

New Results on Vector Meson Production at HERA

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on behalf of the H1 and ZEUS Collaborations

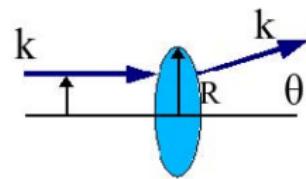
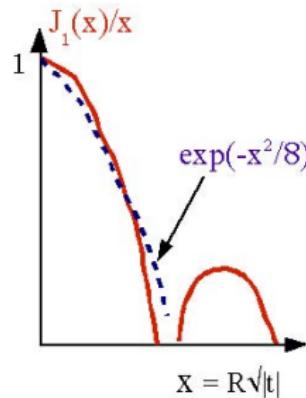
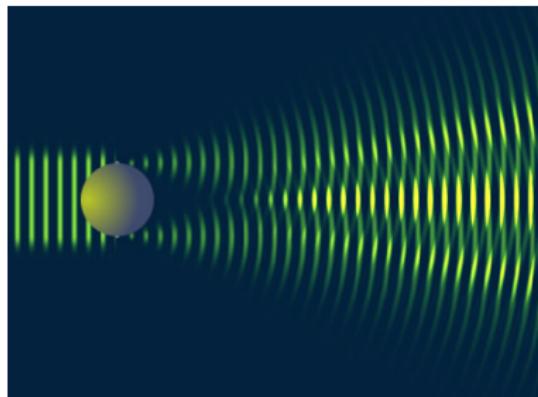
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- 1 Introduction: Vector Meson Production in Exclusive Diffraction in ep Scattering
- 2 New Results on Vector Mesons Production at HERA
 - ZEUS: cross section ratio $\psi(2S)/J/\psi(1S)$ in DIS
 - H1: Exclusive ρ^0 mesons with leading neutron in PHP

Diffraction of Light

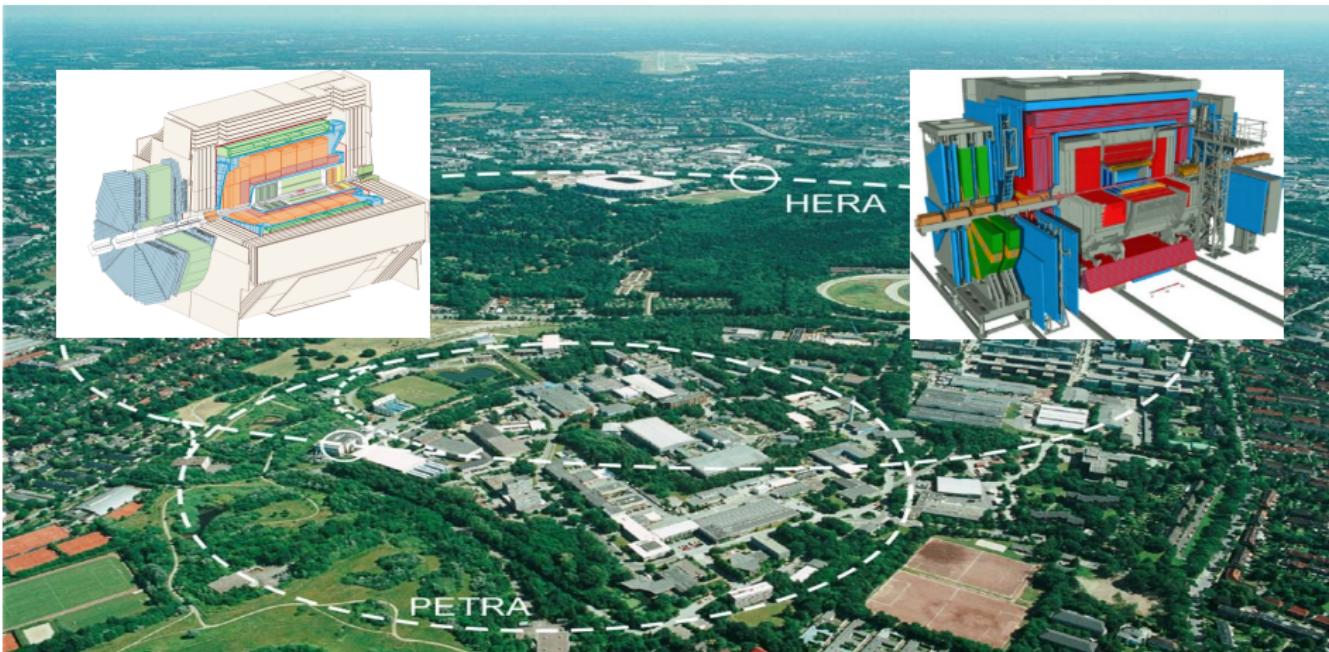


Light Scattering in Fraunhofer approximation (wavelength $\lambda \sim 1/k \ll R$)

- $|t| = 4k^2 \sin^2(\theta/2)$
- $d\sigma/dt \sim e^{-b|t|}$ (first diffractive peak approximated from Bessel function)
- $b = (R/2)^2 \rightarrow$ transverse size of the target
- in the presented studies: target \equiv proton and photon energy $\gg 1$ GeV

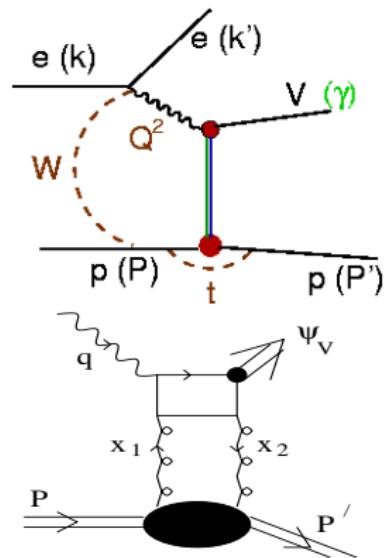
The HERA Accelerator, 1992 – 2007, DESY, Hamburg

World's first and only $e^\pm p$ collider, $E_e = 27.5 \text{ GeV}$, $E_p = 920 \text{ GeV}$ (820, 575, 460 GeV)



Total luminosity: $\int \mathcal{L} \sim 500 \text{ pb}^{-1}$ collected per experiment

Production of Vector Mesons in Exclusive Diffraction in ep Scattering



Kinematics: $M_V^2, Q^2, W, |t|$

M_V^2 - vector meson mass squared

$Q^2 (= -q^2 = -(k - k')^2)$ - the photon virtuality
(emitted by the incoming electron):

- $Q^2 \approx 0$ GeV 2 PHP (Photoproduction)
- larger Q^2 for DIS (Deep Inelastic Scattering)

W - invariant mass of the γp system

Process sensitive to the
gluon density in the proton

$|t|$ - 4-momentum transfer at the proton vertex

$$t = (P - P')^2$$

Kinematics of the exclusive process

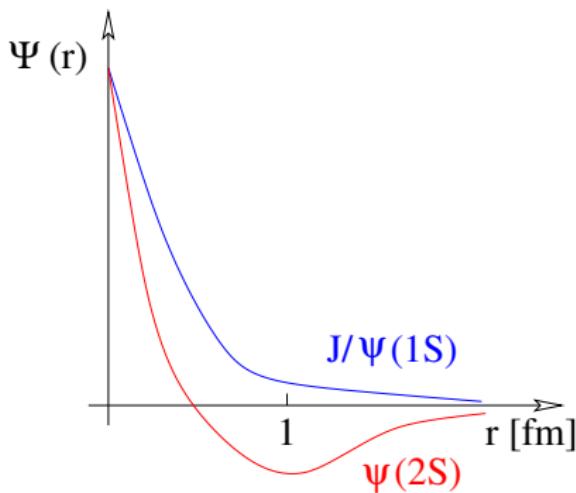
The proton stays intact !

pQCD: M_V^2 and Q^2 - set the scale at which the W and $|t|$ are probed

Outline

- 1 Introduction: Vector Meson Production in Exclusive Diffraction in ep Scattering
- 2 New Results on Vector Mesons Production at HERA
 - ZEUS: cross section ratio $\psi(2S)/J/\psi(1S)$ in DIS
 - H1: Exclusive ρ^0 mesons with leading neutron in PHP

ZEUS: cross section ratio $\psi(2S)/J/\psi(1S)$ in DIS



$$\text{Ratio } R = \frac{\sigma_{\gamma p \rightarrow \psi(2S)p}}{\sigma_{\gamma p \rightarrow J/\psi(1S)p}}$$

- sensitive to radial wave function of charmonium
- provides insight into the dynamics of the hard process

- $J/\psi(1S)$ and $\psi(2S)$ have distinctive wave functions
- $\psi(2S)$ has a node at ≈ 0.4 fm
- $\langle r_{\psi(2S)}^2 \rangle \approx 2 \langle r_{J/\psi(1S)}^2 \rangle$
- pQCD models predict $R \sim 0.17$ in PHP and rise of R with Q^2 in DIS

Investigated channels and data sample

Analyzed channels

- $\psi(2S) \rightarrow J/\psi + \pi^+ \pi^-$; $J/\psi \rightarrow \mu^+ \mu^-$
- $\psi(2S) \rightarrow \mu^+ \mu^-$
- $J/\psi(1S) \rightarrow \mu^+ \mu^-$

HERA II DATA $\mathcal{L} = 354 \text{ pb}^{-1}$ (2003 - 2007)

MC Samples

- Signal: DIFFVM for exclusive VM production (J/ψ and ψ')
- Background: GRAPE for non resonant muon pair production (Bethe-Heitler process)

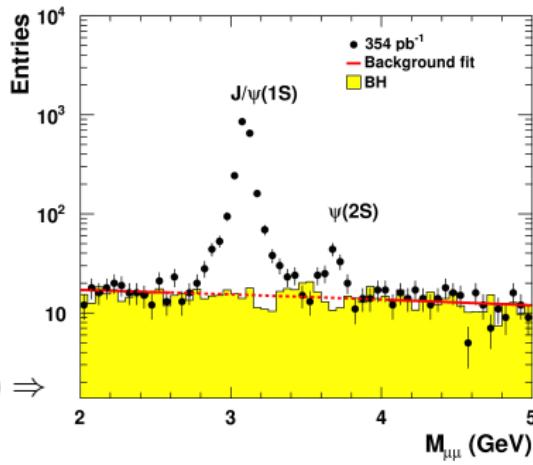
Event selection

- scattered electron $E_{e'} > 10 \text{ GeV}$ in CAL
- 2 (4 for $\psi(2S)$) 4-prongs decay
non-electron tracks from primary vertex,
net charge = 0
- two tracks identified as muons
(CAL, F/B/R/MUO, BAC)
- no other deposits not matched to tracks
(above CAL noise)

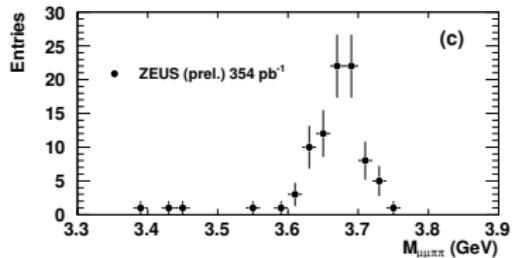
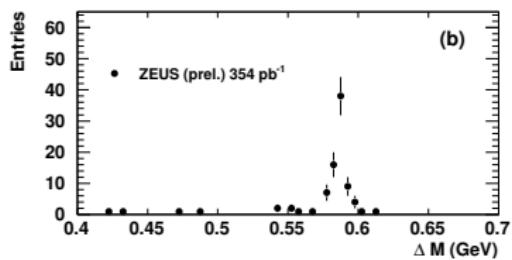
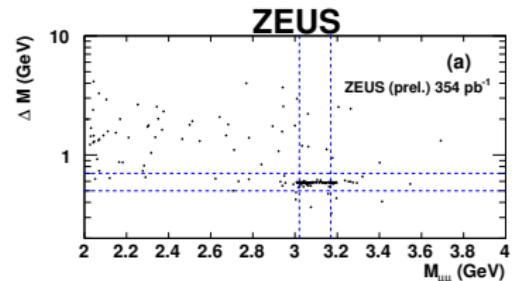
Kinematic range:

- $30 \leq W \leq 210 \text{ GeV}$
- $5 \leq Q^2 \leq 70 \text{ GeV}^2$
- $|t| \leq 1 \text{ GeV}^2$

$$M(\mu^+ \mu^-) \Rightarrow$$



Selection specific for $\psi(2S) \rightarrow \mu^+\mu^-\pi^+\pi^-$ channel



- ΔM vs. $M_{\mu^+\mu^-}$

$\Delta M = M(\mu^+\mu^-\pi^+\pi^-) - M(\mu^+\mu^-)$
cascade decay of $\psi(2S)$

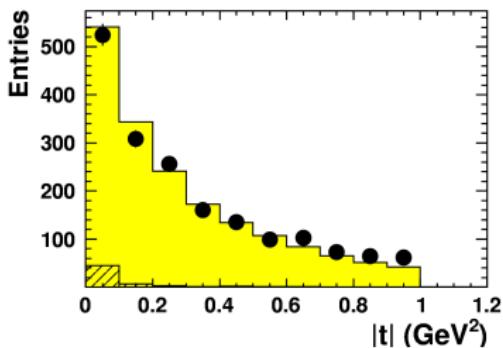
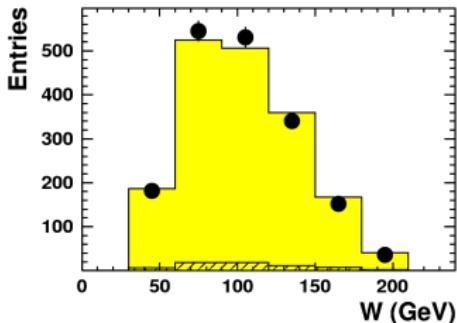
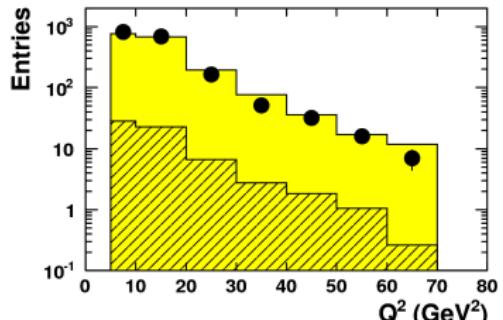
- $0.5 < \Delta M < 0.7 \text{ GeV}$

$3.02 < M_{\mu^+\mu^-} < 3.17 \text{ GeV}$

- $M(\mu^+\mu^-\pi^+\pi^-)$ after cleanup
very clean signature
(≤ 3 background events)

Control plots for $J/\psi \rightarrow \mu^+ \mu^-$ channel

ZEUS



$J/\psi(1S) \rightarrow \mu^+ \mu^-$
● ZEUS (prel.) 354 pb $^{-1}$

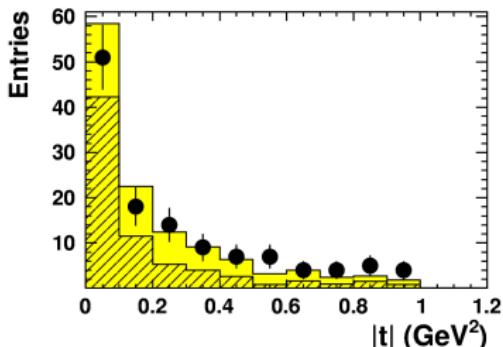
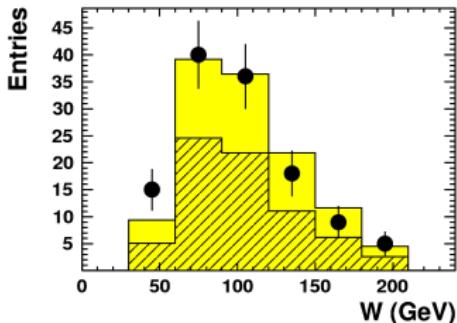
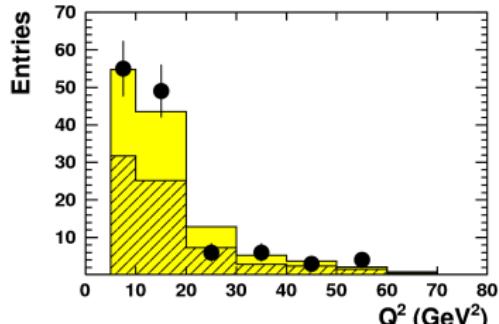
■ DIFFVM + BH

▨ BH

- MC reweighted in Q^2 , $|t|$ and J/ψ decay angles to match the data
- good description → detector efficiency calculation

Control plots for $\psi(2S) \rightarrow \mu^+ \mu^-$ channel

ZEUS



$\psi(2S) \rightarrow \mu^+ \mu^-$

● ZEUS (prel.) 354 pb $^{-1}$

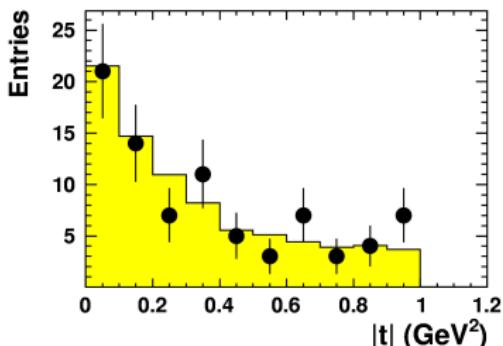
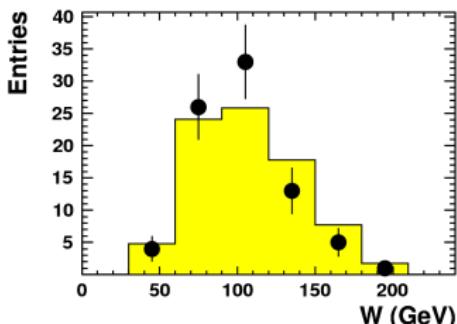
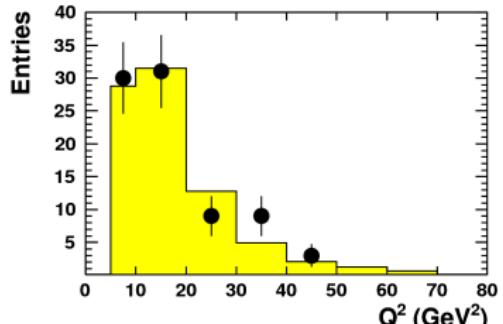
■ DIFFVM + BH

▨ BH

- MC reweighted in Q^2 , $|t|$ and $\psi(2S)$ decay angles using $J/\psi \rightarrow \mu^+ \mu^-$ weights
- good description → detector efficiency calculation

Control plots for $\psi(2S) \rightarrow \mu^+ \mu^- \pi^+ \pi^-$ channel

ZEUS



$\psi(2S) \rightarrow J/\psi(1S) \pi^+ \pi^-$

- ZEUS (prel.) 354 pb^{-1}

- DIFFVM

- MC reweighted in Q^2 and $|t|$ using $J/\psi \rightarrow \mu^+ \mu^-$ weights

- good description → detector efficiency calculation

Cross section ratio $R = \frac{\sigma(\psi(2S))}{\sigma(J/\psi(1S))}$ for full kinematic range

For $30 \leq W \leq 210$ GeV, $5 \leq Q^2 \leq 70$ GeV 2 , $|t| \leq 1$ GeV 2

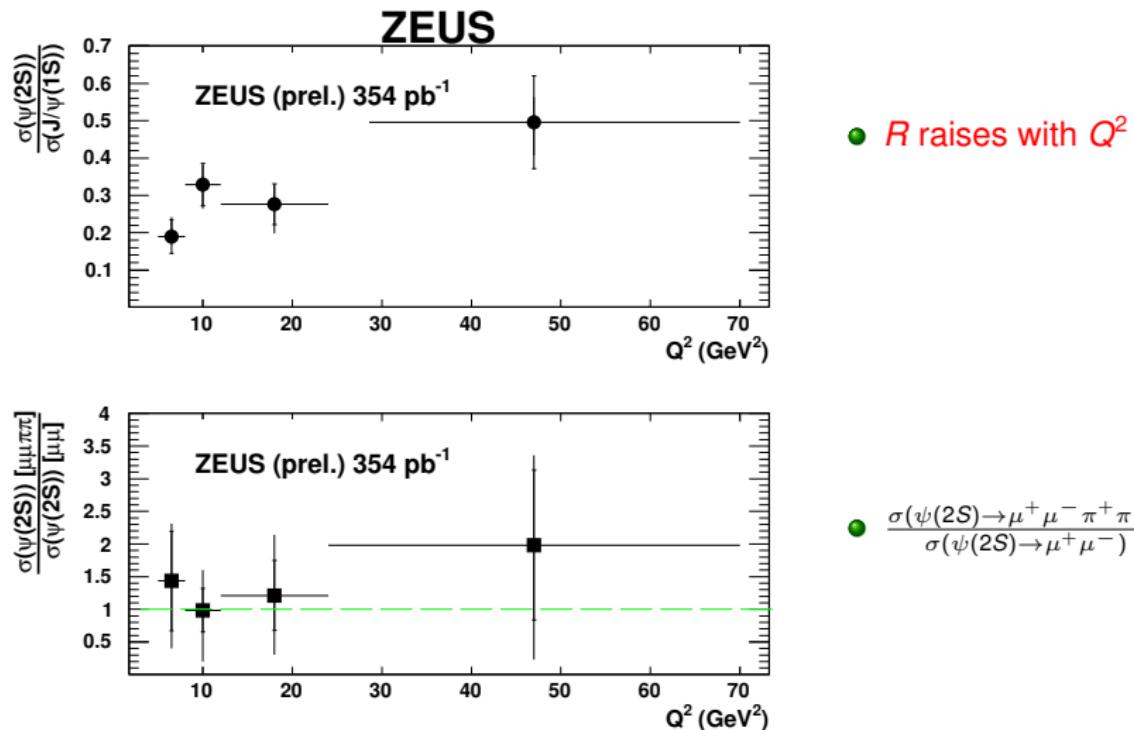
$\psi(2S)$ decay mode	$R = \frac{\sigma(\psi(2S))}{\sigma(J/\psi(1S))}$
$\rightarrow J/\psi(\rightarrow \mu^+ \mu^-) \pi^+ \pi^-$	$0.29 \pm 0.04^{+0.02}_{-0.01}$
$\rightarrow \mu^+ \mu^-$	$0.25 \pm 0.05^{+0.04}_{-0.02}$
combined	$0.28 \pm 0.03^{+0.02}_{-0.01}$

ZEUS Preliminary

- both channels provide consistent results

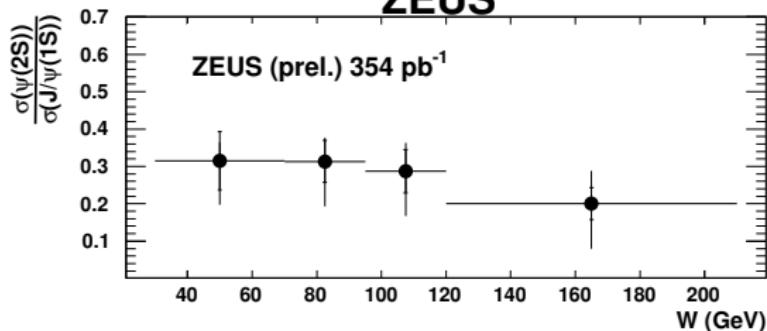
- $R_{\psi(2S) \rightarrow J/\psi \pi^+ \pi^-} = \frac{\sigma_{\psi(2S)}}{\sigma_{J/\psi(1S)}} = \frac{N_{\psi(2S)}}{N_{J/\psi(1S)}} \cdot \frac{Acc_{J/\psi(1S) \rightarrow \mu^+ \mu^-}}{Acc_{\psi(2S) \rightarrow J/\psi \pi^+ \pi^-}} \cdot \frac{1}{BR_{\psi(2S) \rightarrow J/\psi \pi^+ \pi^-}}$
- $R_{\psi(2S) \rightarrow \mu^+ \mu^-} = \frac{\sigma_{\psi(2S)}}{\sigma_{J/\psi(1S)}} = \frac{N_{\psi(2S)}}{N_{J/\psi(1S)}} \cdot \frac{Acc_{J/\psi(1S) \rightarrow \mu^+ \mu^-}}{Acc_{\psi(2S) \rightarrow \mu^+ \mu^-}} \cdot \frac{BR_{J/\psi(1S) \rightarrow \mu^+ \mu^-}}{BR_{\psi(2S) \rightarrow \mu^+ \mu^-}}$
- $Acc_i = \frac{N_i^{reco}}{N_i^{true}}$

Ratio $R = \frac{\sigma(\psi(2S))}{\sigma(J/\psi(1S))}$ vs. Q^2

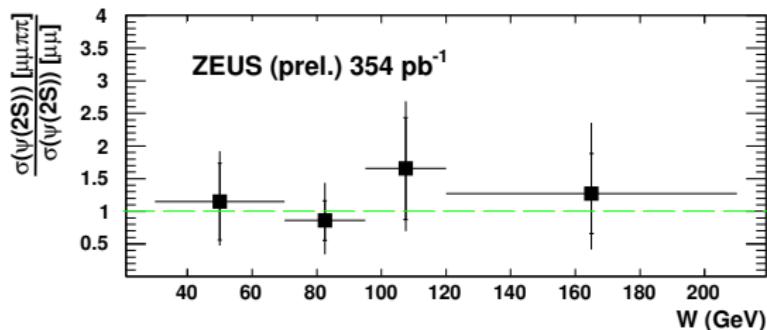


Ratio $R = \frac{\sigma(\psi(2S))}{\sigma(J/\psi(1S))}$ vs. W

ZEUS



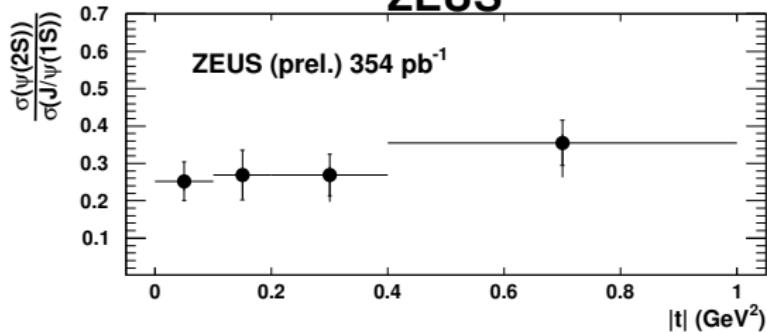
• R independent of W



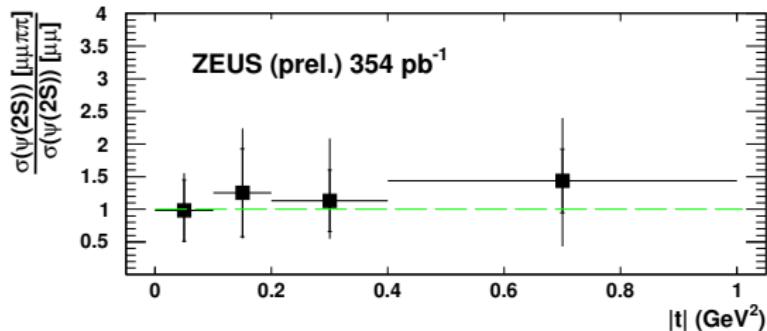
• $\frac{\sigma(\psi(2S) \rightarrow \mu^+ \mu^- \pi^+ \pi^-)}{\sigma(\psi(2S) \rightarrow \mu^+ \mu^-)}$

Ratio $R = \frac{\sigma(\psi(2S))}{\sigma(J/\psi(1S))}$ vs. $|t|$

ZEUS



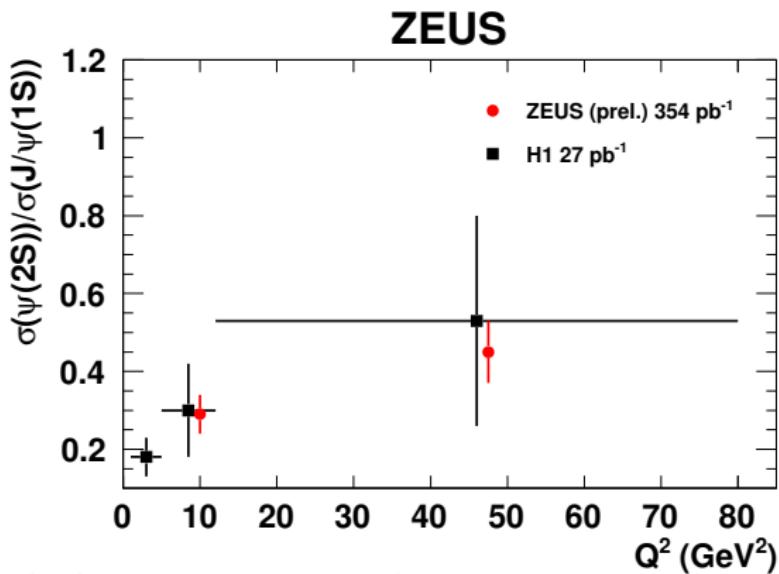
• R independent of $|t|$



• $\frac{\sigma(\psi(2S) \rightarrow \mu^+ \mu^- \pi^+ \pi^-)}{\sigma(\psi(2S) \rightarrow \mu^+ \mu^-)}$

ZEUS to H1 comparison

- cross check: ZEUS data analyzed in Q^2 bins used by H1: [EPJ C10 (1999) 373.]
(5 – 12) GeV^2 and (12 – 80) GeV^2
- $40 < W < 180 \text{ GeV}$ and $1 < Q^2 < 80 \text{ GeV}^2$

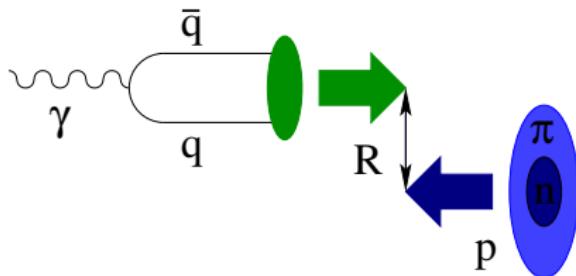


- both measurements are in agreement
- improved accuracy due to the increased statistic of HERA II data

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H1: Exclusive ρ^0 mesons with leading neutron in PHP



- First observation of exclusive VM photoproduction on (virtual) pion

$$\gamma\pi^+ \rightarrow VM + \pi^+$$

- unique opportunity at HERA (γ , π^+ beams existed before, but not the target)
- extends the landscape of Vector Meson production at HERA

- experimental challenge

- trigger: tagged PHP too large W to register the VM;
untagged PHP very large rate, requires prescaling
- limited acceptance for (very) forward π and neutron ($\eta_{LAB} > 6$)

- advantages of H1 during HERA-II run

- improved Forward Neutron Calorimeter (FNC) (identifies and measures n and γ/π^0)
- efficient Fast Track Trigger (FTT) allows to collect untagged soft PHP events

Physics process and analysis phase space

- $\gamma^* + p \rightarrow \rho^0 \pi^+ n, \quad \rho^0 \rightarrow \pi^+ \pi^-$

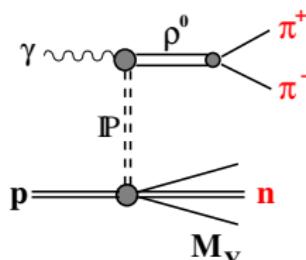
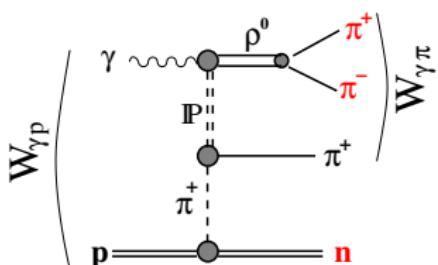
• Key observables

- $x_L = E_n/E_p$ (or $x_\pi = 1 - x_L$, distribution: $\sim f_{\pi/\rho}(x_L)$)
- W dependence: $\sim W^\delta$, nature of exchanged object
- t -slope, $dN/d|t| \sim \exp(-b|t|)$, $b \sim R^2$, size of the interaction region

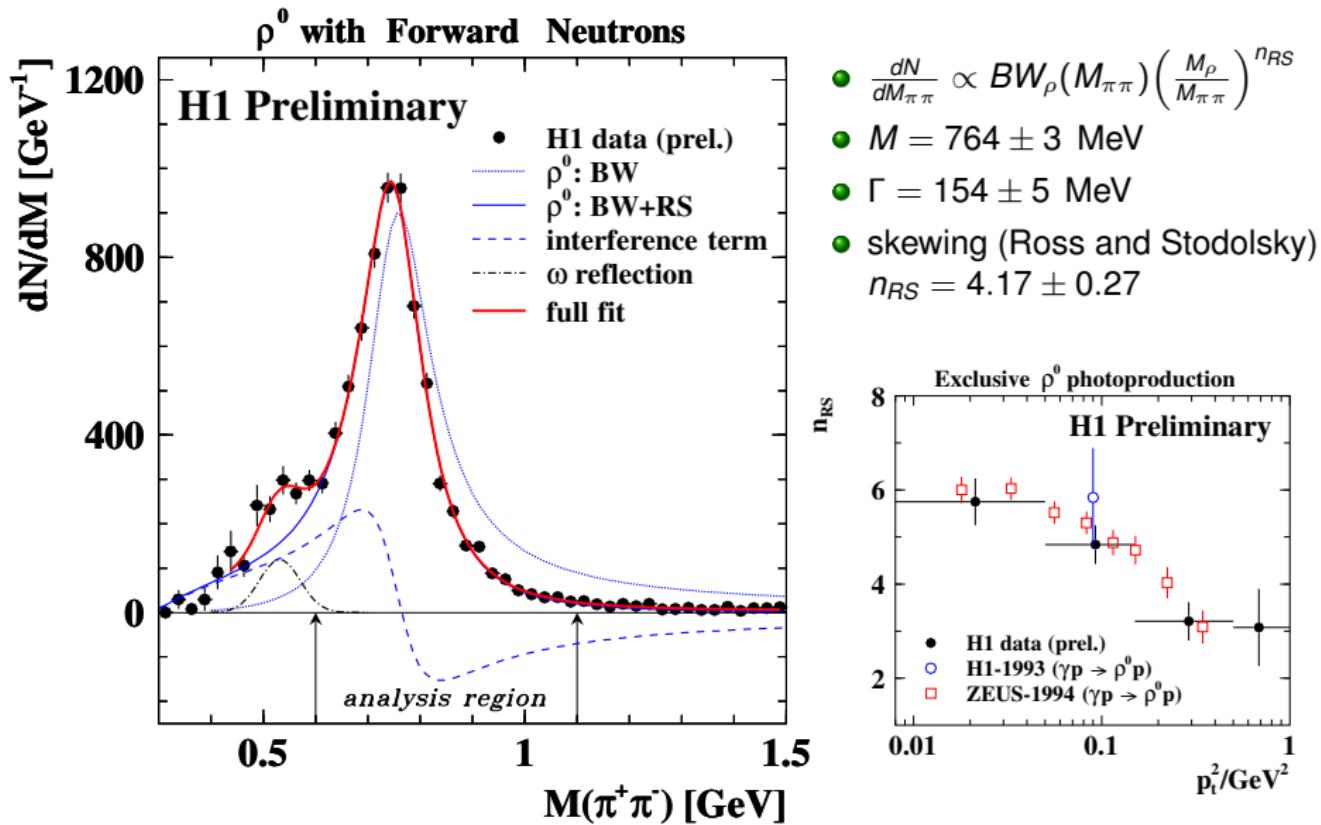
• Kinematics

- Photoproduction: $Q^2 < 2 \text{ GeV}^2$, $\langle Q^2 \rangle = 0.05 \text{ GeV}^2$
- Low p_t : $|t| < 1 \text{ GeV}^2$, $\langle |t| \rangle = 0.20 \text{ GeV}^2$,
 $t = (P_\gamma - P_{\rho^0})^2$ (at top vertex)
- small mass: $0.3 < m_{\pi\pi} < 1.5 \text{ GeV}$
- track acceptance in Central Tracker:
 $20 < W_{\gamma p} < 100 \text{ GeV}$, $\langle W_{\gamma p} \rangle = 48 \text{ GeV}$
 $W_{\gamma p} = \sqrt{2(E - p_z)_\rho E_p}$, $W_{\gamma\pi} = W_{\gamma p} \sqrt{1 - x_L}$
- Forward neutron: $E_n > 120 \text{ GeV}$, $\theta_n < 0.75 \text{ mrad}$

- No hard scale present: Regge framework is most adequate
- Diffractive BG has irreducible component: $M_Y = N^* \rightarrow n \pi^+$

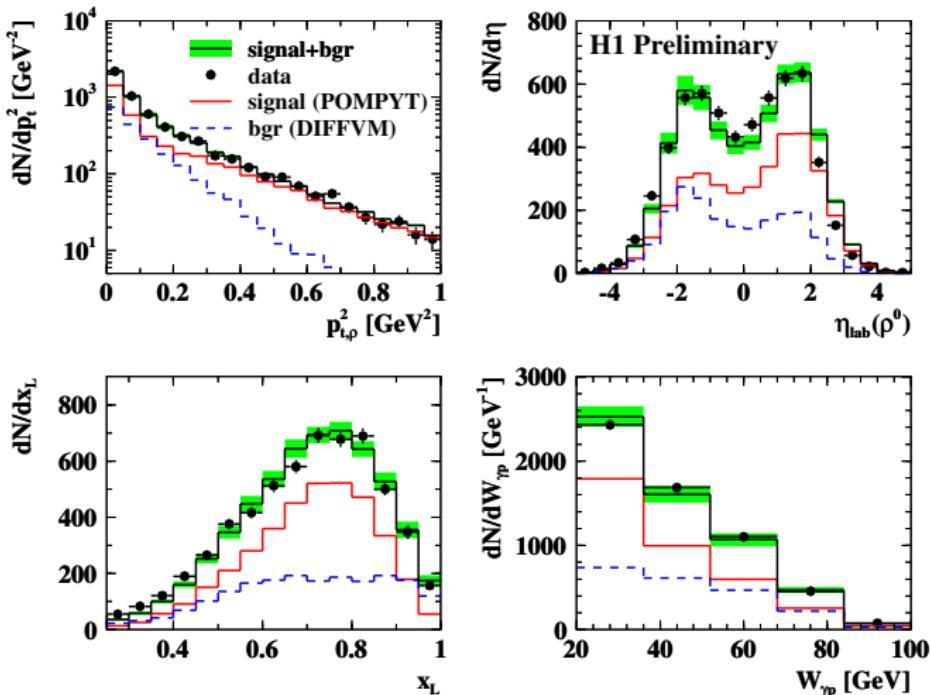


ρ^0 meson shape



ρ^0 photoproduction: control distributions

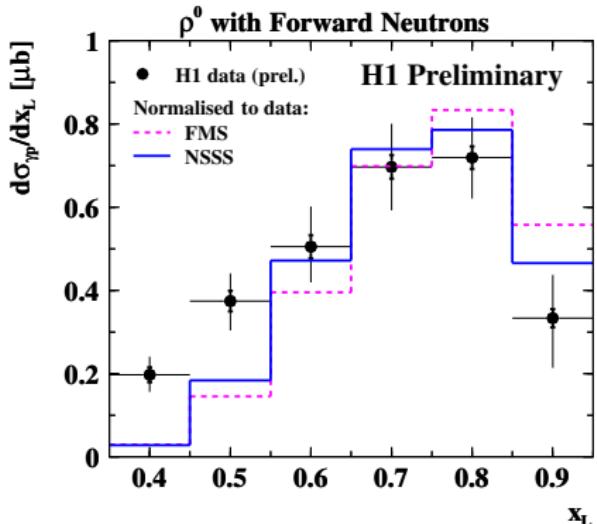
Exclusive photoproduction of ρ^0 with Forward Neutrons



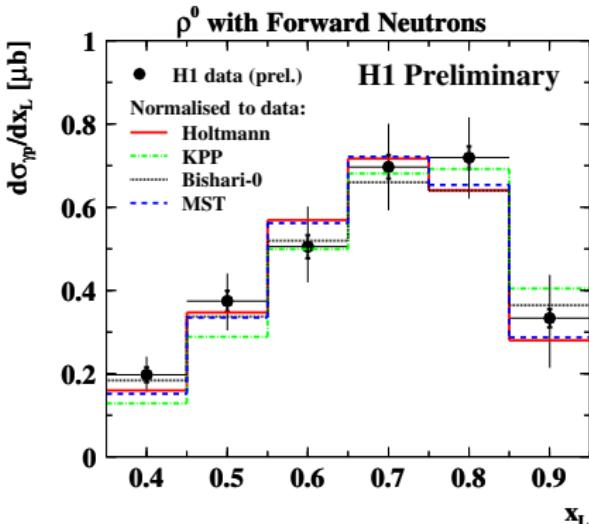
Data points are shown with stat. errors only; green band represents estimated uncertainty on the p.diss BG fraction

Pion fluxes verified by H1 data

Classes of π -fluxes restricted by comparing the x_L distribution



Example of fluxes **excluded by the data** (too soft pions "in the proton")



Fluxes **compatible with H1 data**
($\chi^2 = 2.1$ to 5.5 for 6 points)

Total Cross Sections: H1 Preliminary

$$\bullet \sigma_{\gamma p} = \frac{\sigma_{ep}}{\int f_{\gamma/e}(y, Q^2) dy dQ^2} = \frac{N_{\text{DATA}} - N_{\text{BG}}}{\mathcal{L}(A \cdot \epsilon) \mathcal{F}} \cdot C_p$$

- N_{BG} - proton dissociation background from MC
- \mathcal{L} - integrated luminosity
- $A \cdot \epsilon$ - acceptance and efficiency corrections
- \mathcal{F} - photon flux integrated over $20 < W < 100$ GeV, $Q^2 < 2$ GeV 2
- C_p - correction due to the extrapolation to the full ρ^0 mass range

For $0.35 < x_L < 0.95$, $20 < W < 100$ GeV, $\theta_n < 0.75$ mrad

$$\sigma(\gamma p \rightarrow \rho^0 n(\pi^+)) = (280 \pm 6_{\text{stat}} \pm 46_{\text{syst}}) \text{ nb}$$

$$\bullet \sigma_{\gamma\pi}(\langle W_{\gamma\pi} \rangle) = \frac{\sigma_{\gamma p}}{\int f_{\pi^+/p}(x_L, t) dx_L dt}$$

For $\langle W_{\gamma\pi} \rangle = 22$ GeV

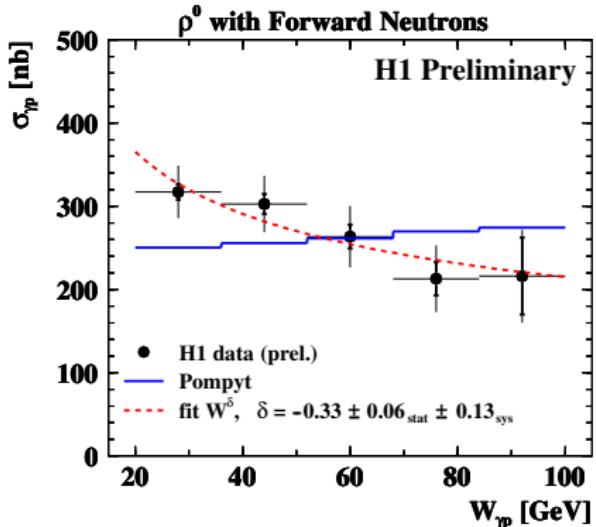
$$\sigma_{el}(\gamma\pi^+ \rightarrow \rho^0\pi^+) = (2.03 \pm 0.34_{\text{exp}} \pm 0.51_{\text{model}}) \mu\text{b}$$

taking interpolated value of $\sigma(\gamma p \rightarrow \rho^0 p) = 9.5 \mu\text{b}$ at corresponding energy

$$r_{el} = \sigma_{\gamma\pi}^{el} / \sigma_{\gamma p}^{el} = 0.21 \pm 0.06 \quad (\text{cf. } r_{\text{tot}} = \sigma_{\gamma\pi}^{\text{tot}} / \sigma_{\gamma p}^{\text{tot}} = 0.32 \pm 0.03 \text{ [ZEUS, Nucl. Phys. B637 (2002) 3.]})$$

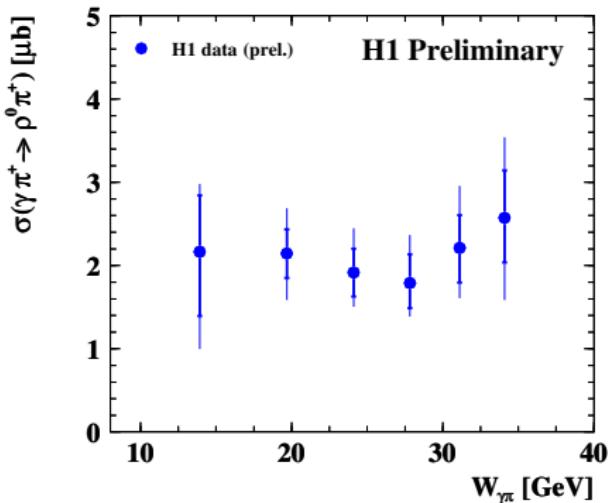
W Dependence of the Total γp and $\gamma \pi$ Cross Sections

inner error bars: statistical uncertainty
outer error bars: $\sqrt{\text{stat}^2 + \text{syst}^2}$



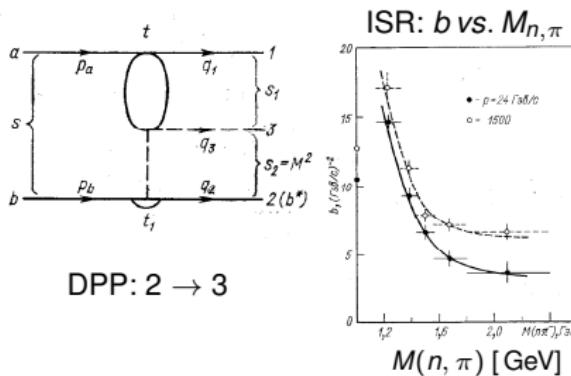
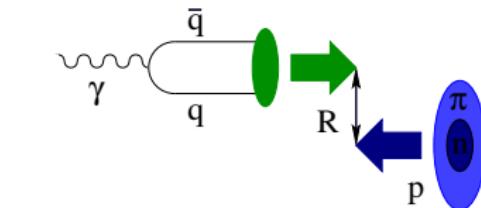
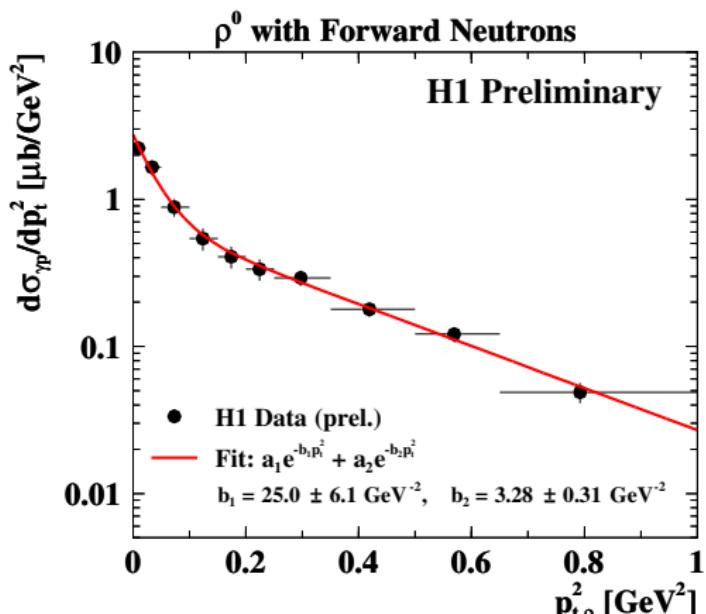
Regge motivated power law fit W^δ yields $\delta < 0$ in qualitative agreement with DPP and in contrast to MC, $\delta_{MC} = 0.08 \pm 0.02$ as expected from purely IP exchange

inner error bars: total experimental uncertainty
outer error bars: $\sqrt{\text{exp}^2 + \text{model}^2}$



Holtmann flux is used for the central values, conservative model uncertainty $\sim 25\%$

Differential cross section in $p_{t,\rho}^2$



- In geometric picture: $\langle r^2 \rangle = 2b_1 \cdot (\hbar c)^2 \simeq 2 \text{ fm}^2 \Rightarrow (1.6R_p)^2 \Rightarrow$ ultra-peripheral process
- DPP model: low mass $\pi^+ n$ state → large slope, high mass → small (less steep) slope

Summary and Conclusions

- Ratio of $\frac{\sigma(\psi(2S))}{\sigma(J/\psi(1S))}$ using HERA II data was measured in the kinematic range:
 $30 \leq W \leq 210$ GeV, $5 \leq Q^2 \leq 70$ GeV 2 , $|t| \leq 1$ GeV 2
 - The ratio increases with Q^2 and is constant as a function of W and $|t|$
 - ZEUS measurement will be extended to $2 < W < 5$ GeV 2 using HERA I data
 - **Theoretical calculations of the ratio $\frac{\sigma(\psi(2S))}{\sigma(J/\psi(1S))}$ are welcome :).**
-
- Cross section for the exclusive PHP of ρ^0 mesons associated with leading neutron has been measured for the first time at HERA
 - Differential cross sections for the reaction $\gamma p \rightarrow \rho^0 n \pi^+$ are compatible with model of Double Peripheral Process (DPP)
 - The elastic photon-pion cross section, $\sigma(\gamma \pi^+ \rightarrow \rho^0 \pi^+)$ was extracted in the One Pion Exchange (OPE) approximation

Thank You For Your Attention

BACKUP PLOTS

BACKUP PLOTS FOLLOWS...

- **Data Sample:** 2006-7 e^+ runs, $\sqrt{s} = 319$ GeV, $\int L = 1.16 \text{ pb}^{-1}$,
~ 6600 events in final sample
- **Tracking:** two opposite charge tracks fitted to event vertex $|z_{\text{vtx}}| < 30$ cm,
 $p_t^{tr} > 0.2$ GeV, $20^\circ < \theta^{tr} < 160^\circ$
effective mass range $M_{\pi\pi} \in (0.6, 1.1)$ GeV,
extrapolated for $\sigma(\rho^0)$ to (0.28, 1.5) GeV
- **FNC:** high energy neutron $E_n > 120$ GeV, within good acceptance $\theta_n < 0.75$ mrad
BG fraction from x_L shape fit: $F_{bg} = 0.36 \pm 0.06$ (subtracted from the data)
- **Exclusivity:** no deposits above detector noise level
except two pions from ρ^0 decay and the leading neutron
- **Monte Carlo modeling**
 - Signal (Double Peripheral Process DPP): POMPYT × PYTHIA6
(π -flux × elastic $\gamma\pi \rightarrow \rho^0\pi$)
 - Background: DIFFVM (elastic, p -diss, γ -diss, double diss)

H1: OPE (One Pion Exchange) and pion fluxes

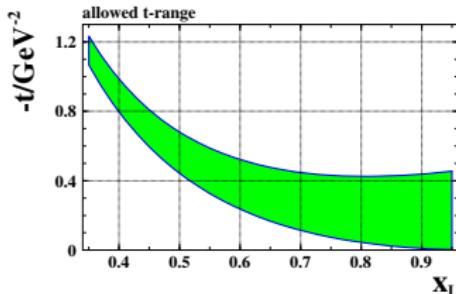
$$\frac{d^2\sigma_{\gamma p}(W^2, x_L, t)}{dx_L dt} = f_{\pi/p}(x_L, t)\sigma_{\gamma\pi}((1-x_L)W^2)$$

$$\frac{d\sigma_{\gamma p}}{dx_L} = \int_{t_0(x_L)}^{t_{min}(x_L)} f_{\pi/p}(x_L, t) dt \cdot \sigma_{\gamma\pi}(W_{\gamma\pi})$$

$$\text{where } t = -\frac{p_{t,n}^2}{x_L} - \frac{(1-x_L)(m_n^2 - m_p^2 x_L)}{x_L}$$

$$\sigma_{\gamma\pi}(W_{\gamma\pi}) = \frac{1}{\Gamma_\pi(x_L)} \frac{d\sigma_{\gamma p}}{dx_L} \text{ and } \overline{\sigma_{\gamma\pi}}(\langle W_{\gamma\pi} \rangle) = \frac{\sigma_{\gamma p}}{\int \Gamma_\pi}$$

$$\text{where } \Gamma_\pi = \int_{t_0(x_L)}^{t_{min}(x_L)} f_{\pi/p}(x_L, t) dt$$



← allowed t -range due to the FNC acceptance

Typical examples of pion fluxes:

$$f_{\pi^+/p}(x_L, t) = \frac{1}{2\pi} \frac{g_{p\pi N}^2}{4\pi} (1-x_L) \frac{-t}{(m_\pi^2 - t)^2} \exp[R_{\pi N}^2 \frac{m_\pi^2 - t}{1-x_L}] \quad \text{-- H. Holtmann et al., Nucl. Phys. A596 (1996) 631.}$$

$$f_{\pi^+/p}(x_L, t) = \frac{1}{2\pi} \frac{g_{p\pi N}^2}{4\pi} (1-x_L) \frac{1-2\alpha'_\pi t}{(m_\pi^2 - t)^2} \exp[R_\pi^2 (m_\pi^2 - t)] \quad \text{-- B. Kopeliovich et al., Z. Phys. C73 (1996) 125.}$$

Very many pion fluxes has been proposed...

