

Tests of the Standard Model in Inclusive Scattering at Large Transverse Momenta

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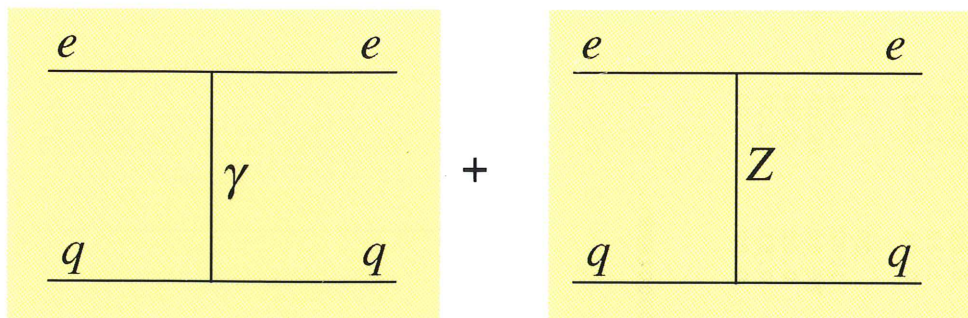
for the H1 and ZEUS Collaborations

HERA results on
Deep Inelastic Scattering

Probing the electroweak model at
 10^{-3} fm in space-like interactions

Boson Exchange at HERA

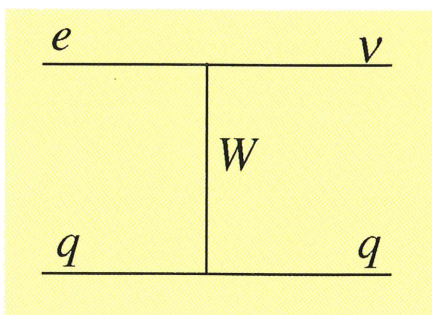
Neutral Current (NC)



$$\frac{d^2 \sigma}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} \phi_{\gamma,Z} = \frac{2\pi\alpha^2}{xQ^4} \left[Y_+ F_2(x, Q^2) \mp Y_- xF_3(x, Q^2) \right]$$

with $Y_{\pm} = 1 \pm (1-y)^2$ for $e^{+/-}$

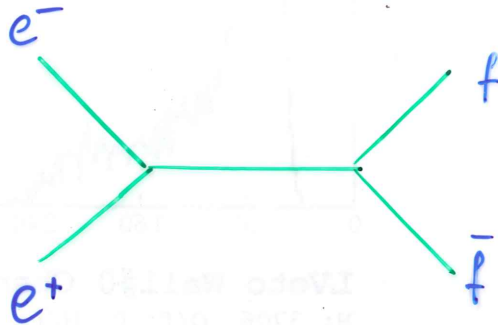
Charged Current (CC)



$$\frac{d^2 \sigma}{dx dQ^2} = \frac{G_F^2}{2\pi x} \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2 \phi_W$$

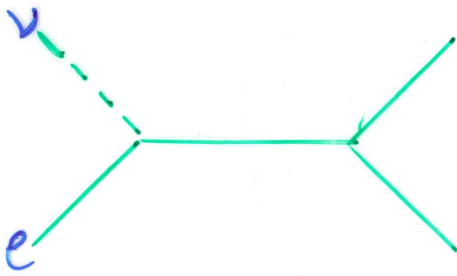
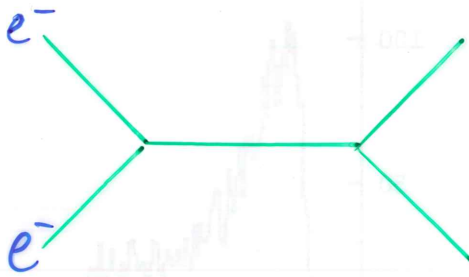
$\phi_{\gamma,Z}, \phi_W$ comprise: structure of the proton
 propagators and
 electroweak couplings
 \approx quark densities

LEP/SLC ANALOGY



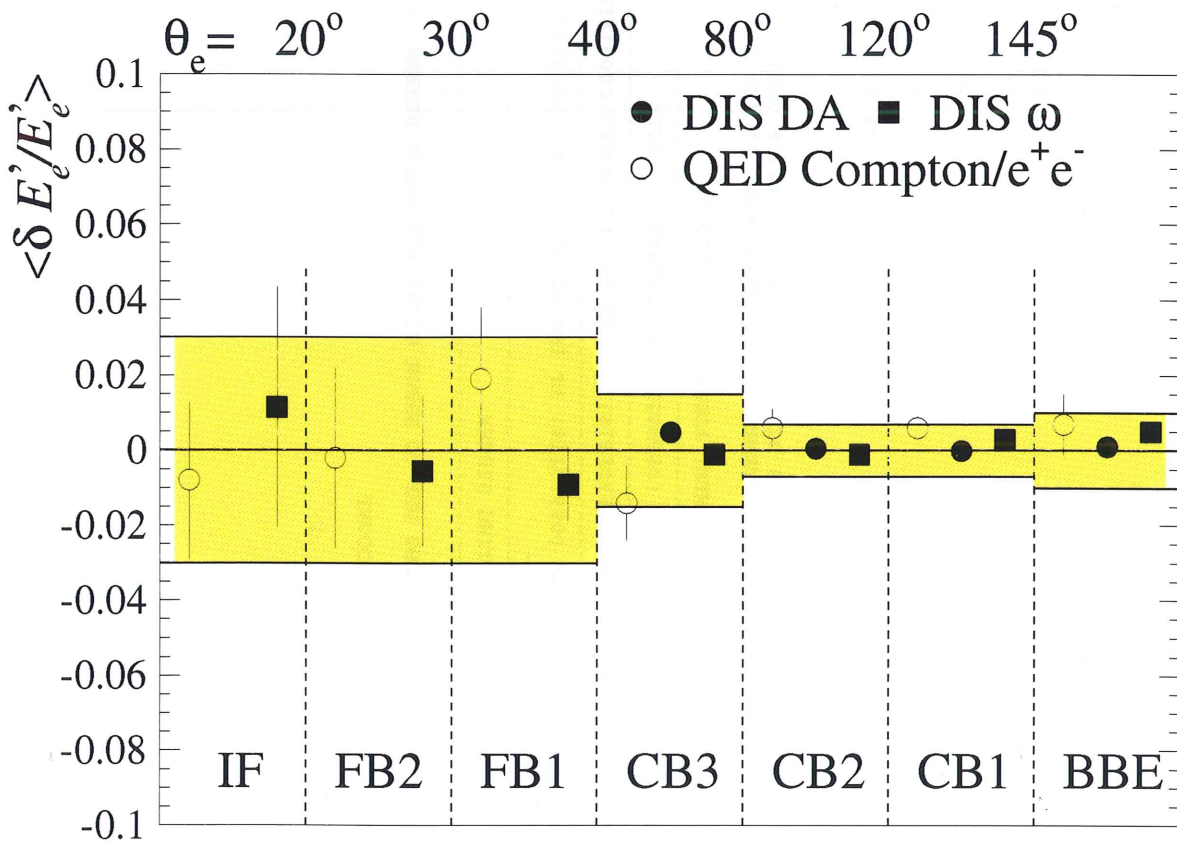
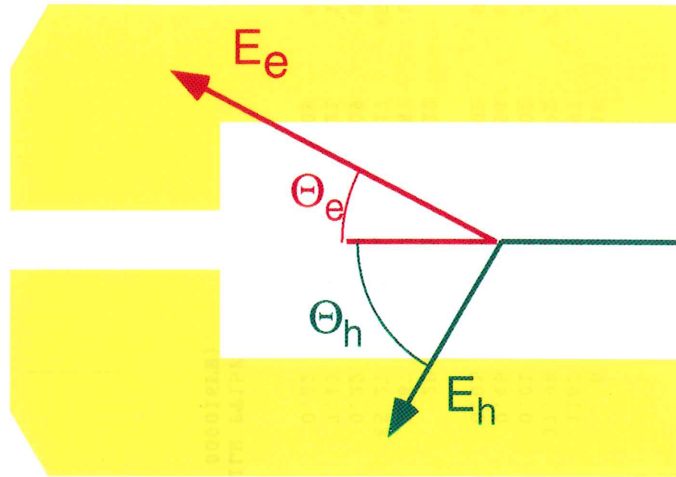
$$\sigma \sim \frac{1}{s^2} R$$

$$R \sim \sum_f \text{coupl} \times \text{propag.} \times (1 + \frac{\alpha_s}{\pi} \dots)$$



Experimental Accuracy

Experimental redundancy in the measurement of the two inclusive scattering parameters provides means for cross calibration.



LAr wheel

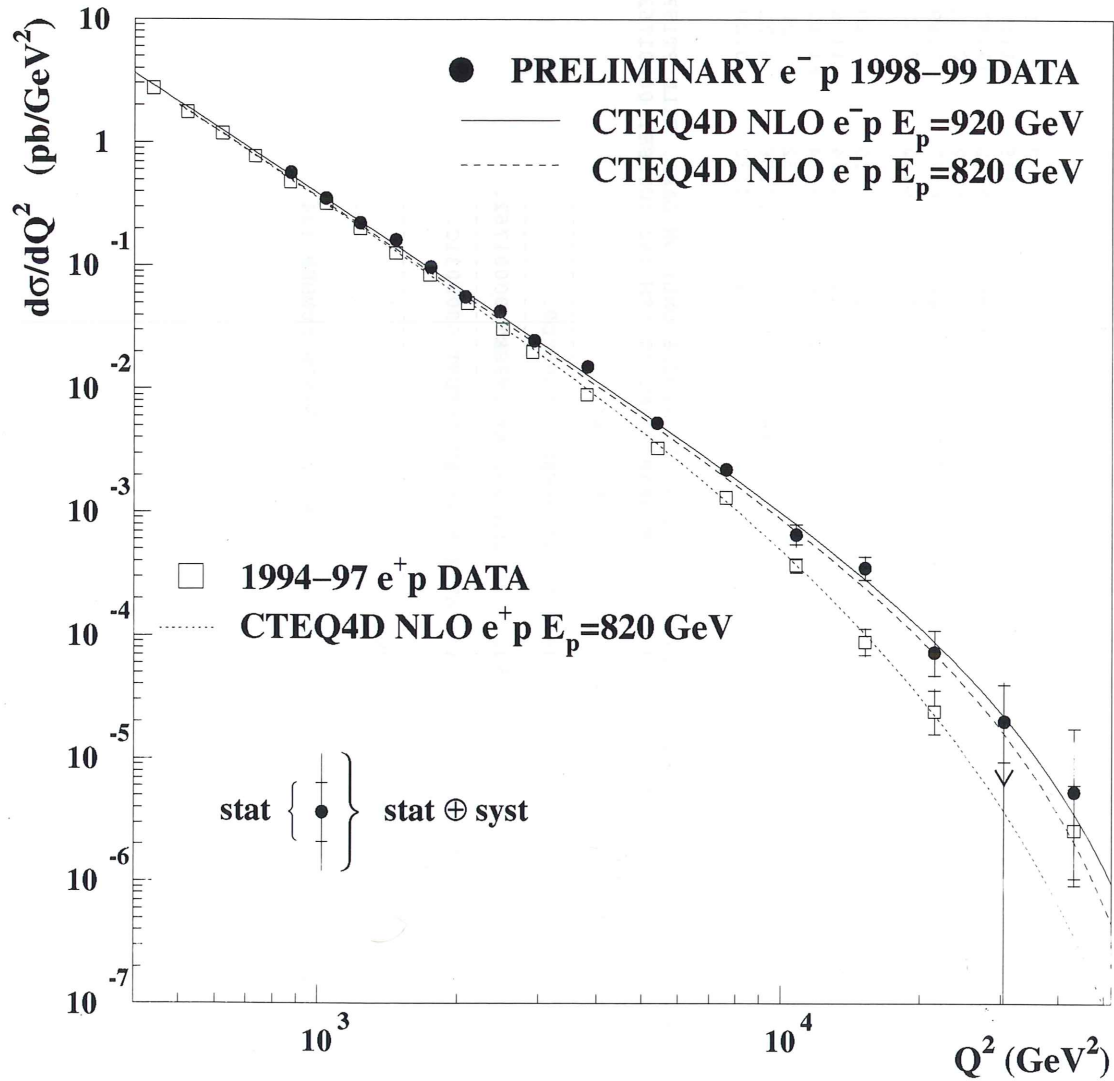
Neutral Currents

$$\frac{d^2}{dx dQ^2} \sigma(e^\pm p) = \frac{2\pi\alpha^2}{xQ^4} \phi_{\gamma,Z}$$

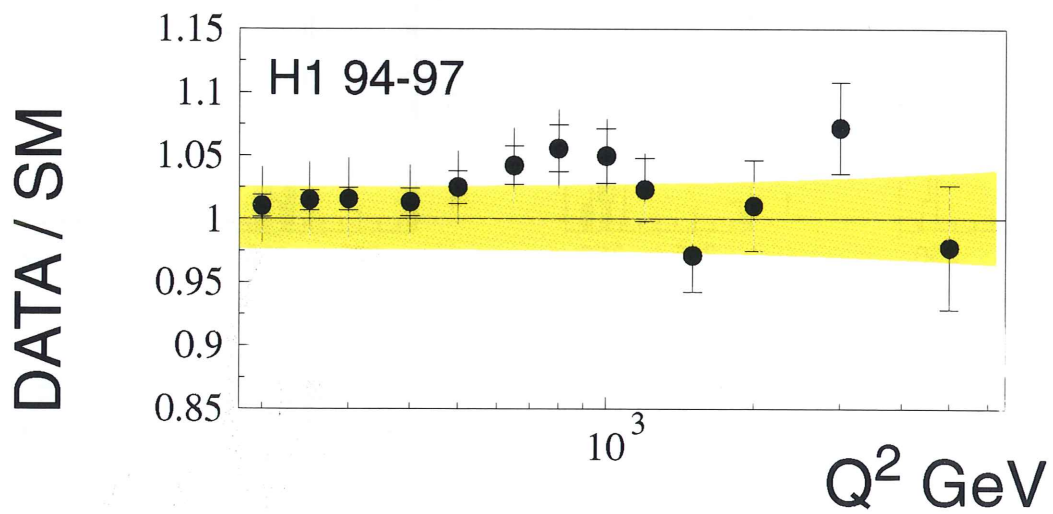
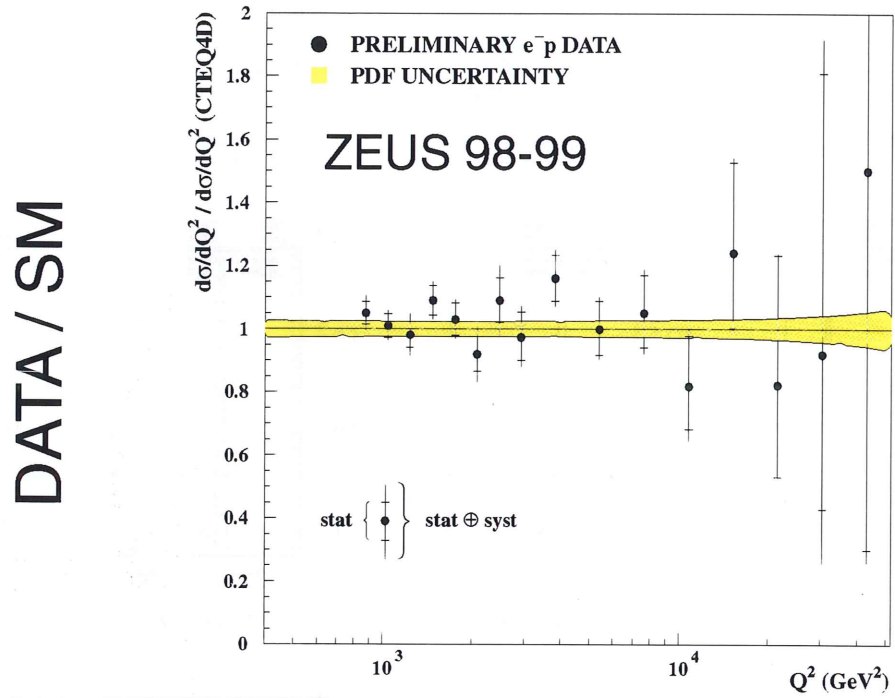
Observe the proper charge dependence:

- γZ interference

ZEUS NC



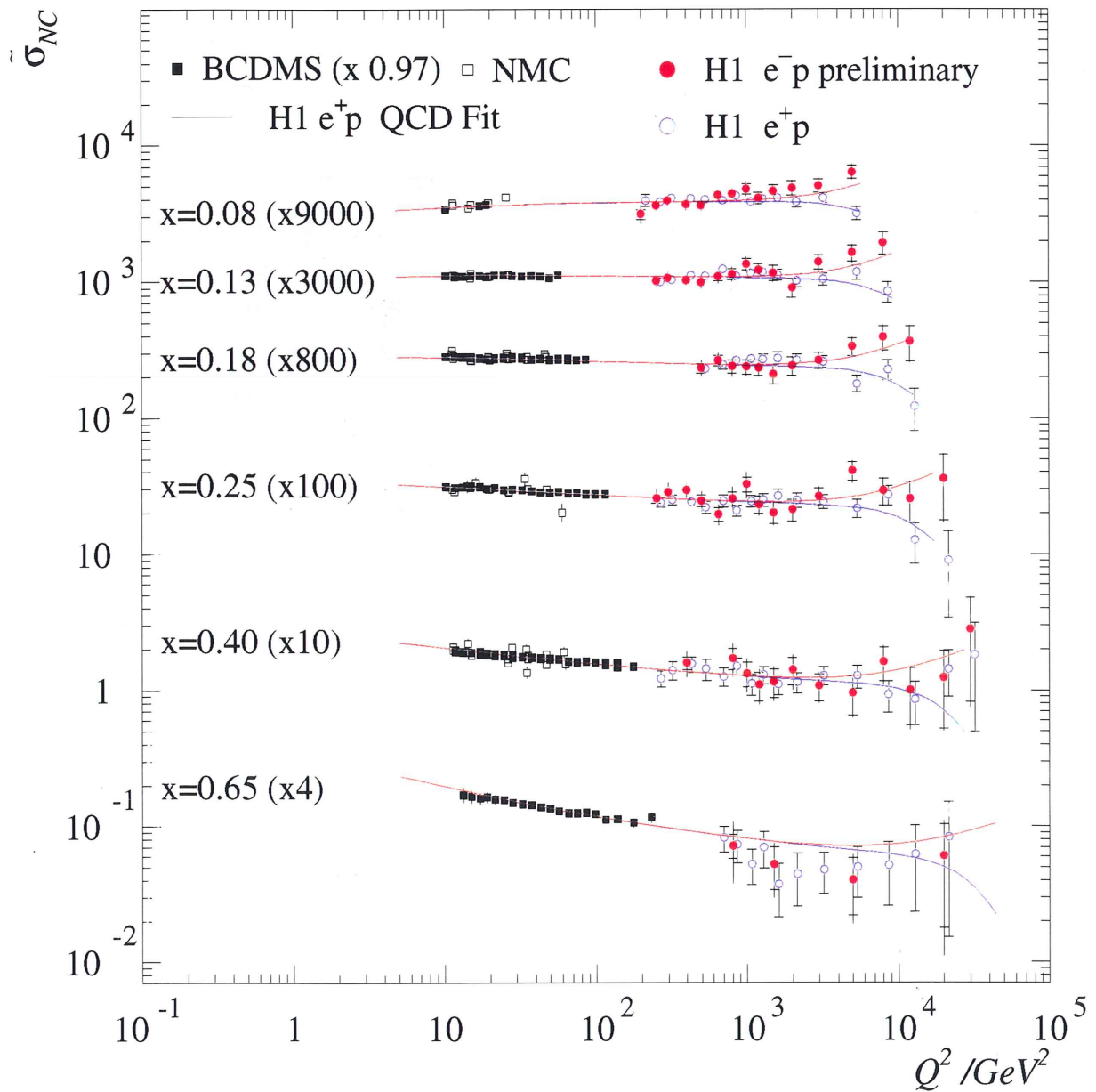
Neutral Currents cont'd



NC at high x

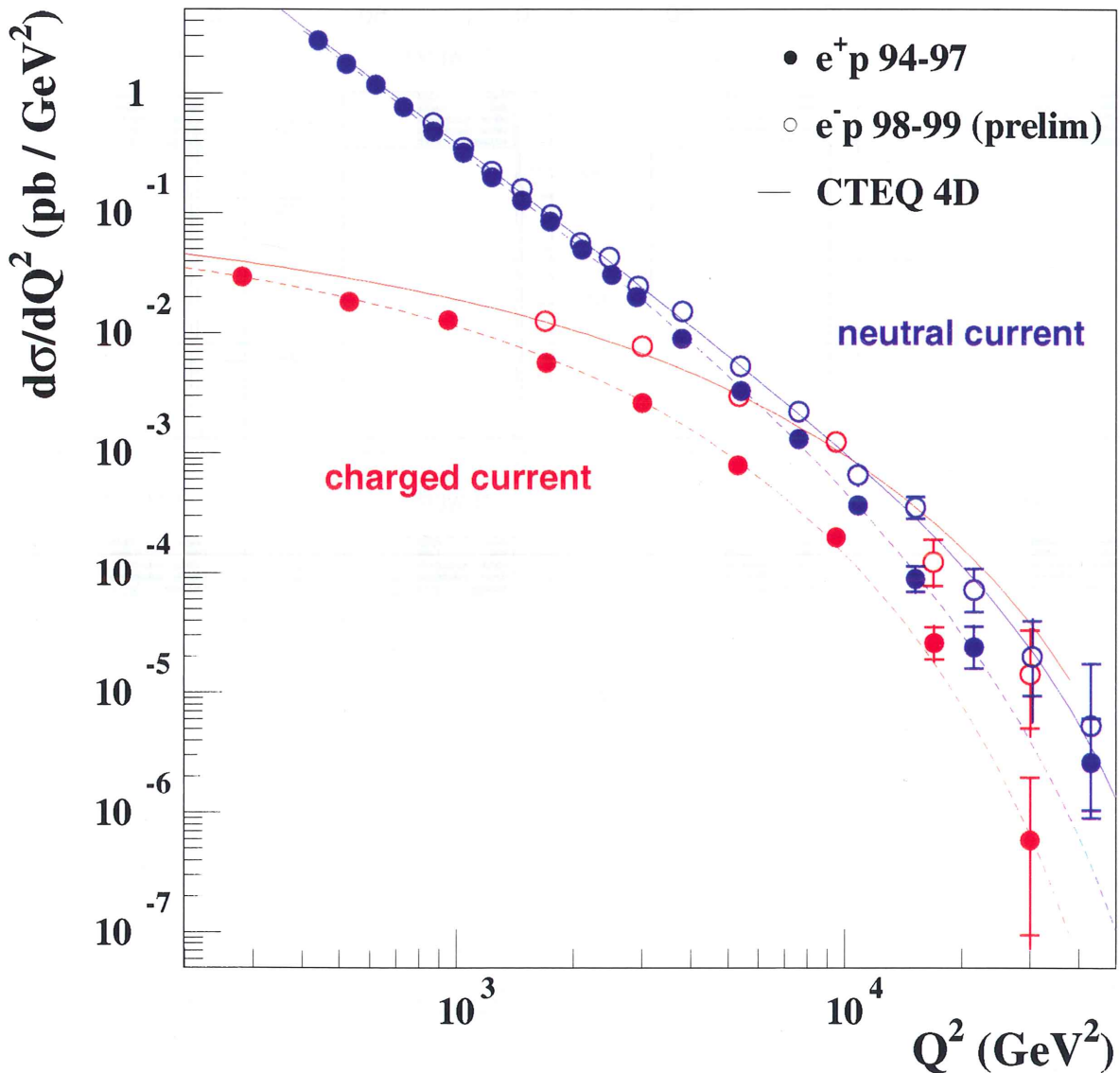
$$\tilde{\sigma} = \left[\frac{2\pi\alpha^2}{xQ^4} Y_+ \right]^{-1} \frac{d^2\sigma}{dx dQ^2}$$

- QCD evolution over 2 orders of magnitude
- high x requires clarification



Electroweak Unification

ZEUS DIS Cross Sections



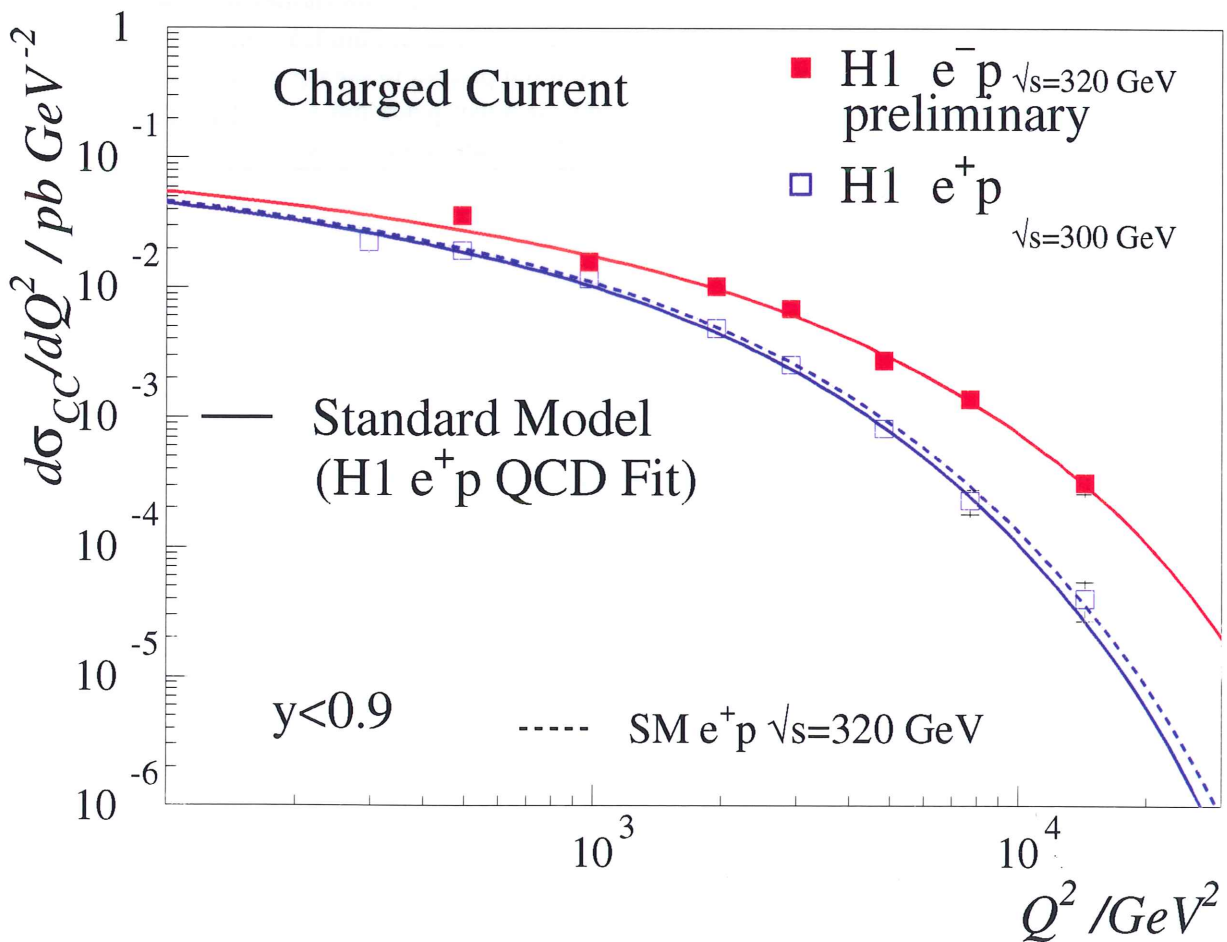
- for large Q^2 neutral and charged current processes have cross sections of similar size
- details depend on quark content and coupling

Charged Current

$$\frac{d\sigma^{CC}}{dQ^2 dx} \sim G_F^2 \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2 \frac{1}{x} \phi_W$$

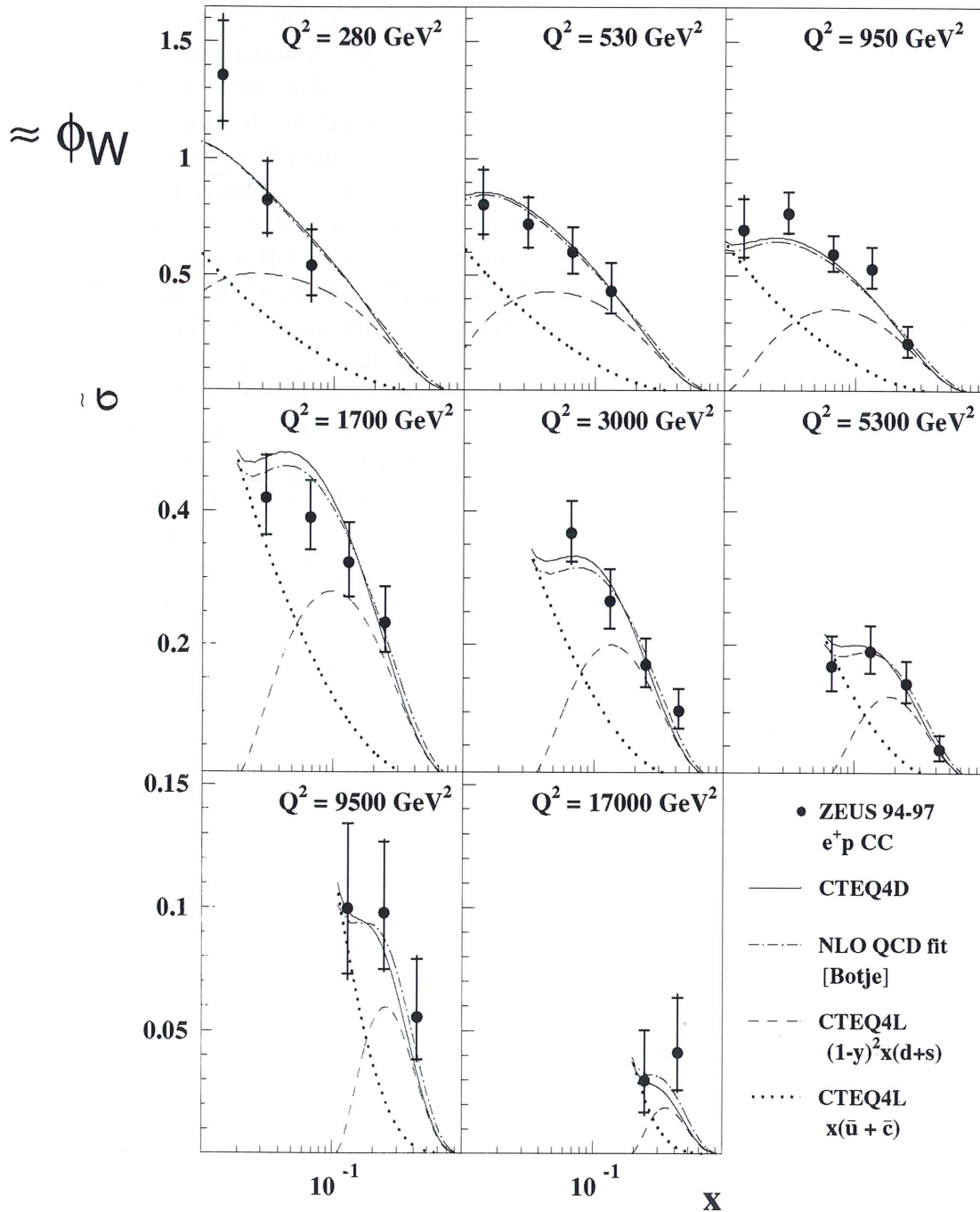
At high Q^2 coupling to valence quarks prevails:

$$\begin{aligned} &\Leftarrow \Rightarrow && \Rightarrow \Rightarrow \\ e^-: \phi_W &\sim xu && e^+: \phi_W \sim (1-y)^2 xd \end{aligned}$$



Charged Current x-Distrib.

Precision of the e^+p data disentangles individual quark distributions in the scattering process.

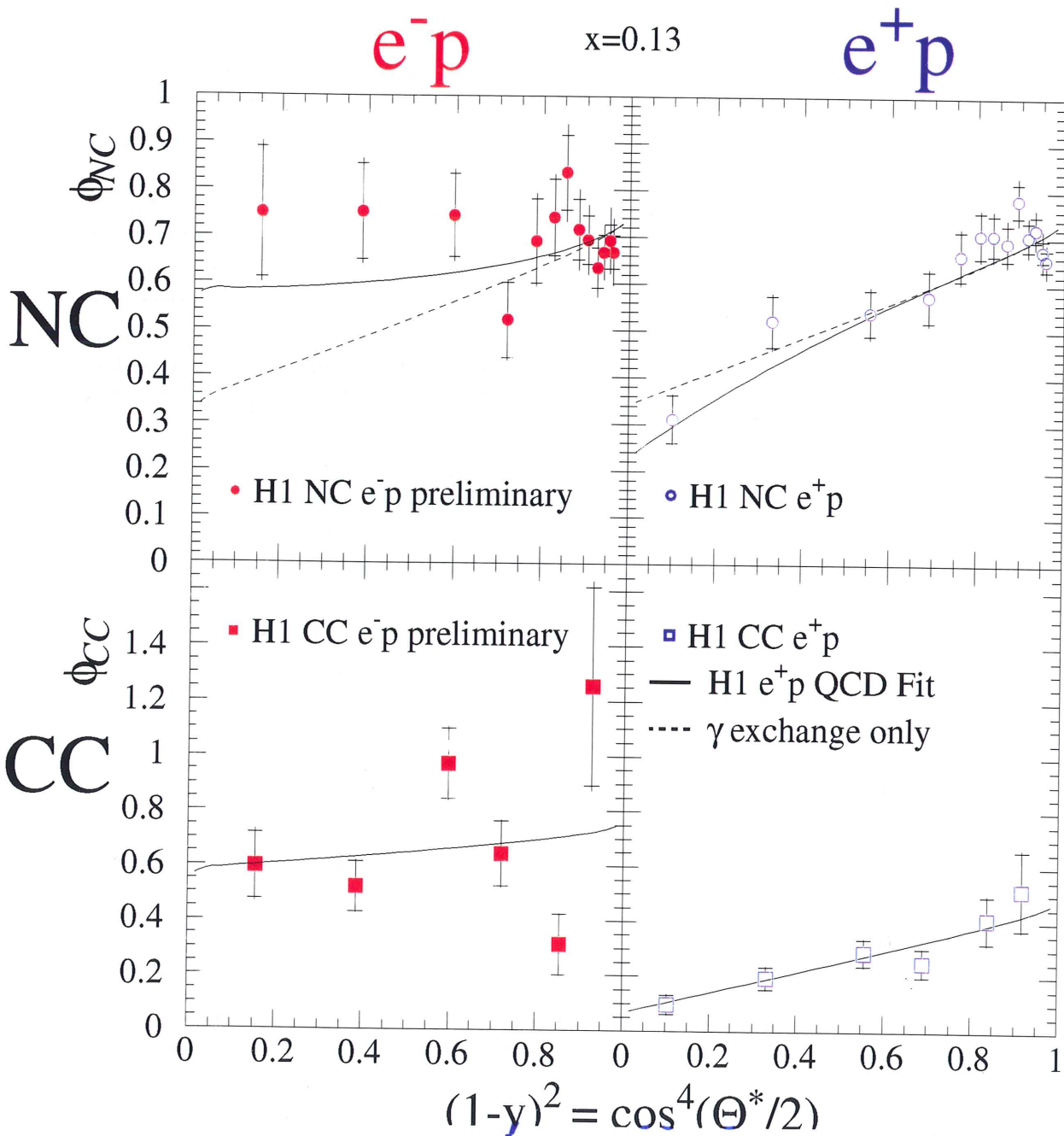


Helicity Structure at Short Distance

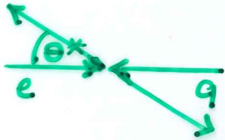
NC: $\phi_{\gamma,Z} \sim (1 + (1-y)^2) \left(\frac{4}{9}(xu + x\bar{u}) + \frac{1}{9}(xd + x\bar{d}) \right)$

CC: $\phi_W \sim xu + (1-y)^2 x\bar{d}$

$\phi_W \sim x\bar{u} + (1-y)^2 xd$



Helicity Structure at Short Distance



NC: $\phi_{\gamma,Z} \sim (1 + (1-y)^2) \left(\frac{4}{9} (xu + x\bar{u}) + \frac{1}{9} (xd + x\bar{d}) \right)$

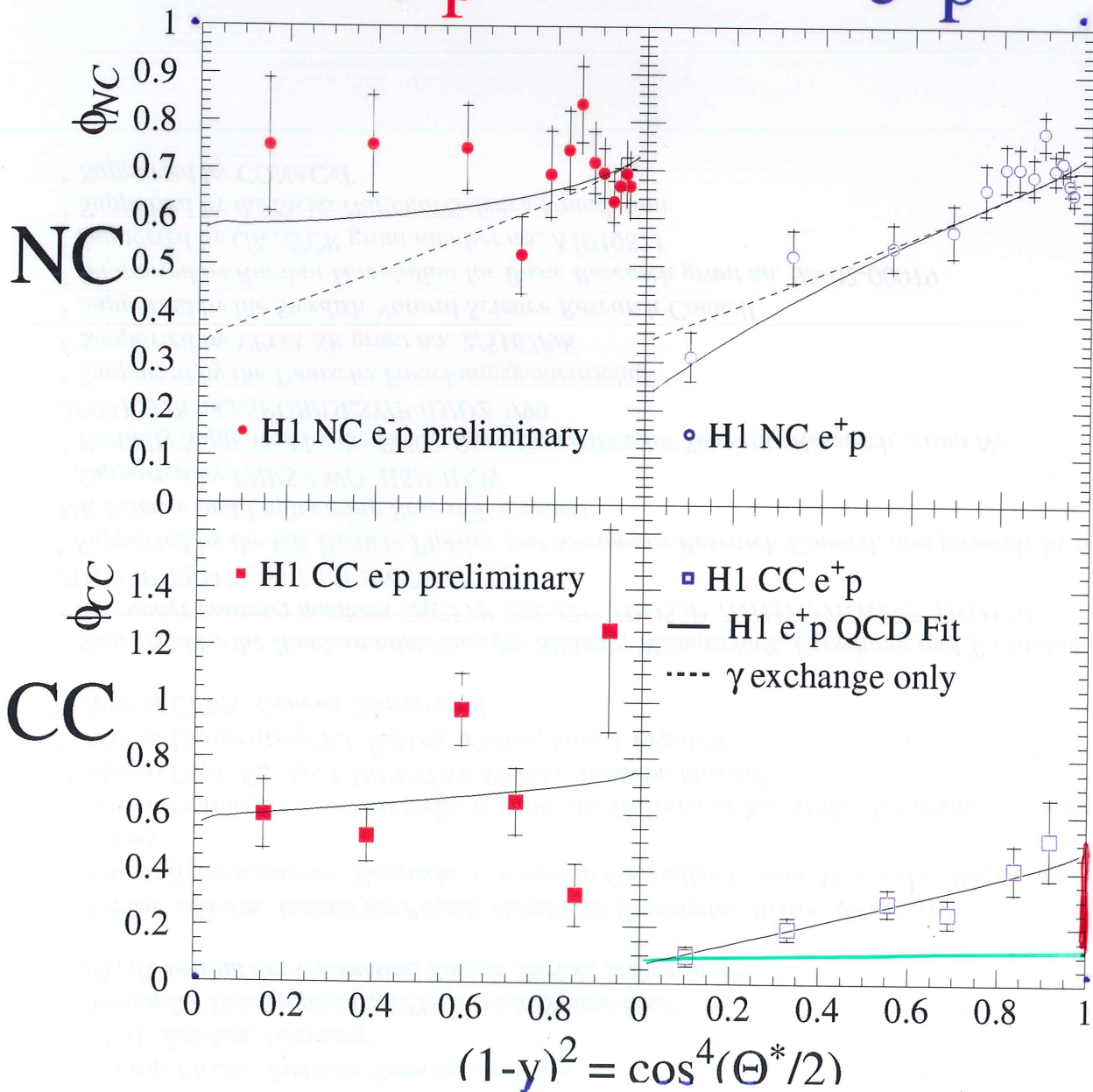
CC: $\phi_W \sim xu + (1-y)^2 x\bar{d}$

$\phi_W \sim x\bar{u} + (1-y)^2 xd$

e^-p

$x=0.13$

e^+p



Electroweak Couplings

The “original” electroweak parameters

$$\alpha, G_F, \sin^2 \Theta_W$$

with

$$\Gamma_\mu = G^2 \frac{m_\mu^5}{192\pi^3} \left(1 - \frac{8m_e^2}{m_\mu^2} \right) \left(1 + \frac{\alpha}{2\pi} \left(\frac{25}{4} - \pi^2 \right) \right)$$

have since the advent of LEP/SLC been replaced by the more accurately known parameters

$$\alpha, G_F, M_Z$$

Since then tests of the electroweak model have been performed through its radiative corrections, which may be quoted as a function of other external quantities such as m_{Higgs} and m_{top} , etc.

Electroweak Tests

- LEP/SLC: $\sigma \propto \text{couplings} \times \left(\frac{M_Z^2}{s - M_Z^2} \right)^2$, time-like
- HERA: $\sigma \propto \text{couplings} \times \left(\frac{M_W^2}{Q^2 + M_W^2} \right)^2 \phi_W$, space-like

Both facilities reach rather similar virtualities.

At HERA Q^2 varies from event to event such that the validity of the model can be simultaneously tested over a large range in Q^2 comparable to the tests in the time-like region at PEP/PETRA, TRISTAN and LEP/SLC.

The purely weak Charged Current process at HERA therefore explores

- the precisely known part at small Q^2 and
- the propagator dependence

$$\frac{d\sigma}{dQ^2} = \sigma^0 \left(\frac{M_W^2}{Q^2 + M_W^2} \right)^2 \phi_W$$

assuming precisely known parton distributions and their Q^2 evolution.

Precision of Space-like Tests

A $\approx 3\%$ cross section uncertainty could be cast into an uncertainty of a single electroweak parameter, such as M_W . With

$$\frac{d\sigma}{dQ^2} \propto G^2 \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2 \phi_W \quad G = \frac{\pi\alpha}{\sqrt{2}M_W^2} \frac{1}{\left(1 - M_W^2/M_Z^2\right)^{1-\Delta r}}$$

the apparent sensitivity is

$$\frac{\Delta\sigma}{\sigma} \approx 10 \times \frac{\Delta M_W}{M_W}$$

$$3\% \Rightarrow 0.3\%$$

.ZEUS get:

$$M_W = 80.50_{-0.25}^{+0.24} \text{ (stat.) }_{-0.16}^{+0.13} \text{ (syst.) } \pm 0.31 \text{ (pdf)}_{-0.16}^{+0.03} (\Delta M_{\text{top}}, \Delta M_H, \Delta M_Z)$$

as expected.

Note $\frac{\Delta G}{G} = 20 \times 10^{-9}$ from τ_μ .

Conclusion

With luminosities of $\approx 50 \text{ pb}^{-1}$ HERA starts to explore the high Q^2 sector of the SM:

- parton distributions extrapolate well into the high Q^2 regime (QCD)
- start to distinguish u and d in a single experiment (flavour separated parton couplings!)
- helicity structure probed at $\approx 10^{-3} \text{ fm}$
- Z^0 contributions observed in NC scattering at high Q^2
- Charged Current consistent with a single propagator
- HERA measures propagator mass to $\approx 2 \text{ GeV}$, consistent with M_W

but

- high x , high Q^2 region needs more data for clarification