

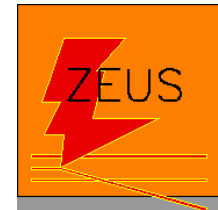
# Charged multiplicities in inclusive and diffractive deep-inelastic ep scattering at HERA

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On behalf of H1 and ZEUS collaborations

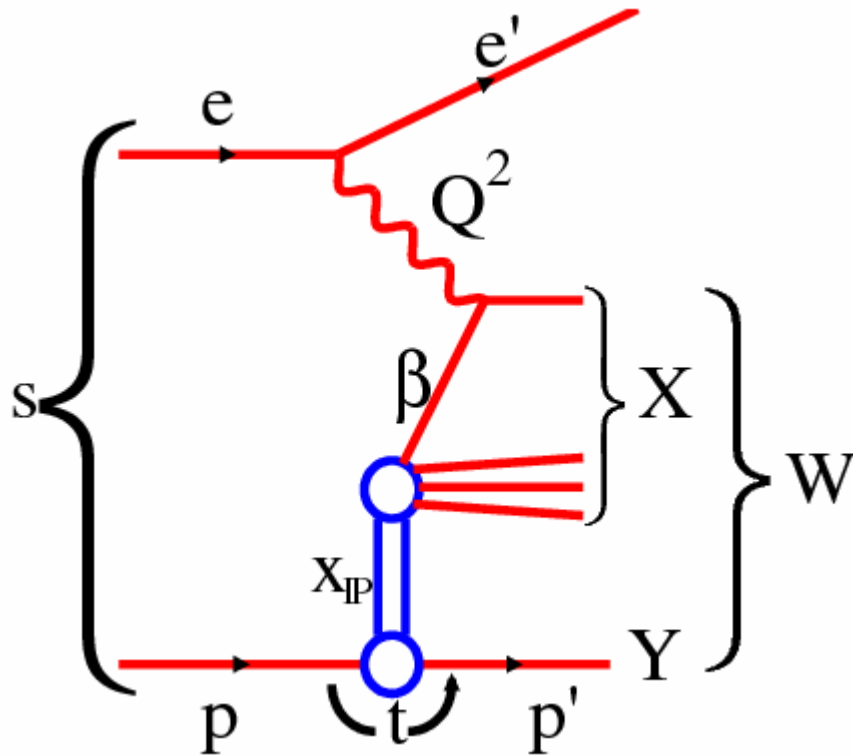
abstracts 377 & 642



## Outline

- H1: Charged particle multiplicity (distributions)
  - Kinematic dependences of  $\langle n \rangle$  in DDIS
  - Comparison of DIS with diffractive DIS (DDIS)
- ZEUS: What is correct energy scale in Breit-frame and in hadronic CM frame for comparisons with other processes?
- Apologies for incomplete coverage

# Diffraction: $ep \rightarrow e' + X + Y$



$M_X$ : inv. mass of diffractive system X

~ 10% of DIS events have a rapidity gap

$$t = (p - p')^2$$

$$\beta = x_{quark}/IP$$

$$x_{IP} = x_{IP}/proton$$

$$M_X^2 = Q^2 \frac{1-\beta}{\beta}$$

$$\text{rapidity-gap} \sim \ln \frac{1}{x_{IP}}$$

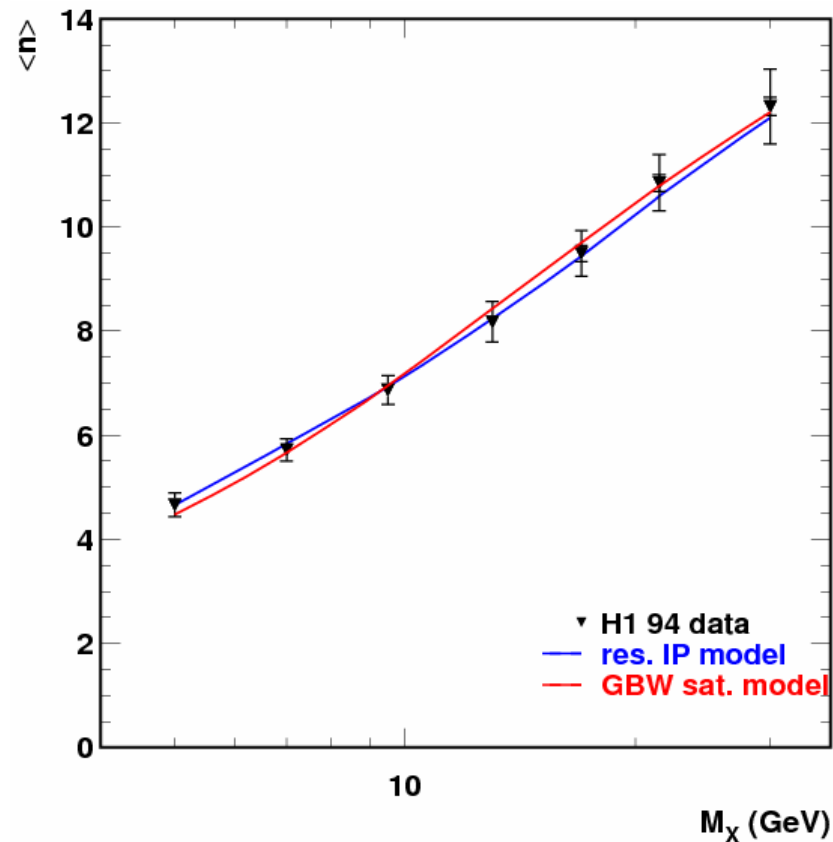
# H1 analysis : Motivation

- Previous H1 analysis on 94 data
  - Dependence of  $\langle n \rangle$  on  $M_X$  only

DDIS:  $W, x, Q^2, \beta, x_{IP}, t, M_X$

Which kinematic variables  
are relevant for multiplicity?

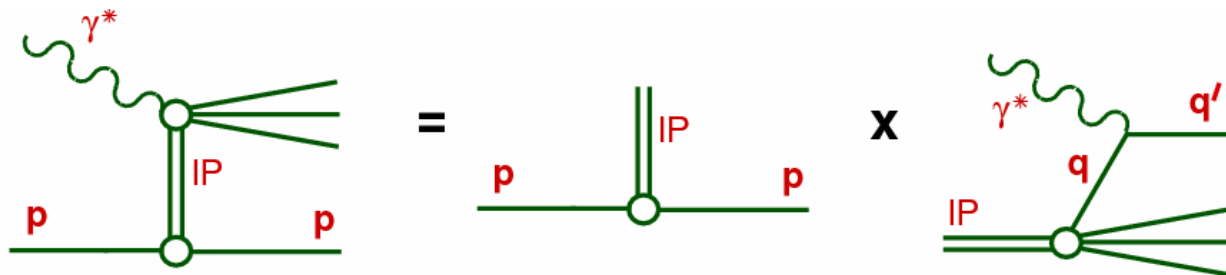
- H1 2000 DDIS data:
  - Large statistics allows more differential study:
    - $W, Q^2, \beta$  dependences at fixed  $M_X$
  - Compare DIS and DDIS



# A model for diffraction

## Combine QCD & Regge theory: resolved Pomeron model

- Proton infinite momentum frame
- Colorless **Pomeron (IP)** is built up of quarks/gluons  $\rightarrow$  diffractive PDF's
- Based on QCD and **Regge factorization**: naïve, probably incorrect but works...!



$$F_2^D(x_{IP}, t, \beta, Q^2) = f_{IP/p}(x_{IP}, t) F_2^{IP}(\beta, Q^2)$$

- **Regge factorization implies diffractive final state is independent of proton (fractional) energy loss  $x_{IP}$**
- Need sub-leading **Reggeon (IR)** component besides Pomeron to fit the  $F_2^D$  Diffractive Structure Function data

# H1: Data selection DIS and DDIS

2000 nominal vertex data:  $46.65 \text{ pb}^{-1}$

Data corrected via Bayesian unfolding procedure:

-DIS MC: DJANGO 1.3, proton pdf CTEQ5L

-DDIS MC: RAPGAP resolved pomeron

## DIS selection:

- Good reconstruction of scattered electron
- Kinematic cuts:
  - $0.05 < \gamma_{av} < 0.65$
  - $5 < Q^2 < 100 \text{ GeV}^2$
  - $80 < W < 220 \text{ GeV}$

## DDIS selection:

- Rapidity gap:
  - No activity in the forward detectors
  - $\eta_{\max} < 3.3$
- Kinematic cuts:
  - $4 < M_X < 36 \text{ GeV}$
  - $x_{IP} < 0.05$

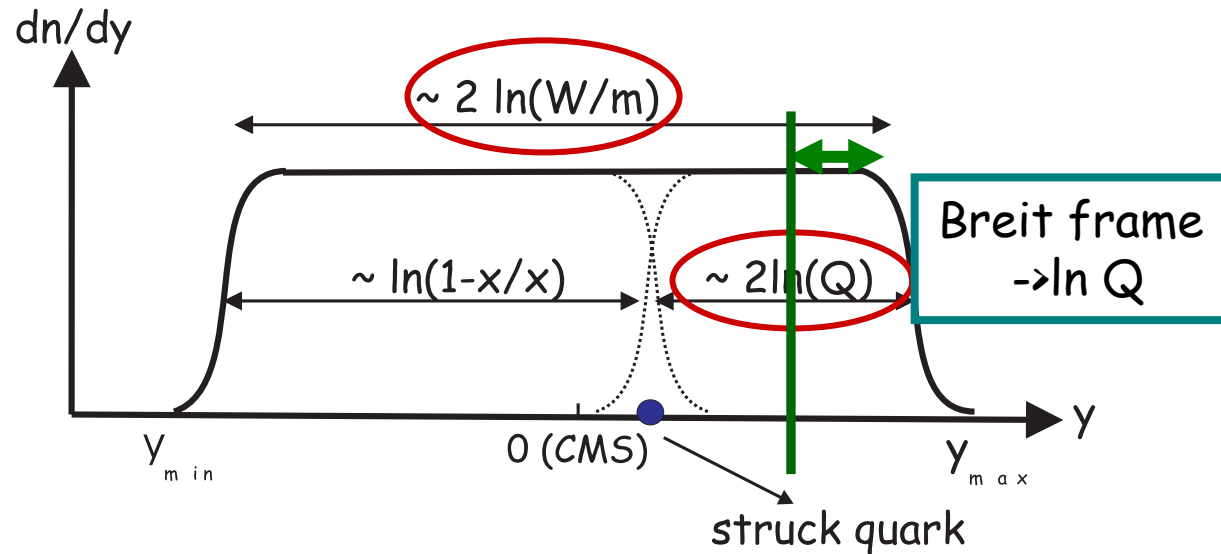
Use charged particles with  $\eta^* > 1$  (acceptance  $> 90\%$ )  
in  $\gamma^*p$  CMS frame

# Kinematics: Bjorken Plot

**DIS**

$$W^2 \sim Q^2/x$$

$$y_{\max} = \ln(W/m_\pi)$$

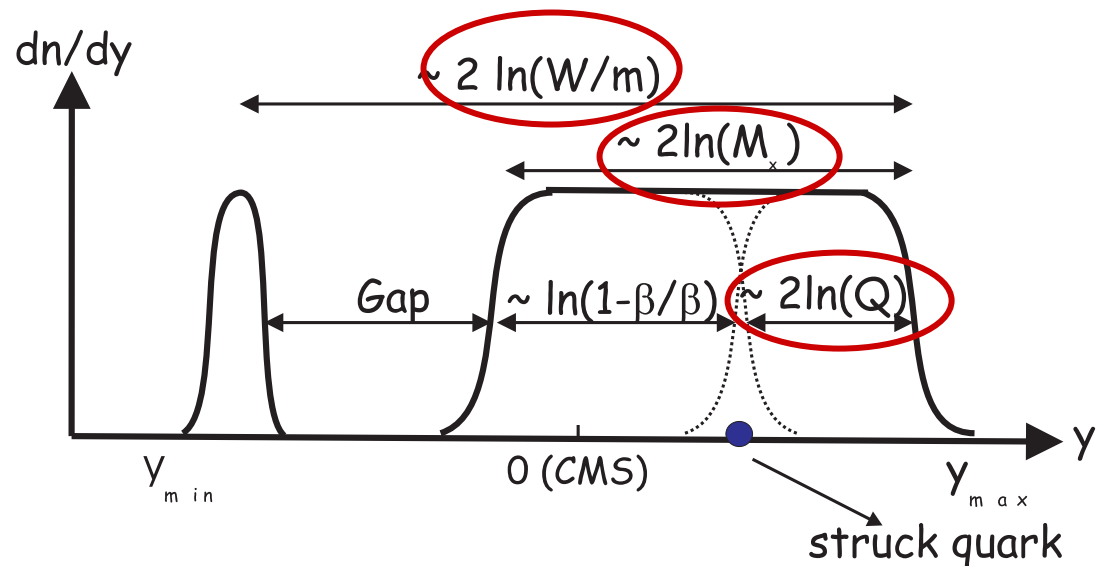


In DDIS  $\beta$  plays role of  $x$  in DIS

**DDIS**

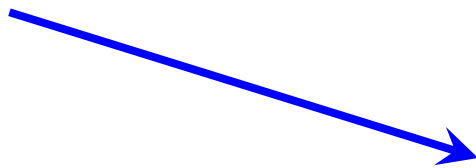
$$M_X^2 = Q^2 \frac{1-\beta}{\beta}$$

$$\text{gap} \sim \ln \frac{1}{x_{\text{IP}}}$$



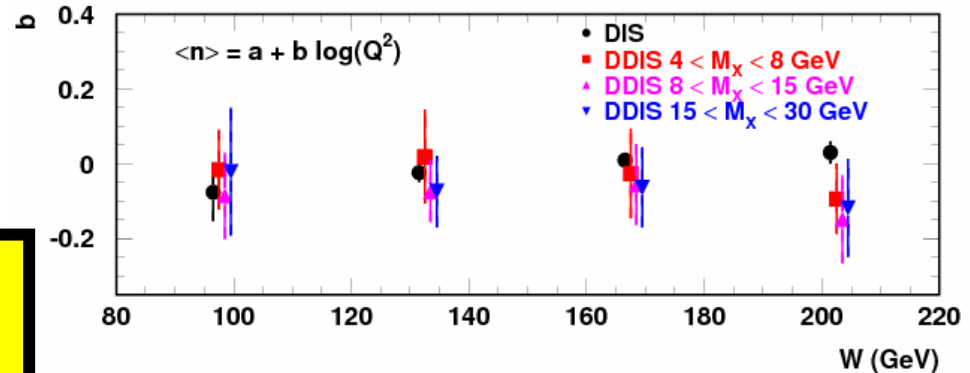
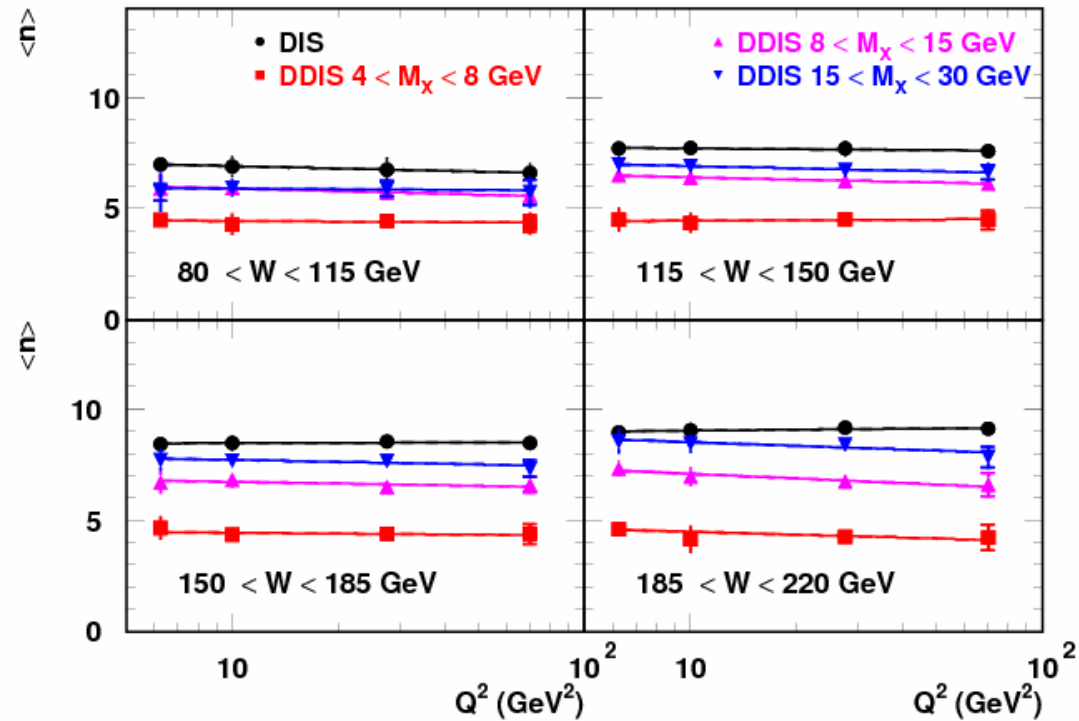
# H1: $\langle n \rangle (Q^2)$ in DIS & DDIS at fixed W

- DIS data
- DDIS data (fixed  $M_x$ )
- No stat. signif. dependence on  $Q^2$
- Weak W-dependence in DDIS
- Fit  $\langle n \rangle$  to  $\langle n \rangle = a + b \log(Q^2)$

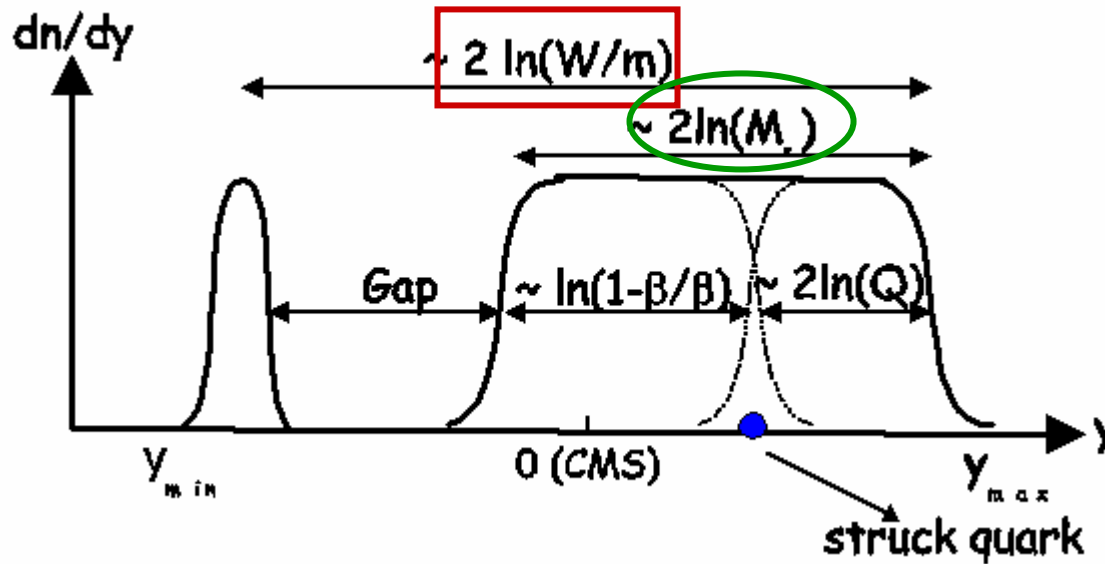


Further: Rapidity spectra show very weak  $Q^2$  dependence

H1 prel. ( $\eta^* > 1$ )



# W dependence of $\langle n \rangle$ in DDIS ?



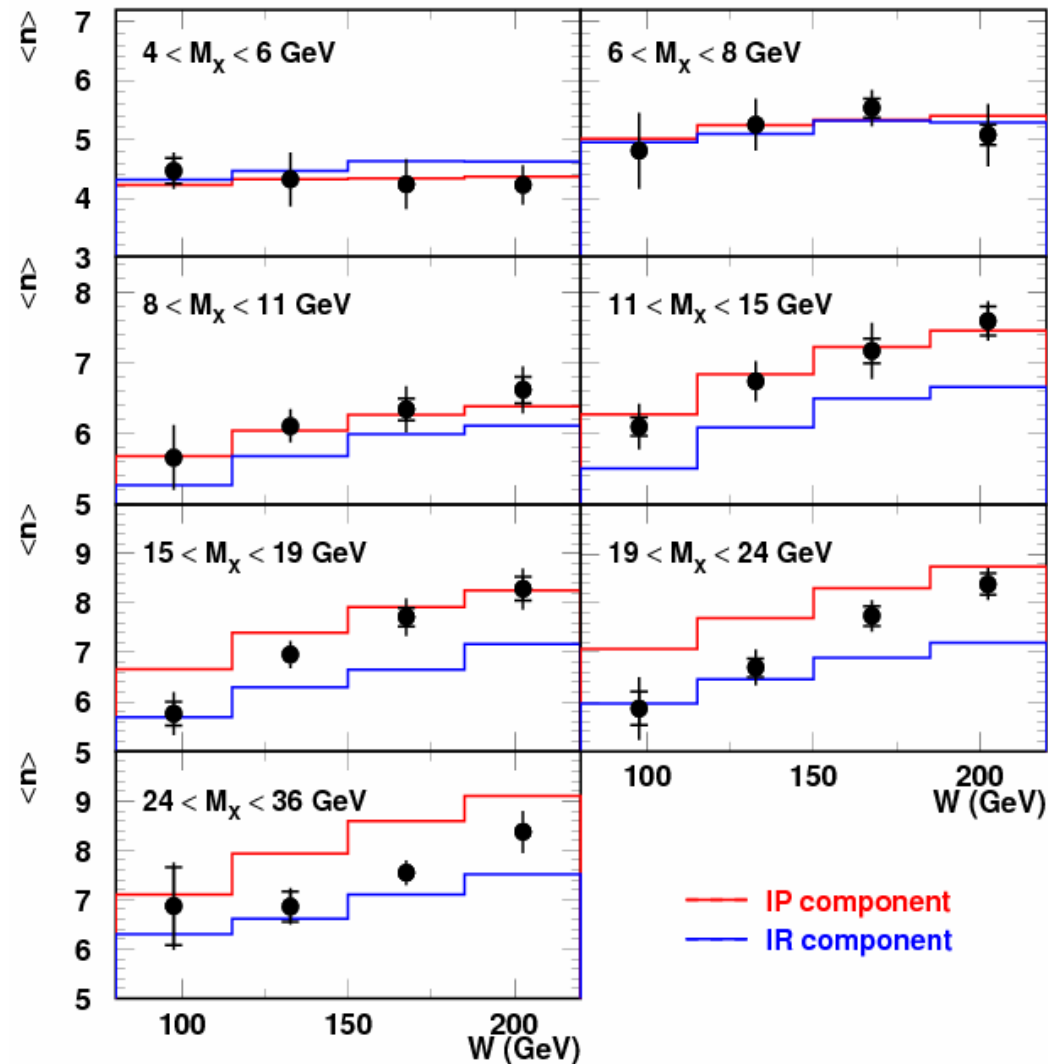
- Changing  $W$  with  $M_x$  fixed = changing gap width
- Gap  $\sim \ln(1/x_{IP})$  thus ... effectively  
investigate dependence on  $x_{IP}$  of  $\langle n \rangle$



# W dep. of $\langle n \rangle$ at fixed $M_x$ in DDIS

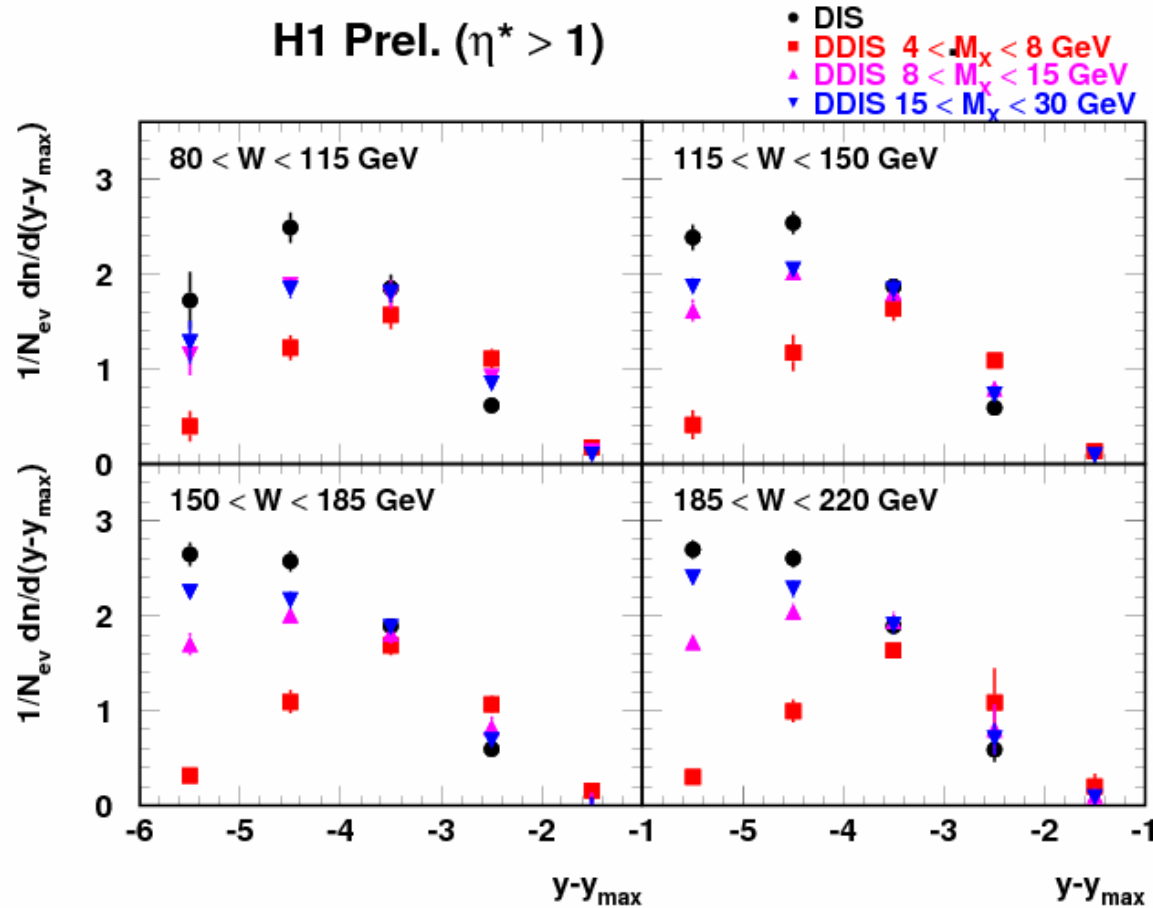
H1 Prel. DDIS ( $\eta^* > 1$ ) **All  $Q^2$**

- At fixed  $M_x$ : changing  $W$  means changing gap and  $x_{IP}$
- Regge factorization means diffractive pdf's AND Final state properties independent of  $x_{IP}$
- **W-dependence = breaking of Regge factorization**
- In resolved Pomeron model: pomeron + reggeon
- **Large  $M_x$ : Data move from Reggeon to Pomeron as  $W$  grows**



# Particle Density in $\gamma$ : DIS $\leftrightarrow$ DDIS

- $(1/N) \frac{dn}{d(y-y_{\max})}$
- Central region:  
particle density similar  
for DIS and high  $M_X$   
DDIS

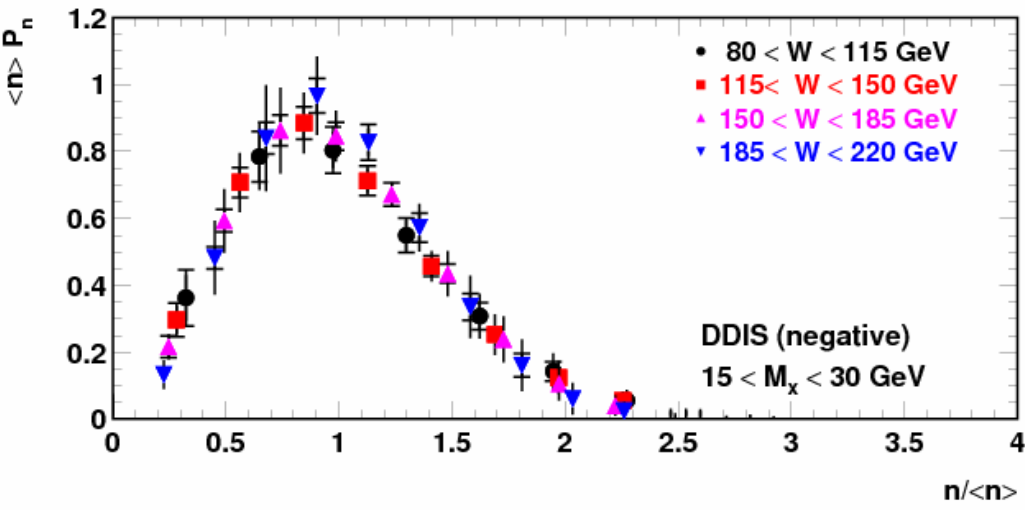
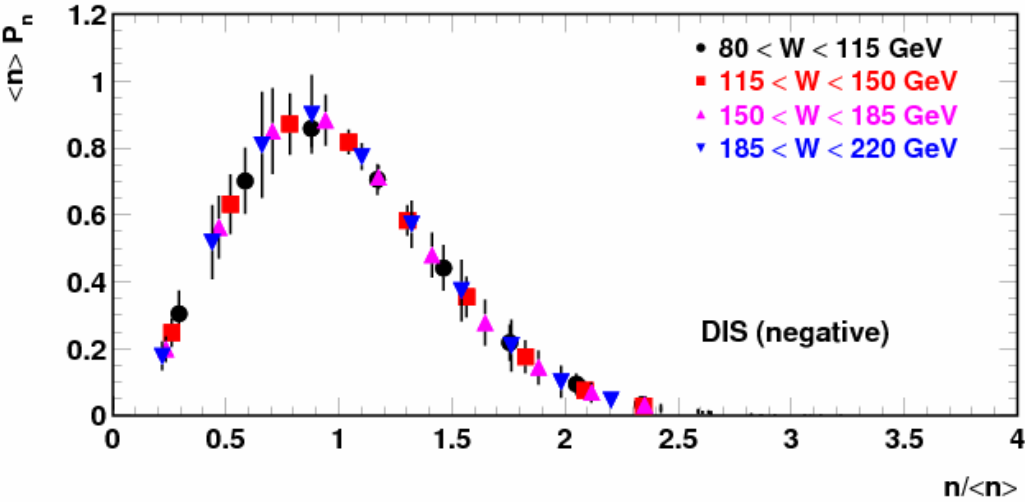


DDIS & DIS particle density not much different  
although  $M_X \ll W$

DDIS = gluon-rich system  $\rightarrow$  higher multiplicity

# Comparison DIS & DDIS: KNO scaling

H1 prel. ( $\eta^* > 1$ )



- Negative particles
- $\psi(z) = \langle n \rangle P_n$   
vs  $z = n/\langle n \rangle$
- Approximate KNO scaling for DIS and DDIS
- Shape of KNO distribution similar for DIS and DDIS
- Implies that correlations are very similar

# H1 : Summary

## Charged particle multiplicity

- studied for DIS and DDIS in ep at HERA

## Kinematic dependences

- $\langle n \rangle$  in DIS: main dependence only on  $W$ , not  $Q^2$  or  $x$  separately
- $\langle n \rangle$  in DDIS: main dependence on  $M_x$  not on  $Q^2$  and  $\beta$  separately (and a bit on  $W$ : Regge factorization breaking),

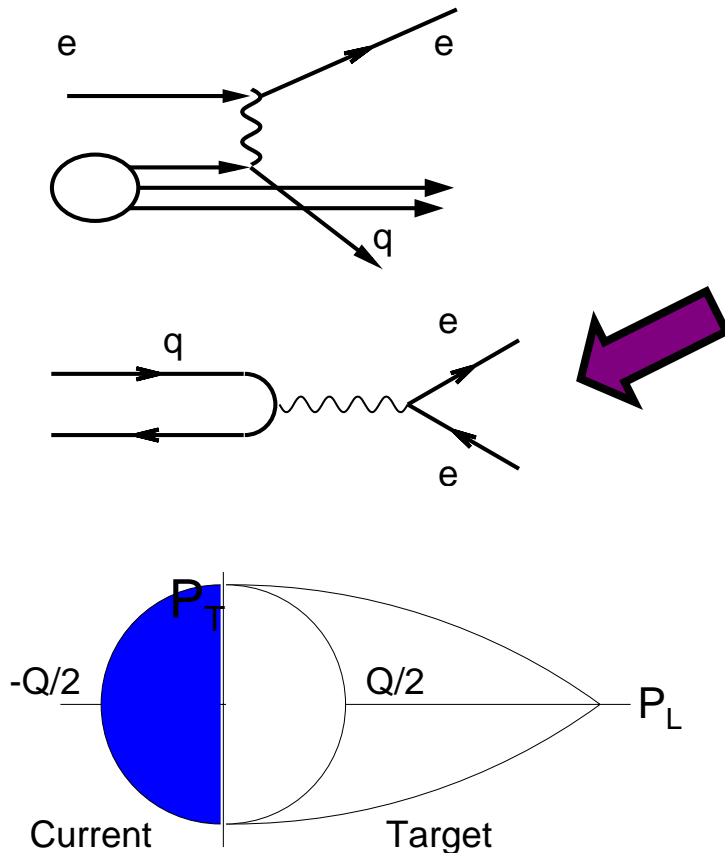
## Comparison DIS and DDIS

- DIS and DDIS: density in rapidity similar at highest  $M_x$
- DIS and DDIS: approximate KNO scaling & same shape.

# ZEUS: $e^+e^-$ & $ep$ : Breit Frame

Breit Frame Breit Frame Lab Frame

## DIS event



- Use Breit frame to compare multiplicity in  $ep$  to (one hemisphere) of  $e^+e^-$

- Breit Frame definition:

$$2xP + q = 0$$

- "Brick Wall frame": incoming quark scatters off photon and returns along same axis.

- Current region (CR) of Breit Frame is analogous to  $e^+e^-$  in  $0^{\text{th}}$  order pQCD = Quark-Parton Model and energy =  $Q/2$

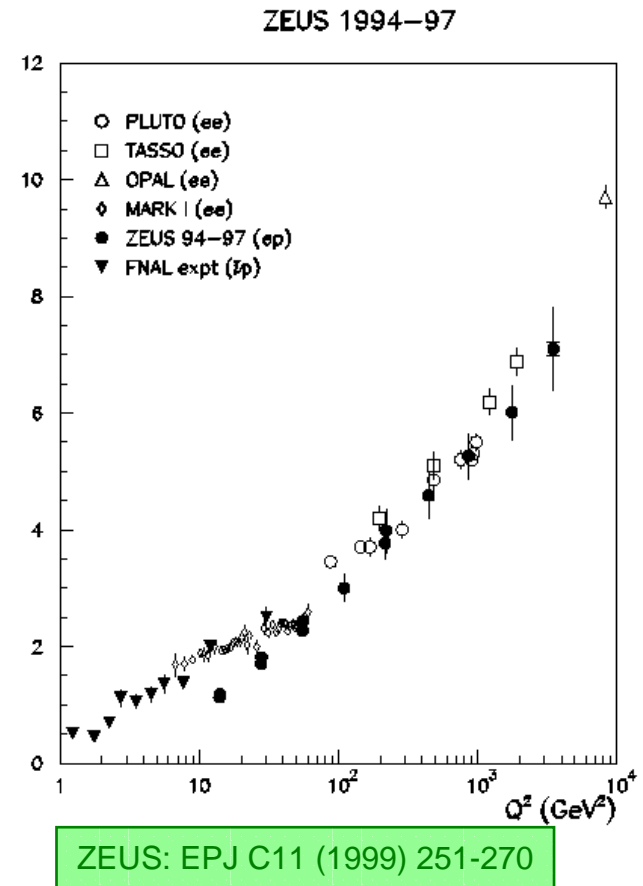
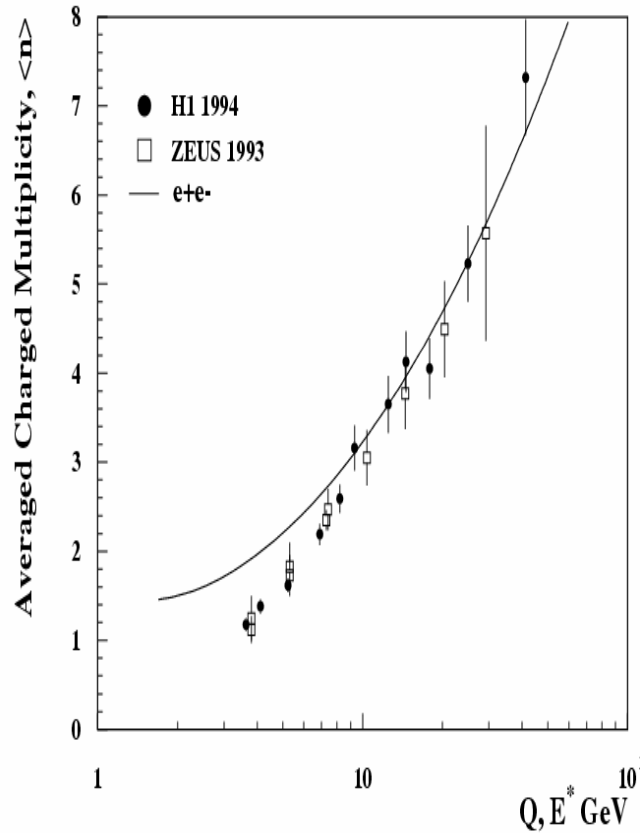
- But: QCD Compton and Boson-Gluon Fusion processes  $\rightarrow$  Particle migration out of current region

K.H.Streng et al. ZPC 2 (1979) 237; S. Chekanov J.Phys. G (1999) 59, hep-ph/9806511; 9810477

- Energy in CR  $< Q/2$

- **ZEUS: use measured energy in CR of Breit Frame as energy scale**

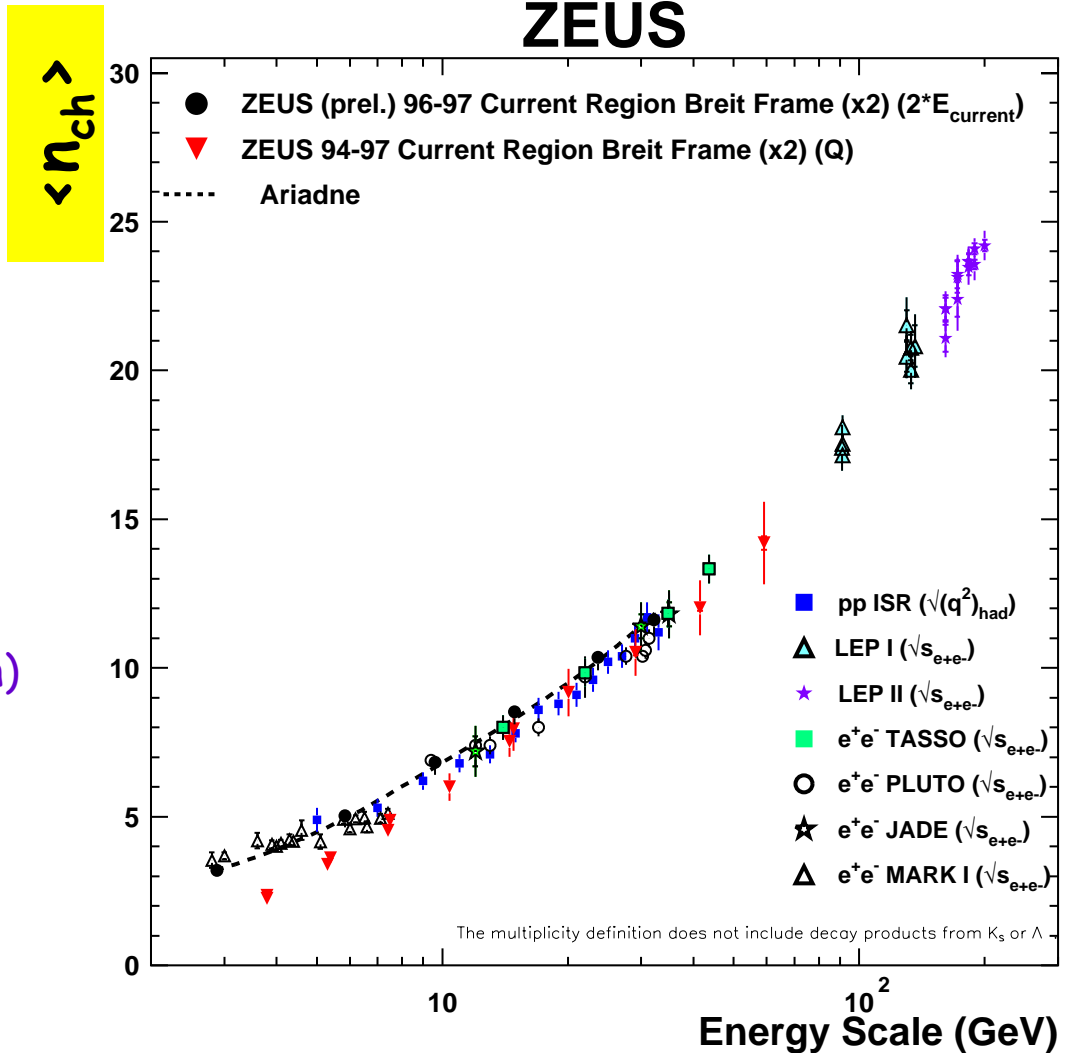
# Early results: $\langle n_{ch} \rangle$ vs. $Q$ in Breit Frame



- Current region Breit frame multiplicity vs.  $Q \leftrightarrow e^+e^-$  data (divided by 2)
- Consistent with  $e^+e^-$  data for high  $Q^2$
- Lower than  $e^+e^-$  at low  $Q^2$

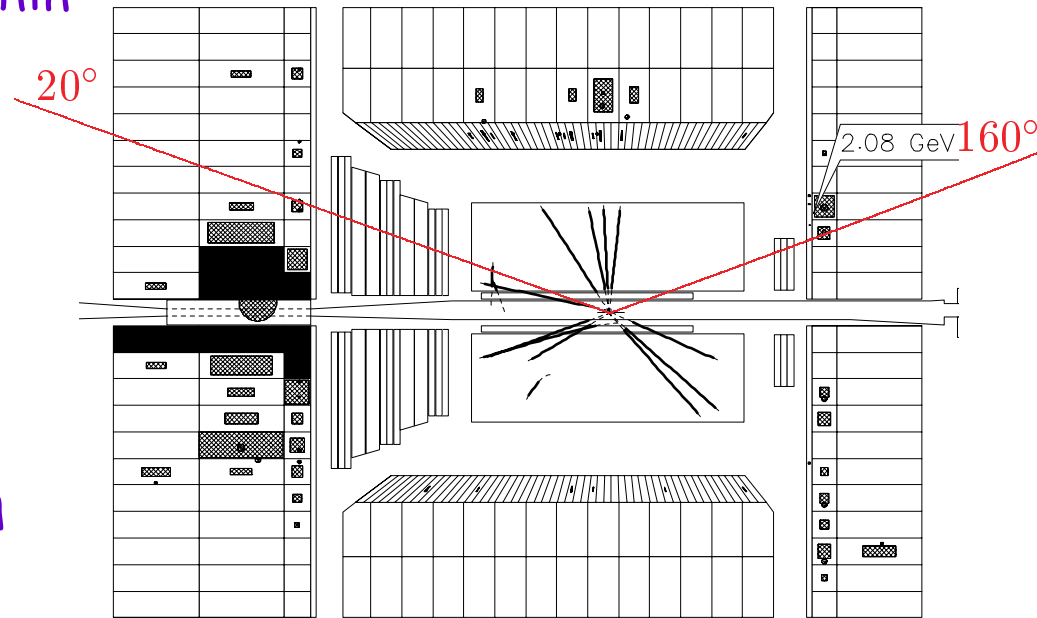
# ZEUS DIS: $\langle n_{ch} \rangle$ vs. $2 \cdot E_{current}$

- NC DIS:  $Q^2 > 25 \text{ GeV}^2$  and  $70 < W < 225 \text{ GeV}$
- Measurement of  $\langle n_{ch} \rangle$  dependence on  $2 \cdot E_{current}$  compared to previous ZEUS measurement vs.  $Q$  (in red), and to  $e^+e^-$  and pp data ( $\langle n_{ch} \rangle$  multiplied by 2 for comparison)
- $2 \cdot E_{current}$  gives better description of multiplicity at low energy



# ZEUS: Effective Mass experimental method

- Measure hadronic final state within  $\Delta\eta$  for best acceptance in the central tracking detector (CTD)
- Measure # charged tracks, reconstruct number of charged hadrons
- Measure invariant mass of the system ( $M_{\text{eff}}$ ) in corresponding  $\Delta\eta$  region
- Energy is measured in the Calorimeter (CAL)



$$M_{\text{eff}}^2 = \left( \sum_{i \neq e} E^i \right)^2 - \left( \sum_{i \neq e} p_x^i \right)^2 - \left( \sum_{i \neq e} p_y^i \right)^2 - \left( \sum_{i \neq e} p_z^i \right)^2$$

Study:  $\langle n_{\text{ch}} \rangle$  vs.  $M_{\text{eff}}$

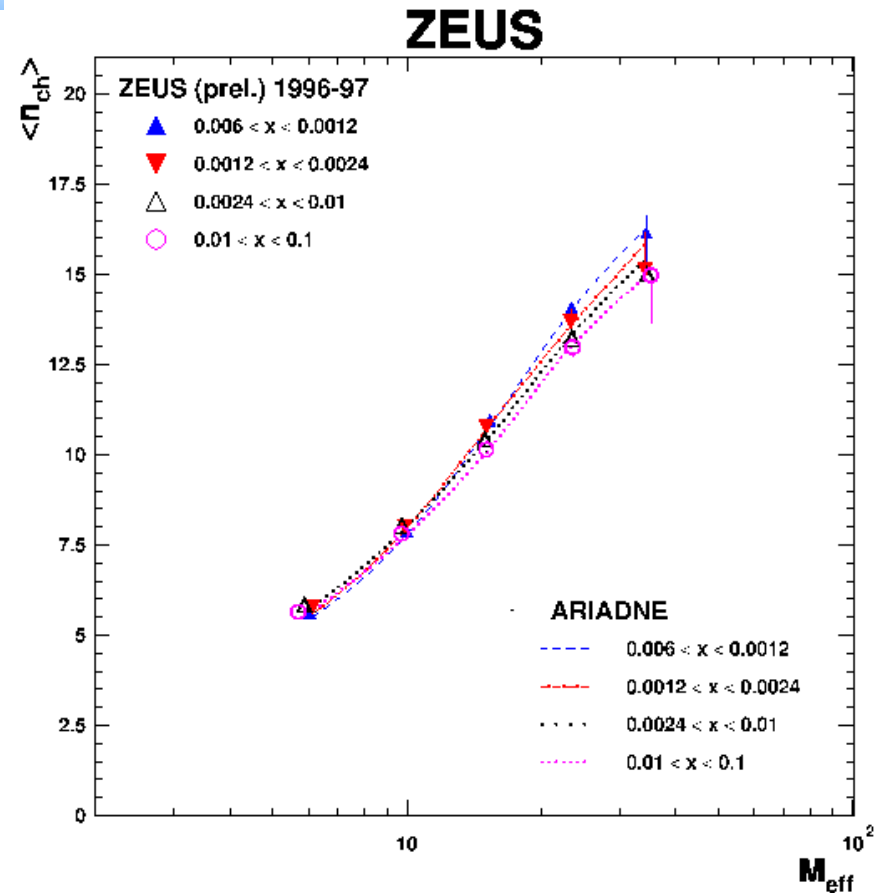
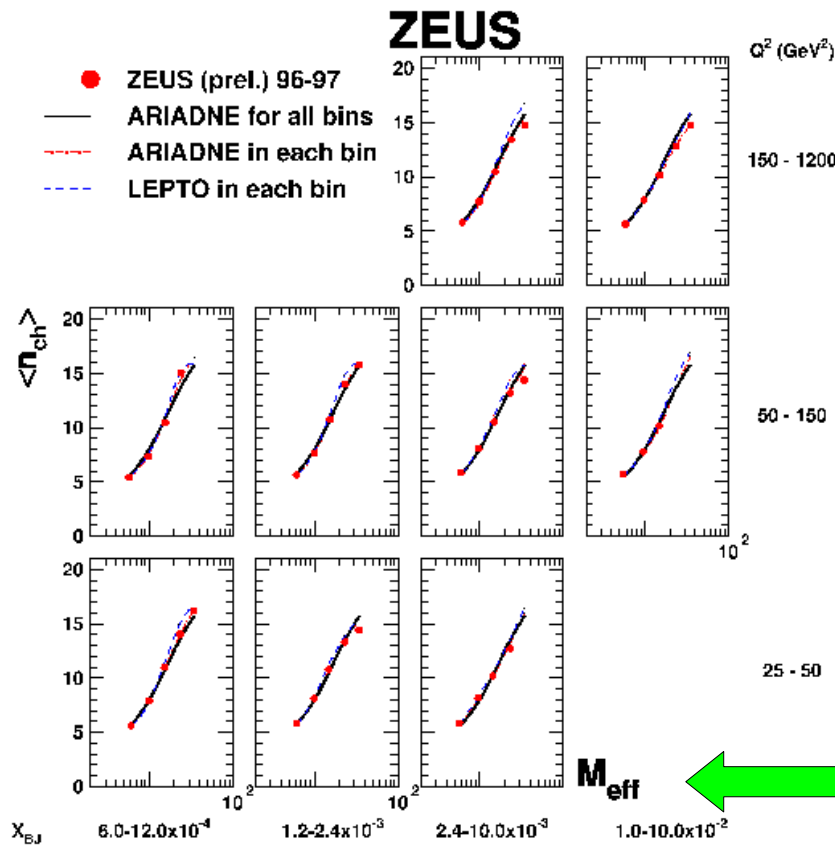
↑  
CTD

↑

CAL within the CTD acceptance



# ZEUS: Lab frame: $\langle n_{ch} \rangle$ vs. $M_{eff}$ in $x$ and $Q^2$ bins



- Data described by ARIADNE
- LEPTO slightly above data
- No  $Q^2$  dependence observed

- Lab frame multipl. vs.  $M_{eff}$ , in  $x_{Bj}$ -bins with Ariadne predictions.
- weak  $x_{Bj}$ -dependence in both data and Monte Carlo.

# ZEUS : Summary

- Measurement in current region of the Breit frame shows similar dependence to  $e^+e^-$  if  $2 \cdot E_{\text{current}}$  is used as the scale (black dots)
- The same dependence is observed for the photon region of the  $\gamma^*p$  frame vs.  $W$ : (blue dots) (not discussed here)
- $\langle n \rangle$  in lab. frame vs.  $M_{\text{eff}}$  shows no  $Q^2$  dependence and weak  $x$ -dependence

