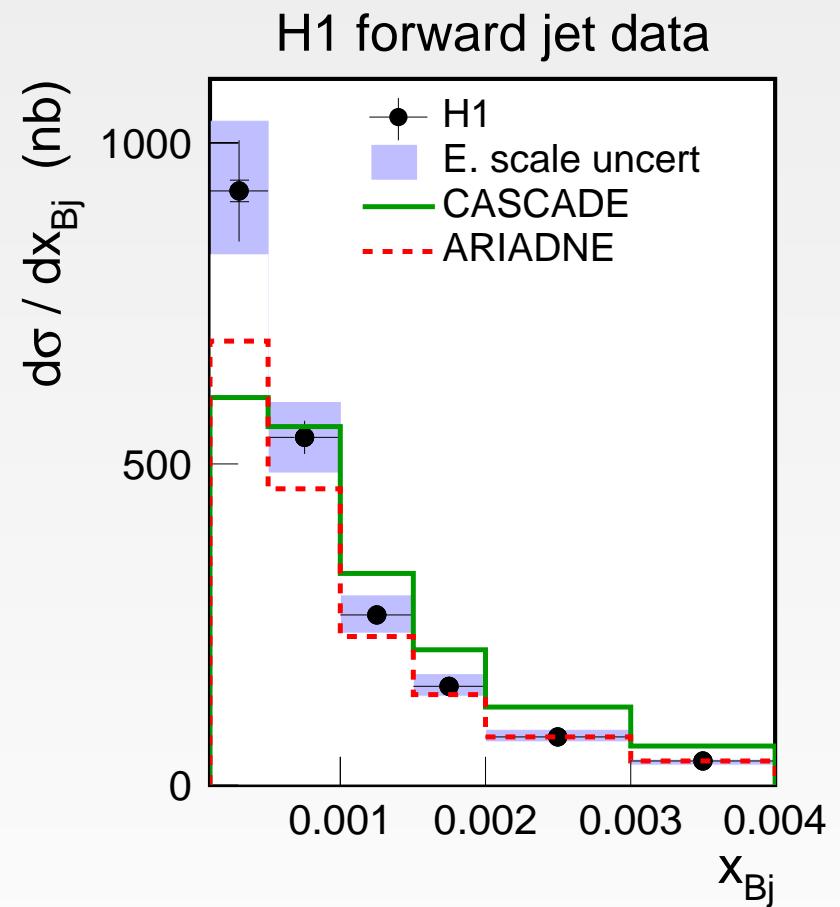
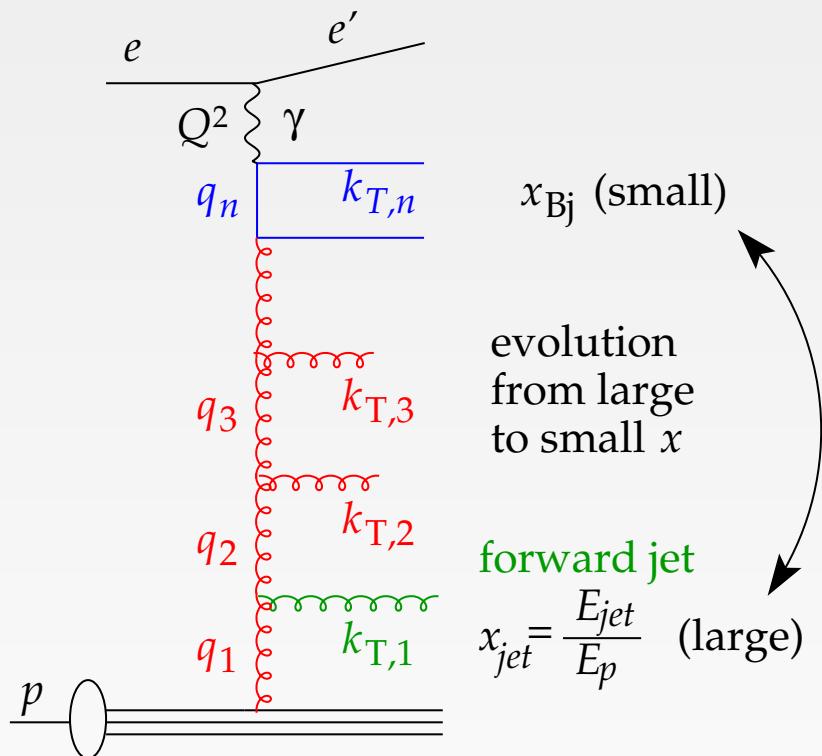


◆ If using DGLAP for parton shower $\Rightarrow k_T$ ordering

$$q_1^2 \ll q_2^2 \cdots \ll q_n^2 \ll Q^2 \Rightarrow k_{T,1} \ll k_{T,2} \cdots \ll k_{T,n}$$



Forward Jets at Low χ



- ◆ Parton shower schemes with different ordering
 - CCFM (MC Cascade)
 - CDM (MC Ariadne)

work better than DGLAP at low x

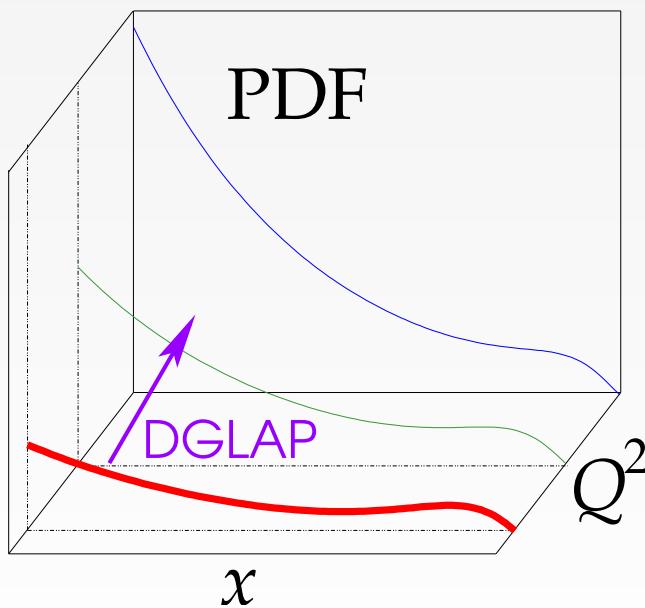
Resummation at Low x

◆ Missing terms?

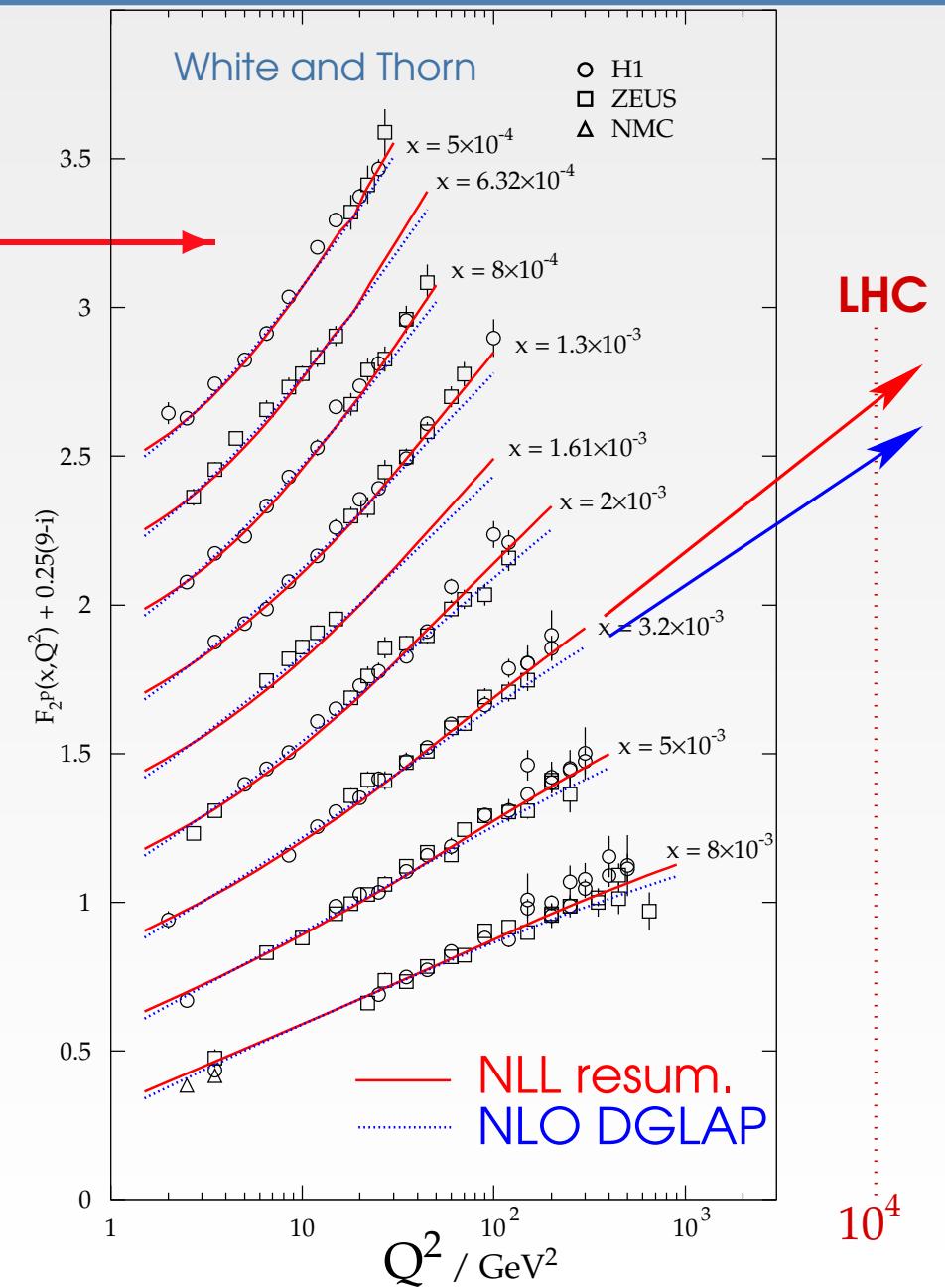
$$\left[\alpha_s \ln \frac{1}{x} \right]^k$$

Attempt to add these to evol. eqs

◆ Difficult to see in inclusive fits
due to free choice of $q(x), g(x)$ at Q_0^2



Critical corner – low x , low Q^2



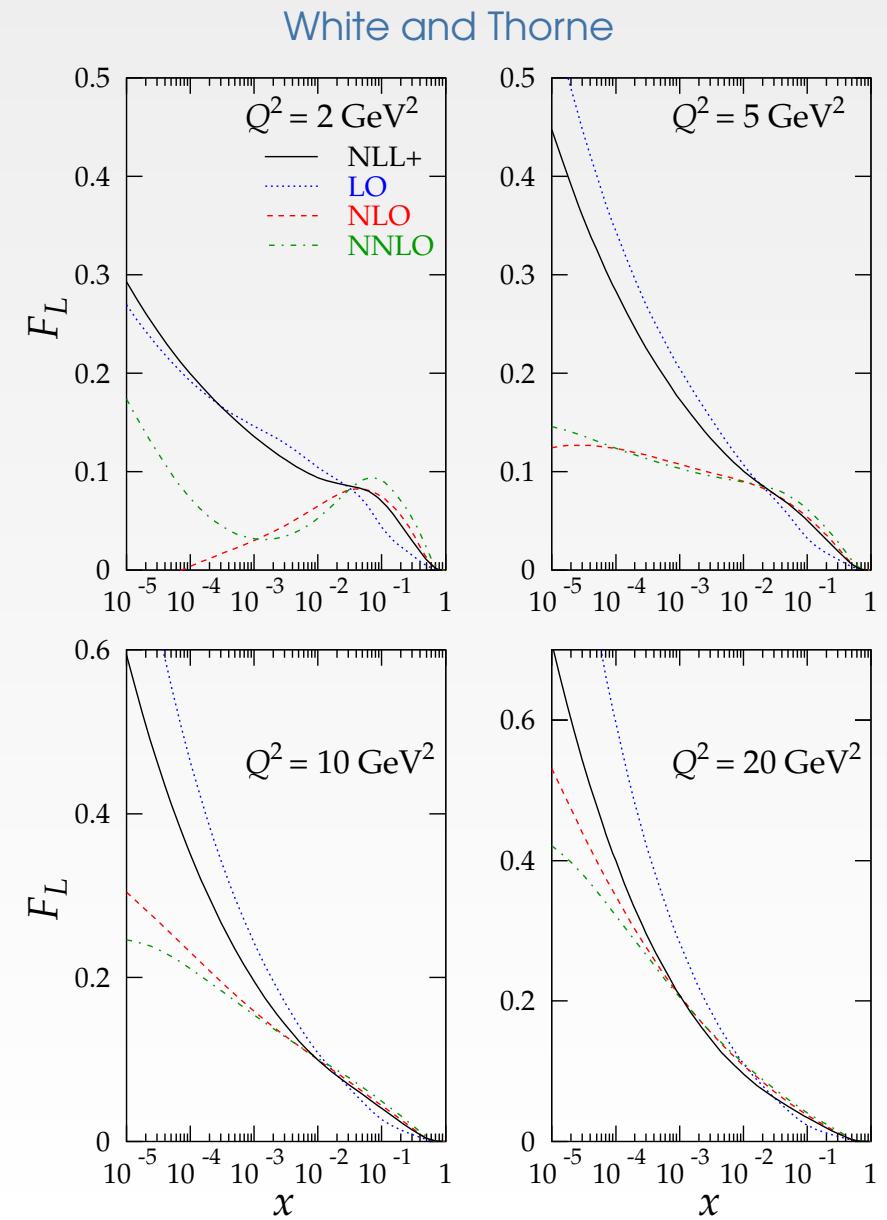
Longitudinal Structure Function F_L

$F_L(x, Q^2)$ – longitudinally polarised photons

- ◆ $F_L = 0$ in Quark Parton Model
- ◆ Sensitive to QCD higher orders (gluon emission)
 $F_L \sim \alpha_s x g$ — **constrains gluon density**

Calculations

- Large spread of calculations for gluon and F_L
- Critical corner – low Q^2 and low x
- Can be used to test resummation approaches



Future Low Energy Run

Cross section:

$$\sigma(x, Q^2) \propto F_2(x, Q^2) - \frac{Q^4}{s^2 x^2} F_L(x, Q^2)$$

Measure σ at the same x, Q^2 for different s

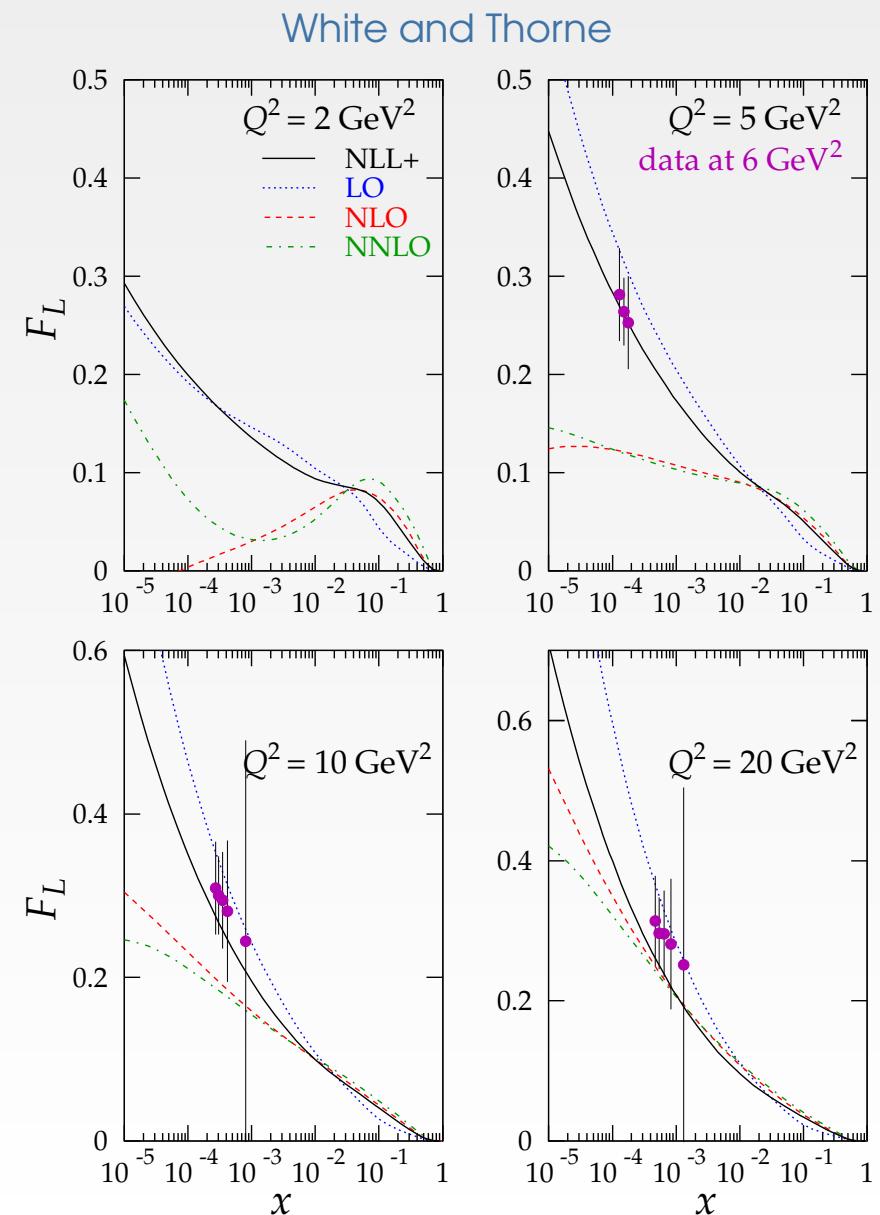
H1 Simulation

30 pb^{-1} at $E_p = 920 \text{ GeV}$

10 pb^{-1} at $E_p = 460 \text{ GeV}$

Can differentiate between calculations

Decision taken. Run in preparation



Last Question

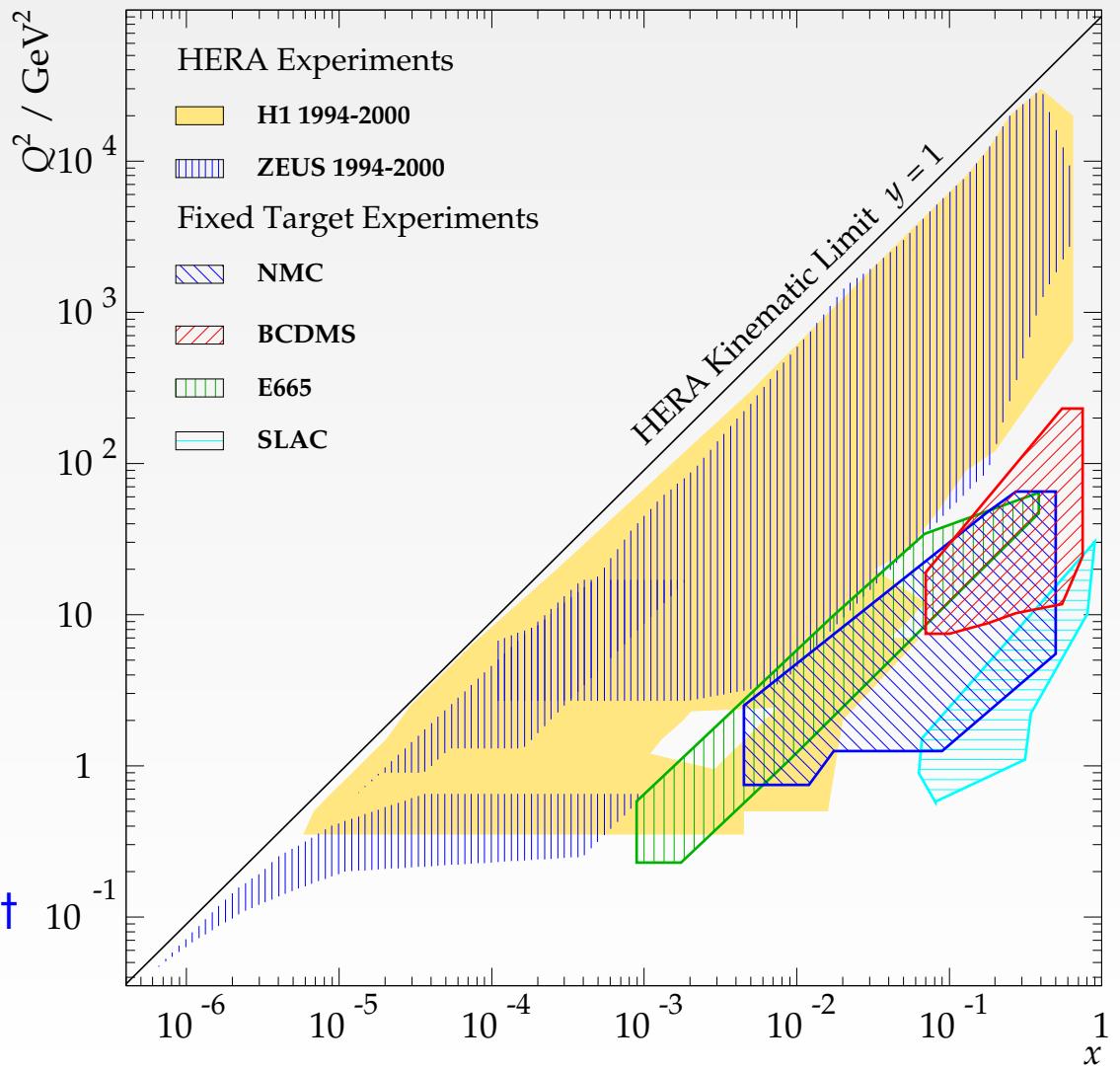
- ◆ How good is our knowledge of PDFs?
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- ◆ Are non-linear QCD effects relevant?

Going to High Parton Densities (Lowest x)

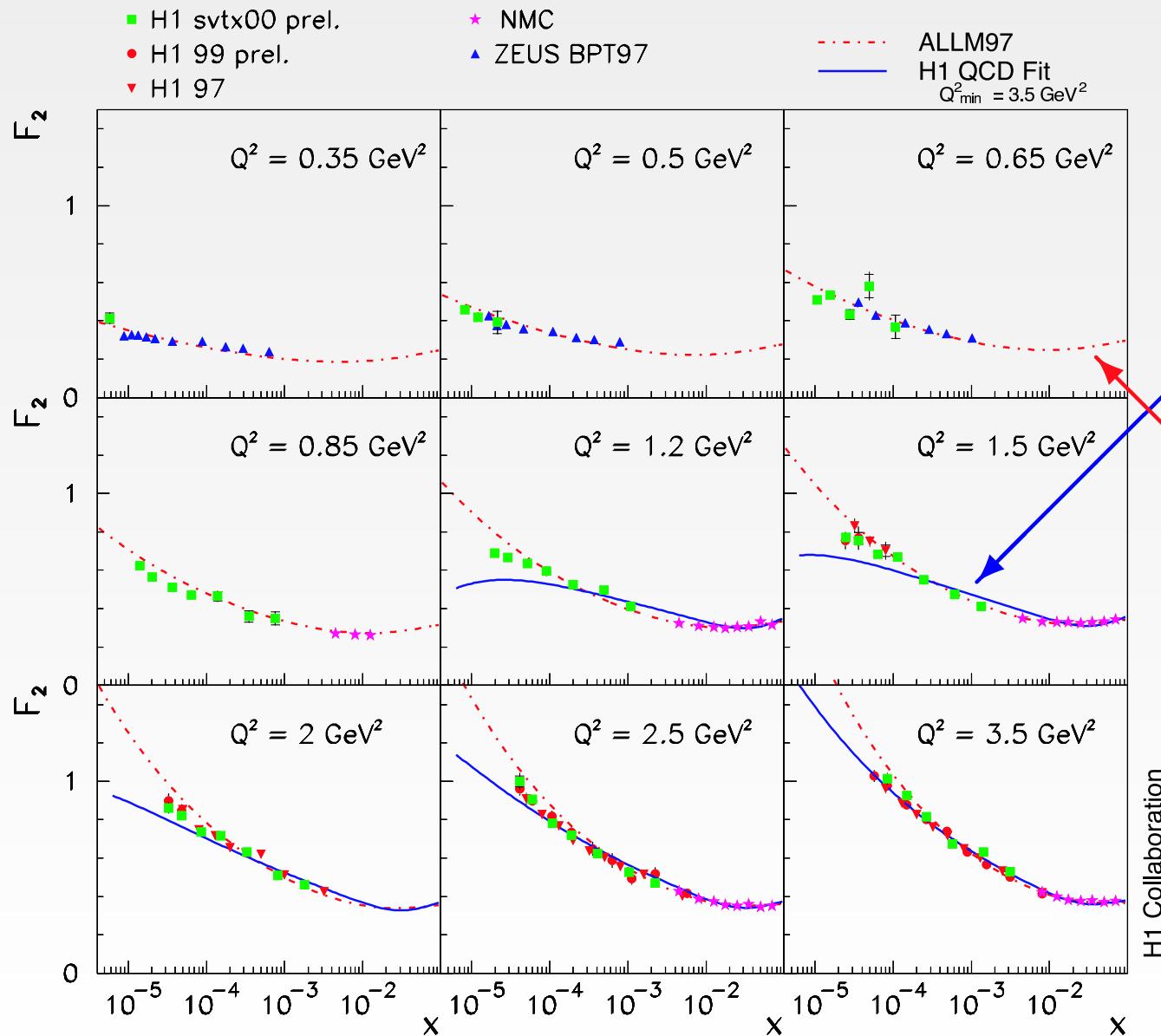
Non-linear QCD dynamics = multi-gluon exchange

Gluon density rises steeply towards low x

Lowest $x \implies$ Lowest Q^2
 $\alpha_s(Q^2)$ becomes large \implies confinement
quarks are no more free



Data in Transition Region



- ◆ At $Q^2 \lesssim 2 \text{ GeV}^2$, pQCD fits fail
- ◆ Only phenomenological models describe data
- ◆ Rise of F_2 at $x < 10^{-2}$ is well described by

$$F_2 \propto x^{-\lambda(Q^2)}$$
 as in pQCD region

Extraction of $\lambda(Q^2)$

$$F_2 = c(Q^2) \cdot x^{-\lambda(Q^2)}$$

► At $Q^2 \gtrsim 2 \text{ GeV}^2$:

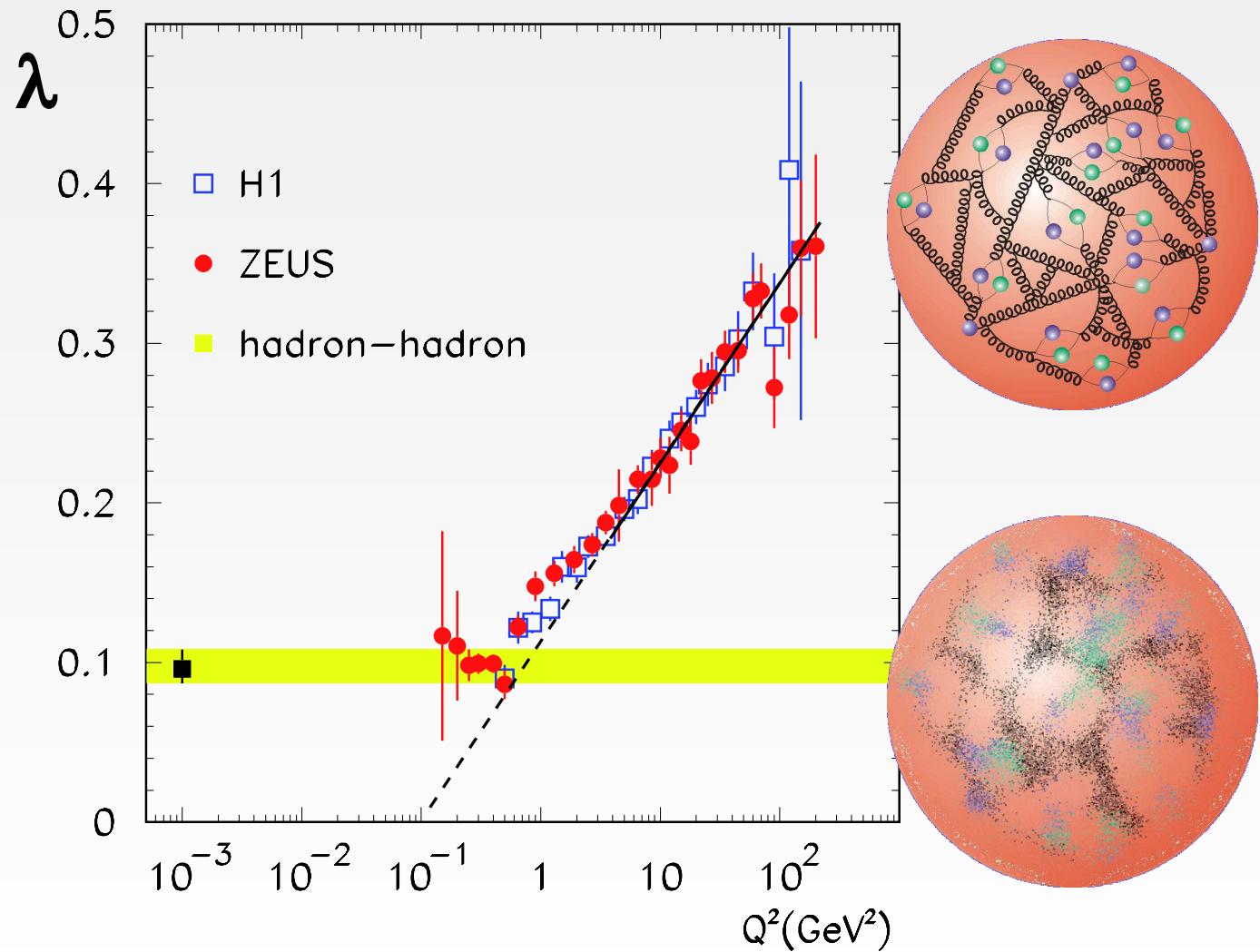
$$\lambda \sim \ln Q^2$$

Partonic degrees of freedom

► At $Q^2 \lesssim 2 \text{ GeV}^2$:

Transition to hadronic d.o.f.

$Q^2 \rightarrow 0$: $\lambda \rightarrow 0.08$ (Regge model)



Extraction of $\lambda(Q^2)$

$$F_2 = c(Q^2) \cdot x^{-\lambda(Q^2)}$$

► At $Q^2 \gtrsim 2 \text{ GeV}^2$:

$$\lambda \sim \ln Q^2$$

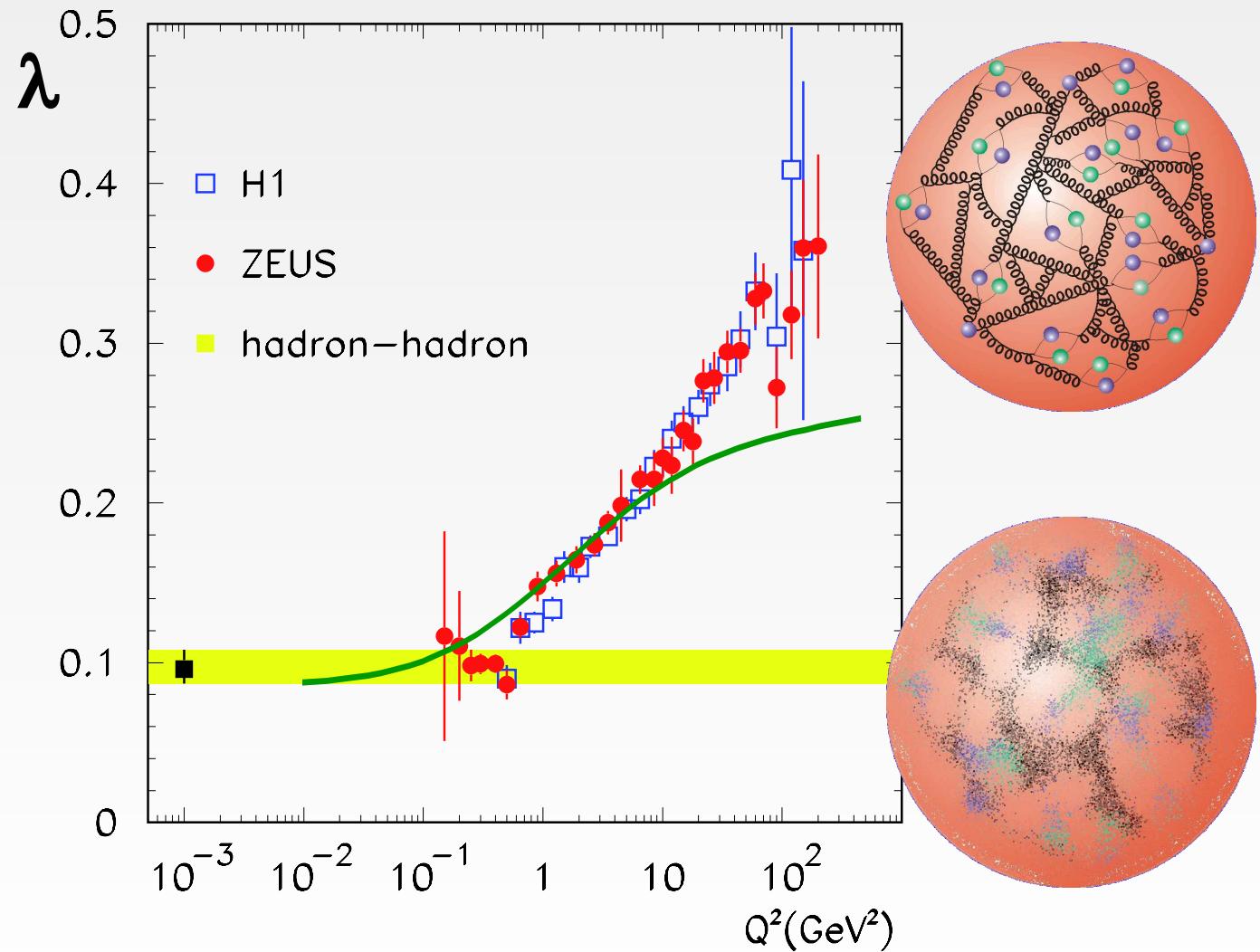
Partonic degrees of freedom

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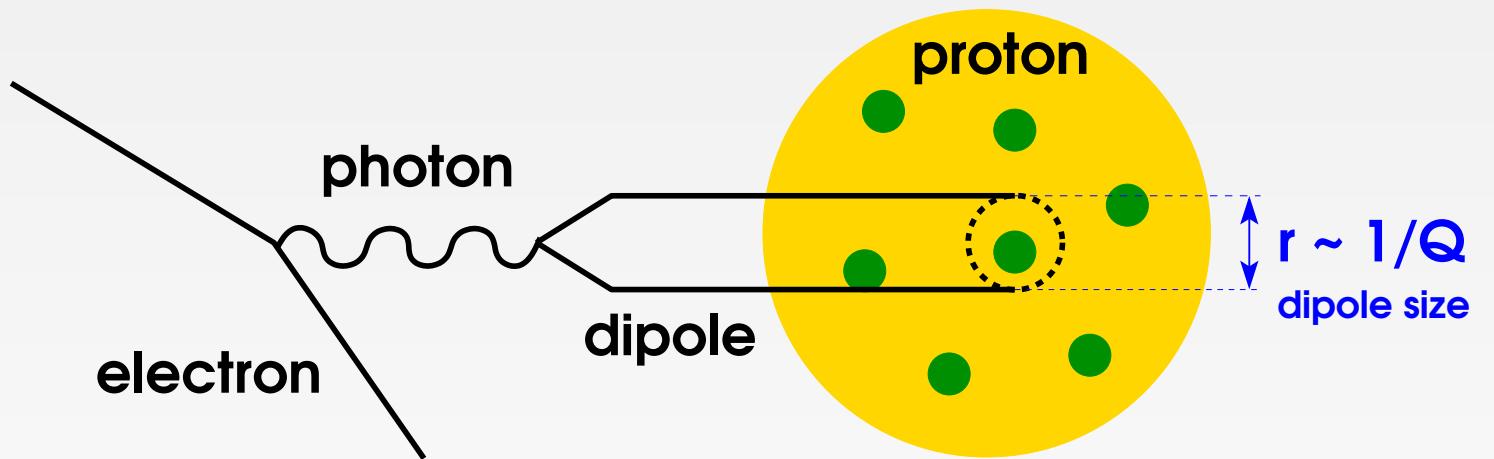
► Example dipole model



Dipole Models

Use colour $q\bar{q}$ dipoles as degrees of freedom

Proton rest frame: Photon fluctuates in $q\bar{q}$ pair which interacts with proton



$$\sigma_{T,L}^{\gamma^* p} = \int d^2 r \dots | \psi_{T,L}(r, \dots) |^2 \hat{\sigma}(r, \dots)$$

photon wave function
(calculable)

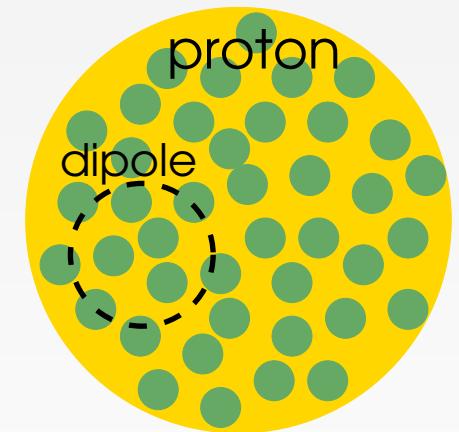
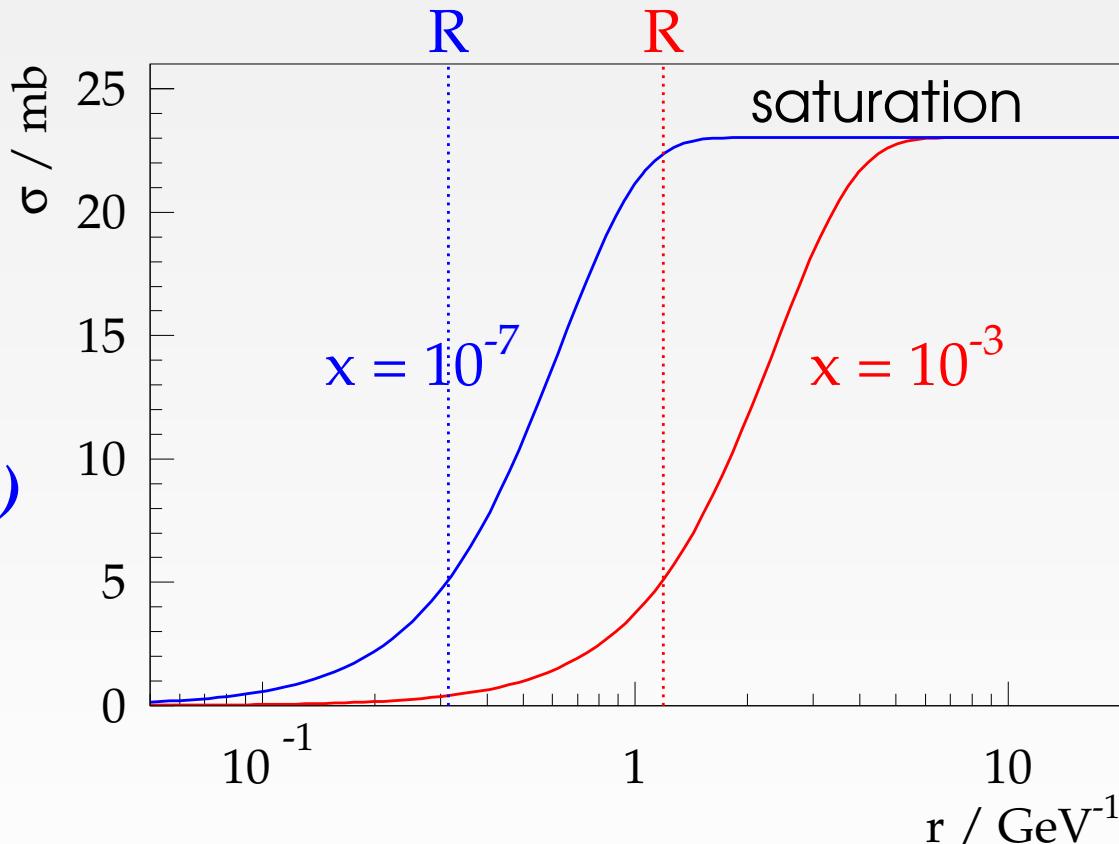
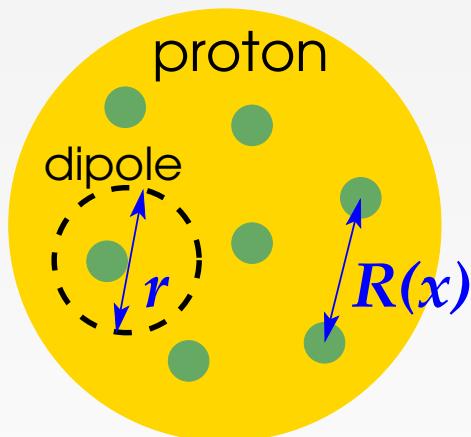
dipole x-section
transverse gluon distribution

Dipole Saturation Model

Photon fluctuates in $q\bar{q}$ pair which interacts with proton

Golec-Biernat, Wüsthoff

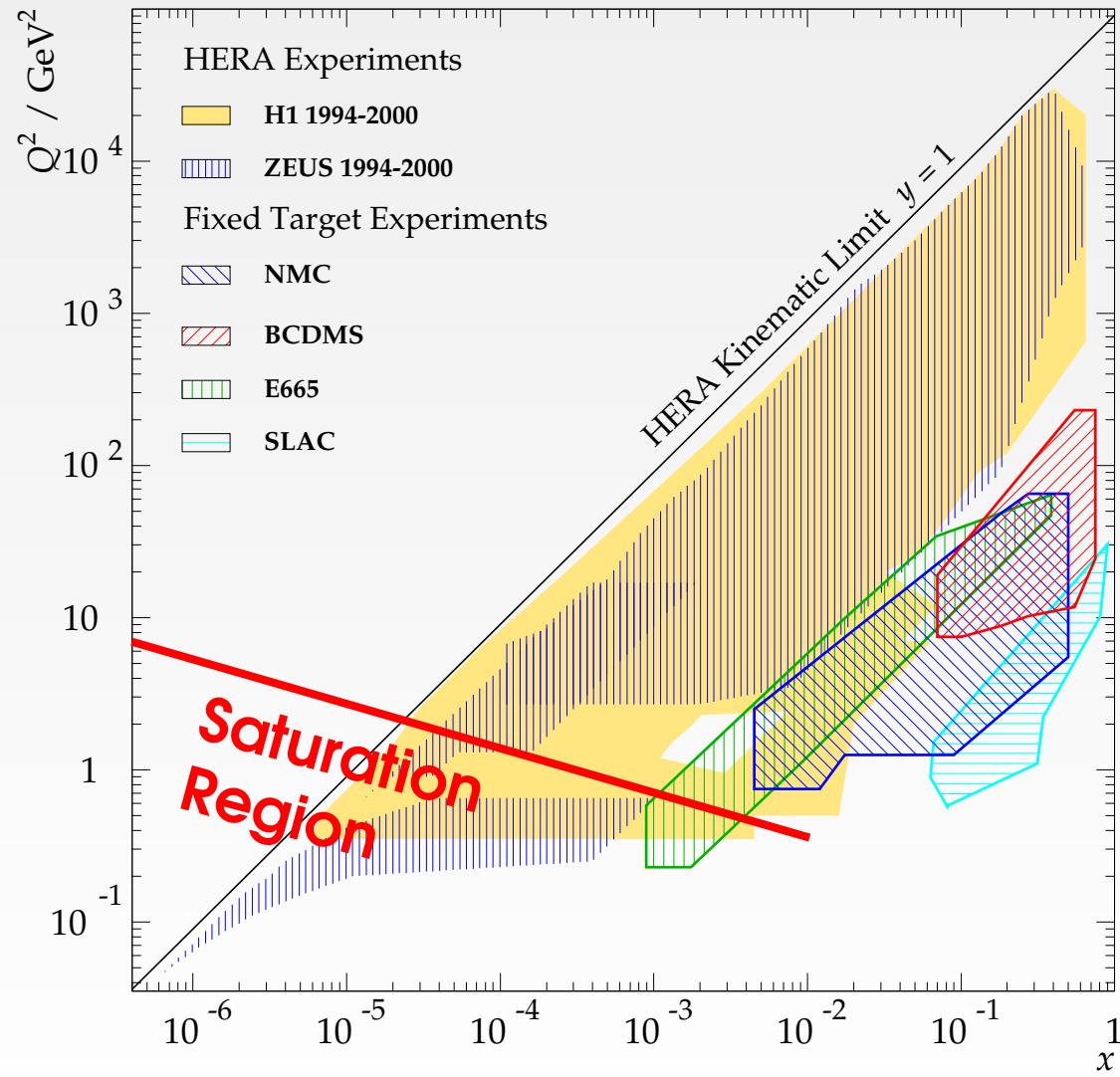
Dipole–proton cross section:



$$\text{At small } r: \hat{\sigma} \sim \frac{r^2}{R(x)^2}$$

At large r : non-linear interactions → **saturation**

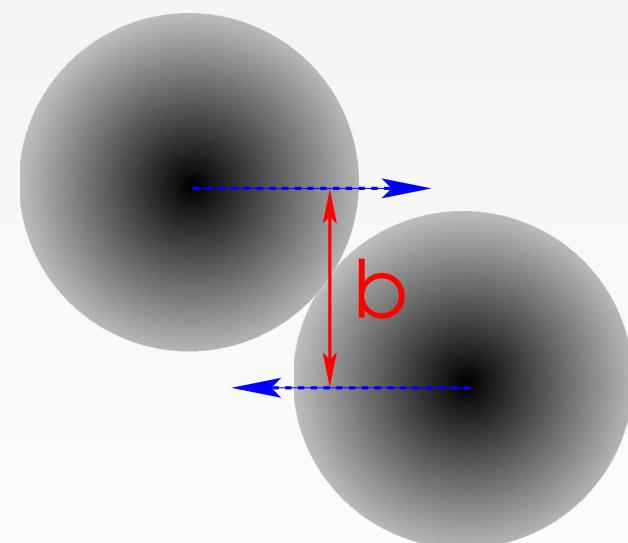
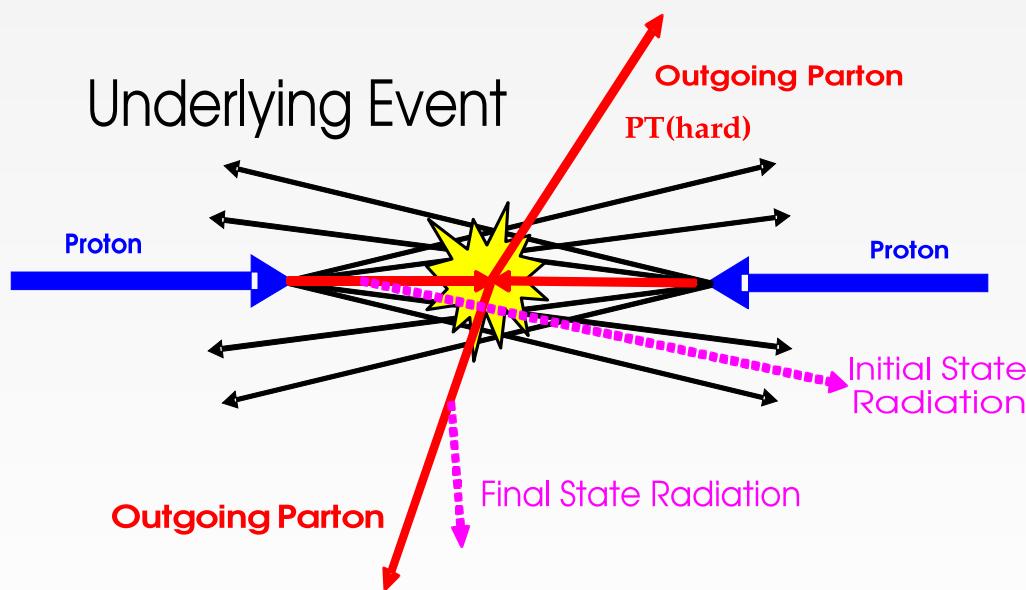
Saturation Region in Dipole Model



- ◆ For $Q^2 \lesssim 1 - 2 \text{ GeV}^2$ saturation model describes transition to soft interactions with only 3 parameters
- ◆ For pQCD Q^2 scales saturation region is beyond HERA reach
- ◆ Still, non-linear effects can affect pQCD evolution at low x and Q^2

Dipole Models

- ◆ Dipole models describe very successfully inclusive diffraction and exclusive channels (light VM, J/ψ , DVCS)
- ◆ They can be used to describe diffraction at pp and multiple interactions (underlying event, minijets ...)



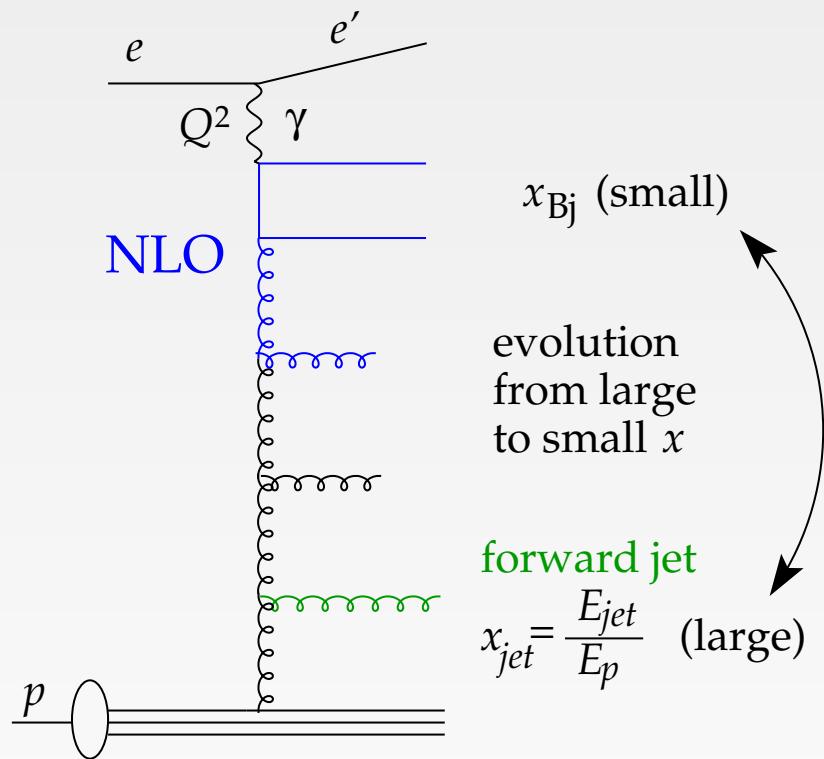
Summary

- ◆ The quest for precision and deeper understanding continues
 - Experiments → statistics, combined data, F_L, \dots
 - Theory → NNLO, resummation, non-linear effects ...
- ◆ DGLAP limitations are clearly visible in semi-inclusive measurements
 - Alternative parton cascade models
- ◆ Models for non-linear dynamics are further developed
 - Transverse picture of the proton
 - Understanding of soft hadron interactions
 - Diffraction, Multiple Interactions

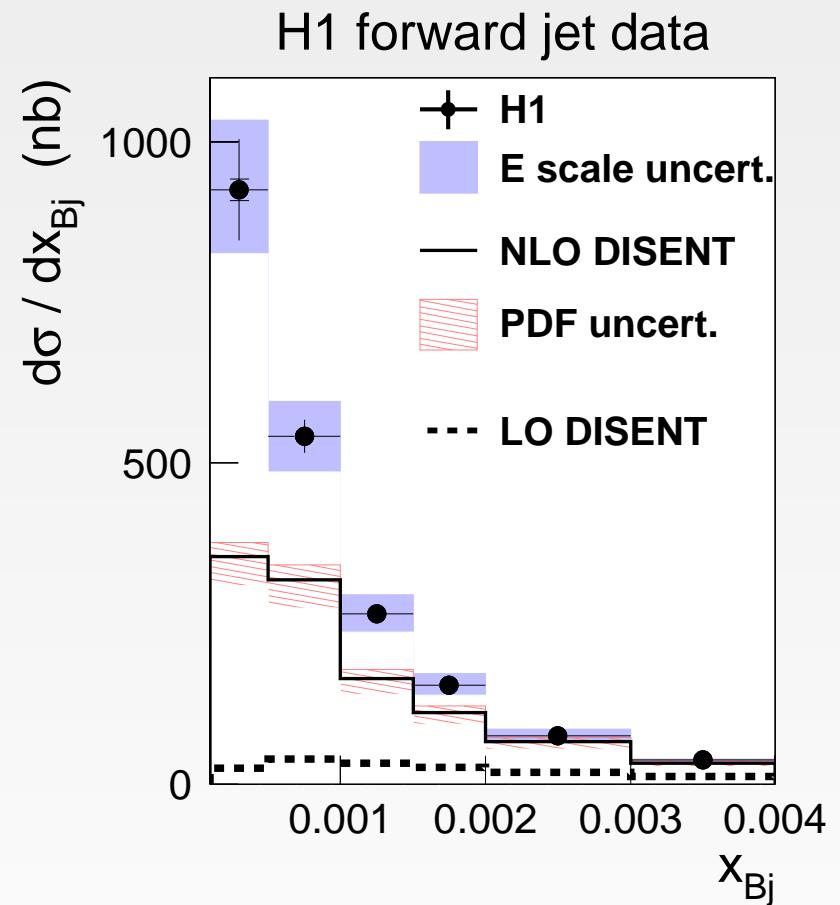
Backup

Additional Information

Forward Jets at Low χ

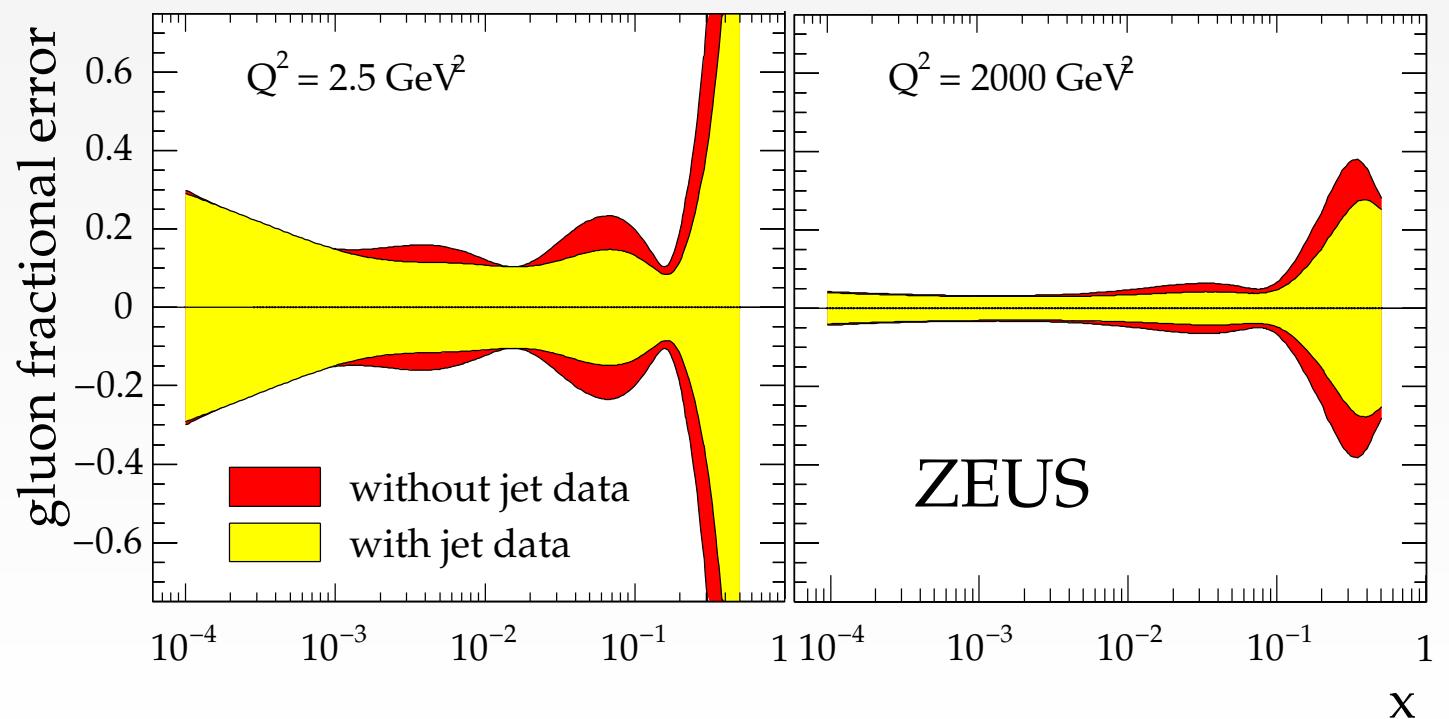
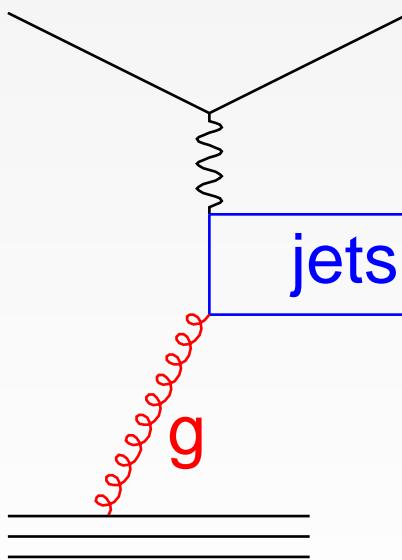


- ◆ Low x – long parton chain
- ◆ Look at forward jet – start of the chain
- ◆ NLO is not sufficient

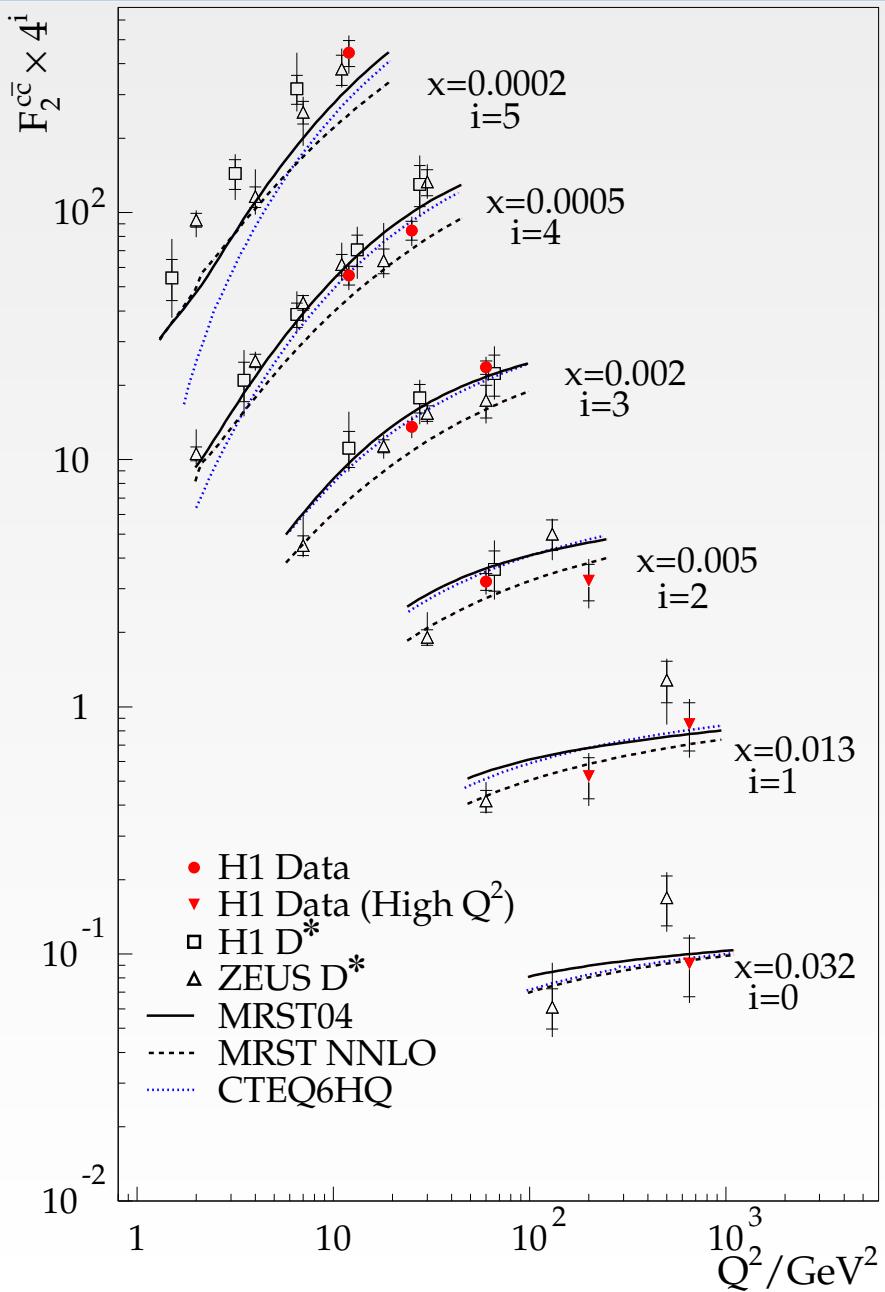


Including Jet Measurements in \mathcal{PDF} Fits

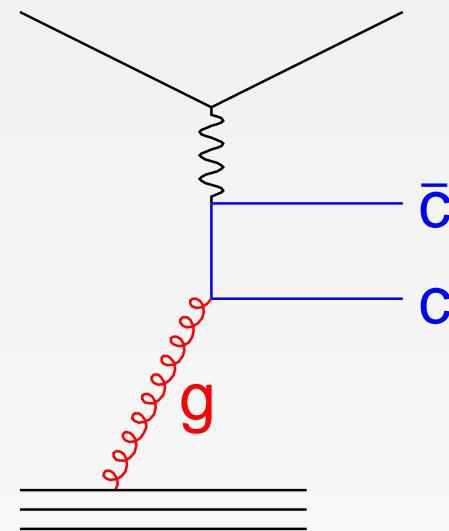
- ◆ HERA data are stat. limited at high x (high Q^2)
- ◆ Theor. group include TeVatron jet data
Large systematics due to jet energy scale
- ◆ ZEUS included its jet data in PDF fits
⇒ Improved gluon at medium–high x



Heavy Flavour Measurements



- ◆ Inclusive c and b based on long lifetime
- ◆ Inclusive D^* production



Possible impact on PDFs

- ▶ Sea decomposition
($F_2^{c\bar{c}}$ already used by some theor. fit groups)
- ▶ Improve gluon distribution?

Large statistics is expected at HERA II

Indirect Extraction of F_L

◆ NC ep cross section

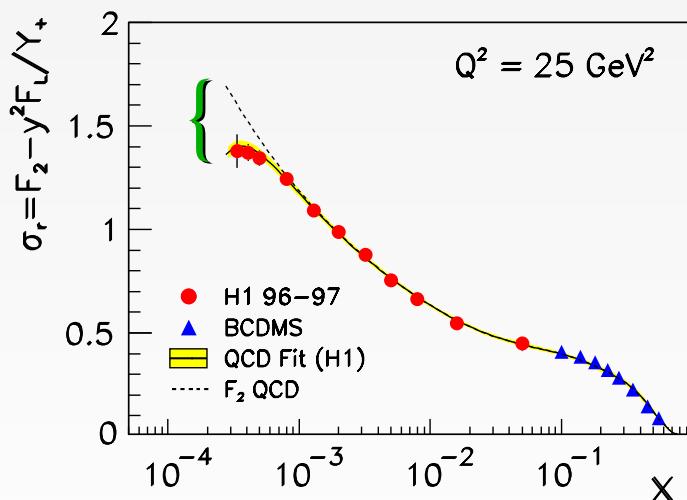
Inelasticity $y = Q^2/(xs)$

$$\frac{d^2\sigma}{dx dQ^2} = \frac{2\pi\alpha^2}{Q^4 x} Y_+ \left\{ F_2(x, Q^2) - \frac{y^2}{Y_+} F_L(x, Q^2) \right\}$$

$$Y_{\pm} = 1 \pm (1 - y)^2$$

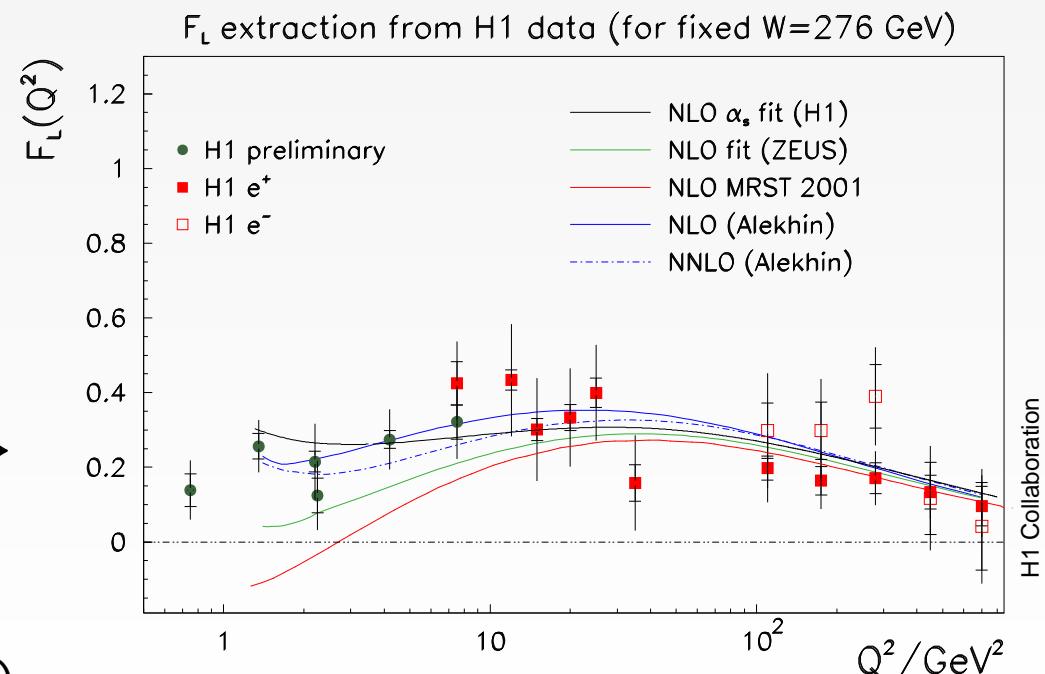
F_L contribution is significant only at high y = at low x

◆ Reduced cross section $\sigma_r = F_2(x, Q^2) - \frac{y^2}{Y_+} F_L(x, Q^2)$

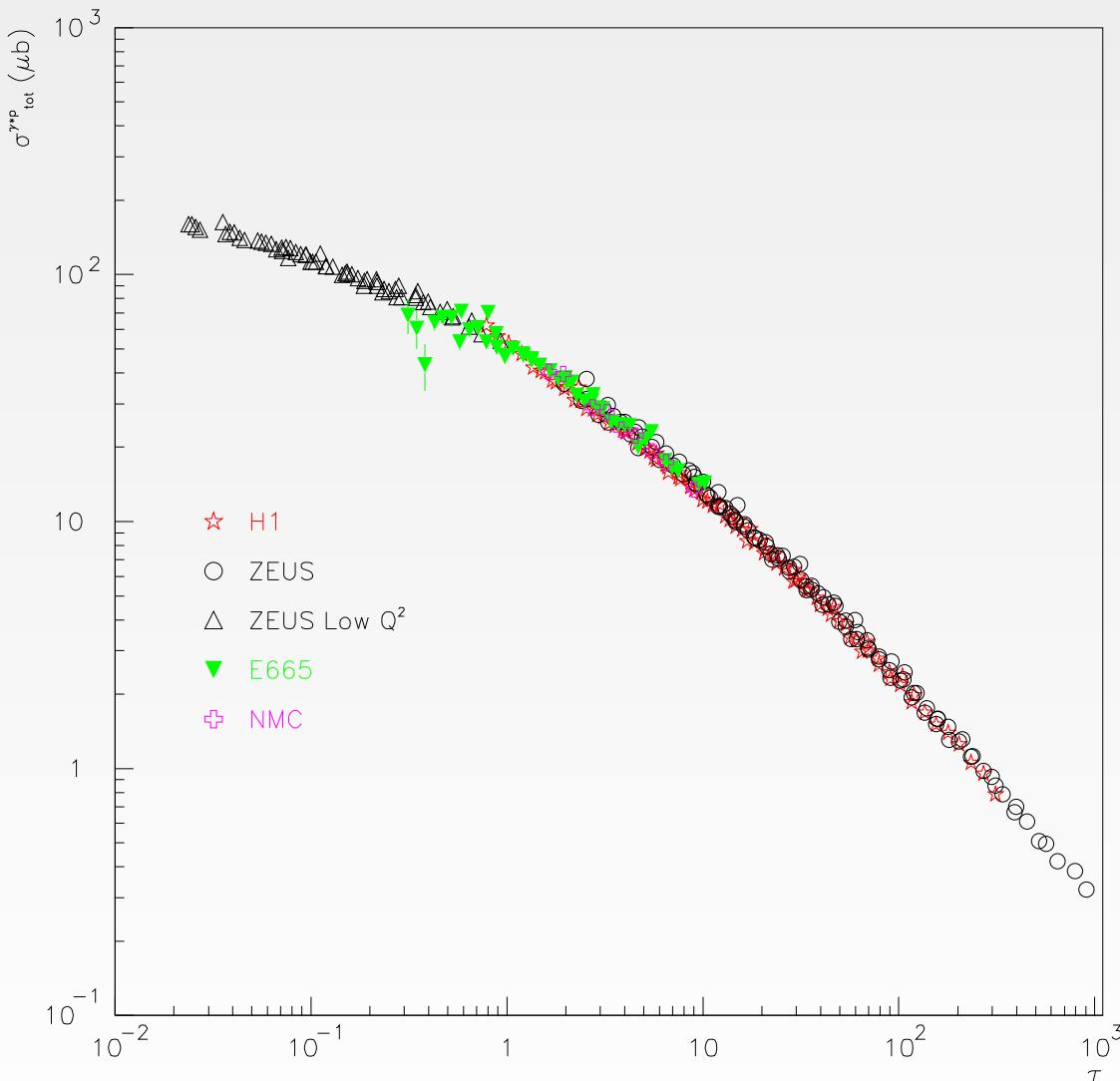


Drawbacks

- No x -dependence
- F_2 -model dependent (esp. at low Q^2)

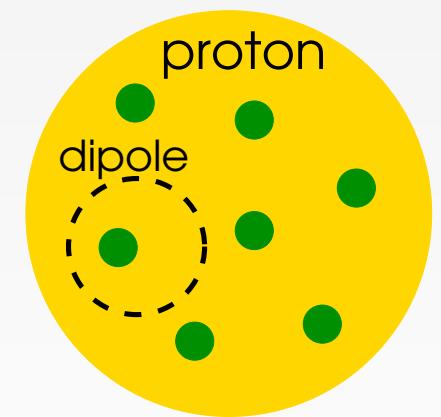
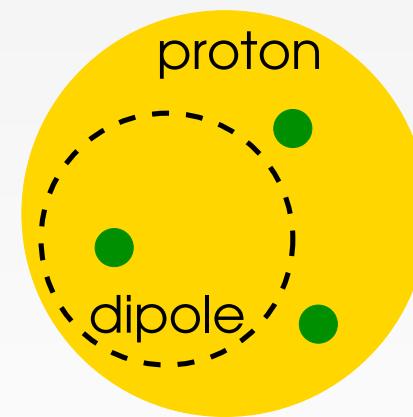


Geometric Scaling at $x < 10^{-2}$



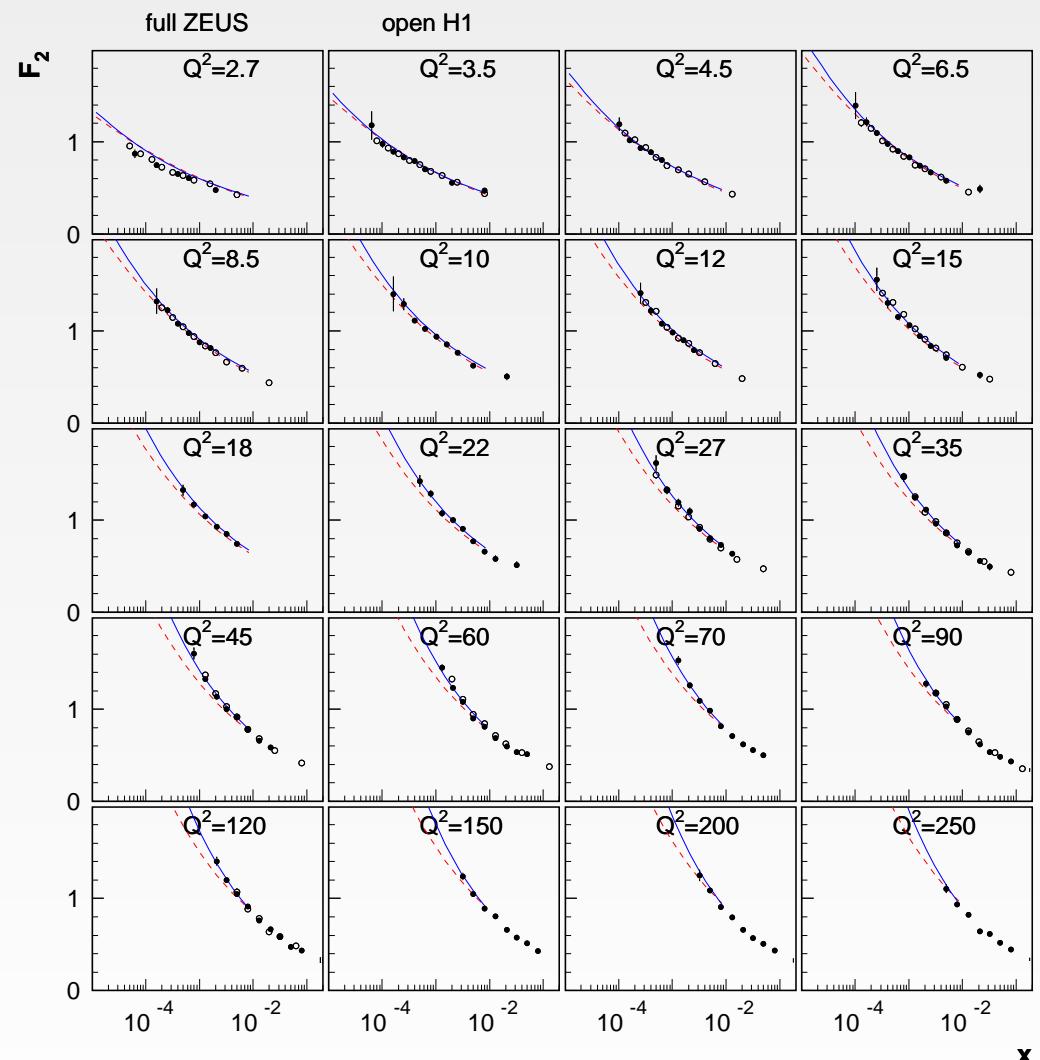
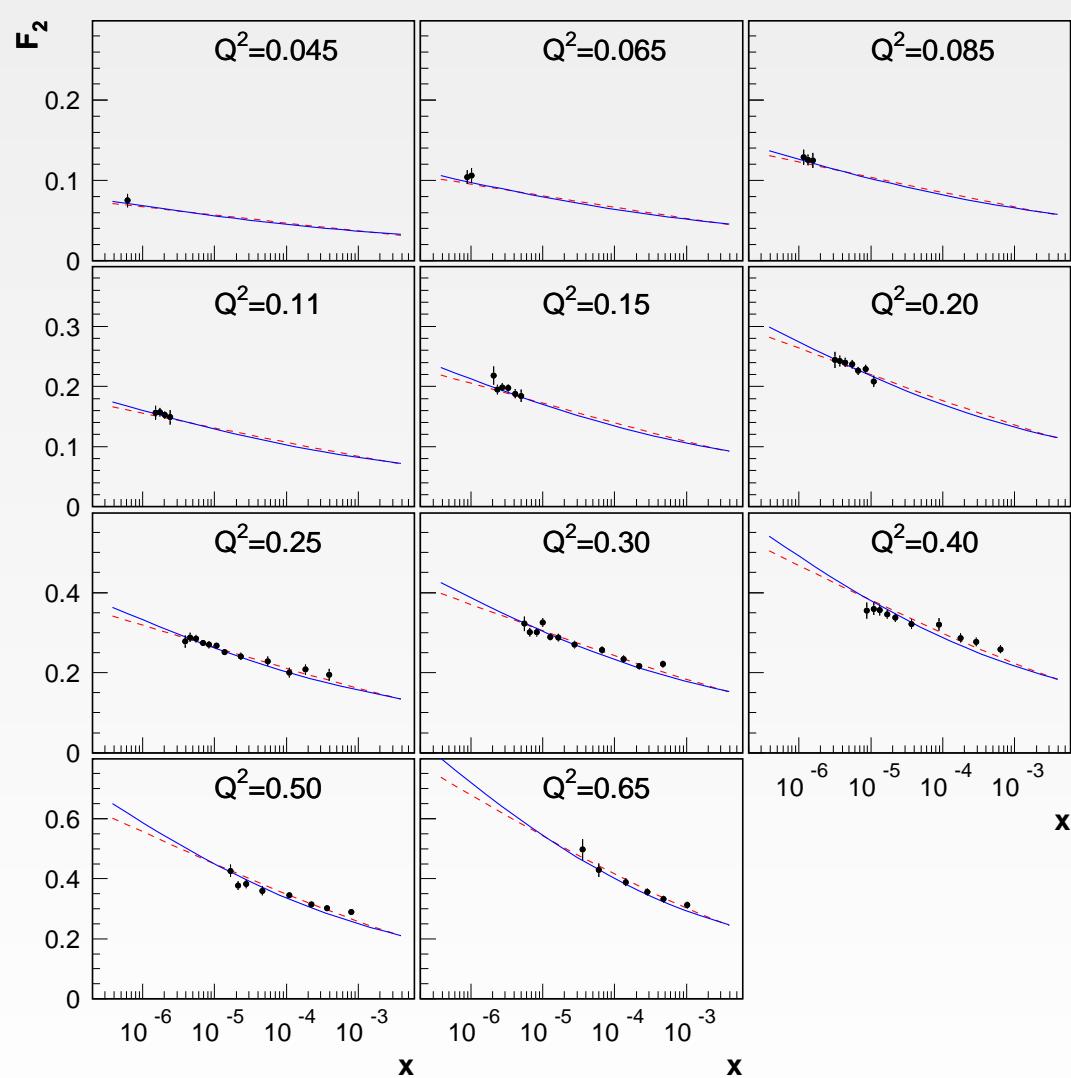
$$\tau = Q^2 R_0^2(x) \sim \frac{R_0^2(x)}{r^2}$$

$\sigma_{\gamma^* p}$ is function of only one variable τ :
Data manifest existence
of saturation scale
as used in dipole model



At high τ : $\sigma \sim 1/\tau$
At lower τ : saturation?

F_2 Description by Saturation Model



Bartels, Golec-Biernat, Kowalski

GBW (+ DGLAP evolution in pQCD region)