

JETS IN DEEP-INELASTIC SCATTERING AT HERA AND DETERMINATIONS OF α_s^*

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Several methods to extract the strong coupling constant α_s by means of highly energetic jets in Deep-Inelastic Scattering are presented. The results from the various methods agree with one another and with the world average. The errors are competitive to those achieved in α_s determinations in other processes such as proton-anti-proton scattering.

1. Introduction

Measurements of the hadronic final state in deeply inelastic ep scattering (DIS) provide precision tests of quantum chromodynamics (QCD). At HERA data are collected over a wide range of the negative four-momentum-transfer Q^2 , and the transverse energy E_T of hadronic final state jets. As sketched in Eq. 1 the jet cross section can be expressed as a power series of α_s combined with a convolution of the hard matrix element, $\hat{\sigma}_{jet}$, and appropriate parton distribution functions of the proton.

$$\sigma_{jet} = \sum \alpha_s^n(\mu_r) \sum \hat{\sigma}_{jet}(\mu_r, \mu_f) \otimes \text{pdf}(\mu_f, \dots), \quad (1)$$

with μ_r and μ_f being mass scales.

Fig. 1 shows diagrams of the leading order, here $O(\alpha_s)$, processes for dijet-production in DIS.

The accurate measurement of jet production, hence, allows for a precise determination of α_s .

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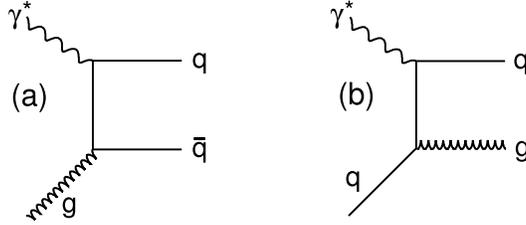


Figure 1. Leading order diagrams for dijet production in ep scattering. (a) photon-gluon fusion and (b) QCD-Compton process.

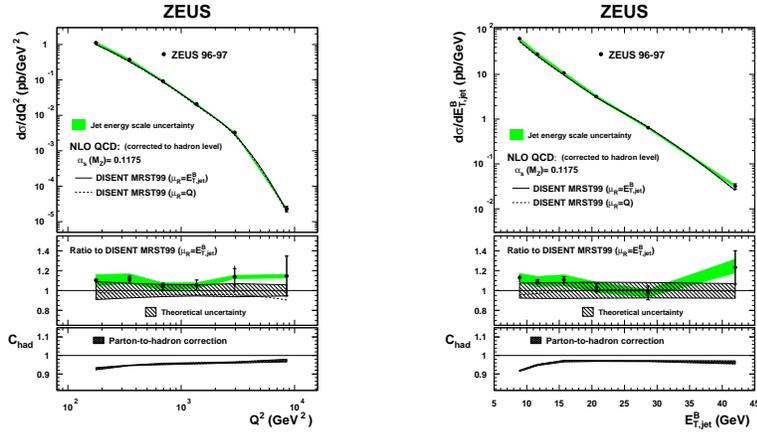


Figure 2. Single inclusive jet cross section as a function of Q^2 (left) and $E_{T,jet}^B$ (right) compared with NLO QCD predictions.

2. Single Inclusive Jet Cross Section

For this analysis it is required to identify at least one jet above a given transverse energy. Fig. 2 shows the measured single inclusive jet cross section as a function of Q^2 and the transverse energy of the jet as measured in the Breit-Frame compared with NLO-QCD calculations ¹.

The data have a typical experimental uncertainty of 7 % and are well reproduced by the theoretical predictions at large values of E_T and Q^2 . The NLO-QCD calculations, currently only available at the parton level, are corrected for hadronisation effects using LO models. These corrections are expected to be small at large E_T and Q^2 .

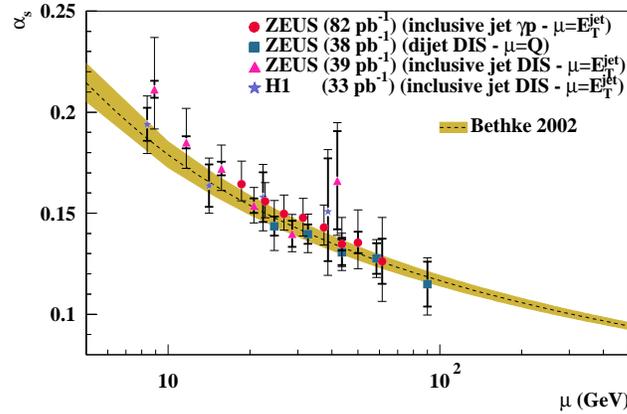


Figure 3. Various results on α_s measurements in DIS as a function of a mass scale μ compared with results from global fits.

2.1. Determination of α_s

The dependency of a generic jet cross section on $\alpha_s(M_Z)$ is parameterized in the corresponding analysis bins with the help of suitable NLO-QCD predictions featuring slightly different values of $\alpha_s(M_Z)$. By comparison of the parameterized cross section with the measured cross section the values for α_s are obtained. The resulting $\alpha_s(M_Z)$ from *e.g.* the single inclusive jet cross section for $Q^2 > 500$ GeV is found in the ZEUS analysis to be

$$\alpha_s(M_Z) = 0.1212 \pm 0.0017(\text{stat.})_{-0.0031}^{+0.0023}(\text{exp.})_{-0.0027}^{+0.0028}(\text{theor.}).$$

The experimental error (exp.) is dominated by the uncertainty on the energy scale for the jet measurement. The largest contribution to the theoretical uncertainty (theor.) is given by a residual dependency on the renormalization scale μ_r which corresponds to uncertainties due to the contributions of terms beyond next-to-leading order.

In Figure 3 the result is displayed as a function of a mass scale μ together with other values of α_s extracted from DIS-jet data ².

The expected running of α_s as a function of μ is clearly visible. In addition the figure demonstrates the compatibility of the resulting α_s values with those obtained in global fits ³.

3. Jet Substructure and Determination of α_s

Jets appear as a collimated spray of particles which are combined to by dedicated algorithms. These particles are the end point of a cascade of successive particle emissions from the hard interaction to the hadronic final state.

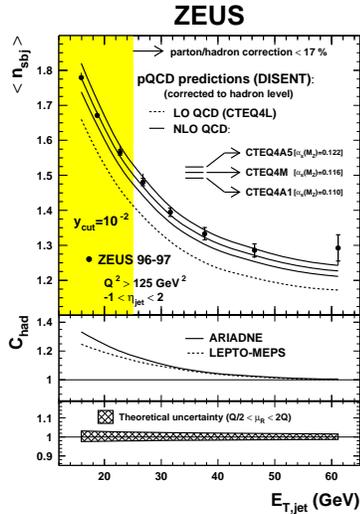


Figure 4. Subjet multiplicites for different transverse energies of jets. Jets were identified in the laboratory frame.

The development of the cascade is governed by the strong coupling constant α_s . An attempt is made to resolve subjets using dedicated algorithms within the identified jets⁴. The amount of subjets is expected to depend on α_s . Fig. 4 shows the number of subjets as a function of the jet- E_T measured in the laboratory frame. The number of subjets decreases as the transverse energy of the jet increases. The transverse energy of the jet sets the scale for α_s . Thus, the probability to radiate partons is small at large transverse energies. The data are well described by NLO-QCD calculations employing different parton pdfs featuring slightly different values of $\alpha_s(M_Z)$. For E_T larger than 30 GeV the hadronization corrections become small allowing for a QCD analysis to determine α_s from subjet multiplicites. The resulting value

is found by ZEUS to be

$$\alpha_s(M_Z) = 0.1187 \pm 0.0017(\text{stat.})_{-0.0009}^{+0.0024}(\text{exp.})_{-0.0076}^{+0.0093}(\text{theor.}).$$

4. 3-jet Cross Sections

3-jet cross sections are well suited for an extraction of α_s because the lowest order contribution to this event class is proportional to α_s^2 . The sensitivity to uncertainties due to proton pdf's can be reduced by building the ratio $R_{3/2}$, *i.e.* the ratio of 3-jet to dijet cross sections. A measurement of this observable is shown in Fig. 5. While a minor sensitivity to variations of the pdf's is observed the ratio is very sensitive to small variations of α_s which underlines the potential of this observable for future determinations of α_s .

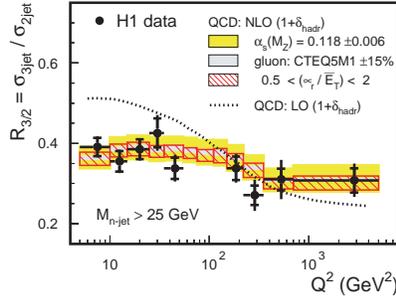


Figure 5. The observable $R_{3/2}$ as a function of Q^2 compared with NLO-QCD predictions.

5. Conclusion and Outlook

The analysis of jet events in DIS allows for precise measurements of the strong coupling constant α_s . A compilation of results is given in Fig. 6. They have a major impact on the current world average value of α_s . Ongoing analysis of HERA I ⁶ data as well as new data expected from HERA II open the possibility for α_s determinations including 3-jet cross sections.

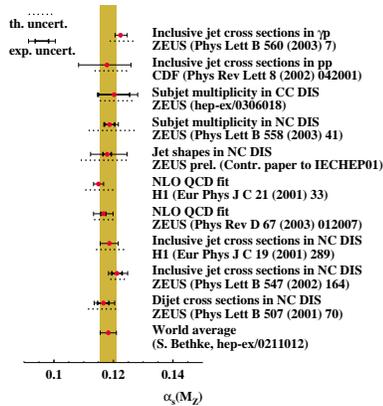


Figure 6. $\alpha_s(M_Z)$ values obtained in DIS together with results from $p\bar{p}$ -collisions and the world average.

References

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