

Diffractive J/psi production at HERA

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Center For Frontiers In Nuclear Science ad-hoc Workshop:

*Target fragmentation and diffraction physics with novel processes:
ultraperipheral, electron-ion, and hadron collisions*

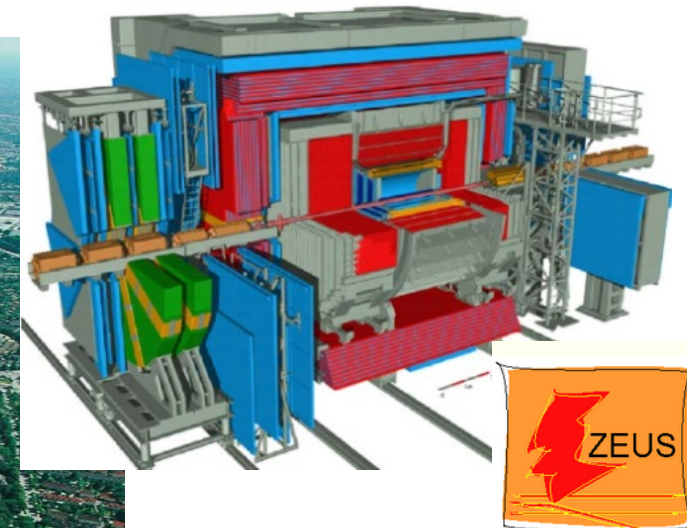
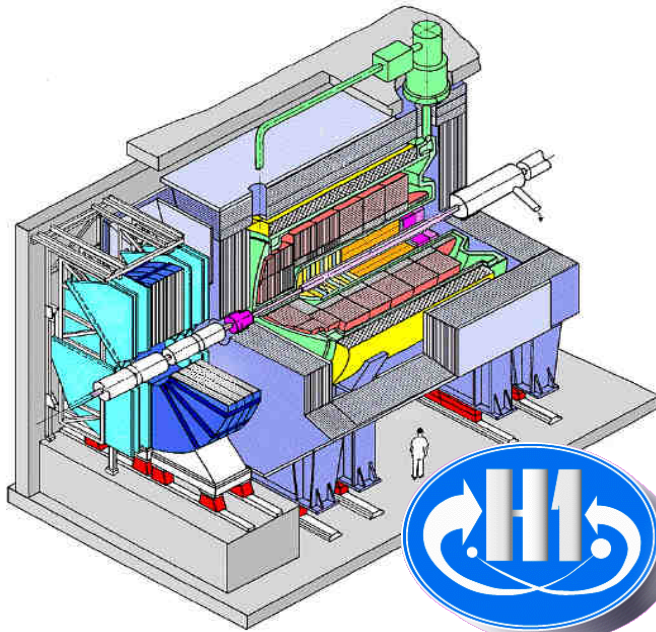
9-11 February 2022

<https://indico.bnl.gov/event/14009/>

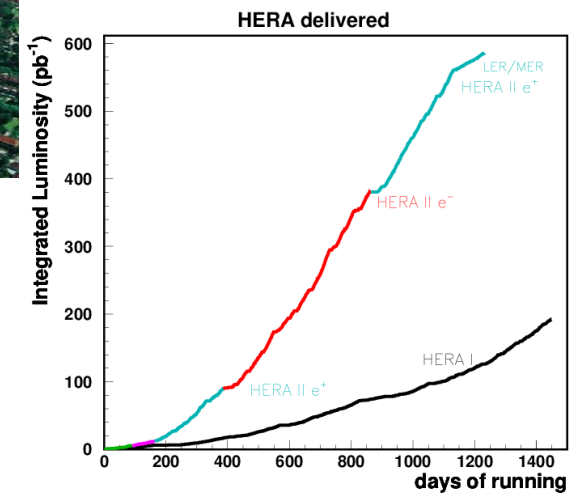


HERA electron/positron-proton collider (1992-2007)

$E_e = 27.6 \text{ GeV}$, $E_p = 920 \text{ GeV}$ (820, 460, 575 GeV)



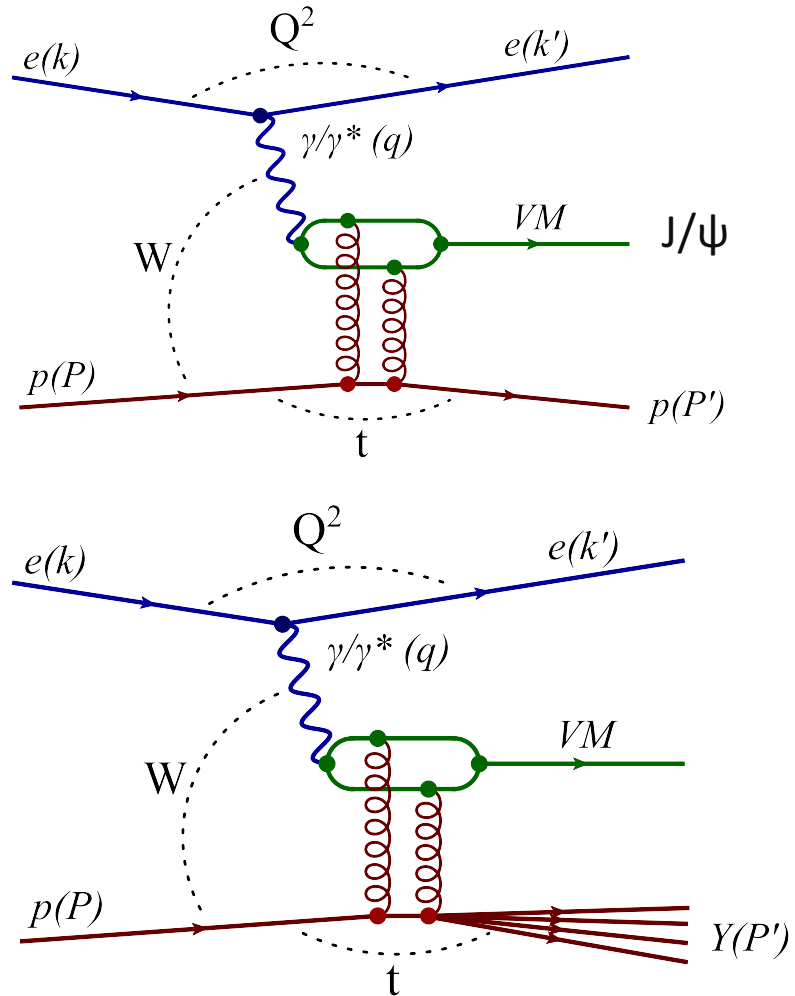
- > total luminosity $\sim 0.5 \text{ fb}^{-1}$ per experiment
- > Center-of-mass energies 318 GeV - 225 GeV



H1 and ZEUS paper on J/ψ , $\psi(2S)$ production

Photoproduction	<p>H1, Elastic and p-diss J/ψ in PHP Eur.Phys.J.C73 (2013) 2466</p> <p>H1, Elastic J/ψ in PHP and DIS Eur.Phys.J.C46 (2006) 585</p> <p>ZEUS, Exclusive PHP of J/ψ mesons Eur. Phys. J. C 24 (2002) 345</p> <p>H1, Elastic J/ψ and ρ^0 in PHP Phys.Lett.B483 (2000) 23</p> <p>H1, Elastic and Inelastic J/ψ in PHP Nucl.Phys.B472 (1996) 3</p>
DIS	<p>ZEUS, Exclusive J/ψ in DIS Nucl. Phys. B 695 (2004) 3</p> <p>H1, Charmonium Production in Deep Inelastic Scattering at HERA, Eur.Phys.J.C10 (1999) 373-393</p> <p>H1, Elastic ρ^0 and J/ψ at large Q^2 Nucl.Phys.B468 (1996) 3</p> <p>ZEUS, Exclusive ρ^0 and J/ψ in DIS Eur. Phys. J. C 6 (1999) 603</p>
High-t	<p>ZEUS, p-dissociative J/ψ in PHP at large t JHEP 05 (2010) 085</p> <p>H1, Diffractive PHP of J/ψ with large t Phys Lett B568 (2003) 205</p>
$\psi(2S)$	<p>$\psi' / J/\psi$ ratio [ZEUS prel-15-003]</p> <p>ZEUS, $R(\sigma\psi(2S)/\sigma J/\psi(1S))$ in DIS Nucl. Phys. B 909 (2016) 934</p> <p>H1, Diffractive PHP of $\psi(2S)$ Phys.Lett.B541 (2002) 251</p>

Kinematics



Exclusive production of Vector Mesons fully described by:

- Q^2 photon virtuality, kinematic regimes:

- photoproduction, $Q^2 \approx 0 \text{ GeV}^2$

- Deep Inelastic Scattering: $Q^2 > 1 \text{ GeV}^2$

- W photon-proton centre-of-mass energy

- $t = (p - p')^2$ – four momentum transfer squared at proton vertex

- x -Bjorken x -fraction of proton's momentum carried by struck quark $x \sim 1/W^2 \gamma p$

Kinematic variable fully reconstructed usually measuring **scattered electron** (in DIS) and vector meson decay products or final photons or jets. **Scattered p** detected with lower acceptance

Diffractive vector meson production studied at HERA as a **function of several scales**: Q^2 , MV , t over a wide range of $W\gamma p$.

Event topology/selection: exclusive di-muon /di-electron production

Very clean signature: J/ψ mesons identified via the decay channels: $\mu^+\mu^-$, e^+e^-

Photoproduction: trigger and selection driven by tracking + muon chambers / electrons in CAL

DIS: trigger given by scattered electron in CAL

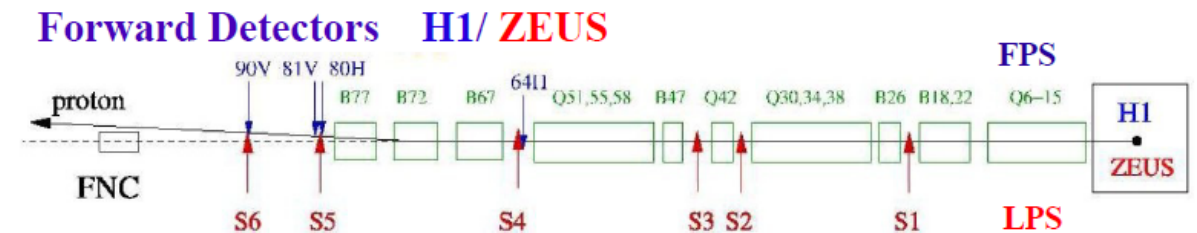
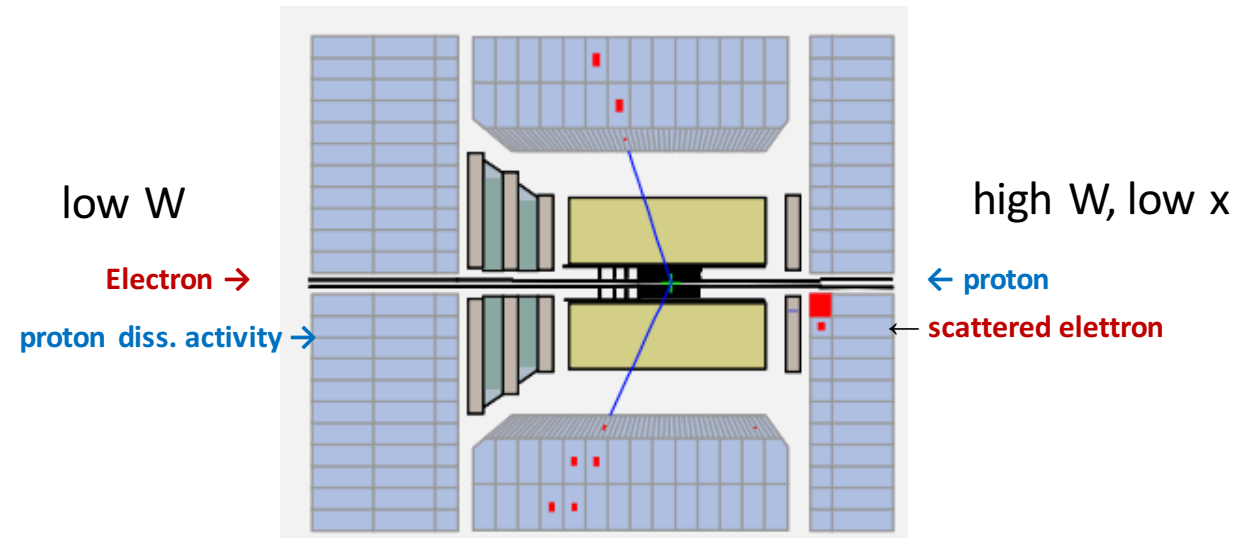
Selection: 2 tracks matched with calorimeter (MIPs or electron clusters), no additional CAL cluster.
For electron channel, 1 track + 1 CAL electron is also possible

Experimental challenge:

- trigger on exclusive events while rejecting non ep background: halo mu and cosmuics

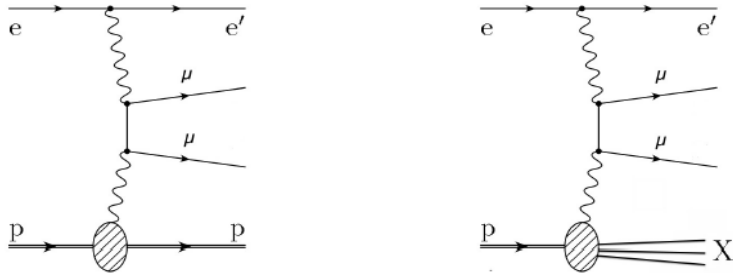
Analysis requires:

- a good understanding of trigger efficiencies
- forward detector to tag/study p-dissociative process



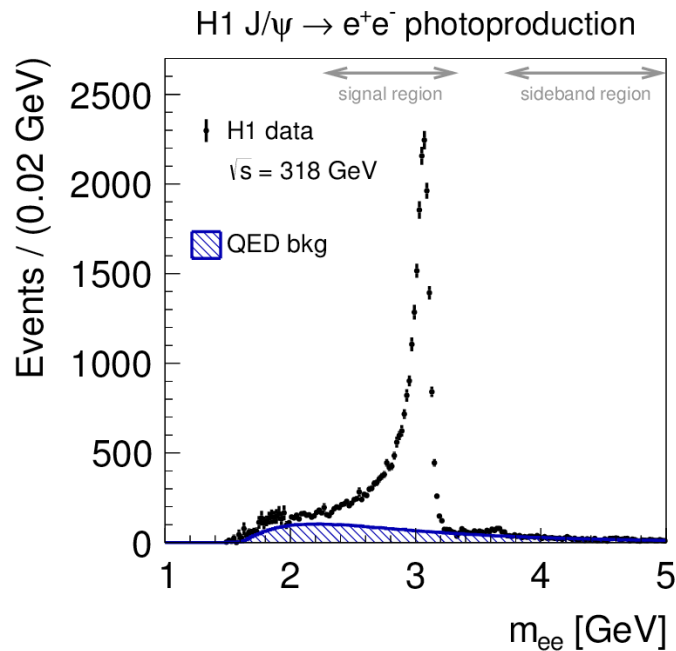
H1: FPS, Very Forward Proton Spectrometer, Forward Tagging System
ZEUS: LPS; Proton Remnant Tagger

Signal extraction from invariant mass distributions



$J/\psi \rightarrow \mu^+\mu^-, e^+e^-$

- only background source from QED non-resonant muon/electron pair production $\gamma\gamma \rightarrow l^+l^-$ (Bethe-Heitler process) subtracted by QED simulation, usually normalized in control region, as it contains p-diss contribution, not perfectly simulated

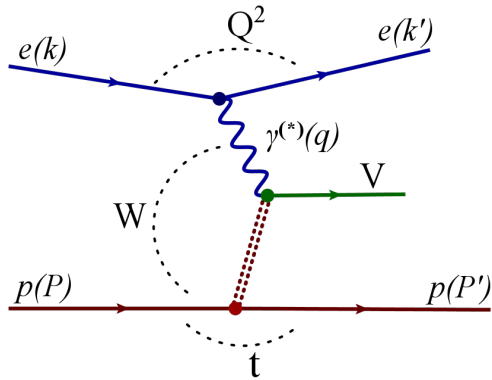


signal events:

- fitted
- or counted in signal region, procedure insensitive to low mass tail due to QED radiation losses and Bremsstrahlung, large for electrons and sizeable for muons; adequate in regions at low and high W where mass resolution is poor.

Theoretical models

- > **Regge phenomenology:** Vector Dominance Model + Soft Pomeron IP exchange



$$\alpha_P(t) = \alpha_0 + \alpha' t$$

$$\alpha_0 = 1.08, \alpha' = 0.25 \text{ GeV}^{-2} \quad (\text{DL})$$

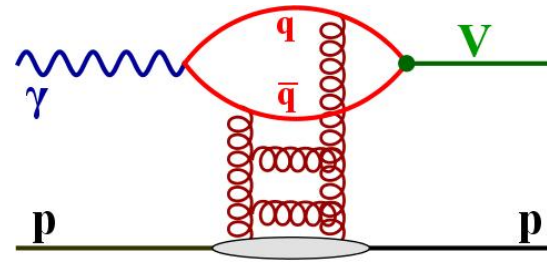
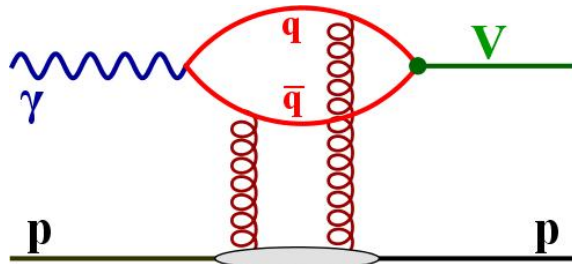
$$\frac{d\sigma}{dt} \propto e^{bt} \left(\frac{W_{YP}}{W_0} \right)^\delta \quad \delta = 4(\alpha_0 - 1)$$

$$b = b_0 + 4\alpha' \ln \left(\frac{W_{YP}}{W_0} \right)$$

Weak energy dependence

$$\sigma \propto W^{\delta_{YP}}$$

- > In the presence of a hard scale (M_{VM}, Q^2, t) calculations in **pQCD** are possible.
- > Exchange of a gluon ladder. **Fast increase of the cross section with energy**

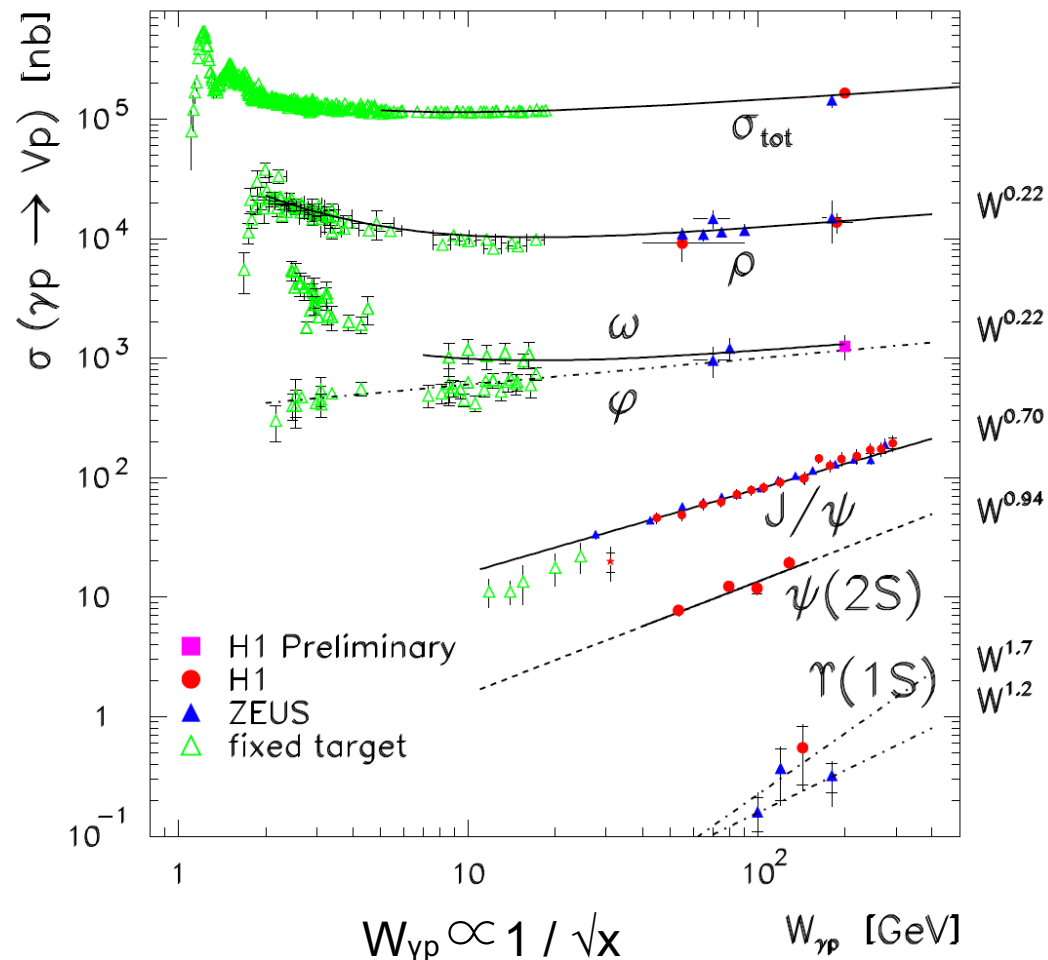


$$\begin{cases} \sigma \propto [x g(x, \mu^2)]^2 \\ x = \mu^2 / W^2 \\ \mu^2 \propto (Q^2 + M_V^2) \end{cases}$$

- > GPDs relevant here. Calculations performed now at NLO.

$W_{\gamma p}$ dependence of Elastic Vector Meson Photoproduction

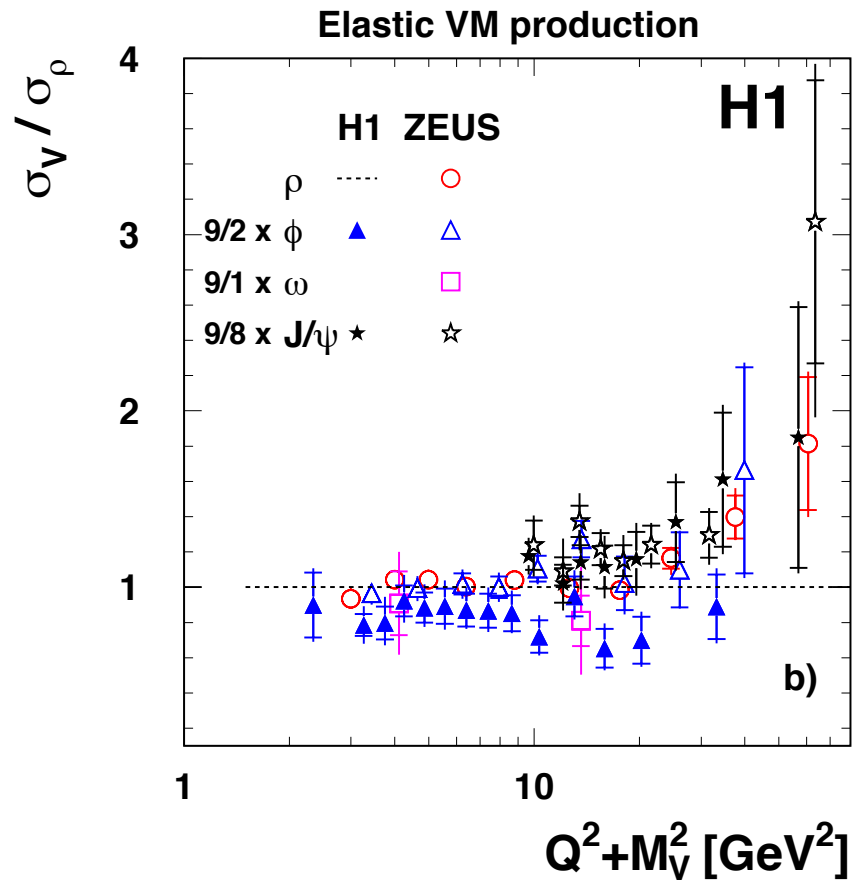
- › The cross section dependence on W can be parameterised as $\sigma \sim W^{\delta}_{\gamma p}$
- › Increasing VM mass, process gets harder



Low mass VM consistent with soft model, $\delta \sim 0.2$

Rapid rise of cross section with $W_{\gamma p}$ related to the increasing gluon density with decreasing of fractional momentum $\sigma \propto [x g(x, \mu)]^2$, with $x \sim 1/W^2_{\gamma p}$

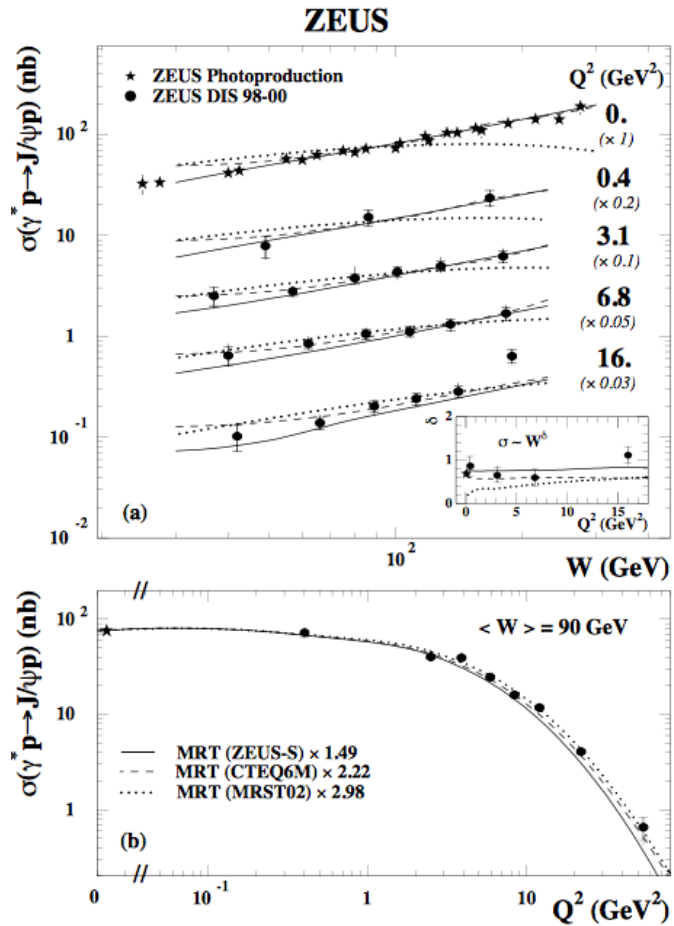
Universality of VM production



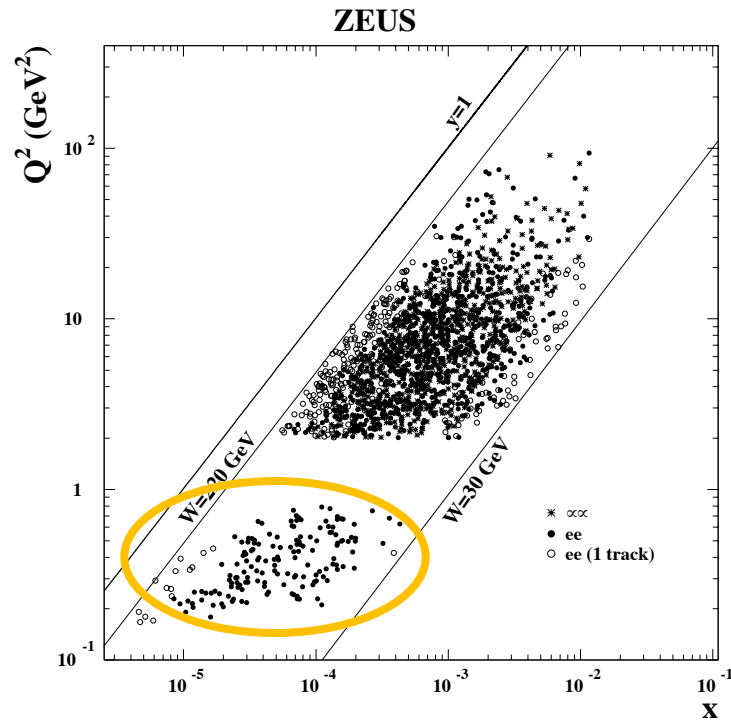
- Despite cross sections of light VM and J/psi differ by order of magnitude, they are close when corrected for factors accounting VM charge (rho : phi : J/psi = 9 : 2 : 8) and plotted vs scaling variable $\mu^2 = (Q^2 + M_V^2) / 4 \approx 3-5 \text{ GeV}^2$
- Ratios, scaled according to quark charge content, close to 1 (up to WF effects) once plotted vs $(Q^2 + M_V^2)$

This supports the dipole approach of VM production at high energy: cross sections are essentially determined by the dipole size

Wyp dependence of J/psi in photoproduction and DIS



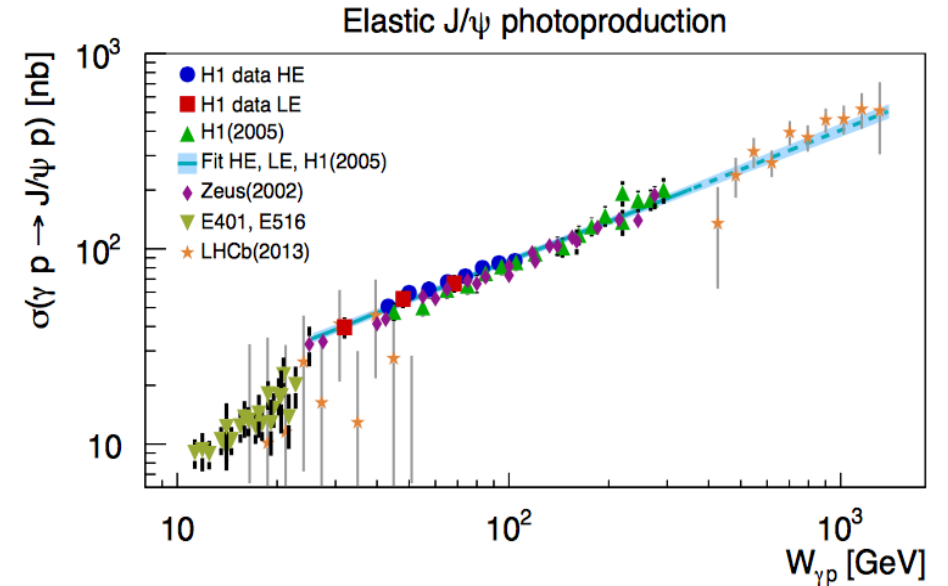
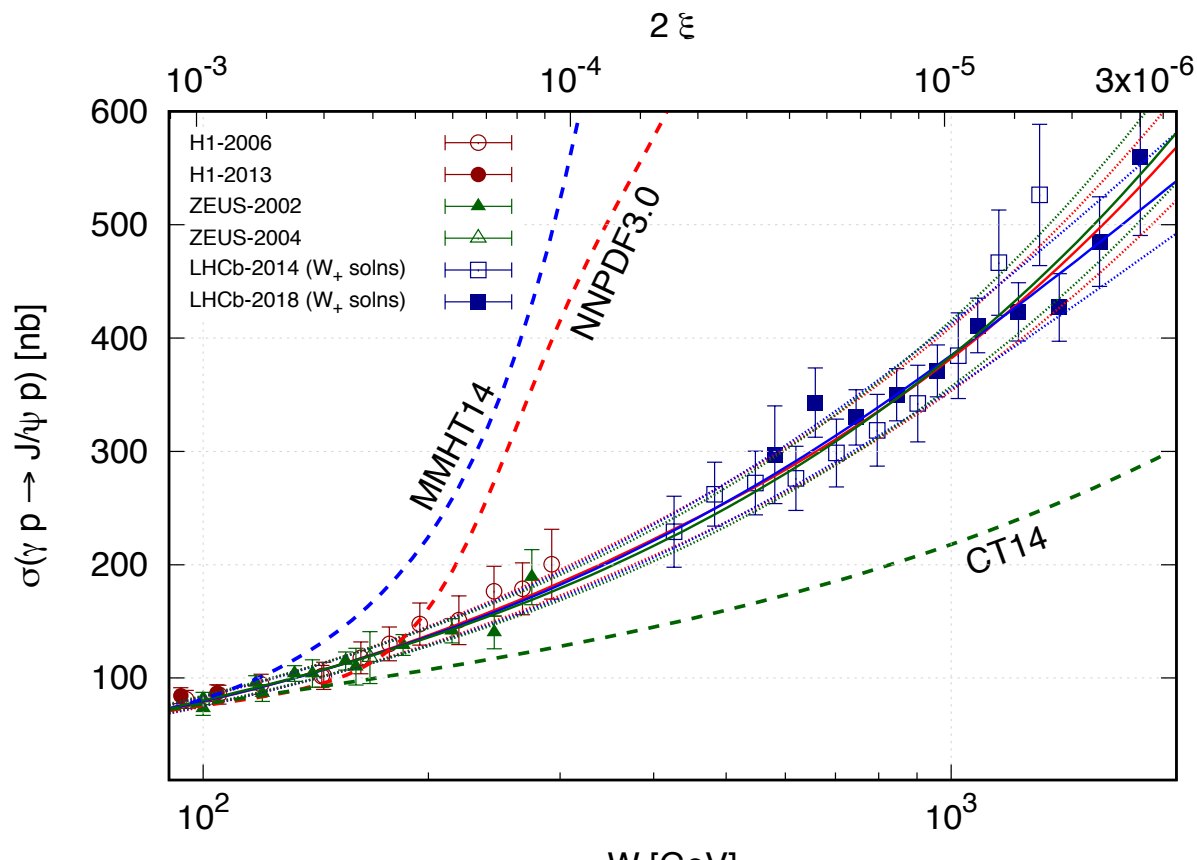
distribution of the J/psi events in the kinematic plane of Bjorken-x and Q^2



lowest values of Bjorken-x reached in photoproduction

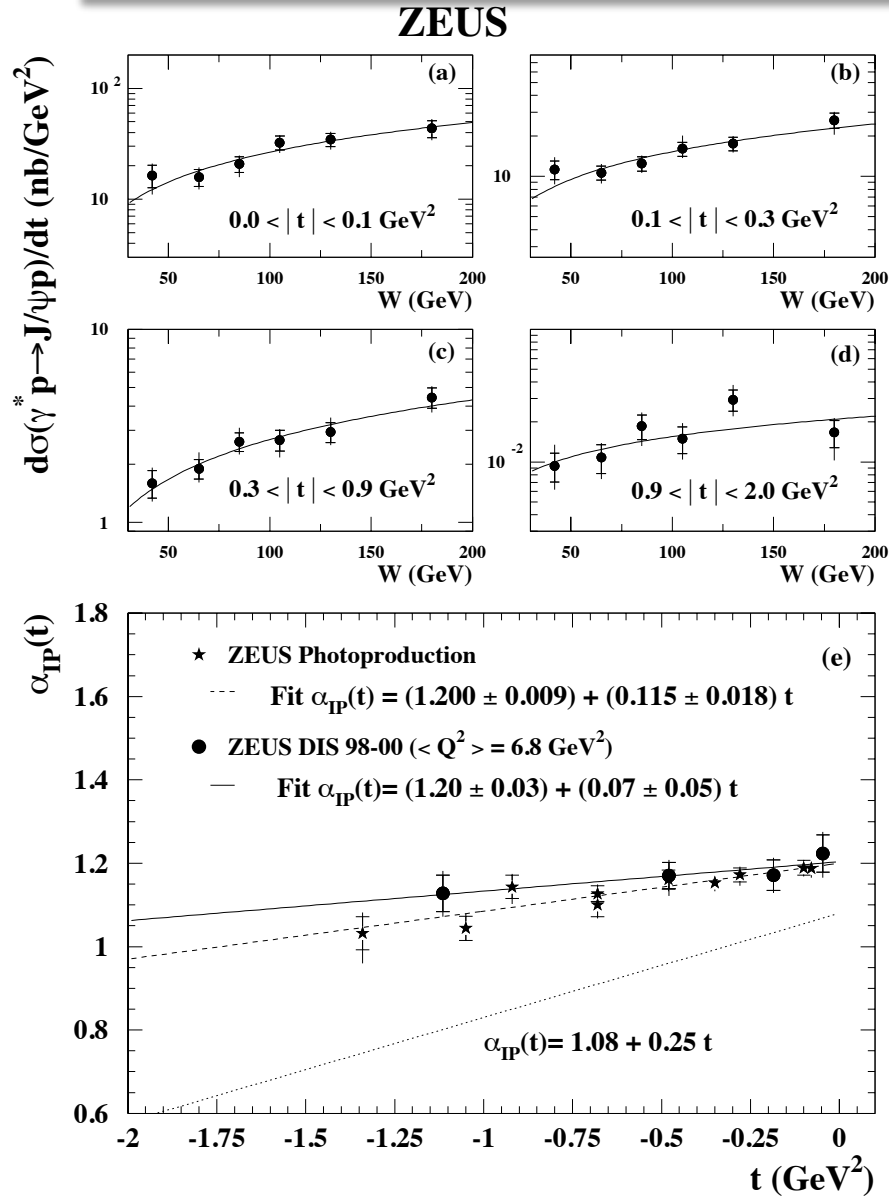
W_γ dependence of J/ψ photoproduction

Since 2015 “data J/ψ p can included in the global PDF fits to determine the gluon in the low x regime” (1507.06942 Jones, Martin, Ryskin, Teubner)



- A. Martin et al Phys. Rev. D 102, 114021 (2020)
- No hint of saturation observed in exclusive J/ψ data at the scale $\mu^2 = 2.4 \text{ GeV}^2$ and x down to 10^{-5} .

Regge models - Pomeron trajectory

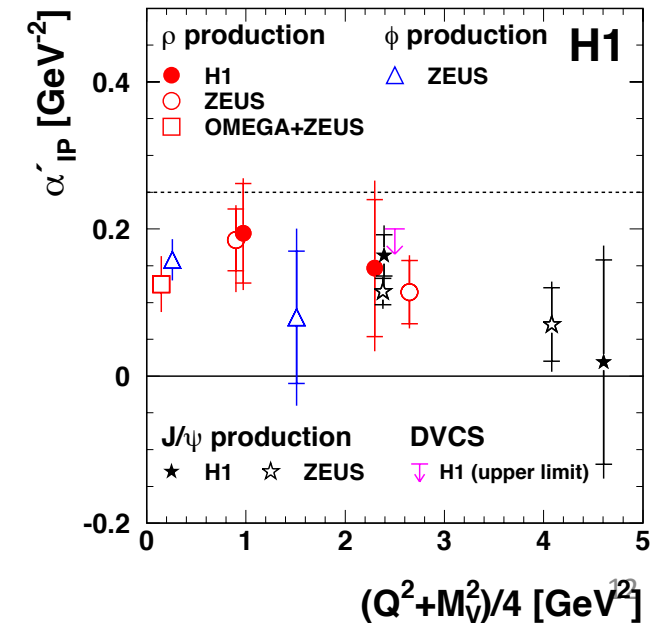
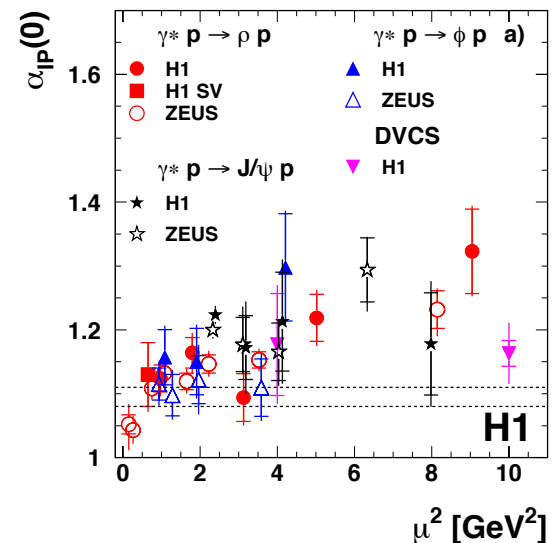


W dependance investigated using parametrisation inspired by Regge theory

Differential cross sections $d\sigma/dt$ as a function of W for fixed ranges of t as $W^4(\alpha_{IP}(t)-1)$

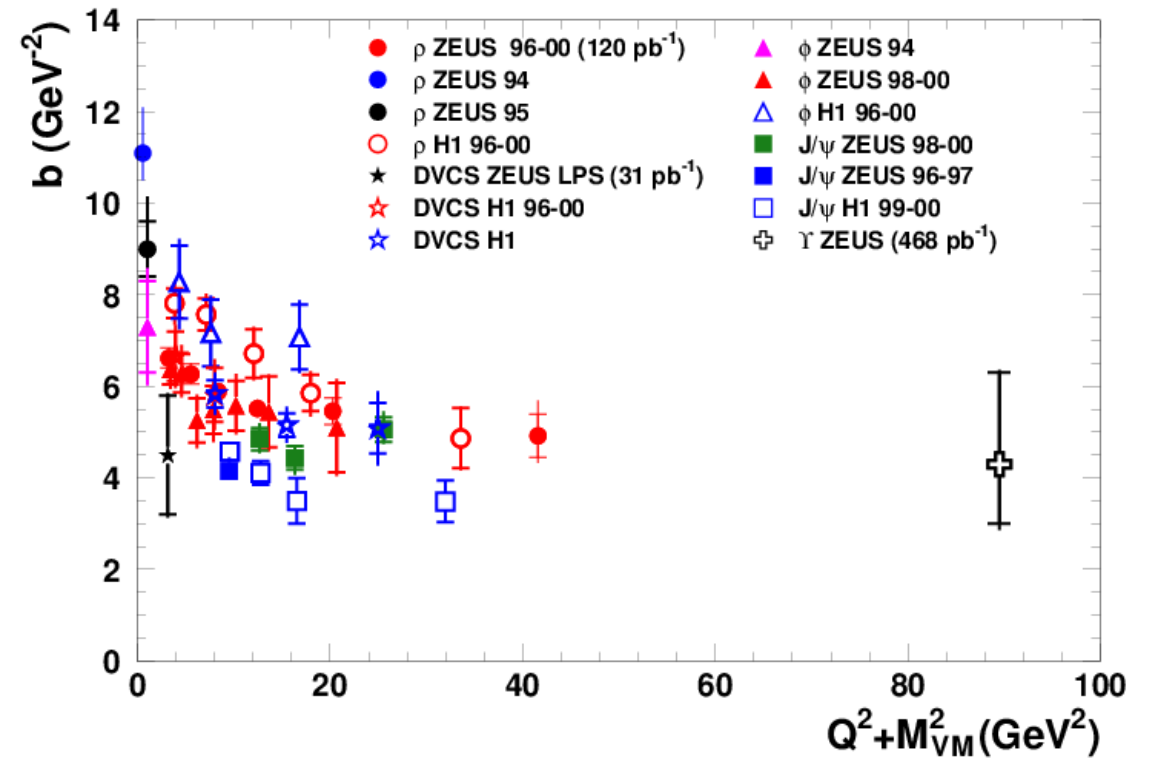
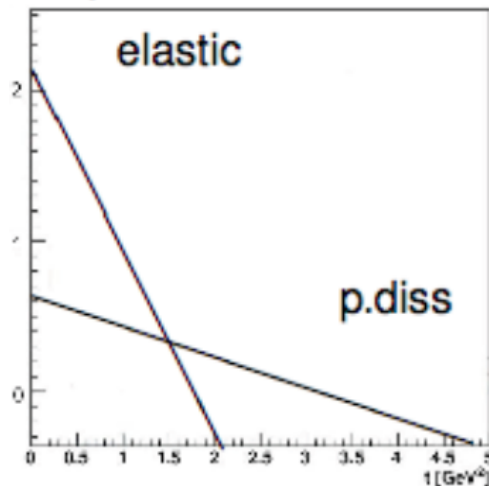
Pomeron trajectory associated to J/psi is not soft

In BFKL approach, α' is related to the average kt of gluons around the ladder in their random walk, expected to be small

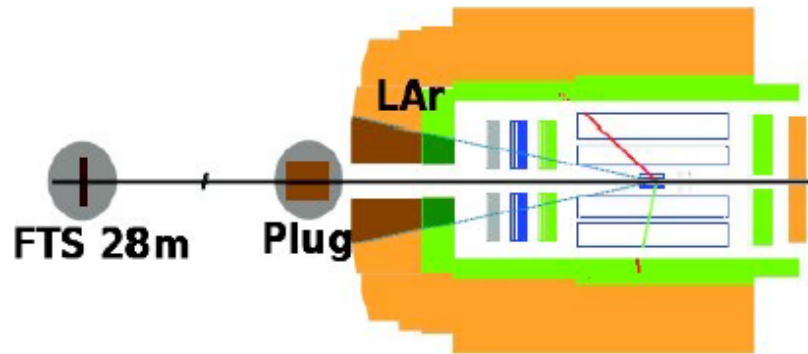


t-dependence of J/psi photoproduction

- ❖ t-dependence of cross section $d\sigma / dt \approx e^{bt}$
- In optical model approach b is related to the quadratic size of interacting objects:
- $b \sim (R_p^2 + R_{VM}^2)/4$
- in p-diss events proton breaks and size is smaller. As σ is similar for exclusive and p-diss. events, p-diss dominates at high-t

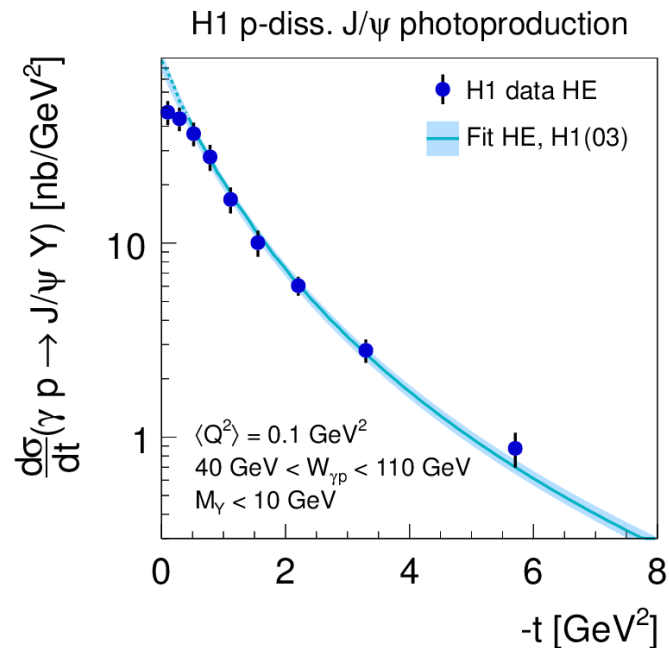
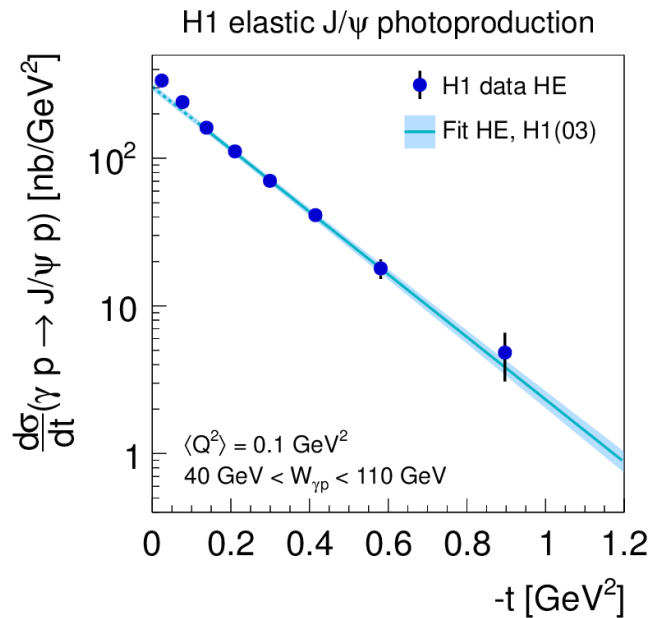


t-dependence of J/psi photoproduction - low t



H1:

- Use forward detectors (FTS, Plug, LAr) to tag proton dissociative process at low $|t|$.
- Use data from HERA low energy run to extend the range to lower $W_{\gamma p}$. Simultaneous measurement of elastic and proton-dissociative process.



Phenomenological fits:

$$d\sigma/dt = N_{el} e^{-b_{el}|t|}$$

$$d\sigma/dt = N_{pd} (1 + (b_{pd}/n)|t|)^{-n}$$

t-dependence of J/psi photoproduction – b slope

H1 PHP. pdiss $m_p < M_\gamma < 10 \text{ GeV}$

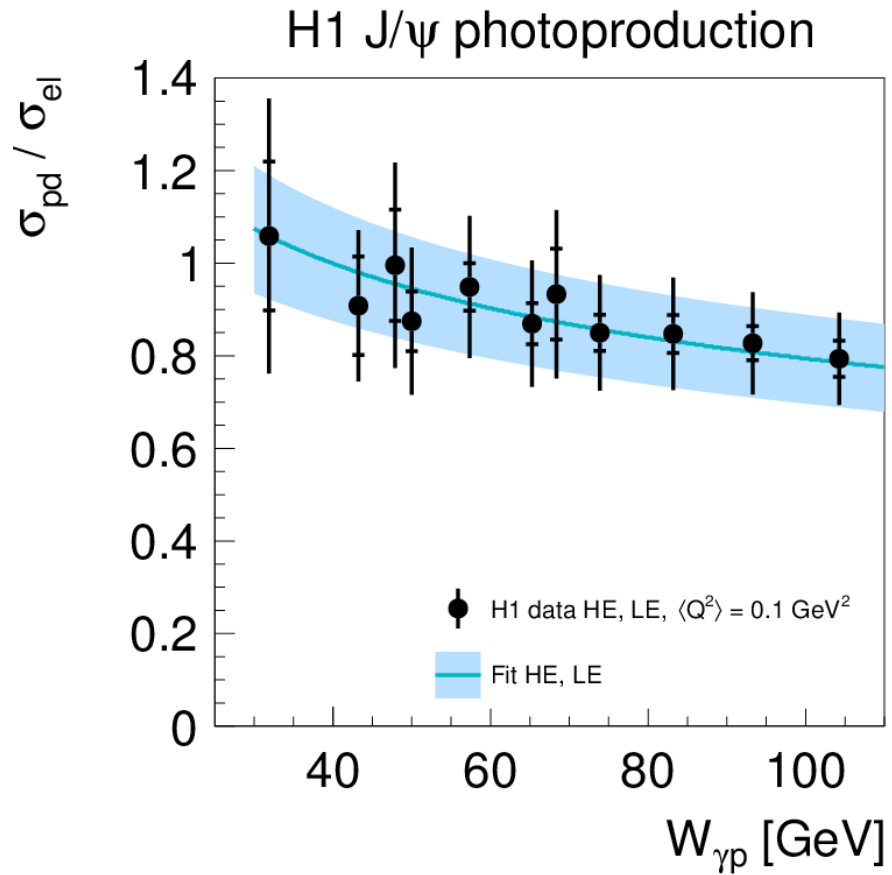
Data period	Process	Parameter	Fit value	Correlation
HE	$\gamma p \rightarrow J/\psi p$	b_{el}	$(4.88 \pm 0.15) \text{ GeV}^{-2}$	$\rho(b_{el}, N_{el}) = 0.50$ $\rho(b_{el}, b_{pd}) = 0.49$ $\rho(b_{el}, n) = -0.21$ $\rho(b_{el}, N_{pd}) = 0.68$
		N_{el}	$(305 \pm 17) \text{ nb/GeV}^2$	$\rho(N_{el}, b_{pd}) = 0.23$ $\rho(N_{el}, n) = -0.07$ $\rho(N_{el}, N_{pd}) = 0.46$
	$\gamma p \rightarrow J/\psi Y$	b_{pd}	$(1.79 \pm 0.12) \text{ GeV}^{-2}$	$\rho(b_{pd}, n) = -0.78$ $\rho(b_{pd}, N_{pd}) = 0.76$
		n	3.58 ± 0.15	$\rho(n, N_{pd}) = -0.46$
		N_{pd}	$(87 \pm 10) \text{ nb/GeV}^2$	
	LE	$\gamma p \rightarrow J/\psi p$	b_{el}	$(4.3 \pm 0.2) \text{ GeV}^{-2}$
N_{el}			$(213 \pm 18) \text{ nb/GeV}^2$	$\rho(N_{el}, b_{pd}) = -0.24$ $\rho(N_{el}, N_{pd}) = -0.10$
$\gamma p \rightarrow J/\psi Y$		b_{pd}	$(1.6 \pm 0.2) \text{ GeV}^{-2}$	$\rho(b_{pd}, N_{pd}) = 0.53$
		n	3.58 (fixed value)	
		N_{pd}	$(62 \pm 12) \text{ nb/GeV}^2$	

- ZEUS Eur. Phys. J. C 24 (2002)
- Finite Mass Sum Rule (from Mueller generalized optical theorem) applied to proton diffraction provide constraints on M_γ spectrum of proton diffractive state, which MC have to respect. It also explain why t-distribution is steeper for resonant part than for the continuum part of proton diffraction
- MC simulation: *In the baryon resonance region, at low M_γ , a resonant component with slope $b = 6.5 \text{ GeV}^{-2}$ was considered. A second component due to non-resonant proton dissociation with slope $b = 0.65 \text{ GeV}^{-2}$ was added. The two components were constrained to satisfy the first moment of the finite-mass sum rule.* Ref: G. Alberi and G. Goggi, Phys. Rev. 74, 1 (1980)

ZEUS; proton–diss, continuum region for for $3.5 < M_\gamma < 30 \text{ GeV}$ and $p_T^2 < 10 \text{ GeV}^2$

- p-diss $b = 0.65 \pm 0.10 \text{ GeV}^{-2}$
- $\sigma \simeq 1/M_\gamma^{\text{beta}}$ with $\text{beta} = 2.6 \pm 0.3$
- elastic $b = 4.15 + 0.05 \text{ (stat)} + 0.30 - 0.18 \text{ (syst)} \text{ GeV}^{-2}$

W-dependence of J/psi photoproduction – p. diss.



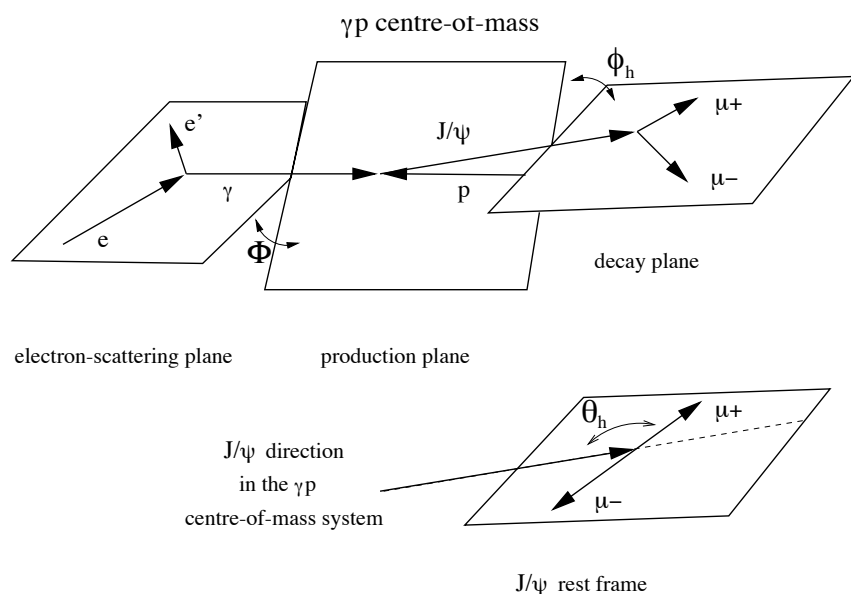
Process	Parameter	Fit value	Correlation
$\gamma p \rightarrow J/\psi p$	δ_{el}	0.67 ± 0.03	$\rho(\delta_{el}, N_{el}) = -0.08$ $\rho(\delta_{el}, \delta_{pd}) = 0.01$ $\rho(\delta_{el}, N_{pd}) = 0.09$
	N_{el}	$81 \pm 3 \text{ nb}$	$\rho(N_{el}, \delta_{pd}) = -0.27$ $\rho(N_{el}, N_{pd}) = -0.18$
$\gamma p \rightarrow J/\psi Y$	δ_{pd}	0.42 ± 0.05	$\rho(\delta_{pd}, N_{pd}) = 0.09$
	N_{pd}	$66 \pm 7 \text{ nb}$	

Ratio of elastic and proton dissociative cross is approximately equal to 1

A slight dependence of this ratio as a function of $W_{\gamma p}$ is observed

Decay angular distribution, helicity frame

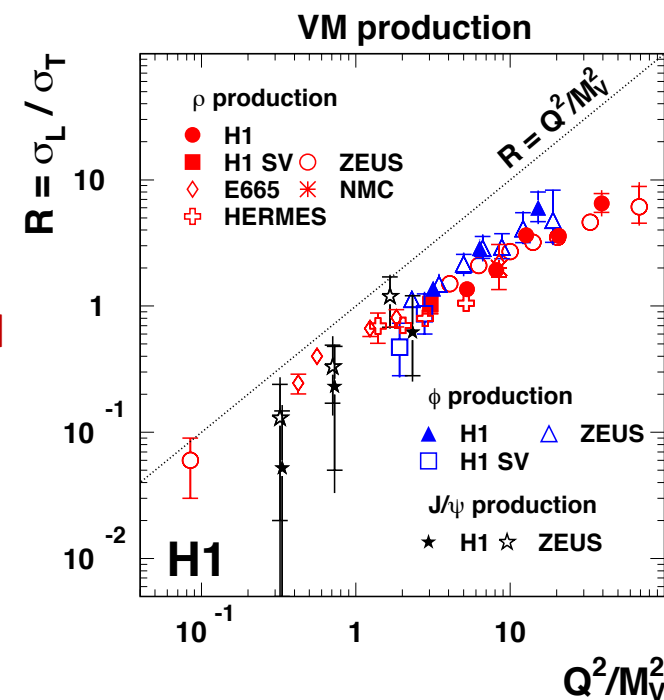
Measurement of VM production and decay angles give access to spin density matrix elements
 Angular distr. \rightarrow spin density matrix elements r_{ji}^{kl} , \rightarrow helicity amplitudes $T\lambda_{VM}\lambda_\gamma$



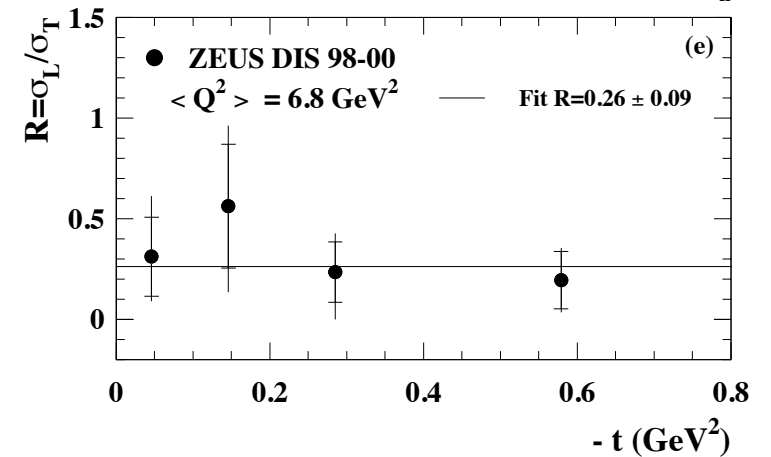
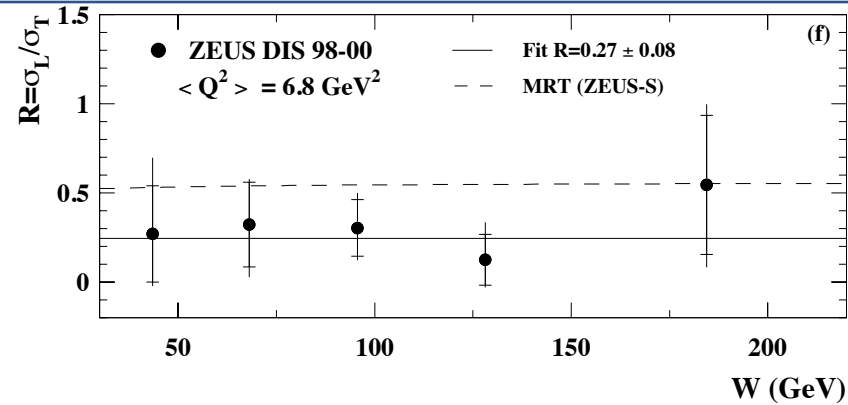
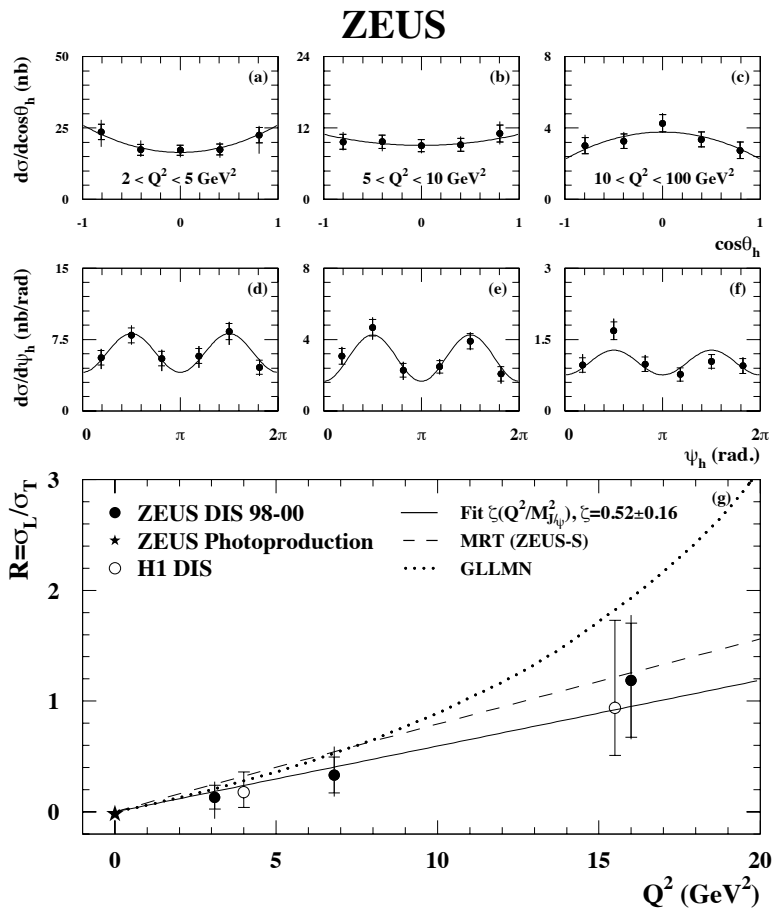
S-channel helicity conservation: the outgoing VM retains the γ helicity

In the accessible Q^2 range, J/ψ production is almost transverse, while for light VM the longitudinal amplitude dominates

θ_h, ϕ_h angles of decay muons in the meson rest frame
 Φ angle between scattering and production plane



Ratio $R = \sigma_L/\sigma_T$ as a function of Q^2 , W and t



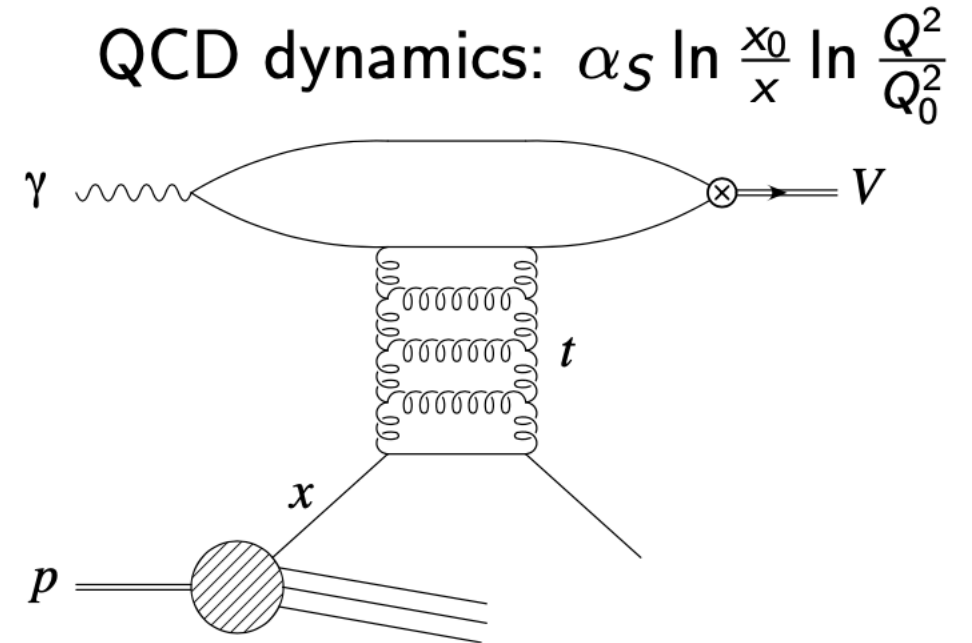
In the accessible Q^2 range, J/psi production is almost transverse
 No dependence from W or t

pQCD: during the interaction, the orbital angular momentum of $q\bar{q}$ can be modified due to the transfer of momentum of the gluons; \Rightarrow the helicity of the outgoing VM differs from the one of the γ , helicity flip between photon and meson is possible

Photoproduction J/ψ at high t – test BFKL dynamics

large values of $|t|$ provide a hard scale for $|t| > 1 \text{ GeV}^2$, process dominated by p-diss production

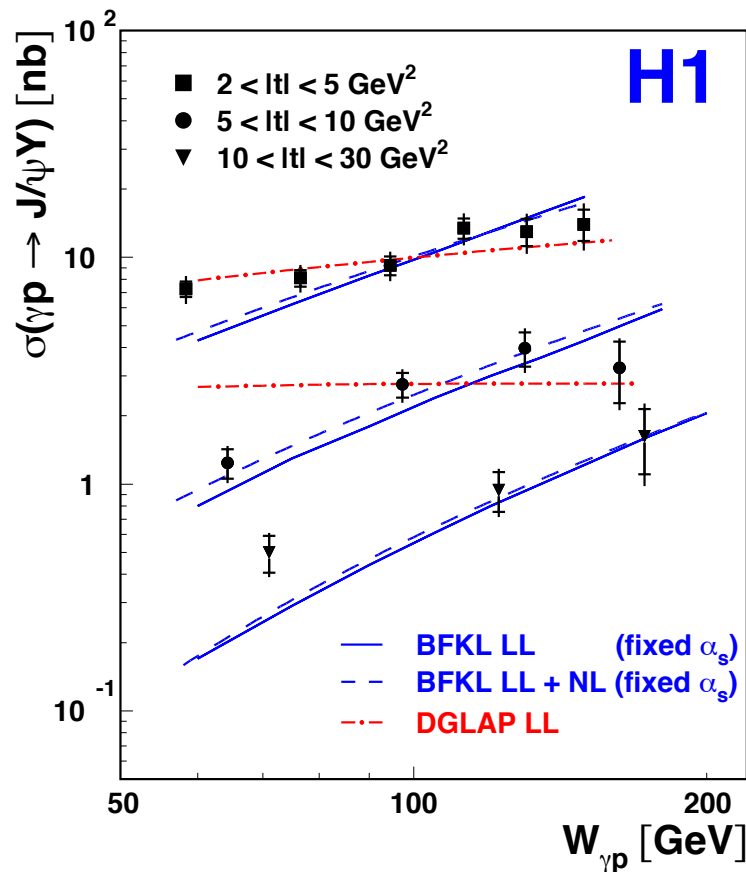
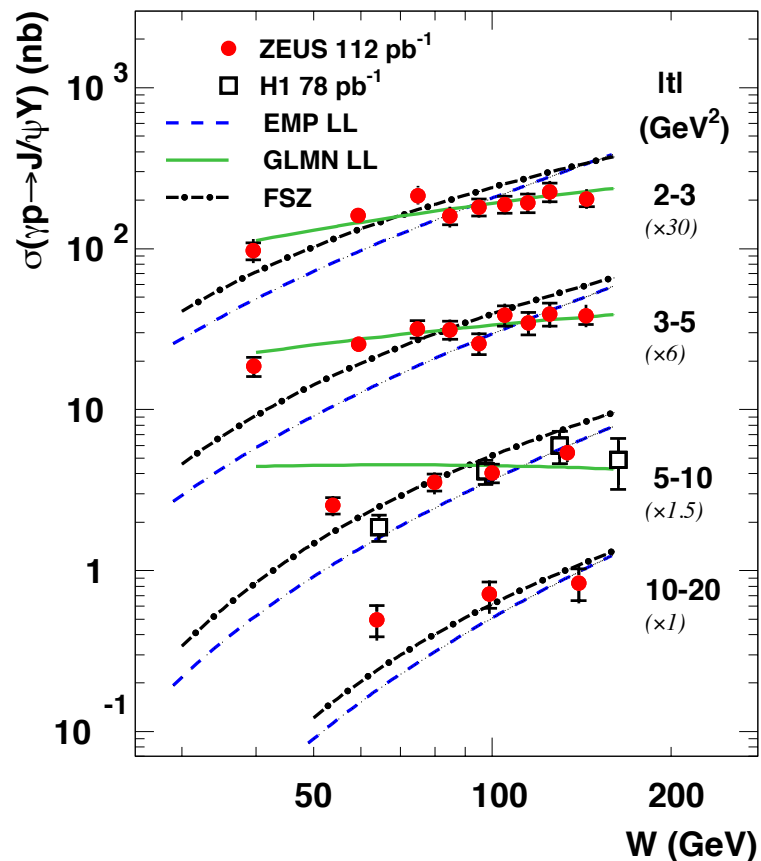
Hard scale at both end of the exchanged gluon ladder, no strong kt ordering is expected (expected in DGLAP), as in typical BFKL evolution



HERA I data, 1996-2000, 112 pb^{-1} of luminosity
 $30 < W < 160 \text{ GeV}$, $2 < |t| < 20 \text{ GeV}^2$
 $z > 0.95$, inelasticity cut

J/ $\psi(1S)$ at large $|t|$, W dependence

ZEUS



H1 and ZEUS results in good agreement in the common kinematic region

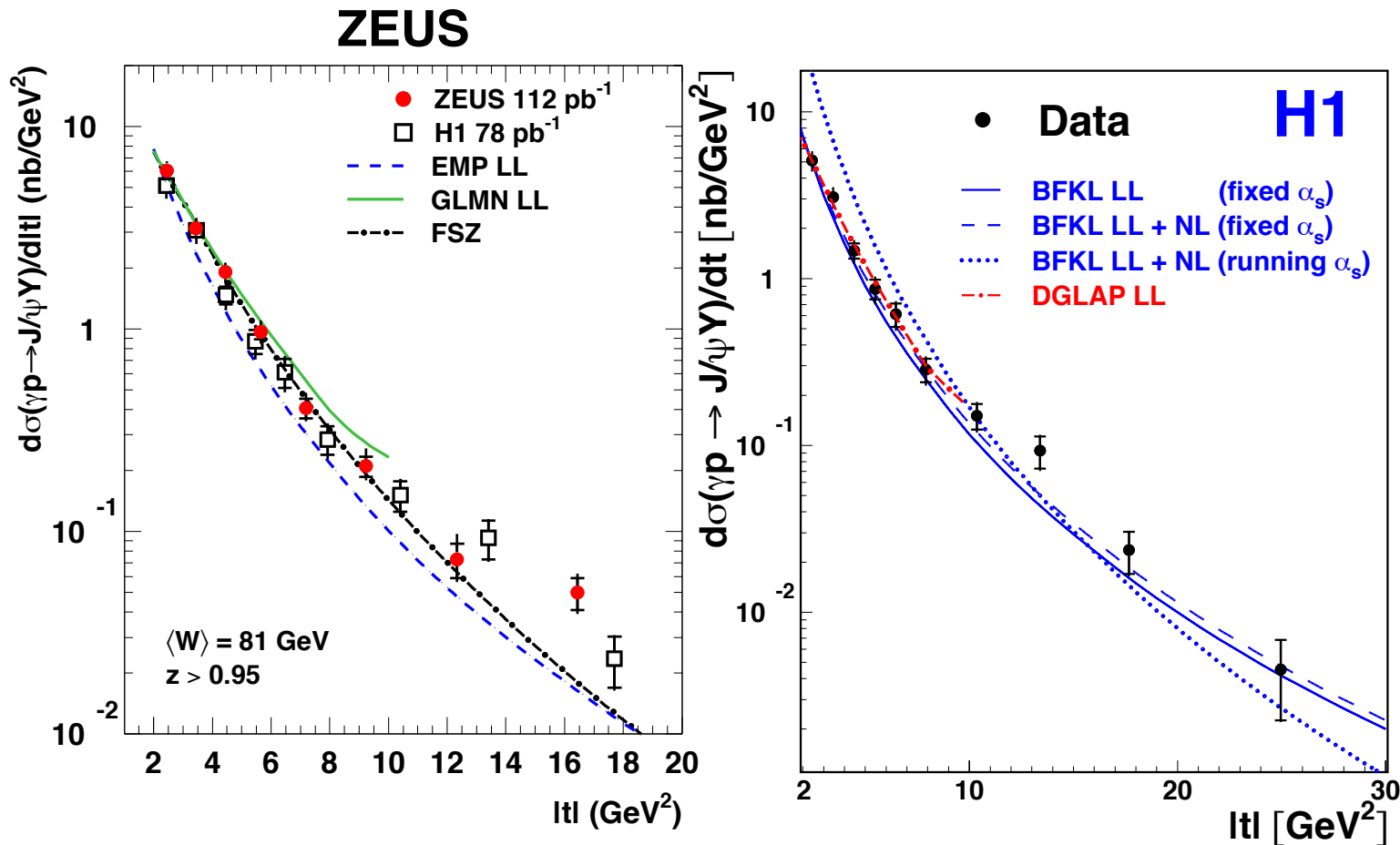
Data rise with W for all t region

BFKL (EMP) predictions too steep

DGLAP (GLMN) approach fails to describe σ rise at low x

FKS: increase of σ due to gluon distribution in the proton

J/ψ(1S) at large |t|; t-dependence



Data cannot be described by a single exponential fit $d\sigma/dt \approx e^{bt}$ neither by a single power $d\sigma/dt \approx |t|^n$

Models predict n dependent on t

Models able to describe data but not for the full t (x) range

EMP (BFKL) below data

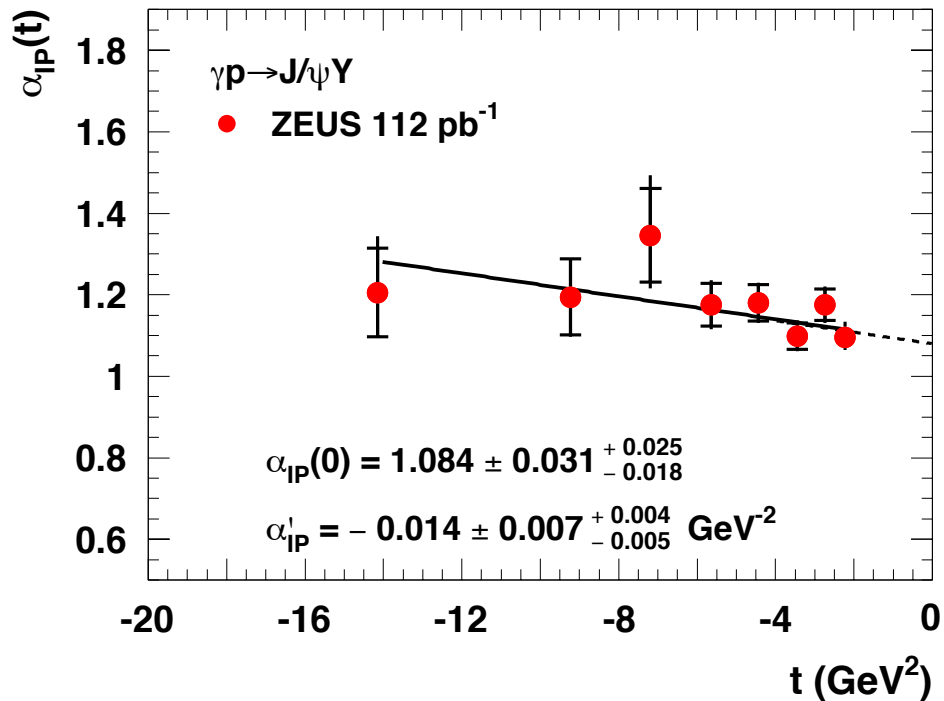
GLMN (DGLAP) validity range of $|t| < M_{J/\psi}$

FSZ describe data up to 12 GeV²

J/ψ(1S) at large |t|, W dependence as a function of t

Effective Pomeron trajectory $d\sigma/dt = F(t) \cdot W^{4(\alpha_P(t)-1)}$

ZEUS



Pomeron intercept consistent with soft $\alpha_{PI}(0) = 1.0808$

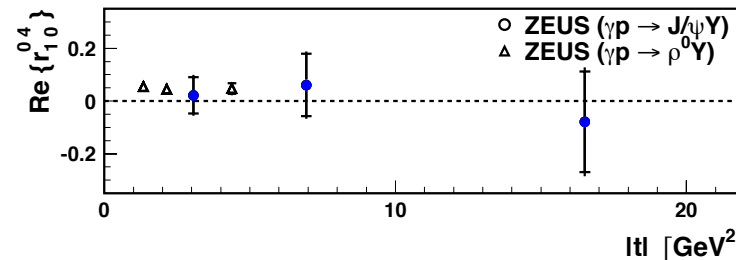
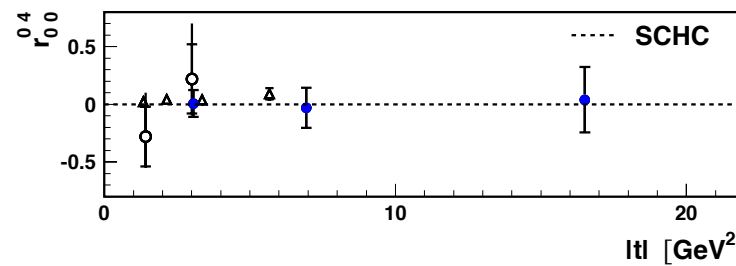
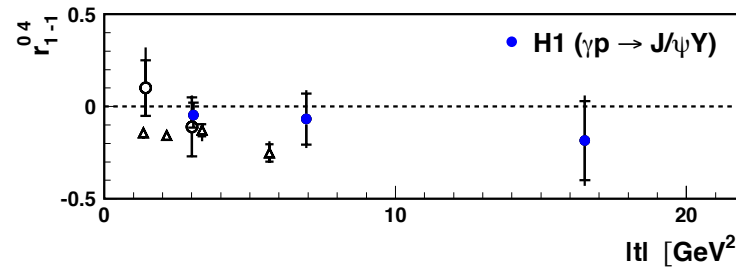
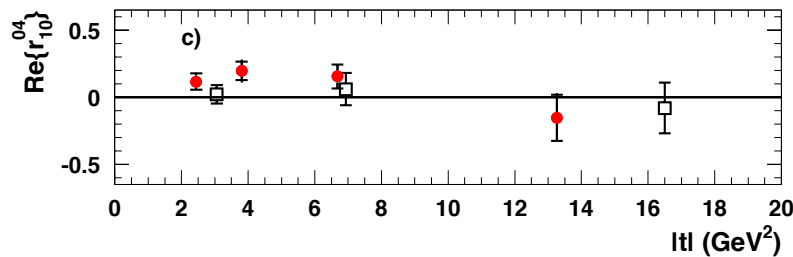
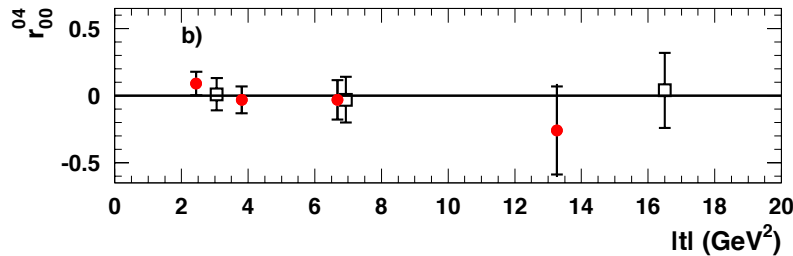
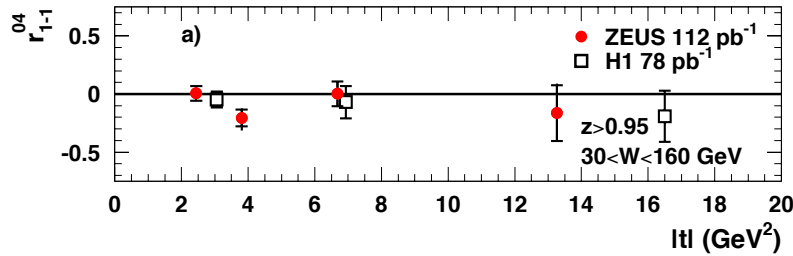
Pomeron slope consistent with BFKL Pomeron

H1 $\alpha_{PI}(0) = 1.167 \pm 0.048(\text{stat}) \pm 0.024(\text{syst})$

H1 $\alpha' = -0.0135 \pm 0.0074(\text{stat}) \pm 0.0051(\text{syst})$

Helicity spin density matrix elements as a function of $|t|$

ZEUS



θ_h, ϕ_h angles

angles of decay muons in the meson rest frame estimated in different t bins
Spin density matrix elements are extracted from fit to the angular distributions

r_{1-1}^{04} is related to interference between non-flip and double flip amplitude

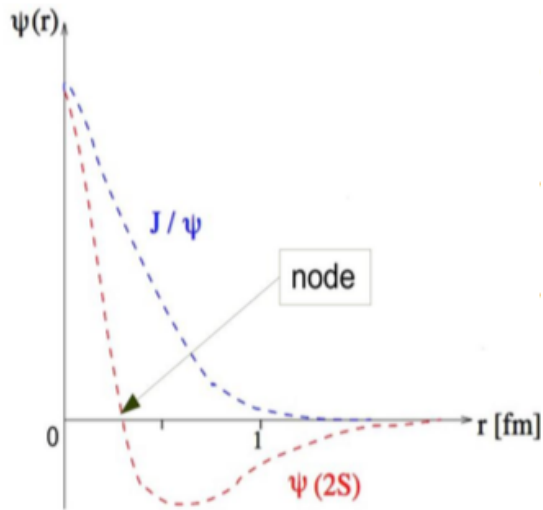
r_{00}^{04} represents the probability that $J\psi$ has 0 helicity

$\text{Re}(r_{00}^{04})$ is proportional to the single flip amplitude 10

These spin density elements expected to be 0 in SCHC

$\psi(2S)$ in photoproduction and DIS

$\psi(2S)$ and $J/\psi(1S)$ have same quark content but different wave function



$$\text{Ratio } R = \sigma_{\gamma p \rightarrow \psi(2S)p} / \sigma_{\gamma p \rightarrow J/\psi(1S)p}$$

- sensitive to radial wave function of charmonium
- provides insight into the dynamics of the hard process

HIKT, Hufner et al.: dipole model, dipole-proton constrained by inclusive DIS data

AR, Armesto and Rezaeian: impact parameter dependent CGC and IP-Sat model
KMW, Kowalski Motyka Watt: QCD description and universality of quarkonia production
FFJS, Fazio et al.: two component Pomeron model

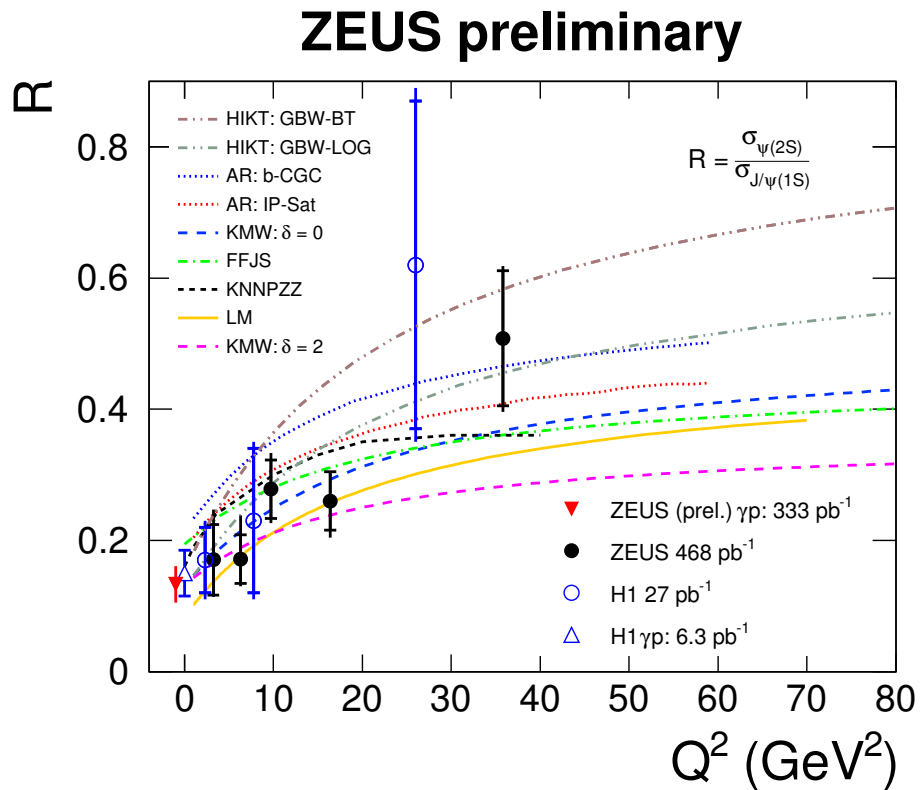
KNNPZZ, Nemchik et al.: color-dipole cross section derived from BFKL generalised eq.
LM, Lappi and Mäntysaari: dipole picture in IP-Sat model

$$\langle r^2 \psi(2S) \rangle \sim \langle 2 r^2 J/\psi(1S) \rangle$$

$\psi(2S)$ has a radial node at ~ 0.4 fm

models predicts $\sigma_{\gamma p \rightarrow \psi(2S)p}$ suppressed w.r.t. $\sigma_{\gamma p \rightarrow J/\psi(1S)p}$

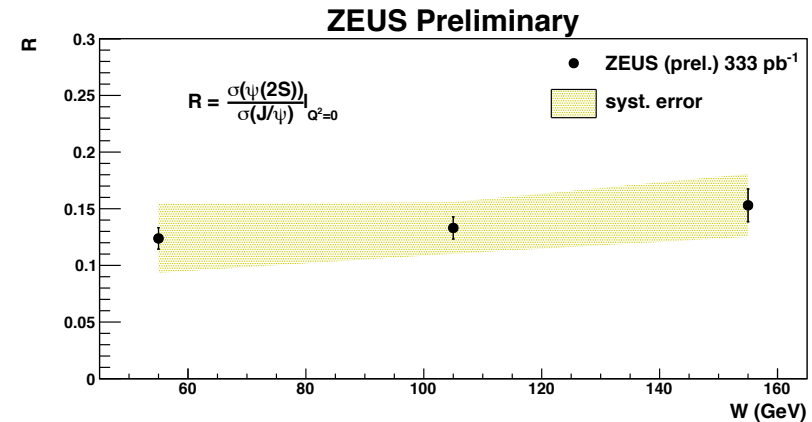
$$R = \sigma_{\gamma p \rightarrow \psi(2S)p} / \sigma_{\gamma p \rightarrow J/\psi(1S)} \text{ vs } Q^2, t, W$$



$$R = \sigma_{\gamma p \rightarrow \psi(2S)p} / \sigma_{\gamma p \rightarrow J/\psi(1S)p}$$

ZEUS prel: highest precision at $Q^2 = 0 \text{ GeV}^2$
 $R = 0.1332 \pm 0.0065(\text{stat.}) \pm 0.0270(\text{syst.})$

- Increase in Q^2
- no dependence in t
- no or moderated dependence on W



pQCD model calculations predict $R \sim 0.17$ in PHP rising with Q^2 reaching plateau at $Q^2 \gg M^2 \psi$

Summary

- Exclusive $J/\psi(1S)$ production extremely interesting for theory, specially pQCD
- It is a clean experimental process, but precise measurement requires a strategy for trigger and dedicated detectors to tag proton-dissociation

Waiting for the discussion!

H1 and ZEUS paper on VM production

H1 Topic	Journal	ZEUS Topic	Journal
<p>Exclusive $\pi^+\pi^-$ and ρ^0 in PHP</p> <p>Exclusive ρ^0 with Leading n in PHP</p> <p>Elastic and p-diss J/ψ in PHP</p> <p>Diffractive ρ^0 and ϕ in DIS</p> <p>Diffractive PHP of ρ^0 with large t</p> <p>Elastic J/ψ in PHP and DIS</p> <p>Diffractive PHP of J/ψ with large t</p> <p>Diffractive PHP of $\psi(2S)$</p> <p>Helicity structure of ρ^0 in DIS</p> <p>Elastic ϕ in DIS</p> <p>Elastic J/ψ and Υ in PHP</p> <p>Elastic ρ^0 in DIS</p> <p>Quasi-elastic ($z > 0.95$) $\psi(2S)$ in PHP</p> <p>P-diss. ρ^0 and Elastic ϕ in DIS</p> <p>Elastic and Inelastic J/ψ in PHP</p> <p>Elastic ρ^0 and J/ψ at large Q^2</p> <p>Elastic Rho0 in PHP</p>	<p>Eur.Phys.J.C80 (2020), 1189</p> <p>Eur.Phys.J.C76 (2016) 1, 41</p> <p>Eur.Phys.J.C73 (2013) 2466</p> <p>JHEP05 (2010) 032</p> <p>Phys.Lett.B 638 (2006) 422</p> <p>Eur.Phys.J.C46 (2006) 585</p> <p>Phys Lett B568 (2003) 205</p> <p>Phys.Lett.B541 (2002) 251</p> <p>Phys.Lett.B539 (2002) 25</p> <p>Phys.Lett.B483 (2000) 360</p> <p>Phys.Lett.B483 (2000) 23</p> <p>Eur.Phys.J.C13 (2000) 371</p> <p>Phys.Lett.B421 (1998) 385</p> <p>Z.Phys.C75 (1997) 607</p> <p>Nucl.Phys.B472 (1996) 3</p> <p>Nucl.Phys.B468 (1996) 3</p> <p>Nucl.Phys.B463 (1996) 3</p>	<p>$R(\sigma_{\psi(2S)}/\sigma_{J/\psi(1S)})$ in DIS</p> <p>Exclusive Electroproduction of 2π</p> <p>$\Upsilon(1S)$ in PHP (t-dependence)</p> <p>P-dissociative J/ψ in PHP at large t</p> <p>Exclusive PHP of Υ Mesons</p> <p>Exclusive ρ^0 in DIS</p> <p>Exclusive ϕ in DIS</p> <p>Exclusive J/ψ in DIS</p> <p>P-dissociative VM in PHP at large t</p> <p>Exclusive PHP of J/ψ mesons</p> <p>Exclusive ω in DIS</p> <p>Diffractive PHP of VM at large t</p> <p>Spin-Density ME of Exclusive ρ^0 in DIS</p> <p>Exclusive ρ^0 and J/ψ in DIS</p> <p>Elastic Υ Photoproduction</p> <p>Elastic and p-Dissociative ρ^0 in PHP</p> <p>Elastic J/ψ in PHP</p> <p>Elastic ω in PHP</p> <p>$\gamma^*p \rightarrow \phi p$ in DIS</p> <p>Elastic ϕ in PHP</p> <p>Elastic ρ^0 in PHP</p> <p>Exclusive ρ^0 in DIS</p>	<p>Nucl. Phys. B 909 (2016) 934</p> <p>Eur.Phys.J. C 72 (2012) 1869</p> <p>Phys.Lett. B 708 (2012) 14</p> <p>JHEP 05 (2010) 085</p> <p>Phys. Lett. B 680 (2009) 4</p> <p>PMC Physics A 1, 6</p> <p>Nucl. Phys. B 718 (2005) 3</p> <p>Nucl. Phys. B 695 (2004) 3</p> <p>Eur. Phys. J. C 26 (2003) 389</p> <p>Eur. Phys. J. C 24 (2002) 345</p> <p>Phys. Lett. B 487 (2000) 273</p> <p>Eur. Phys. J. C 14 (2000) 213</p> <p>Eur. Phys. J. C 12 (2000) 393</p> <p>Eur. Phys. J. C 6 (1999) 603</p> <p>Phys. Lett. B 437 (1998) 432</p> <p>Eur. Phys. J. C 2 (1998) 247</p> <p>Z. Phys. C 75 (1997) 215</p> <p>Z. Phys. C 73 (1996) 73</p> <p>Phys. Lett. B 380 (1996) 220</p> <p>Phys. Lett. B 377 (1996) 259</p> <p>Z. Phys. C 69 (1995) 39</p> <p>Phys. Lett. B 356 (1995) 601</p>

• Ratio:

- Independent of W
- Independent of t
- Increase with Q^2

