

$\psi(2S)$ to $J/\psi(1S) \rightarrow \mu^+ \mu^-$ exclusive xsection ratio with ZEUS (case study)

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(ZEUS Collaboration)

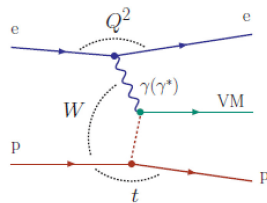
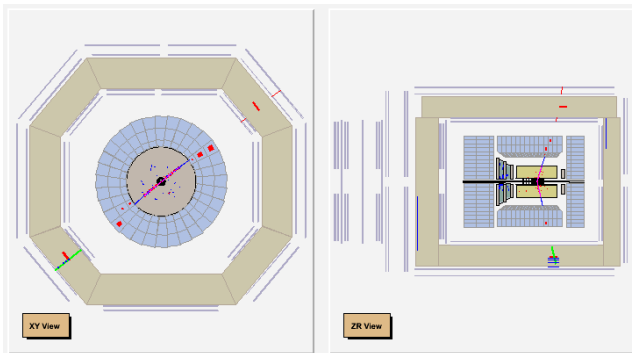
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CFNS Workshop: Hera 4 EIC, (online), 8 – 10 June 2022

- **HERA II DATA:** $\mathcal{L} = 373 \text{ pb}^{-1}$ (2003 - 2007)
- **Investigated decay channels:**
 - $\psi(2S) \rightarrow \mu^+ \mu^-$, $\psi(2S) \rightarrow J/\psi + \pi^+ \pi^-$, $J/\psi(1S) \rightarrow \mu^+ \mu^-$
- **exclusive (elastic) photoproduction sample**
- **kinematic range (analysis phase space):**
 - $30 < W < 180 \text{ GeV}$ (acceptance of muon chambers: $< \sim 300 \text{ GeV}$, HERA c.m.s.)
 - $|t| < 1.0 \text{ GeV}^2$ (elastic to proton dissociative ratio)
 - $Q^2 < 1 \text{ GeV}^2$ (median $Q^2 \approx 3 \times 10^{-5} \text{ GeV}^2$) (e' acceptance)
- differential measurement in 5 W bins and 5 $|t|$ bins
→ specific issues at the edge of the phase space

Example of Final State Topology for $ep \rightarrow J/\psi p$, $J/\psi \rightarrow \mu^+ \mu^-$

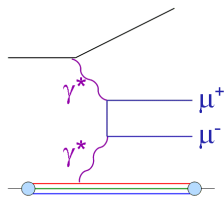


Exclusive process,
reaction mediated by
exchange of **colorless**
object;
proton stays intact.

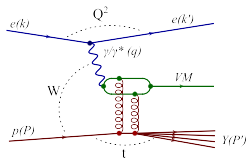
J/ψ and $\psi(2S)$ are detected in the 2- or 4-prong final states ($\mu^+ \mu^-$ or $\mu^+ \mu^- \pi^+ \pi^-$)
very clean final state topology:

Photoproduction ($Q^2 < 1 \text{ GeV}^2$): two or four charged particles and nothing else
 \Rightarrow **experimental challenge: triggering on soft muons**
(Electroproduction ($Q^2 > 1 \text{ GeV}^2$): scattered electron also visible in the detector)

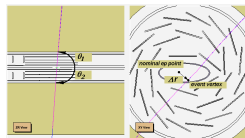
- QED di-muons (like $\gamma^* \gamma^* \rightarrow \mu^+ \mu^-$) from the Bethe-Heitler process



- J/ψ and $\psi(2S)$ mesons production with the dissociation of the proton



- Cosmic muons can mimic $\mu^+ \mu^-$ pairs when passing close to the interaction point



- **Signal MC:** DIFFVM (exclusive/elastic)
- **Background MC:** DIFFVM (proton dissociation)
- **Background MC:** GRAPE
(Bethe-Heitler continuum $\mu^+\mu^-$ pairs)

DIFFVM – A Monte Carlo Generator for Diffractive Processes in ep Scattering.

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- soft diffractive processes in the Regge framework and Vector Dominance Model

- $\frac{d\sigma}{dQ^2} \propto \frac{1}{(1+Q^2/M_Y^2)^{1.5}}$

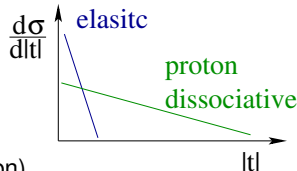
- $\frac{d\sigma}{d|t|} \propto W_{\gamma p}^{4\epsilon} e^{-b|t|}$ ($4\epsilon = \delta$) (elastic)

- $\frac{d^2\sigma}{d|t|dM_Y^2} \propto W_{\gamma p}^{4\epsilon} e^{-b'|t|} M_Y^{-\beta}$ (p.diss)

- $\frac{d\sigma}{dM_Y^2} \sim \frac{f(M_Y^2)}{M_Y^{2(1+\epsilon)}}$ for $M_Y^2 < 3.6 \text{ GeV}^2$ (p resonance region),

- $\frac{d\sigma}{dM_Y^2} \sim \frac{1}{M_Y^{2(1+\epsilon)}}$ for $M_Y^2 \geq 3.6 \text{ GeV}^2$ (continuum region)

- assuming SCHC: s-channel helicity conservation



Signal Monte Carlo

- DIFFVM OK for efficiency/acceptance corrections
- requires tuning of many free parameters
- oversimplified parameterization (enough for moderate $|t|$ in ep scattering)
- more realistic description of diffraction pattern will be welcome :).

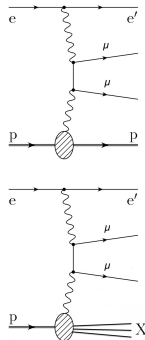
GRAPE-Dilepton (Version 1.1)

A generator for dilepton production in ep collisions

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- based on the **exact matrix elements** in the electroweak theory at tree level via $\gamma\gamma$, γZ^0 , $Z^0 Z^0$ and via photon internal conversion (QED Compton)
- **Feynman amplitudes** are generated by the automatic calculation system **GRACE**
- **proton vertex** covers the whole kinematical region
- interface to PYTHIA and SOPHIA
→ complete hadronic final state
- **covers elastic, quasi-elastic and DIS processes**



Bethe-Heitler continuum di-muon pairs

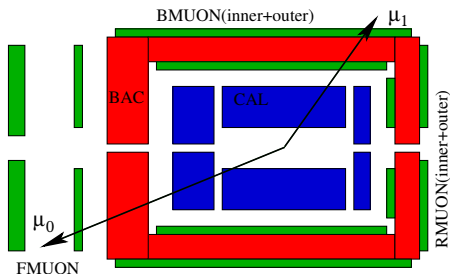
- LO MC (like GRAPE) sufficient for ep scattering
- not sure for eA collisions
(stronger electric field around heavy nuclei)
- NLO, multi-photon effect important (?)

- **Muon efficiency corrections**

Single muon efficiency corrections: TAG and PROBE

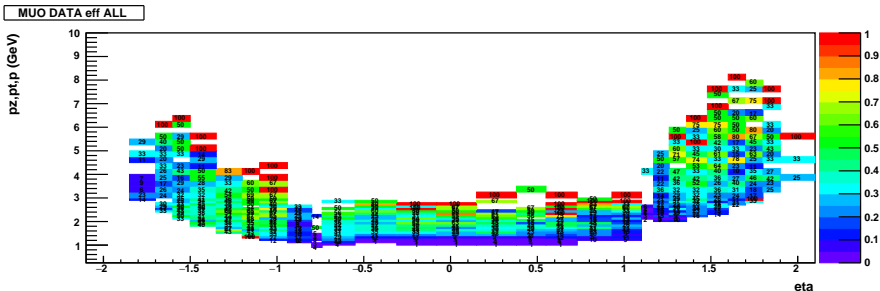
- TAG: “the triggering” muon
- PROBE: “the tested” muon
- effic in given $(p_{t_{eff}}^j, \eta^j)$ bin:

$$\epsilon = N_{PROBE}^j / N_{TAG}^j$$



- **one step correction** for (FLT and SLT and TLT and off-line REC)
- separate maps for F/B/R/MUO, BAC and CAL (off-line only)
- evaluated for single muon in $(p_{t_{eff}}, \eta)$ bins, where as $p_{t_{eff}}$ is used:
(motivated by the CAL/BAC geometry and scaling of the muon path length)
 - p in Forecap
 - p_t in Barrel
 - p_z in Rearcap
- proper identification of **the triggering muon** is crucial
- → the DATA/MC ratio delivers the correction weight: $\epsilon_x = \frac{\epsilon_{DATA}}{\epsilon_{MC}}$

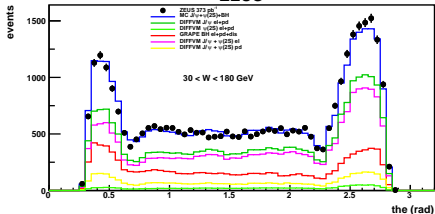
Muon correction maps: (p_z, p_t, p vs. η) - DATA



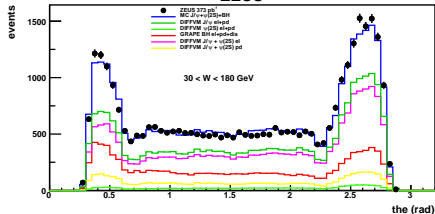
- muon tomography (+ muon range)
- probability (%) to fire FLT-SLT-TLT-REC chain by single muon on ($p_z, p_t, p; \eta$) grid
- X-axis (along eta): Rear-MUO, Barrel-MUO, Forward-MUO detectors
- only events with $M(\mu^+, \mu^-) < 6$ GeV
(ie. in the phase space range of di-muon mass fits)
- p_z, p_t, p grid: 100 MeV per bin ($p_{eff} < 3$ GeV), 250 MeV per bin ($p_{eff} > 3$ GeV)
- size of the grid is subject to systematics

θ and η distribution for μ^+ and μ^-

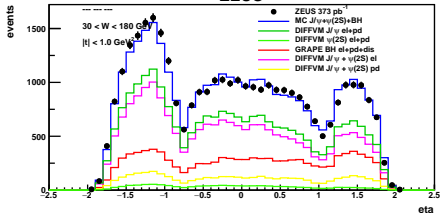
ZEUS



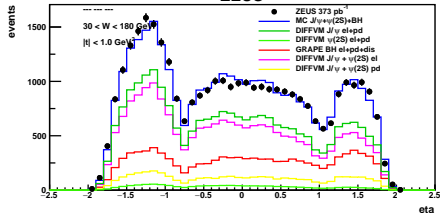
ZEUS



ZEUS

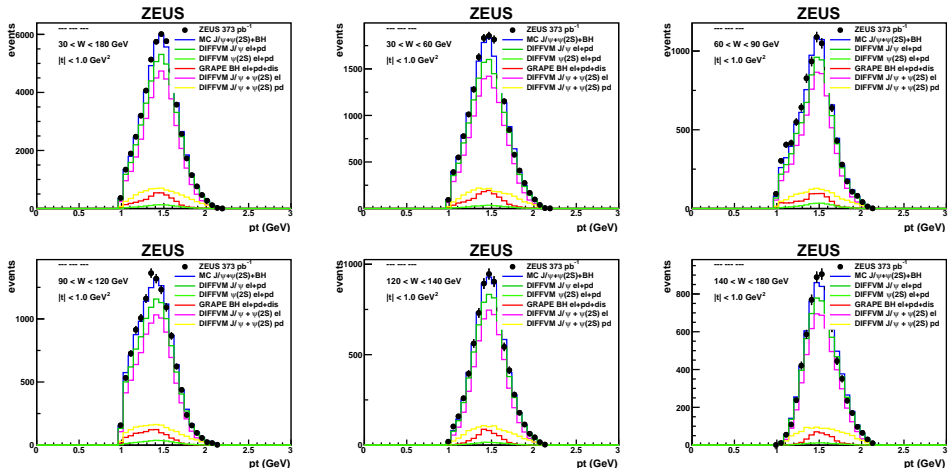


ZEUS



- θ and η distribution for muon tracks (μ^+ and μ^- , $\eta \in (-2, 2)$)
- after MC tuning and muon corrections
- well reproduced acceptance and efficiency of the detector

2-prongs: JPSI mass window: $p_t^{\mu^\pm}$ in W bins



- ALL events and in W_1, W_2, W_3, W_4, W_5 bins

- **Jacobian peak around 1.5 GeV ($M_{JPSI}/2$)**

- almost no JPSI (and PSI2S) event loss due to the $p_t^{\mu^\pm} > 1$ GeV cut

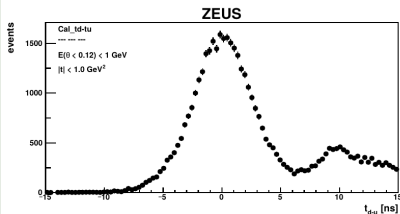
Triggering on soft muons ($p_T \sim 1 \div 2$ GeV)

- working on the steep edge of the trigger acceptance curve
- muon efficiency corrections extracted from data are crucial (trigger and off-line corrections)
- design the trigger to identify **the triggering muon** on all trigger levels

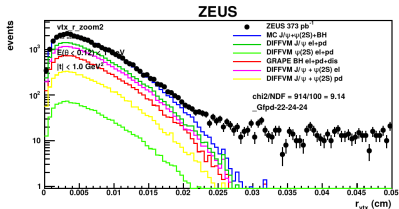
- **W distributions: 2-prongs**

Lesson learned:

CAL timing: Down-Up MIP clusters



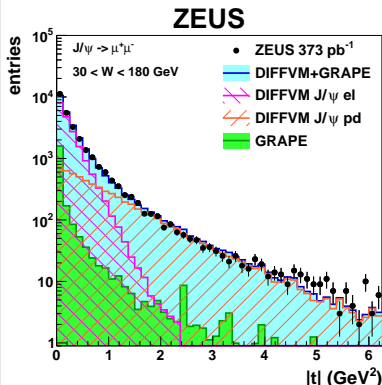
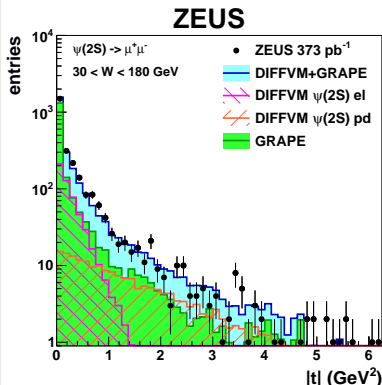
transverse distance to the vertex



Reducing the Cosmic muon background

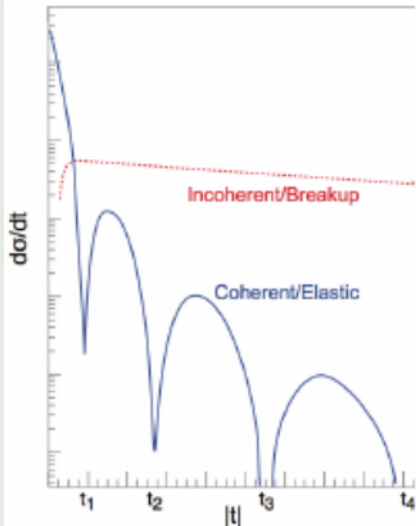
- angular tracking resolution (co-linearity cut !)
- timing cut on CAL MIP clusters

- $|t|$ distributions: 2-prongs

$|t|: 2.8 < M(\mu^+\mu^-) < 3.4$ GeV $|t|: 3.4 < M(\mu^+\mu^-) < 4.0$ GeV

- single exp() parameterization of elastic and p.diss (DIFFVM MC)
- proton dissociation dominates for $|t| > 1.0$ GeV²
- proton dissociative fraction: $f_{p,diss} = 0.16 \pm 0.01$ ($|t| < 1.0$ GeV²) from t -spectra fit
- $\psi(2S)$ channel dominated by Bethe-Heitler continuum- $\mu^+\mu^-$ background

Elastic vs. p-diss t -spectra



Modeling $|t|$ -distribution

- single $\exp(-b|t|)$ fit sufficient for ep data and moderate $|t|$
- b -slopes extracted from data...
- for eA expected more complex structure (diffraction pattern):
 - $t = -p_{t,VM}^2$: tracker resolution !
 - suppression of p-diss BG
- proper description of Bethe-Heitler BG is important

- **Modeling of nucleon resonance states**
- (low M_Y proton dissociation)

- $\frac{d\sigma}{dM_Y^2} \sim \frac{1}{M_Y^{2(1+\epsilon)}}$ for $M_Y^2 \geq 3.6 \text{ GeV}^2$ (continuum region)
- $\frac{d\sigma}{dM_Y^2} \sim \frac{f(M_Y^2)}{M_Y^{2(1+\epsilon)}}$ for $M_Y^2 < 3.6 \text{ GeV}^2$ (p resonance region)
- $f(M_Y^2)$ from the **fit the the p.diss cross section on deuterium:**
 $pD \rightarrow YD$ (Phys. Rep. **101** (3) (1983), 169)
- for $M_Y < 1.9 \text{ GeV}$ several resonances are included
(Pomeron carries quantum numbers of the vacuum ($l=0$, $G = P = C = +$)
only N^{*+} states with $J^P = \frac{1}{2}^+, \frac{3}{2}^-, \frac{5}{2}^-, \dots$)
- $N^{*+} = N(1440), N(1520), N(1680), N(1700)$
- N^{*+} decays into: $N\pi, \Delta\pi, N\rho, N\pi\pi$ included (BR from PGD 1992)
- N^{*+} decays isotropically in their rest frame
- dissociation in the continuum state carried by JETSET
(splitting proton into $q - qq$ system, q couples to \mathbb{P} , leaving qq spectator)

- $d\sigma \sim L_{\mu\nu} W^{\mu\nu}$

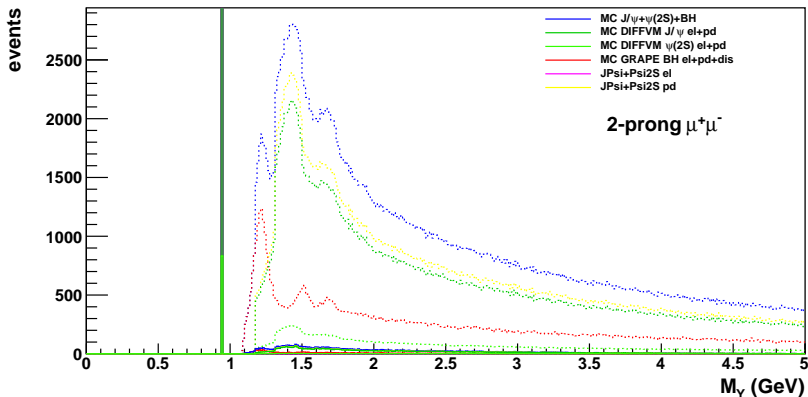
- hadron tensor:

$$W^{\mu\nu} = W_1 \left(-g^{\mu\nu} + \frac{q^\mu q^\nu}{q^2} \right) + W_2 \frac{1}{M_P^2} \left(p_P^\mu - \frac{p_P \cdot q}{q^2} q^\mu \right) \left(p_P^\nu - \frac{p_P \cdot q}{q^2} q^\nu \right)$$

- $W_{1,2}(Q_P^2, M_{had})$ are proton electromagnetic structure functions
- for $M_{had} < 2 \text{ GeV}$ $W_{1,2}$ parameterized by Brasse et al. (Nucl. Phys. **B 110** (1976) 413.) (resonance region)
- for $M_{had} > 2 \text{ GeV}$ $W_{1,2}$ parameterized by ALLM97 (hep-ph/9712415) (continuum)
- both parameterizations from **fits to experimental total $\gamma^* p$ cross sections**
- exclusive hadronic final state generated by SOPHIA
- (plus DIS di-leptons diagrams, in the framework of QPM, using PDF's)

MC generator level: zoom at low $M_Y < 5$ GeV (lin scale)

M_Y gener before and after cuts



- different structure of nucleon resonances between **GRAPE** and **DIFFVM** (!?)
- which is right ?
- how much it is important for MC based p.diss BG subtraction ?

Modeling of nucleon resonances N^* , etc...

- more realistic, unified treatment of nucleon resonances will be required
- feedback from data (?)

- **Modeling of nucleon resonance states**
- (b -slope for low $M_Y N^*$ states)

dependence of p.diss t -slopes on M_Y

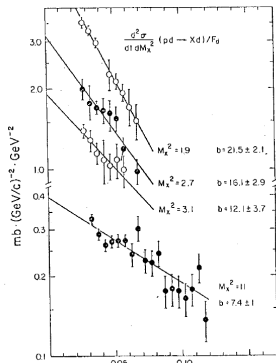
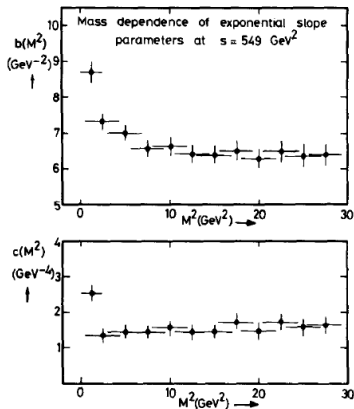


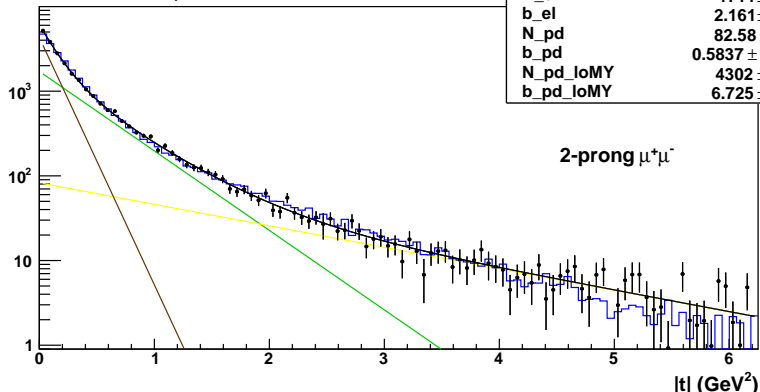
Figure 2 - Differential cross-sections vs. t for $pp \rightarrow Xp$, extracted from $pd \rightarrow Xd$, at $P_{\text{lab}} = 275 \text{ GeV}/c$, for $M_X^2 = 1.9, 2.7, 3.1$, and 11 GeV^2 .

$$\frac{d^2\sigma}{dt dM_X^2} (pd \rightarrow Xd)$$

- Fermilab and ISR data: Phys. Rev. **D 14**, 3148 (1976), Nucl. Phys. **B 108**, 1 (1976)
- steeper proton dissociative t -spectra for very low M_Y and $|t|$ ($b_{pd} \approx 10 \div 20 \text{ GeV}^{-2}$)
- effect observed for pp (and pd) collisions (old language: $M_X \rightarrow M_Y$)
- is it also present for γp p-dissociative interactions ?
- if yes, it has big impact on the estimation of p.diss fraction (extrapolated from model dependent MC)

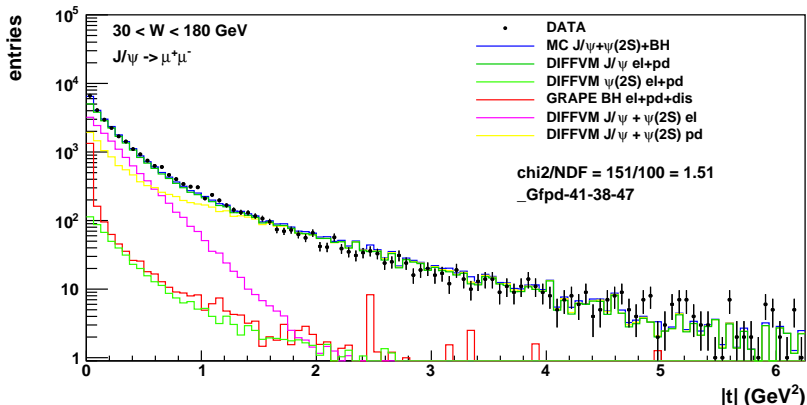
J/ψ : ALL W

events



- J/ψ mass window, BH background subtracted
- **unconstrained triple exp() fit**, all $M_Y < \sim 6$ GeV (elastic selection)
- there is some room for third, steeper component $b_{pd-loMY} \approx 6.5 \text{ GeV}^{-2}$
- but the elastic b_{el} slope becomes very unrealistic: $b_{el} \approx 2 \text{ GeV}^{-2}$

Introducing third $\exp()$ component on MC level

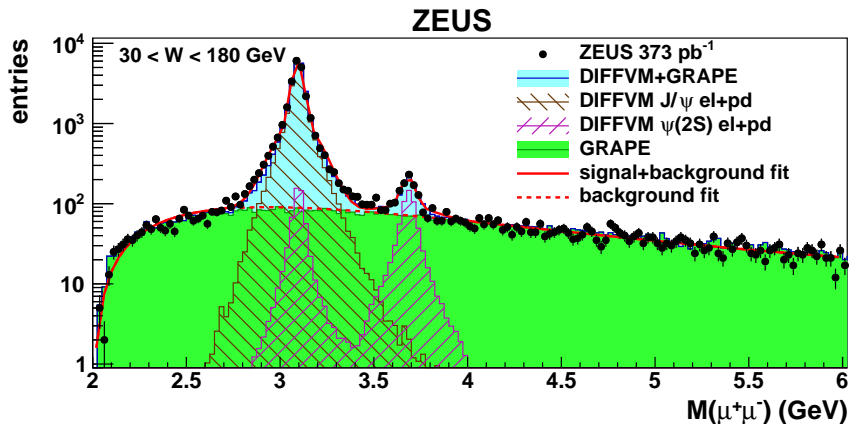


- **constraining the low M_Y component**
- reweighting the resonances $M_Y < 1.9 \text{ GeV}$ to $b_{pd-loM_Y} = 6.5 \text{ GeV}^{-2}$ (yellow histogram)
- **DATA/MC agreement is as good as for two components** (see next page)
- ... but the p.diss fraction increases from ~ 20 to $\sim 40\%$ (for $|t| < 1 \text{ GeV}^2$)

Modeling of t -slope for nucleon resonances N^* , etc...

- is the steeper t -slope dependence observed for ep and eA data ?

- **Signal Extraction, Mass spectra**

$M(\mu^+\mu^-)$ 

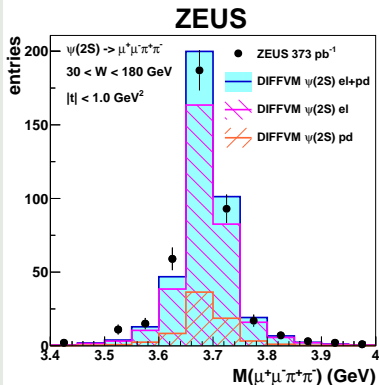
- full phase space: $30 < W < 180$ GeV, $|t| < 1.0$ GeV²
- events yield: $\sim 23\,000$ J/ψ and ~ 700 $\psi(2S)$
- resonant background under J/ψ peak

Tracker resolution

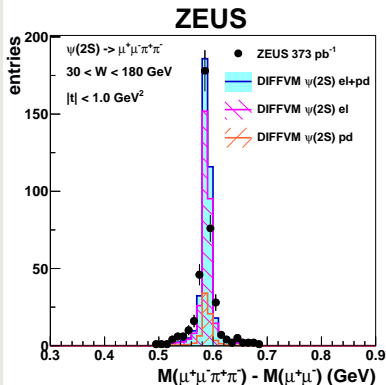
- instrumentation (tracking) close to the beam-pipe very important for high and low W measurement

4-prongs: mass spectra: $\psi(2S) \rightarrow \mu^+ \mu^- \pi^+ \pi^-$

$$M(\mu^+ \mu^- \pi^+ \pi^-)$$



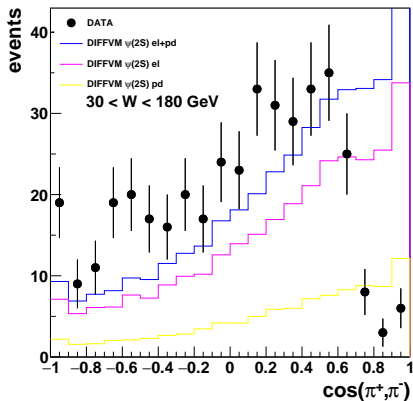
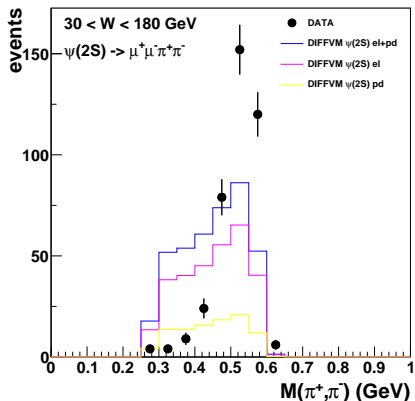
$$M(\mu^+ \mu^- \pi^+ \pi^-) - M(\mu^+ \mu^-)$$



- events yield: $\sim 400 \psi(2S)$
- better resolution on mass difference \rightarrow cascade decay of $\psi(2S)$
- **almost background free channel, but: ...**
- **TAKE CARE: SLOW PIONS are VERY FRAGILE !**

- **(SLOW) Pions phase space reweighting**

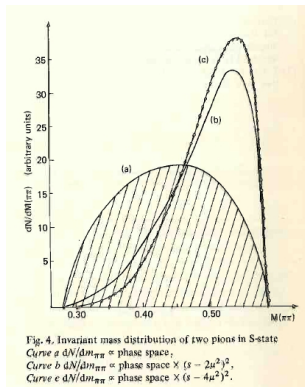
4-PRONGS: $M(\pi^-, \pi^+), \cos(\pi^-, \pi^+)$



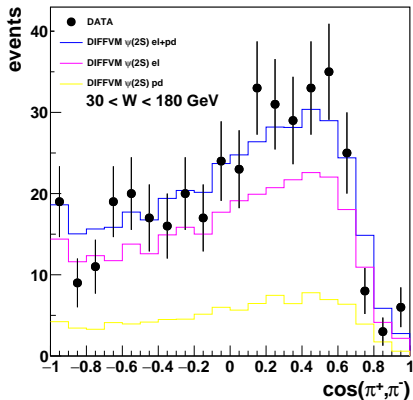
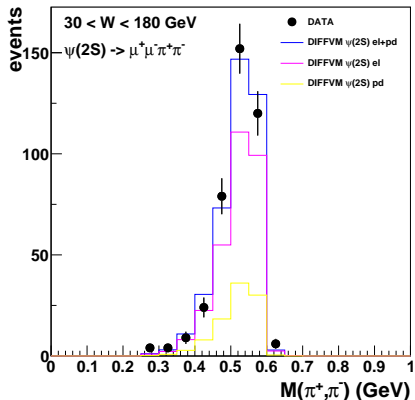
- $\psi' \rightarrow J/\psi + \pi^+ \pi^-$
- $M(\pi^-, \pi^+), \cos(\pi^-, \pi^+)$
- DIFFVM MC **before** pions phase space reweighting

Pions phase space reweighting (DIFFVM 4-prongs)

- $weight = (M(\pi^-, \pi^+)^2 - 4.0 * M_\pi^2)^2$
- ref: Phys_Lett_B61_1976_183.pdf
- final $\pi^+\pi^-$ interaction is not in pure S-state
- → for the impact of this correction see next page



4-PRONGS: $M(\pi^-, \pi^+), \cos(\pi^-, \pi^+)$



- $\psi' \rightarrow J/\psi + \pi^+ \pi^-$
- $M(\pi^-, \pi^+), \cos(\pi^-, \pi^+)$
- DIFFVM MC **after** pions phase space reweighting

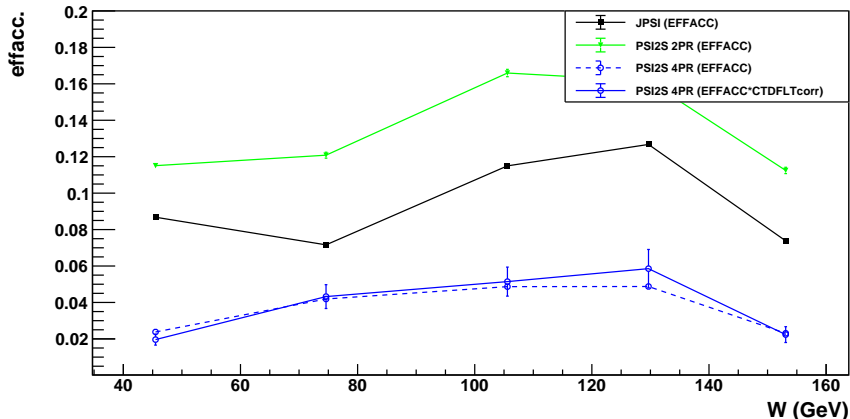
SLOW PIONS PHASE SPACE: $\psi(2S) \rightarrow \mu^+ \mu^- \pi^+ \pi^-$

- trust and check :)
- after long time wisdom of old generations sometimes disappear...

- **selection efficiency**

Acceptance*efficiency in W bins: elastic

EFFACC (el) of JPSI, PSI2S-2PR, PSI2S-4PR vs. W

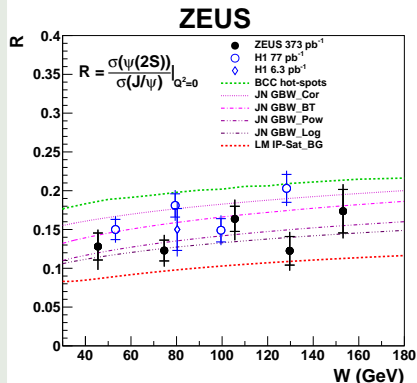


- JPSI, PSI2S 2- and 4-prong (2 ÷ 16%)
- Higher di-muon acceptance for higher mass state (PSI2S)
- similar angular coverage for final state di-muons
- second W bin (W_2) is the “dip” for di-muon acceptance

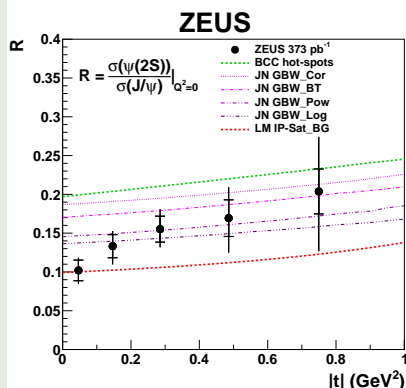
- **Cross section ratio**


cross section ratio $R = \sigma_{\psi(2S)}/\sigma_{J/\psi(1S)}$: Final Results

R vs. W



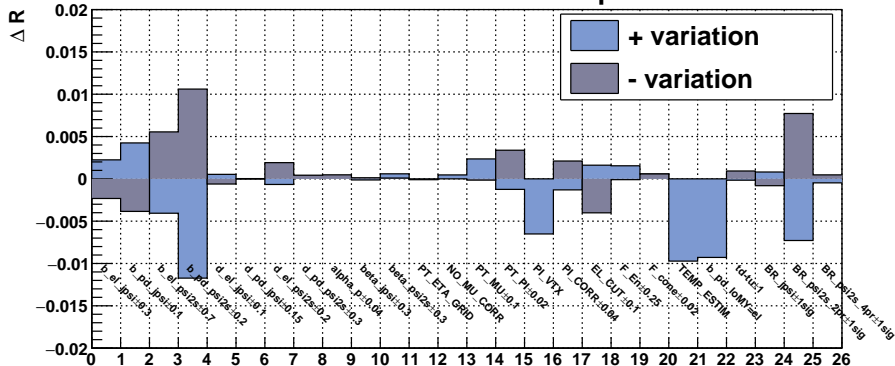
R vs. $|t|$



- for R vs. W ZEUS (full dots) and H1 (open markers) results are compared
- **no W dependence observed, moderate increase with $|t|$**
- good agreement between data and theoretical models (see next page)
- errors at high- $|t|$ points dominated by systematics (\rightarrow proton dissociative fraction) 

h_SYST_TOT_0

R GLOBAL: SYST error components



● biggest contributions from:

- b -slope variation of t -dependence (esp. for b_{pd} of $\psi(2S)$) → **p.diss fraction**
- event number estimator (MC templates fit instead of Gaussian fit) → **BH bckgr.**
- slow pions vertexing → **very soft tracks reconstruction**
- $BR(\psi(2S) \rightarrow \mu^+ \mu^-)$ → **PDG world average...**

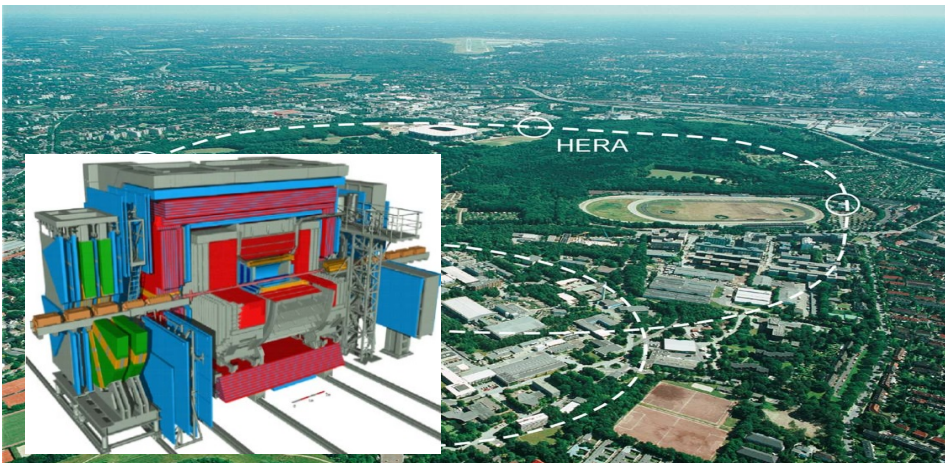
- case study on exclusive $J/\psi \rightarrow \mu^+ \mu^-$
- **DIS vs. photoproduction** (trigger challenge)
- Monte Carlo(s) models : tuning of parameters
- **better measurements** of MC parameters needed (t -slopes, ...)
- mass resolution vs. W (forward and rear trackers acceptance)
- **proton dissociation** → **role of forward taggers**
- proton dissociation → t -spectra study (are two $\exp()$ sufficient at low $|t|$?)
- to make progress: **precision is needed !**
- **statistic** → more luminosity
- **systematics** → **better instrumentation** (forward detectors, taggers)
- **systematics** → **better theoretical models** (MC), ...

- Backup plots follow...

- HERA accelerator, ZEUS experiment
- case study based on new analysis
(J/ψ to $\psi(2S)$ cross section ratio in exclusive photoproduction, $\mu^+\mu^-$ channel)
- Physics Motivation \longrightarrow see backup plots
- DIS vs. PHP (trigger strategy)
- Data sample, signal and background
- proton dissociation
- cosmic muons
- Monte Carlo, tuning of parameters,
controlling efficiency/acceptance corrections
- extracting t -slopes from data
- proton dissociation (N^* nucleon resonances, etc.)
- Summary

HERA and ZEUS: 1992 – 2007, DESY, Hamburg

HERA: world's first and only $e^\pm p$ collider, $E_e = 27.5 \text{ GeV}$, $E_p = 920 \text{ GeV}$ ($\sqrt{s} = 318 \text{ GeV}$)

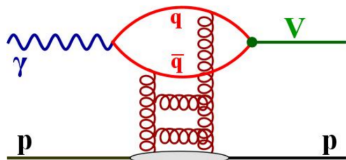


ZEUS: multipurpose, hermetic detector (MVD, CTD, CAL, F/B/RMUON, BAC, ...)

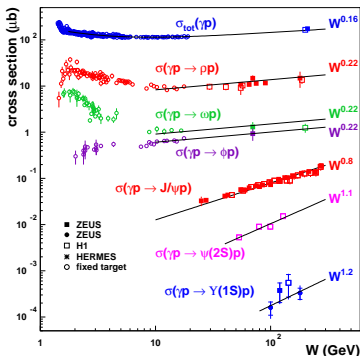
Total luminosity: $\int \mathcal{L} \sim 500 \text{ pb}^{-1}$ collected during HERA I + II running periods

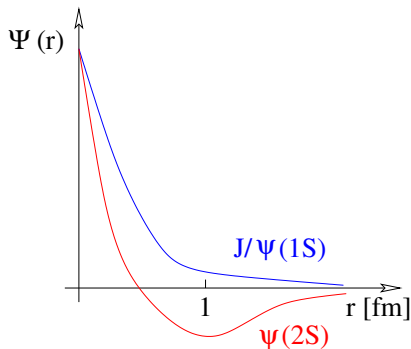
- **Physics Motivation**

W-dependence of the VM exclusive cross section for PHP: $\sigma(\gamma p \rightarrow Vp)$



- process driven by gluons
 - $\sigma \propto [xg(x, \mu^2)]^2$
 - $x = \mu^2/W^2$: high $W \rightarrow$ small x
 - $\mu^2 \propto (Q^2 + M_V^2)$
- total cross sections σ_{VM} for VM photoproduction spans over 6 orders of magnitude ! (for higher VM masses \rightarrow smaller transverse size of $q\bar{q} \rightarrow$ "color screening")
 - σ_{VM} rises with γp c.m. energy W as W^δ ($x \sim 1/W^2$: high $W \rightarrow$ small x) (for heavy VMs expected from the gluon behavior in the proton \rightarrow probing small x)
 - power δ rises with M_V^2 from "soft" ($\delta \approx 0.22$) to "hard" ($\delta \approx 1.2$) processes



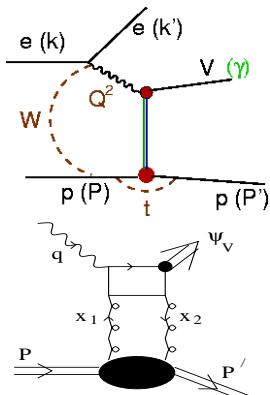


$$\text{Ratio } R = \frac{\sigma_{\gamma p \rightarrow \psi(2S)p}}{\sigma_{\gamma p \rightarrow J/\psi(1S)p}}$$

- sensitive to radial wave function of charmonium
- provides insight into the dynamics of the hard process
- (some systematics cancel out)

- $J/\psi(1S)$ and $\psi(2S)$ have the same quark composition but distinctive wave functions
- $\psi(2S)$ has a node at ≈ 0.4 fm
- $\langle r_{\psi(2S)}^2 \rangle \approx 2 \langle r_{J/\psi(1S)}^2 \rangle$
- pQCD models predict $R \sim 0.17$ in PHP and **rise of R with Q^2 in DIS**
- $\psi(2S)$ **cross section is expected to be suppressed w.r.t. the J/ψ production**
- (Both Vector Mesons masses are much smaller than the γp center-of-mass energy)

Production of Vector Mesons in Exclusive Diffraction in ep Scattering



Exclusive process: proton stays intact
 Proton dissociation also possible \rightarrow
 background

pQCD: M_V^2 and Q^2 - set the scale at which the W and $|t|$ are probed
 Process sensitive to the **gluon density** in the proton

Kinematics: $M_V^2, Q^2, W, |t|$

M_V^2 - vector meson mass squared

$Q^2 (= -q^2 = -(k - k')^2)$ - the photon virtuality
 (emitted by the incoming electron):

- $Q^2 \approx 0$ GeV² PHP (*Photoproduction*)
- larger Q^2 for DIS (*Deep Inelastic Scattering*)

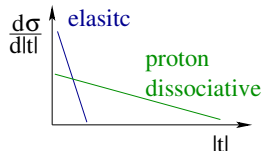
$W = (q + P)^2$ - invariant mass of the γp system

$$W \approx \sqrt{2E_P(E - p_z)}_V$$

$|t|$ - 4-momentum transfer at the proton vertex

$$t = (P - P')^2$$

$$t \approx -p_{T,V}^2$$



- **DATA sample**
- (selection of exclusive di-muon events)

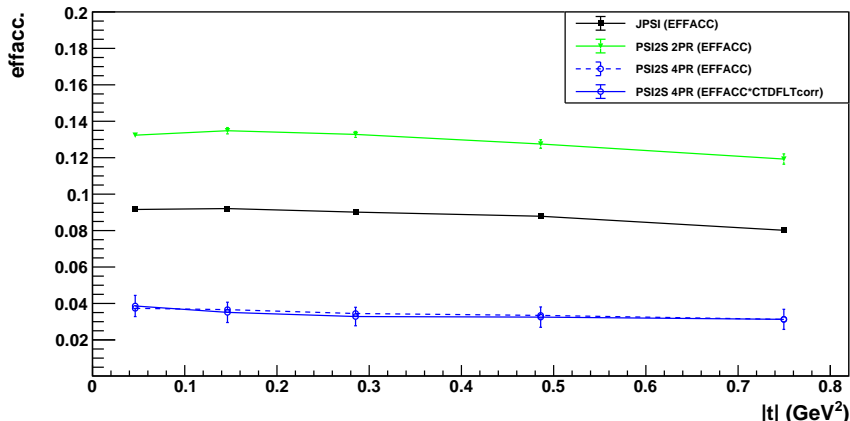
- **Exclusive Muon Triggers** (F/B/R/MUON or BAC)
- **Tracking and Vertex**
 - $N_{track} = 2$, oppositely charged tracks matched to the primary vertex ($\eta \in (-1.9, 1.9)$)
 - **both tracks identified as a muon in CAL, at least one in F/B/RMUON or BAC**
 - $p_T > 1.0$ GeV of each track
 - anti-COSMIC cuts (CAL timing, acolinearity: $\cos(\mu^+, \mu^-) < -0.985$)
- **Elasticity/Exclusivity and Photoproduction cuts** (on CAL Energy)
 - no scattered electron found in CAL
 - $E_{clu} < 0.5$ GeV for clusters not matched to muons (or pions)
(corresponds to an effective cut on $Q^2 < 1$ GeV²)
 - $E(\theta < 0.12rad) < 1$ GeV
the sum of the energy in the FCAL cone around the beam-pipe;
to suppress proton-dissociative events, $ep \rightarrow e + VM + Y$
(corresponds to a requirement for $M_Y \lesssim 5$ GeV)
- **Kinematic range (analysis phase space):**
 - $30 < W < 180$ GeV
 - $|t| < 1.0$ GeV²
 - $Q^2 < 1$ GeV² (median $Q^2 \approx 3 \times 10^{-5}$ GeV²)

Event Selection: 4-prongs sample: $\psi(2S) \rightarrow \mu^+ \mu^- \pi^+ \pi^-$

- (only differences w.r.t. the 2-prong channel)
- $N_{track} = 4$, (two oppositely charged pairs, sorted by p_T)
- **highest momentum pair: muon candidates**
lowest momentum pair: pion candidates
- no anti-COSMIC cuts
- transverse momentum of pion candidates: $p_T^\pi > 0.12$ GeV;
- $2.8 < M(\mu^+ \mu^-) < 3.4$ GeV (J/ψ window)
- $M(\mu^+ \mu^- \pi^+ \pi^-) - M(\mu^+ \mu^-)$ in (0.5 – 0.7) GeV window
(**cascade decay of $\psi(2S)$**)

Acceptance*efficiency in $|t|$ bins: elastic

EFFACC (el) of JPSI, PSI2S-2PR, PSI2S-4PR vs. $|t|$



- JPSI, PSI2S 2- and 4-prong (4 ÷ 14%)
- Higher di-muon acceptance for higher mass state (PSI2S)
- flat in $|t|$ (no angular correlations to $|t|$)
- dashed line after CTD FLT corrections

2-prongs: Signal extraction: fit parameterization

- **Double Gaussian** shape: $G(x)$ or $g(x) = N \cdot \Delta \cdot \frac{1}{\sqrt{2\pi}\sigma} \exp\left(-\frac{(x-m)^2}{2\sigma^2}\right)$
where: N – number of events, Δ – mass bin width,
 m – mean value, σ – RMS

- for J/ψ : $N_1 \cdot G_1(x) + N_2 \cdot G_2(x)$

- for ψ' : $N'_1 \cdot g_1(x) + N'_2 \cdot g_2(x)$

- introducing: $N = N_1 + N_2$, $N' = N'_1 + N'_2$, $R = \frac{N'}{N}$

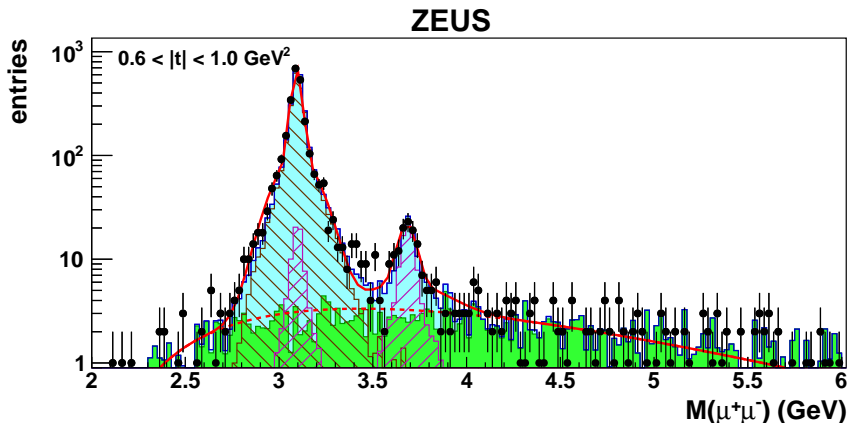
- **with additional constrains**: $m_1 = m_2$, $m'_1 = m'_2$,
 $\frac{\sigma'_1}{\sigma_1} = \frac{\sigma'_2}{\sigma_2} = \alpha$, $\xi = \frac{N_1}{N} = \frac{N'_1}{N'}$ (**scaling of the mass resolution**)

- final formulae:

$$F(x) = N \cdot ((\xi \cdot G_1(x) + (1 - \xi) \cdot G_2(x)) + R \cdot (\xi \cdot g_1(x) + (1 - \xi) \cdot g_2(x))) + BG(x)$$

- **background function**: $BG(x) = A \cdot (x - B)^C \cdot \exp(-D(x - B) - E(x - B)^2)$
where A, B, C, D, E are fit parameters, B fixed ($= 2p_{t,min}^\mu$)

$M(\mu^+\mu^-)$



- **t5 bin:** $30 < W < 180 \text{ GeV}$, $0.6 < |t| < 1.0 \text{ GeV}^2$
- higher $|t|$, small Bethe-Heitler continuum $\mu^+\mu^-$ contribution
- **BUT: high contamination proton dissociative events** $\rightarrow t$ -spectra

- **Cross section ratio**

Cross section ratio $R = \frac{\sigma(\psi(2S))}{\sigma(J/\psi(1S))}$, full kinematic range

$30 < W < 180$ GeV, $|t| < 1.0$ GeV², $Q^2 < 1.0$ GeV²

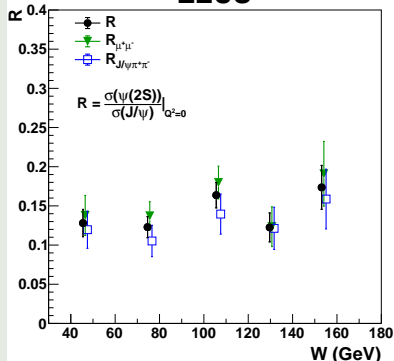
$\psi(2S)$ decay mode	$R = \frac{\sigma(\psi(2S))}{\sigma(J/\psi(1S))}$
$\mu^+ \mu^-$	0.154 ± 0.012
$J/\psi(\rightarrow \mu^+ \mu^-) \pi^+ \pi^-$	0.125 ± 0.019
combined	$0.146 \pm 0.010^{+0.016}_{-0.020}$

- $$R_{J/\psi \pi \pi} = \frac{N_{\psi(2S)}}{N_{J/\psi(1S)}} \cdot \frac{Acc_{J/\psi(1S) \rightarrow \mu^+ \mu^-}}{Acc_{\psi(2S) \rightarrow J/\psi \pi^+ \pi^-}} \cdot \frac{1}{BR_{\psi(2S) \rightarrow J/\psi \pi^+ \pi^-}} \cdot \frac{1 - f_{p.diss}^{\psi(2S)}}{1 - f_{p.diss}^{J/\psi(1S)}}$$
- $$R_{\mu\mu} = \frac{N_{\psi(2S)}}{N_{J/\psi(1S)}} \cdot \frac{Acc_{J/\psi(1S) \rightarrow \mu^+ \mu^-}}{Acc_{\psi(2S) \rightarrow \mu^+ \mu^-}} \cdot \frac{BR_{J/\psi(1S) \rightarrow \mu^+ \mu^-}}{BR_{\psi(2S) \rightarrow \mu^+ \mu^-}} \cdot \frac{1 - f_{p.diss}^{\psi(2S)}}{1 - f_{p.diss}^{J/\psi(1S)}}$$
- $Acc_i = \frac{N_i^{reco}}{N_i^{true}}, f_{p.diss}^i$ - fraction of proton dissociative events
- $BR(\psi(2S) \rightarrow J/\psi \pi^+ \pi^-) = (34.68 \pm 0.3)\%$, $BR(\psi(2S) \rightarrow \mu^+ \mu^-) = (0.80 \pm 0.06)\%$,
 $BR(J/\psi \rightarrow \mu^+ \mu^-) = (5.961 \pm 0.033)\%$, $BR(\psi(2S) \rightarrow \mu^+ \mu^- \pi^+ \pi^-) = (2.07 \pm 0.02)\%$ (PDG 2020)
- both channels have similar precision and provide consistent results**

cross section ratio $R = \sigma_{\psi(2S)}/\sigma_{J/\psi(1S)}$: 2- and 4-prongs comparison

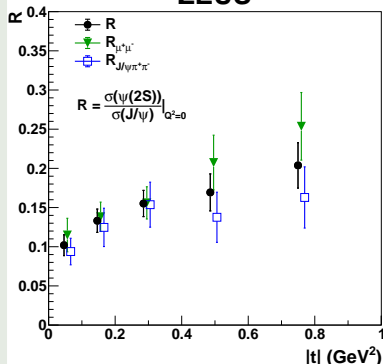
R vs. W

ZEUS



R vs. $|t|$

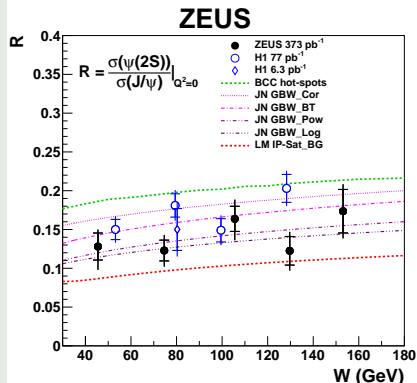
ZEUS



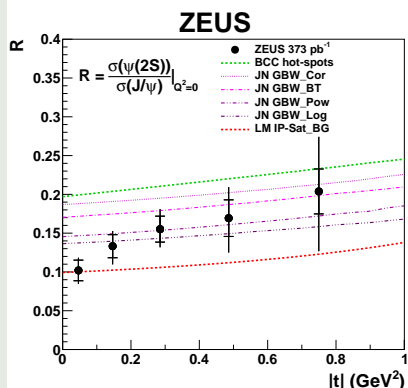
- $R_{\mu\mu}$ (2-prongs channel), $R_{J/\psi\pi\pi}$ (4-prongs channel) and combined R (full dots)
- statistical errors only
- **good agreement between two channels**

cross section ratio $R = \sigma_{\psi(2S)}/\sigma_{J/\psi(1S)}$: Final Results

R vs. W



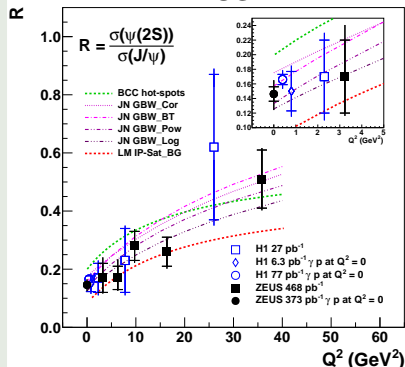
R vs. $|t|$



- for R vs. W ZEUS (full dots) and H1 (open markers) results are compared
- **no W dependence observed, moderate increase with $|t|$**
- good agreement between data and theoretical models (see next page)
- errors at high- $|t|$ points dominated by systematics (\rightarrow proton dissociative fraction)

R vs. Q^2

ZEUS



theoretical models:

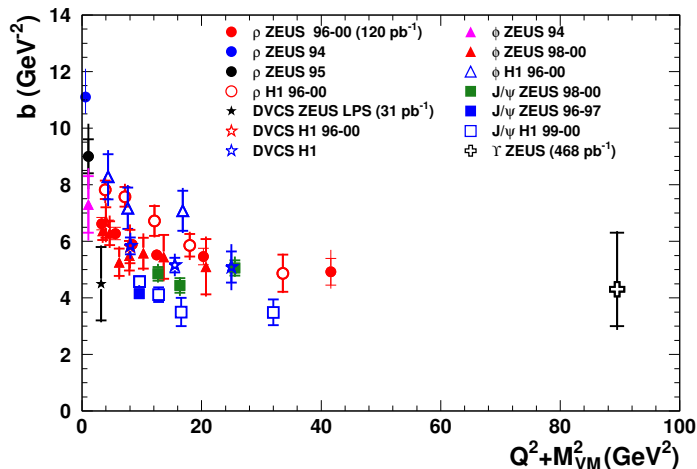
- Bendova, Cepila and Contreras (BCC hot-spots):
 - Phys. Rev. **D 99**, 034025 (2019).
- Jan Nemchik et al. (JN):
 - Eur. Phys. J. **C 79**, 154 (2019).
 - Eur. Phys. J. **C 79**, 495 (2019).
 - Phys. Rev. **D 103**, 094027 (2021).
- Lappi and Mäntysaari (LM IP-Sat):
 - Phys. Rev. **C 83**, 065202 (2011).
 - Phys. Rev. **D 87**, 034002 (2013).
 - PoS (DIS2014), 069 (2014).

- ZEUS (full dot) and H1 (open markers) photoproduction results plotted at $Q^2 \sim 0$
- DIS results are also presented vs. Q^2 : ZEUS (full squares) and H1 (open squares)
- good agreement between data and theoretical models (→ backup plots, page ??)

- **Tuning of DIFFVM Monte Carlo**

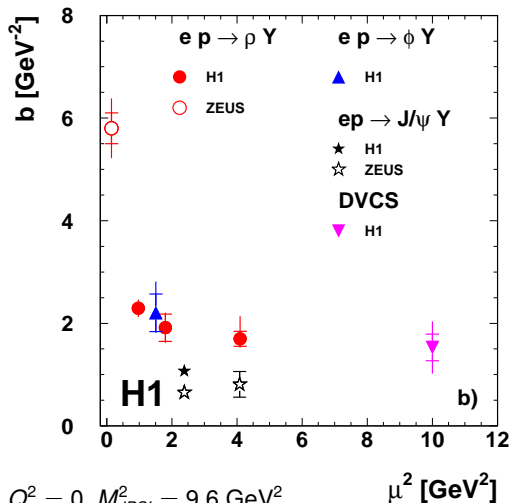
- Reweighting of MC sample at generator level
- **|t| dependence**: $\sim \exp(-b|t|)$, generated with $b_{el} = 4.0$, $b_{pd} = 1.0$
reweighted to:
 $b_{el} = 4.6 \pm 0.3$, $b_{pd} = 1.0 \pm 0.1$ (JPSI)
 $b_{el} = 4.3 \pm 0.7$, $b_{pd} = 0.7 \pm 0.2$ (PSI2S)
- shrinkage added by reweighting: $b = b_0 + 4.0\alpha' \log(W/W_0)$;
 $\alpha' = 0.12 \pm 0.04 \text{ GeV}^{-2}$, $W_0 = 90 \text{ GeV}$ (elastic only)
- **W dependence**: $\sigma \sim W^\delta$,
generated with $\delta = 0.88$ for both elastic and p.diss
reweighted to:
 $\delta_{el} = 0.67 \pm 0.10$, $\delta_{pd} = 0.42 \pm 0.15$ (JPSI)
 $\delta_{el} = 1.10 \pm 0.20$, $\delta_{pd} = 0.70 \pm 0.30$ (PSI2S)
- **M_Y dependence**: $\sim \frac{1}{M_Y^\beta}$, generated with $\beta = 2.5$
reweighted to $\beta = 2.4 \pm 0.3$ (both JPSI and PSI2S, p.diss only)
- **all parameters are subject to systematics checks**

Summary of elastic b -slopes



- PHP: $Q^2 = 0$, $M_{J\psi}^2 = 9.6$ GeV²
- significant difference for different measurements...

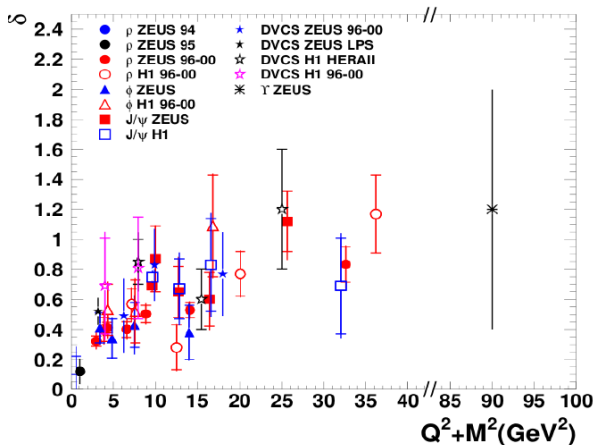
Summary of p.diss b -slopes



- PHP: $Q^2 = 0$, $M_{J/\psi}^2 = 9.6 \text{ GeV}^2$
- for J/ψ : $b_{pd} \sim 1.0 \text{ GeV}^{-2}$

μ^2 [GeV²]

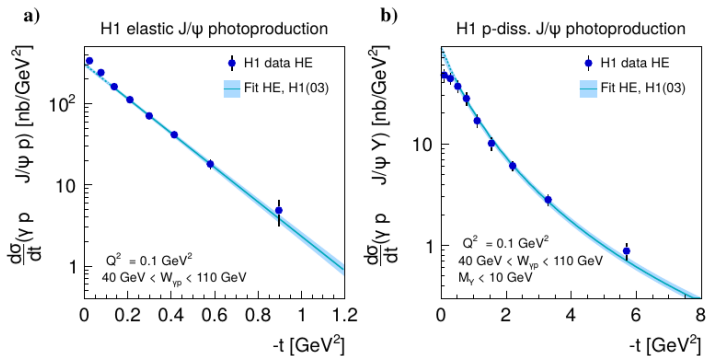
Summary of elastic δ -powers



- PHP: $Q^2 = 0$, $M_{J\psi}^2 = 9.6 \text{ GeV}^2$
- significant difference for different measurements...

- **Extracting of the model parameters**
- (example for the t -spectra slopes)

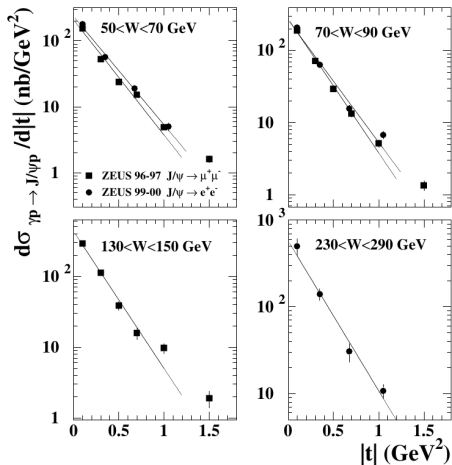
Example t -spectra: H1 results



- (Eur. Phys. J. **C 73** (2013), 2466)
- two **features common for many exclusive VM analyses**:
- **excess of elastic events for very small $|t|$**
(BH background ? p.diss BG for low M_Y - steeper b -slope ?)
- **deficit of p-diss. events for very small $|t|$** (threshold effect ?)

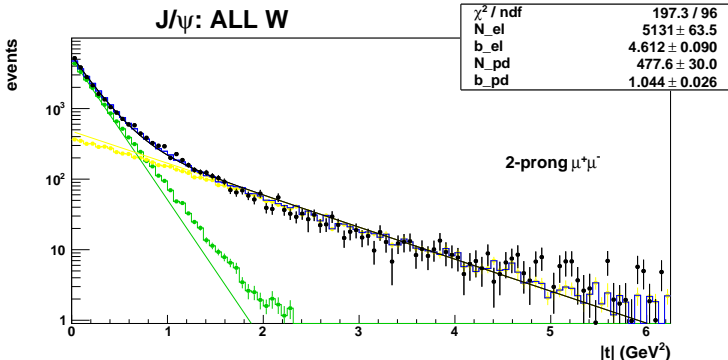
- **elastic t -slope:** $\sim \exp(-b_{el}|t|)$
- **H1:** (HERA-II JPSI paper) : $b_{el} = 4.88 \pm 0.15 \text{ GeV}^{-2}$
(Eur. Phys. J. **C 73** (2013), 2466)
- **H1:** (HERA-I PSI2S paper) : $b_{el} = 4.99 \pm 0.13 \pm 0.39 \text{ GeV}^{-2}$
(Phys. Lett. **B 541** (2002), 251-264)
- **ZEUS:** (HERA-I JPSI paper) : $b_{el} = 4.15 \pm 0.05^{+0.30}_{-0.18} \text{ GeV}^{-2}$
(Eur. Phys. J. **C 24** (2002) 3, 345-360)
- **some tension between ZEUS and H1 results**
- **H1:** unfolding t -spectra to correct for bin-to-bin migration, simultaneous fit to elastic + p.diss component
- **ZEUS:** single $\exp()$ fit to uncorrected data but sophisticated p.diss subtraction
- (this analysis : $b_{el} = 4.6 \pm 0.3 \text{ (stat. } \oplus \text{ syst.) GeV}^{-2}$)

ZEUS



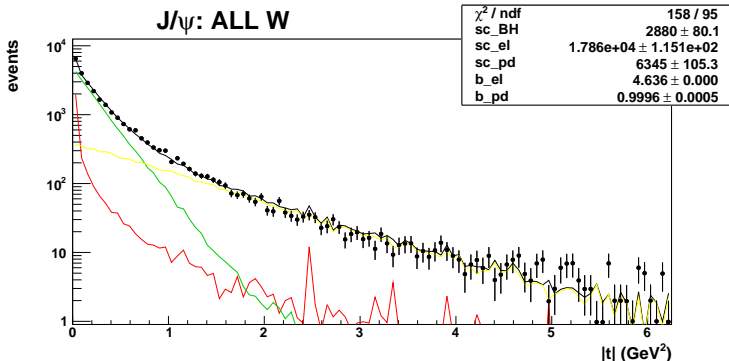
- single $\exp()$ fits for $|t| < 1.2 \text{ GeV}^2$
- p.diss dominates for $|t| > \sim 1.0 \text{ GeV}^2$

Figure 6: The differential cross-section $d\sigma_{\gamma p \rightarrow J/\psi p}/d|t|$ for exclusive J/ψ photoproduction for representative bins of W and for the decay channels, $J/\psi \rightarrow \mu^+\mu^-$ (squares) and $J/\psi \rightarrow e^+e^-$ (points). The vertical bars indicate the statistical uncertainties only. The full lines represent the results of a fit of the form $d\sigma/dt = d\sigma/dt|_{t=0} \cdot e^{bt}$ performed in the range $-t < 1.2 \text{ GeV}^2$ for the muon channel and in the range $-t < 1.25 \text{ GeV}^2$ for the electron channel.



- Green: elastic, Yellow: p.diss components
- double $\exp()$ fit to uncorrected data (MC superimposed on top of the fit)
- **Bethe-Heitler background subtracted** (via MC templates)
- $b_{el} = 4.6 \pm 0.1$ (stat.), $b_{pd} = 1.0 \pm 0.03$ (stat.)
- some smearing visible for MC on detector level (see next pages)

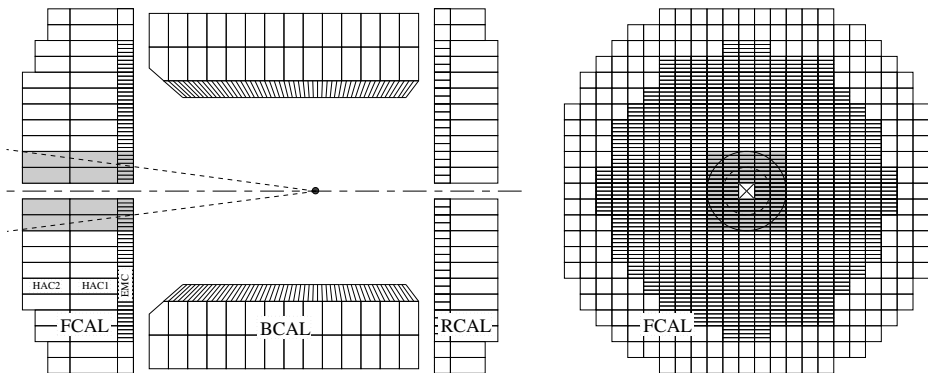
another approach: t -spectra template fit



- triple MC template fit to uncorrected data
- simultaneous fit of **BH+el+p.diss** components
- both b -slopes reweighted on generator “true” level
- $b_{el} = 4.6$, $b_{pd} = 1.0$ (problems with MINUIT errors, limited MC statistic)
- smearing and other “detector/cuts” effect included in templates

- **extracting fractions of proton dissociation**

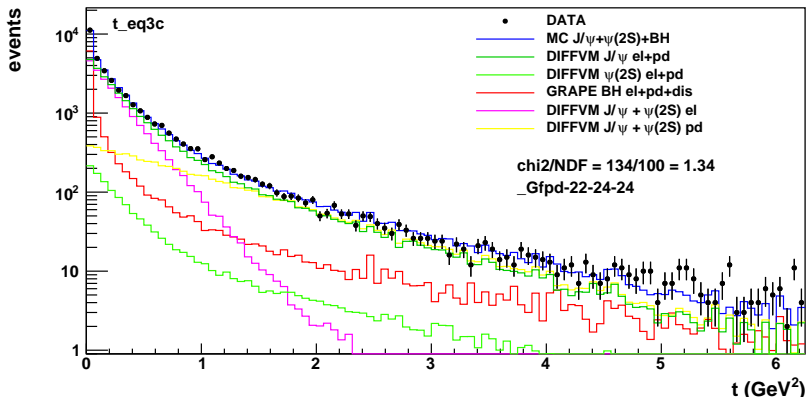
Proton dissociation taggers



- Energy in forward cone to **suppress p.diss events**: $\theta_{max} = 0.12 \text{ rad}$
- using EFO : “Energy Flow Objects” (trackers + CAL info):

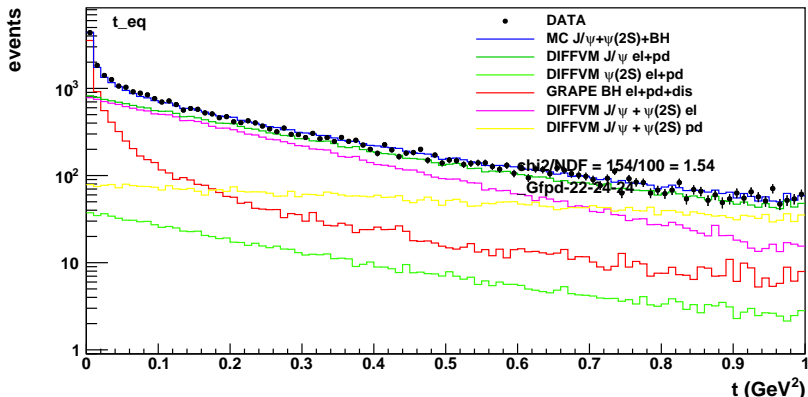
$$\left(\sum_{EFOs} E(\theta_{EFO} < \theta_{max}) \right) < 1 \text{ GeV}$$

2-prongs: $|t|$ distribution: all 2-prong events

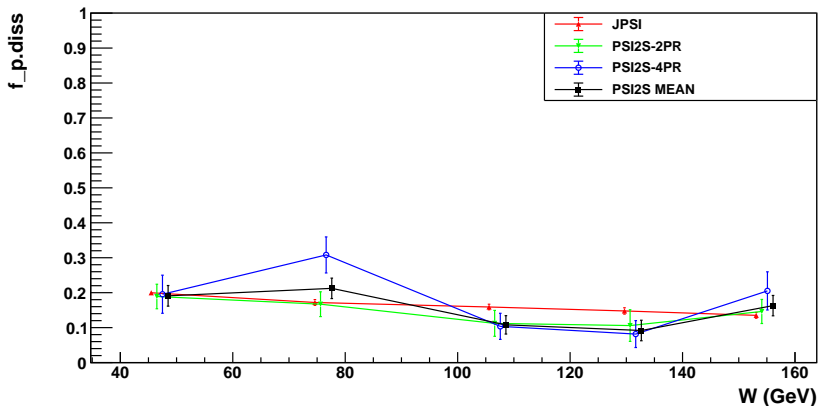


- spectra like this are used to evaluate the p.diss fractions (use longer “lever arm” then integrate it up to $|t| = 1.0 \text{ GeV}^2$)
- using root package TFractionalFitter
- fitted $f_{p.diss} = 0.22$ and $= 0.24$ (JPSI and PSI2S, BH subtracted)
- p.diss take over elastic around $\sim 1 \text{ GeV}^2$ (yellow and magenta histos)

2-prongs: $|t|$ distribution: all 2-prong events ($|t| < 1 \text{ GeV}^2$)

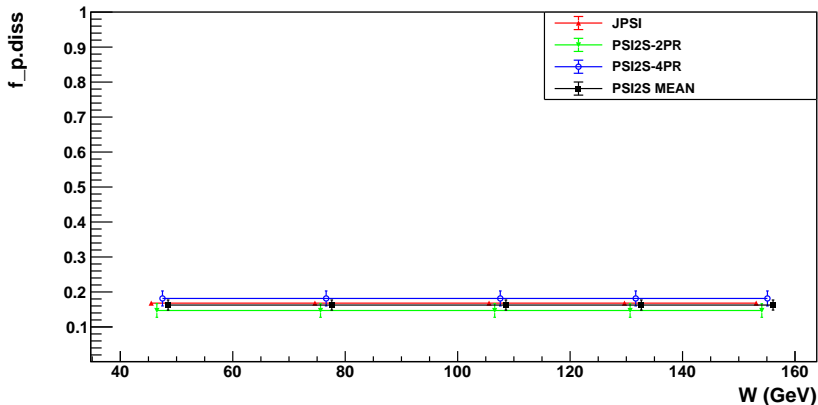


- zoom for $|t| < 1 \text{ GeV}^2$
- fitted $f_{p,diss} = 0.22$ and $= 0.24$ (JPSI and PSI2S, BH subtracted)
- $b_{el} = 4.6 \text{ GeV}^{-2}$, $b_{pd} = 1.0 \text{ GeV}^{-2}$ (JPSI)
- $b_{el} = 4.3 \text{ GeV}^{-2}$, $b_{pd} = 0.7 \text{ GeV}^{-2}$ (PSI2S)
- good MC modeling of the spectra is crucial !

fraction $f_{p,diss}$: JPSI and PSI2S 2PR, 4PR vs. W 

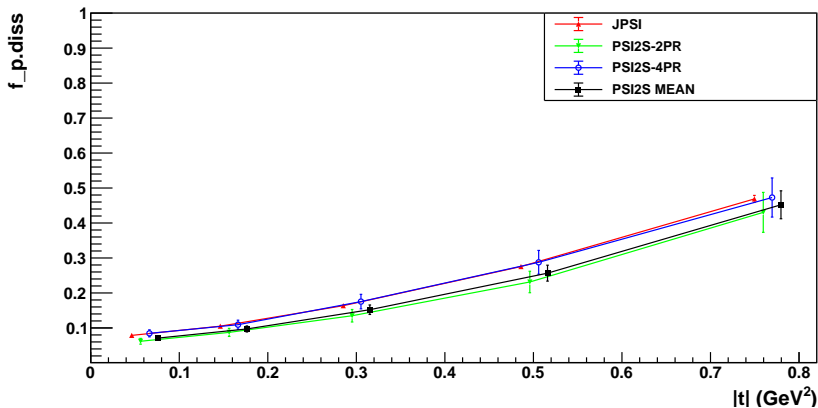
- average value $\sim 17\%$ JPSI and $\sim 16\%$ PSI2S (mean)
- compatible results for 2- and 4-prong channels, no W dependence
- **black**: weighted mean for PSI2S 2- and 4-prong
- bigger fluctuations for PSI2S 2- and 4-prongs

fraction $f_{p,diss}$: JPSI and PSI2S 2PR, 4PR vs. W



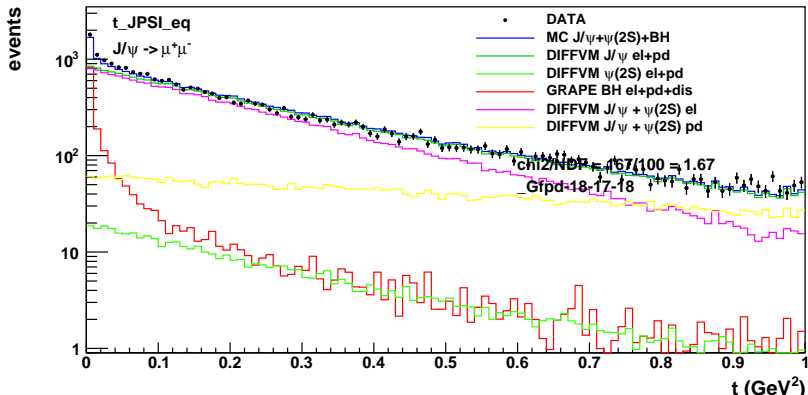
- average value $\sim 17\%$ JPSI and $\sim 16\%$ PSI2S (mean)
- **black**: weighted mean for PSI2S 2- and 4-prong (used in analysis)
- **for R analysis: the same mean value is used for all W bins**
- no significant impact on final ratio R ($f_{p,diss}$ fractions cancels out)

fraction $f_{p,diss}$: JPSI and PSI2S 2PR, 4PR vs. $|t|$



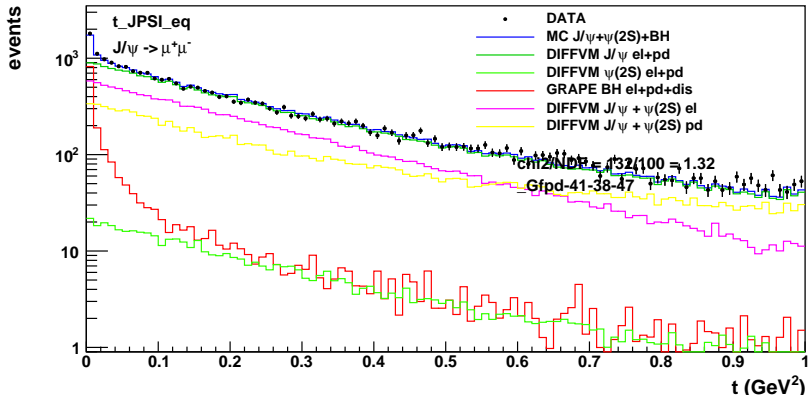
- compatible results for 2- and 4-prong channels
- **black**: weighted mean for PSI2S 2- and 4-prong (used in analysis)
- negligible effect on final R analysis ($f_{p,diss}$ fractions cancels out)
- bigger impact on systematics for large $|t|$ due to the b -slope variation !

Comparison with two $\exp()$ model (zoom for $|t| < 1.0 \text{ GeV}^2$)



- “standard” two components model: single $\exp()$ for elastic and p.diss
- yellow : p.diss with $b_{pd} = 1.0 \text{ GeV}^{-2}$
- some indication for excess of events at very low $|t| < 0.1 \text{ GeV}^2$ (?)
- could be imperfect description of BH background as well...
- $f_{p.diss} = 18\%$ ($|t| < 1 \text{ GeV}^2$)

Introducing third $\exp()$ component on MC level (zoom for $|t| < 1.0 \text{ GeV}^2$)

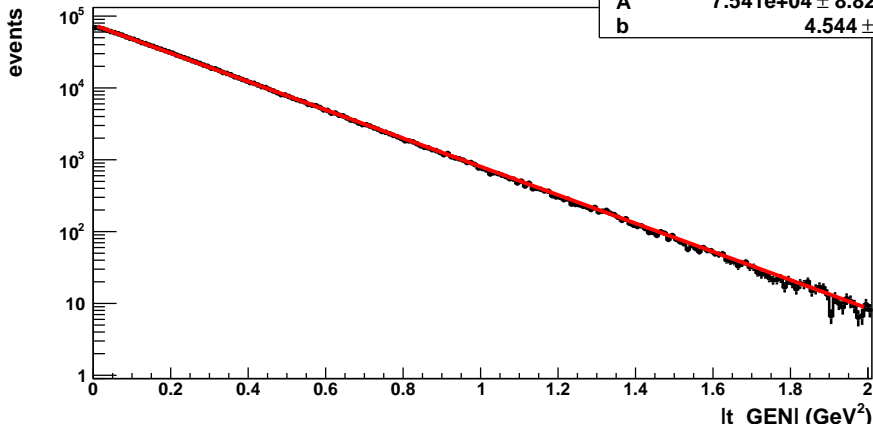


- reweighting the resonances $M_Y < 1.9 \text{ GeV}$ to $b_{pd-loMY} = 6.5 \text{ GeV}^{-2}$ (yellow histogram)
- a bit better agreement in terms of χ^2 ...
- $f_{p.diss} = 41\%$ ($|t| < 1 \text{ GeV}^2$) !?
- which scenario is realized by Nature ?
- → selection of clean p.diss, low M_Y events needed (incl. triggering on low M_Y)

- **MC exercise:** generator level \rightarrow detector level
- elastic and p.diss t -spectra

JPSI: MC gen. level ALL: b_el slope fit

χ^2 / ndf	740.3 / 198
A	$7.541\text{e}+04 \pm 8.821\text{e}+01$
b	4.544 ± 0.003



- elastic JPSI: generator level “true” $|t|$ -distribution before cuts, (calculated using gen-level 4-vectors including scattered e')
- reweighted to: $b_{el} = 4.6 \text{ GeV}^{-2}$ (realistic example)

JPSI: MC det. level: b_el slope fit

 χ^2 / ndf

226.5 / 158

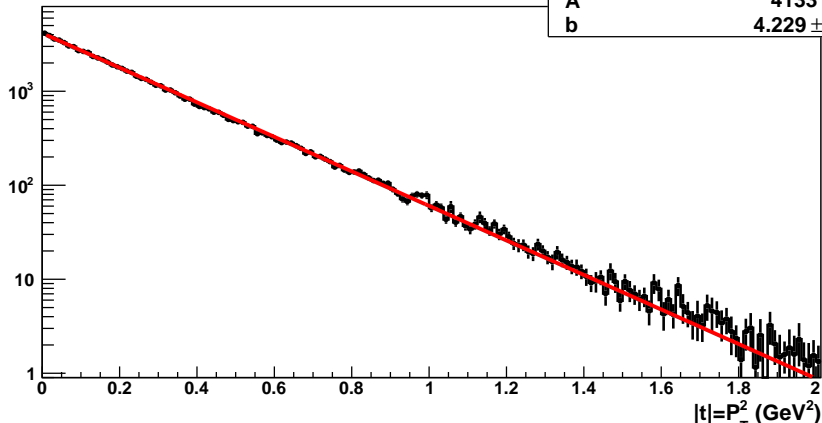
A

 4133 ± 17.9

b

 4.229 ± 0.012

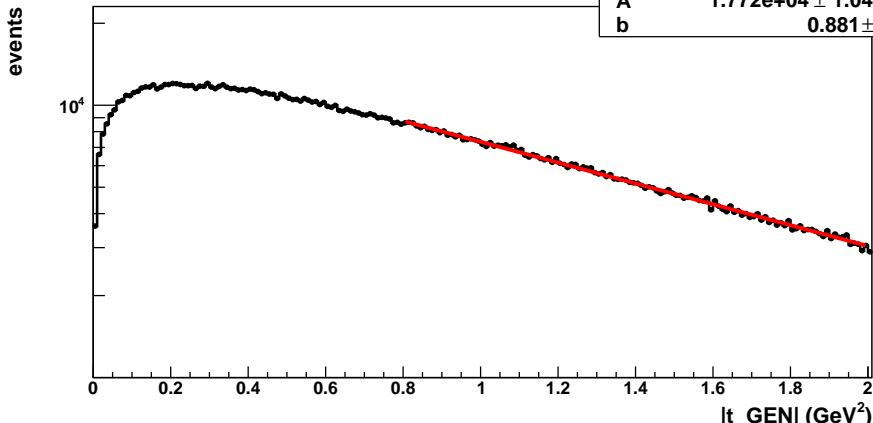
events



- elastic JPSI: **detector level** $|t|$ -distribution **after selection cuts**, **smearing** due to the finite resolution (bin to bin migration), calculated using **approximate formulae** $|t| \approx p_{T,VM}^2$
- **systematic shift** in the t -slope from 4.6 to 4.2 GeV^{-2}

JPSI ALL M_Y: MC gen. level ALL: b_pd slope fit

χ^2 / ndf	147 / 118
A	$1.772\text{e}+04 \pm 1.044\text{e}+02$
b	0.881 ± 0.004



- p.diss t -distribution is non-exponential already at generator level !
- **kinematical threshold effect**
(minimal 4-momentum transfer needed to excite/break-up the proton)

- t' variable corrected for threshold effect:

```

amprot = PMASS;
mdiff  = M_Y_gen;
q2g    = Q2_gen;
mppg   = mass_VM_gen;
Wg     = W_gen;

t_min = ((amprot*amprot-mdiff*mdiff+q2g+mppg*mppg) / (2*Wg)) *
        ((amprot*amprot-mdiff*mdiff+q2g+mppg*mppg) / (2*Wg))
        - (sqrt(((Wg*Wg + amprot*amprot+q2g) / (2*Wg)) *
                ((Wg*Wg + amprot*amprot+q2g) / (2*Wg)) - amprot*amprot))
        - sqrt(((Wg*Wg+mdiff*mdiff-mppg*mppg) / (2*Wg)) *
                ((Wg*Wg+mdiff*mdiff-mppg*mppg) / (2*Wg)) - mdiff*mdiff)) *
        (sqrt(((Wg*Wg + amprot*amprot+q2g) / (2*Wg)) *
                ((Wg*Wg + amprot*amprot+q2g) / (2*Wg)) - amprot*amprot))
        - sqrt(((Wg*Wg+mdiff*mdiff-mppg*mppg) / (2*Wg)) *
                ((Wg*Wg+mdiff*mdiff-mppg*mppg) / (2*Wg)) - mdiff*mdiff));

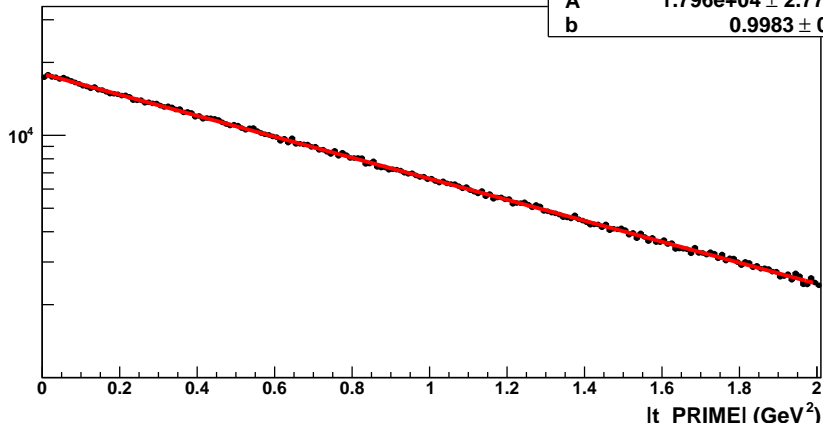
t_prime = t_gen - abs(t_min);

```

JPSI: MC gen. level ALL: b_pd slope fit

χ^2 / ndf	197.8 / 198
A	$1.796\text{e}+04 \pm 2.774\text{e}+01$
b	0.9983 ± 0.0018

events

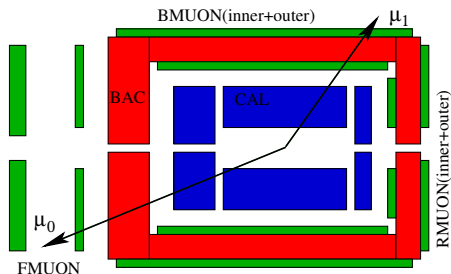


- p.diss t' -distribution is exponential
- reweight p.diss MC using t' !
- correct fitted b_{el} slope for detector level bias

Muon effic. corrections: TAG and PROBE

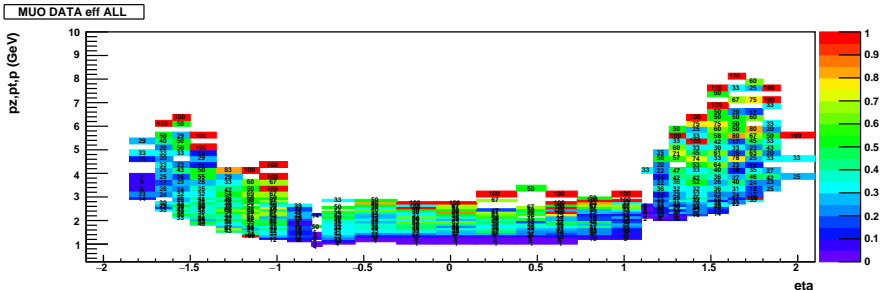
- TAG: “the triggering” muon
- PROBE: “the tested” muon
- effic in given $(p_{t_{eff}}^i, \eta^j)$ bin:

$$\epsilon = N_{PROBE}^i / N_{TAG}^i$$



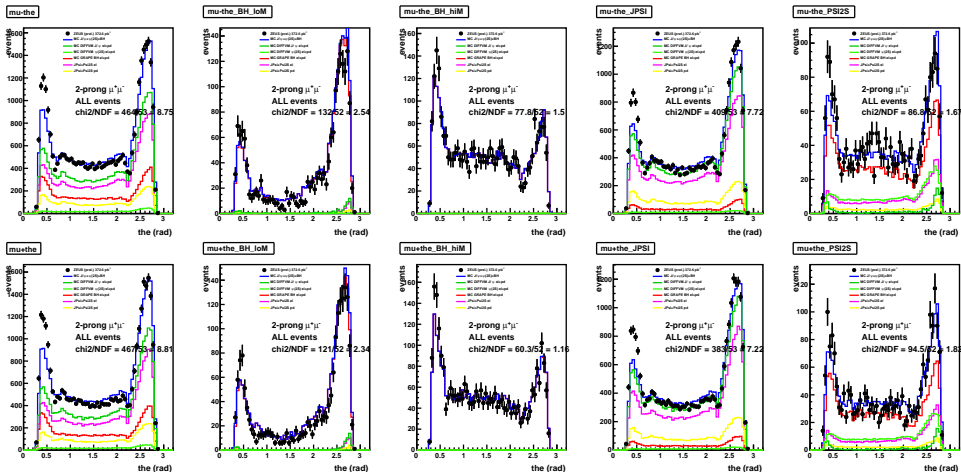
- **one step correction** for (FLT and SLT and TLT and off-line REC)
- separate maps for F/B/R/MUO, BAC and CAL (off-line only)
- evaluated for single muon in $(p_{t_{eff}}, \eta)$ bins, where as $p_{t_{eff}}$ is used:
(motivated by the CAL/BAC geometry and scaling of the muon path length)
 - p in Forecap
 - p_t in Barrel
 - p_z in Rearcap
- proper identification of **the triggering muon** is crucial
- → the DATA/MC ratio delivers the correction weight: $\epsilon_x = \frac{\epsilon_{DATA}}{\epsilon_{MC}}$

Muon correction maps: (p_z, p_t, p vs. η) - DATA



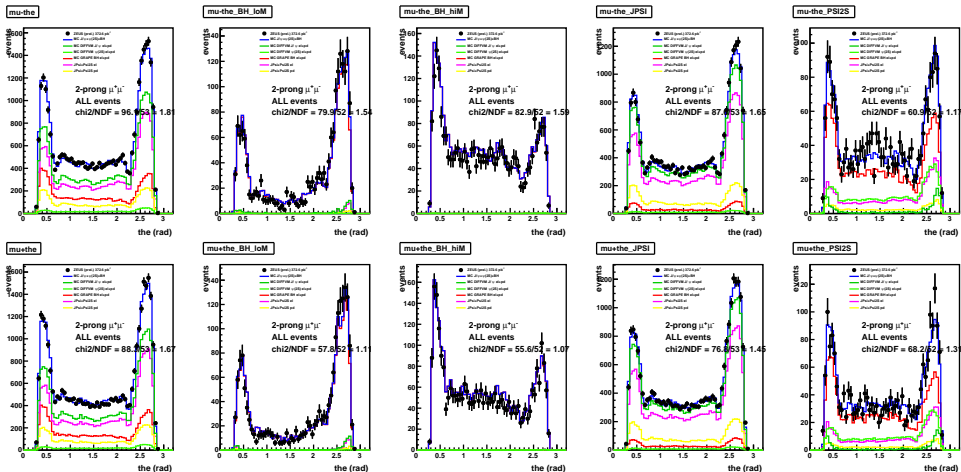
- muon tomography
- probability (%) to fire FLT-SLT-TLT-REC chain by single muon on ($p_z, p_t, p; \eta$) grid
- X-axis (along eta): Rear-MUO, Barrel-MUO, Forward-MUO detectors
- only events with $M(\mu^+, \mu^-) < 6$ GeV (ie. in the phase space range of di-muon mass fits)
- current choice for p_z, p_t, p grid: 100 MeV per bin ($p_{eff} < 3$ GeV), 250 MeV per bin ($p_{eff} > 3$ GeV)
- size of the grid is subject to systematics

2-prongs: before muon corrections: θ_{μ^\pm} in mass bins



- top: θ_{μ^-} , bottom: θ_{μ^+}
- ALL events, BH-loM, BH-hiM, J/ψ peak, ψ' peak
- standard (production version) of detector MC: "as it is"
- MC deficit in forward direction or MC over-efficient in rear

2-prongs: after muon corrections: θ_{μ^\pm} in mass bins



- ALL events, BH-loM, BH-hiM, J/ψ peak, ψ' peak
- good agreement in all mass windows (validation of muon corrections)
- (different processes, different μ^\pm angular/momentum distributions, different boost)