

Exclusive Diffraction and Leading Baryons at HERA

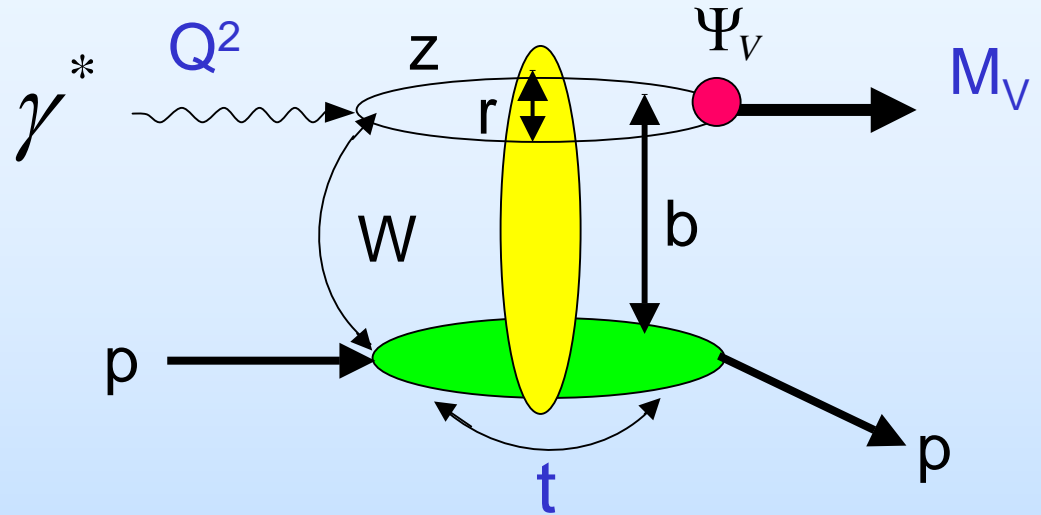
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representing the H1 and ZEUS Collaboration

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Exclusive Vector Meson Production

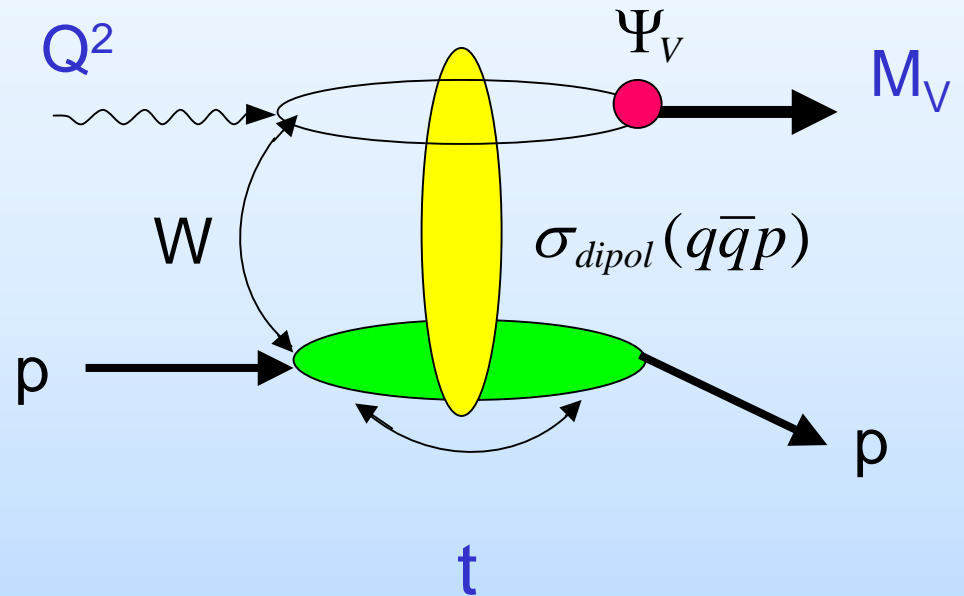


Factorization:

- $\gamma^* \rightarrow q\bar{q}$ $\Psi_\gamma(z, r)$ QED
- dipole–proton interaction
 $\text{Ampl} \sim \Psi_\gamma(r, z) \otimes \sigma_{\text{dip}}(r, z, b) \otimes \Psi_V(z, r)$
- $q\bar{q} \rightarrow V$ Ψ_V model
 parton-hadron duality

Continuous transition **soft** \rightarrow **hard physics**

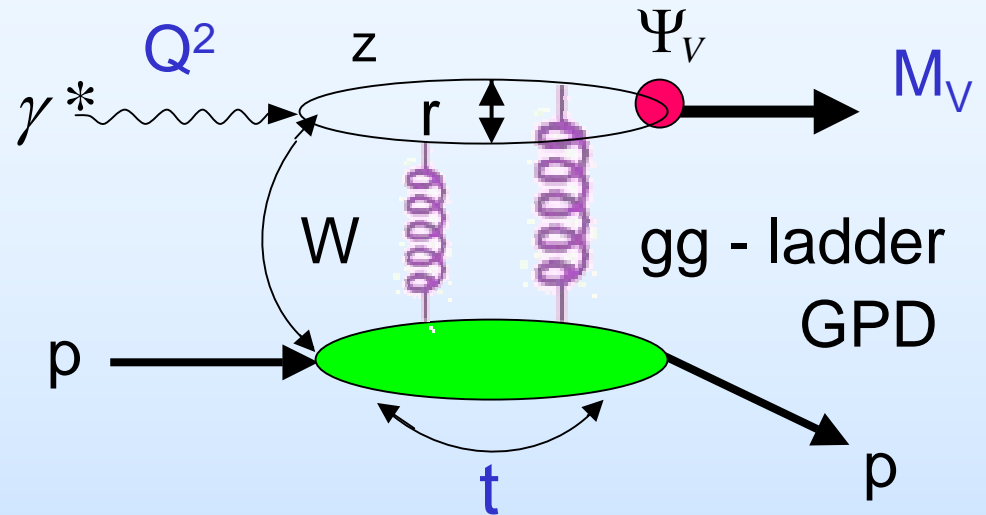
Dipole scattering:



- small Bjorken x : factorization
- valid also at low Q^2
- σ_{dipol} universal: applicable to **DIS**, **DDIS** and **VM** production
- Saturation included

LO 2-gluon exchange

pQCD



Expectation:

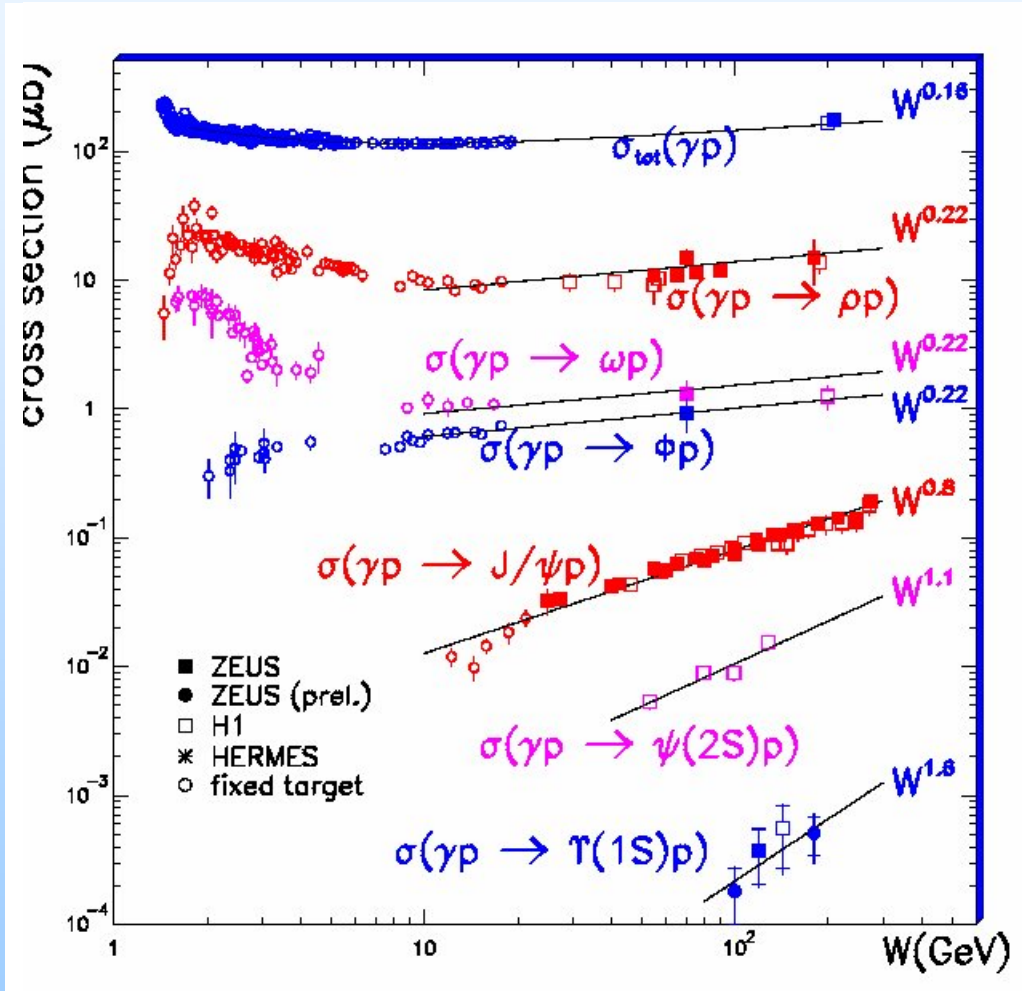
- steep rise with W : $\sigma \sim (xg(x, Q^2))^2$ $x \approx Q^2 / W^2$
 steep rise of $g(x)$ with x decreasing $\rightarrow \sigma \sim W^\delta$
 δ increases with M_V, Q^2

- r decreases with Q^2, M_V
 $\bar{Q}^2 = z(1-z)(Q^2 + M_V^2)$ in perturbative domain:
 A_L $z \approx 1/2$ scale variable $\bar{Q}^2 = 1/4(Q^2 + M_V^2)$
 A_T : contribution at $z = 0, 1 \Rightarrow$ scaling delayed

Hard scale: M_V



Photoproduction: $\sigma \sim W^\delta$



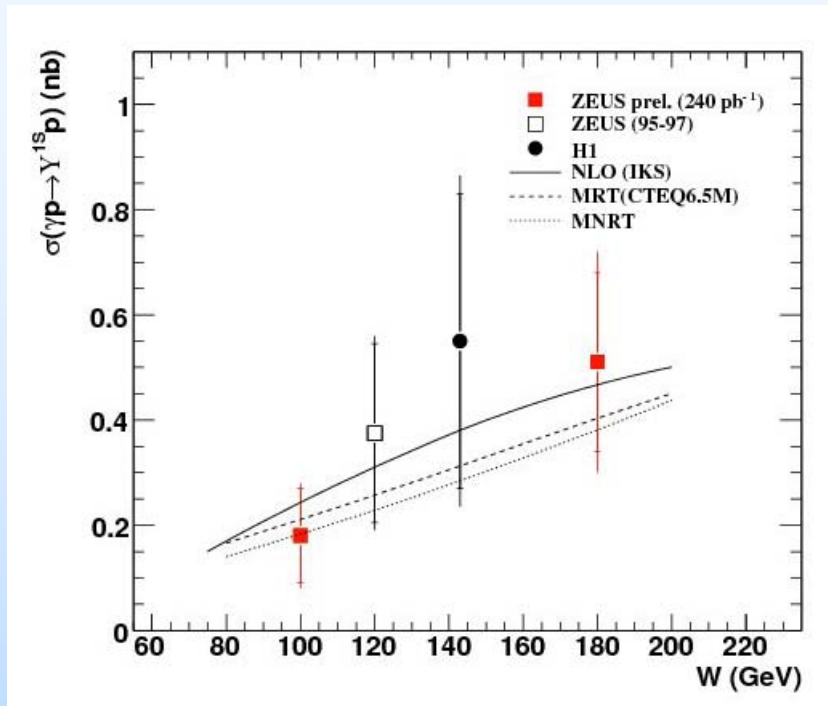
↓ Regge model
↑ pQCD

$$d\sigma/dt \sim (W)^{4(\alpha(t)-1)}$$
$$\alpha(t) = 1.08 + \alpha' t$$

M_V hard scale for
 $\Psi(1S), \Psi(2S), \Upsilon(1S)$

$\Psi(2S)$ special case: zero of wave function \rightarrow smaller dipole

Comparison with models: $\gamma p \rightarrow \Upsilon(1S) p$



$$\sigma \sim W^\delta, \delta \approx 1.8$$

M_Ψ, M_Υ hard scale

Models:

Ivanov, Krasnikov, Szymynowski (IKS): NLO, GPD

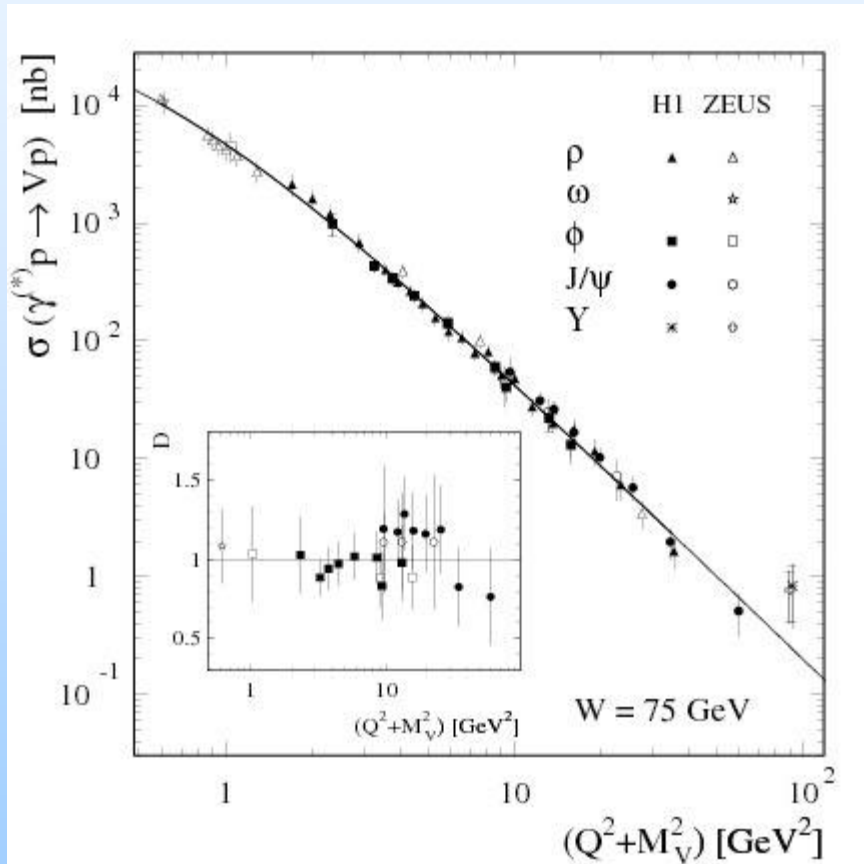
Martin, Ryskin, Teubner (MRT): NLO, skewed gluons
CTEQ6.5M gluon

MRT + Nockles (MNRT): gluons from $\Psi(1S)$ data

Hard scale: Q^2



Scaling of vector meson elastic cross sections:



SU(4) flavor factors considered

Universal scale: $Q^2 + M_V^2$

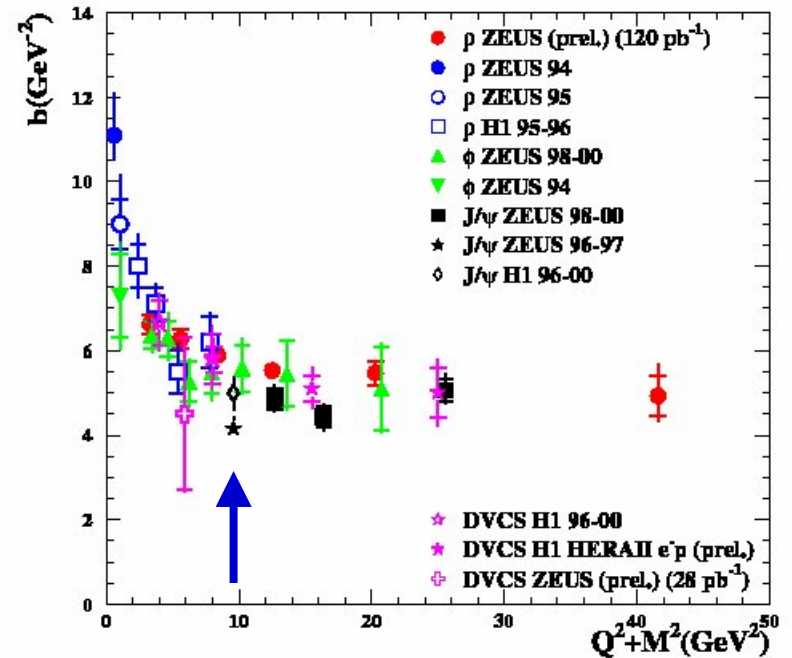
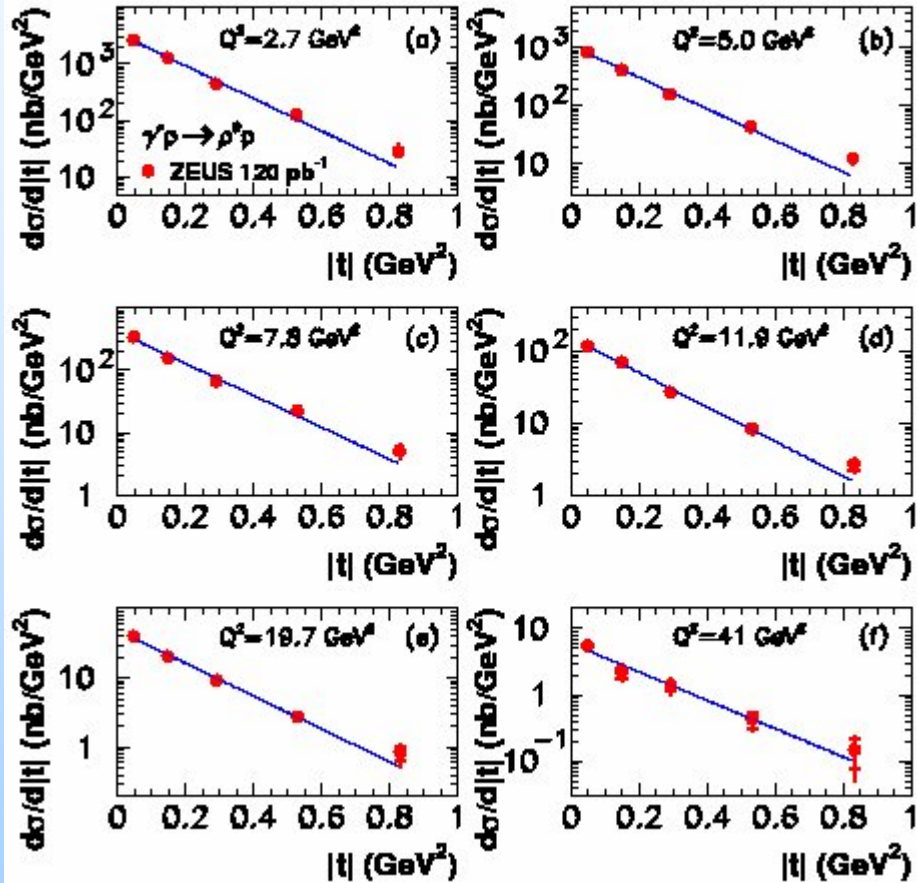
Transition to hard scale:

$$\gamma^* p \rightarrow V p$$

$$d\sigma/dt \sim \exp(-b t)$$

$$b = b_{\text{dip}} \oplus b_{\text{nucleon}}$$

ZEUS



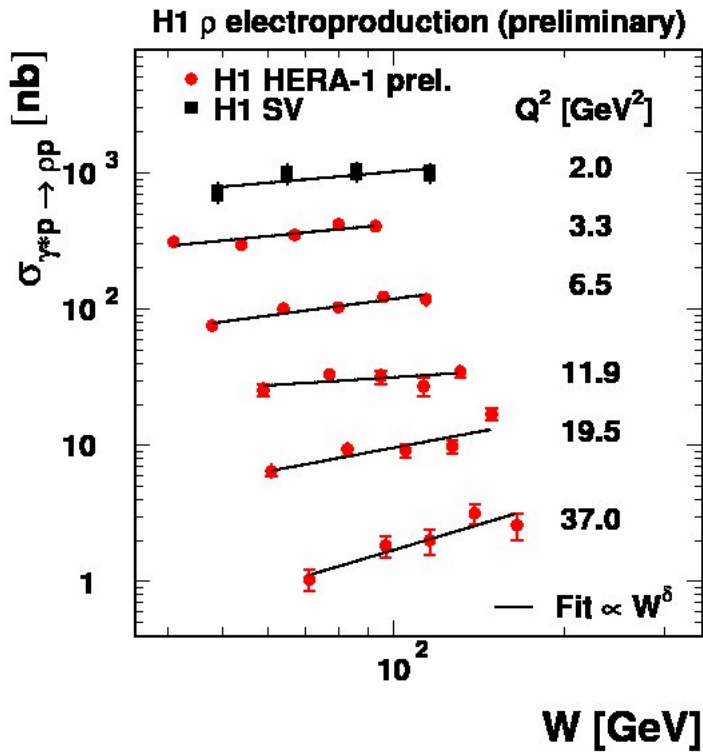
$Q^2 > 10$ GeV² hard scale:

point like dipole probes gluon cloud of proton

$$\gamma^* p \rightarrow \rho^0 p$$

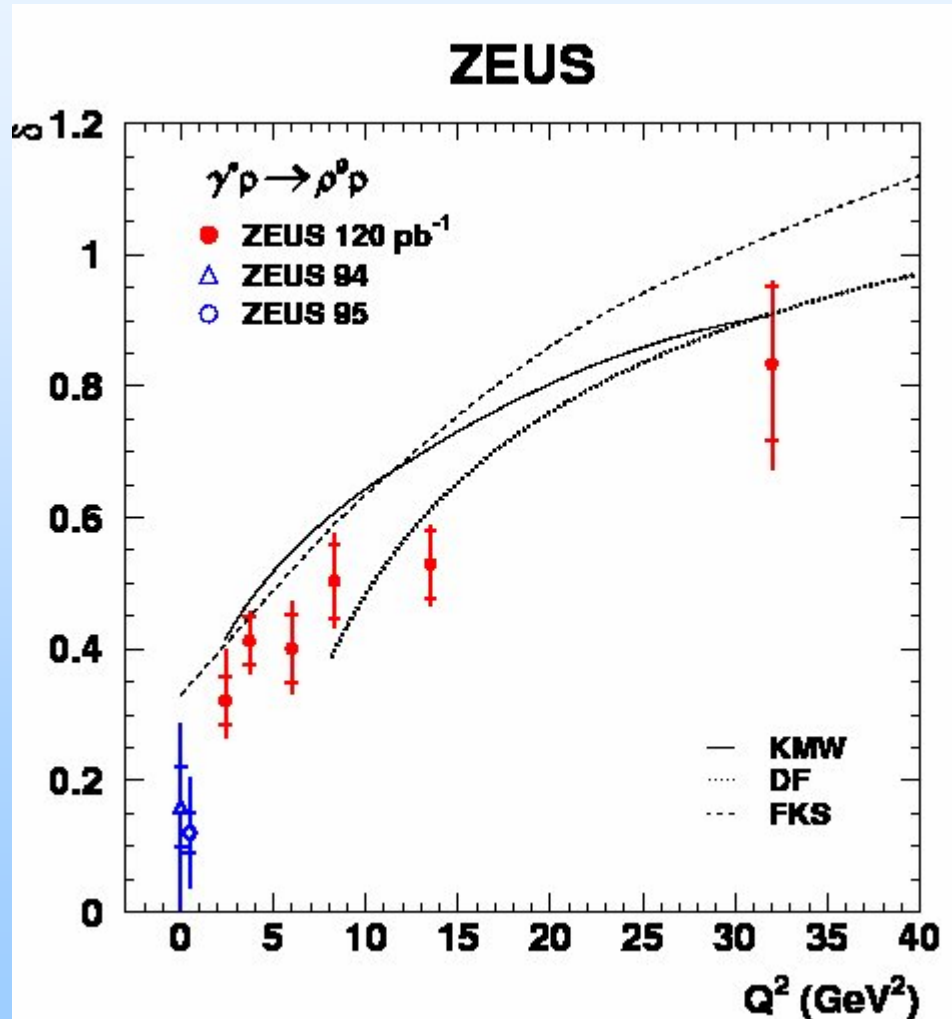
$$\sigma \sim W^\delta, \delta = \delta(Q^2)$$

- Harder with increase of Q^2



Model predictions:

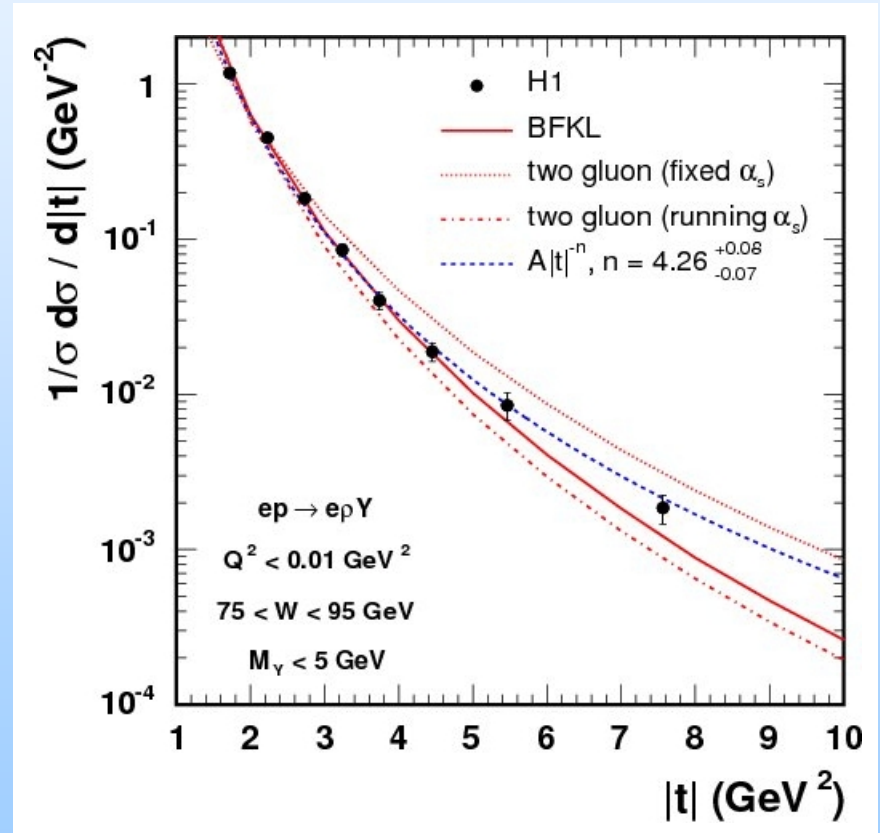
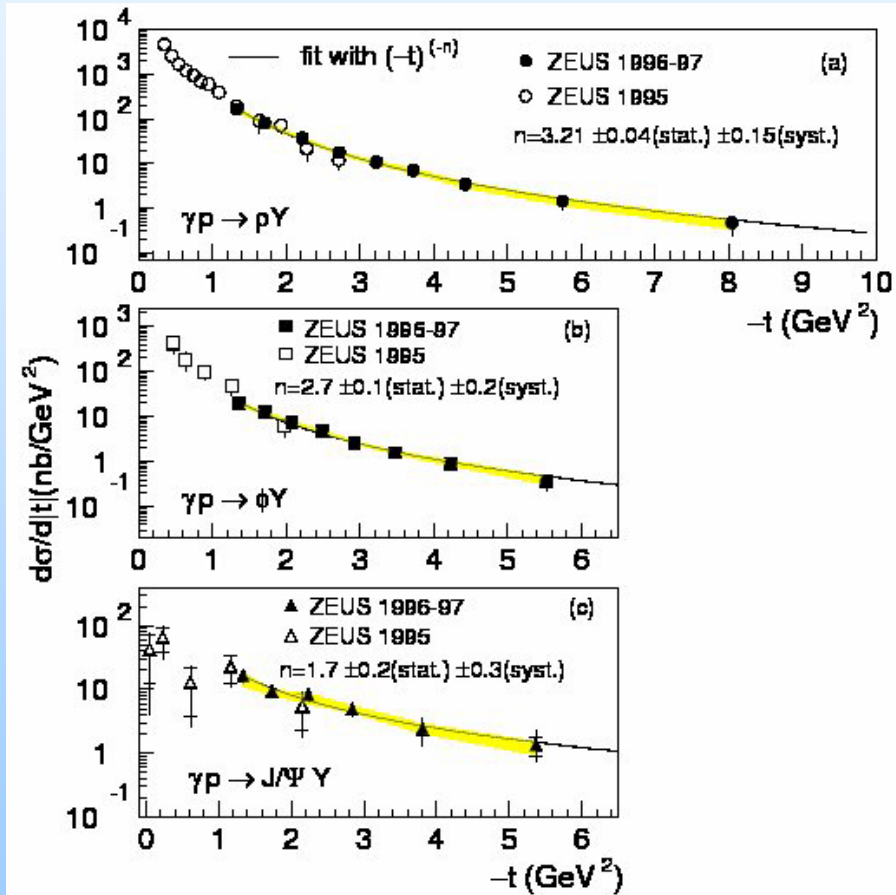
- FKS: 2-gluon exchange, gluons from fit to DIS
- KMW: saturation model, b dependence, DGLAP
- DF: σ_{dip} Wilson loop



Hard scale t:

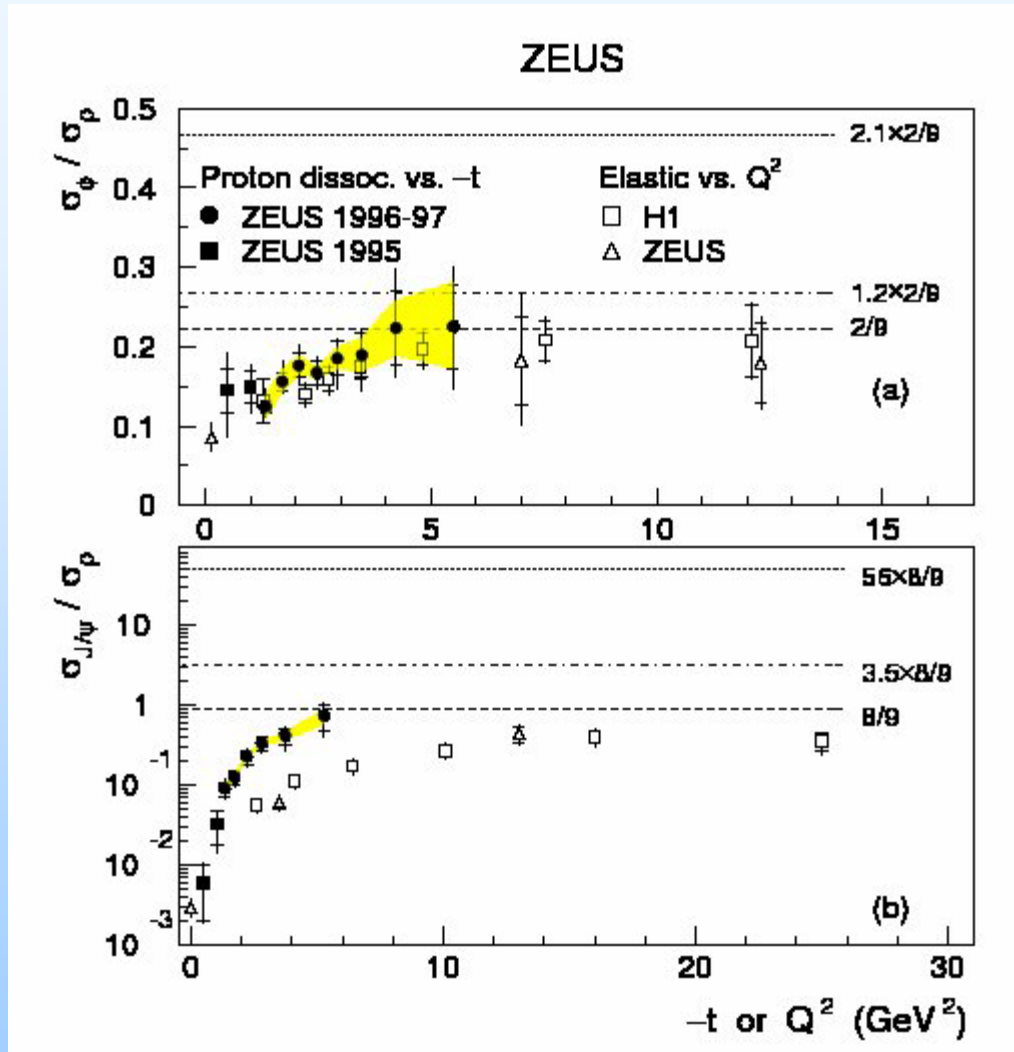
$$\gamma p \rightarrow V Y$$

Reminder: high p_t physics in pp reactions:
power law behavior of $d\sigma/dp_t^2 \sim p_t^{-2n}$



Assume vertex factorization

Flavor restoration at similar values of t and Q^2



- t distribution
proton vertex
 $\gamma^* p \rightarrow V Y$

- Q^2 distribution
photon vertex

$$\gamma^* p \rightarrow V p$$

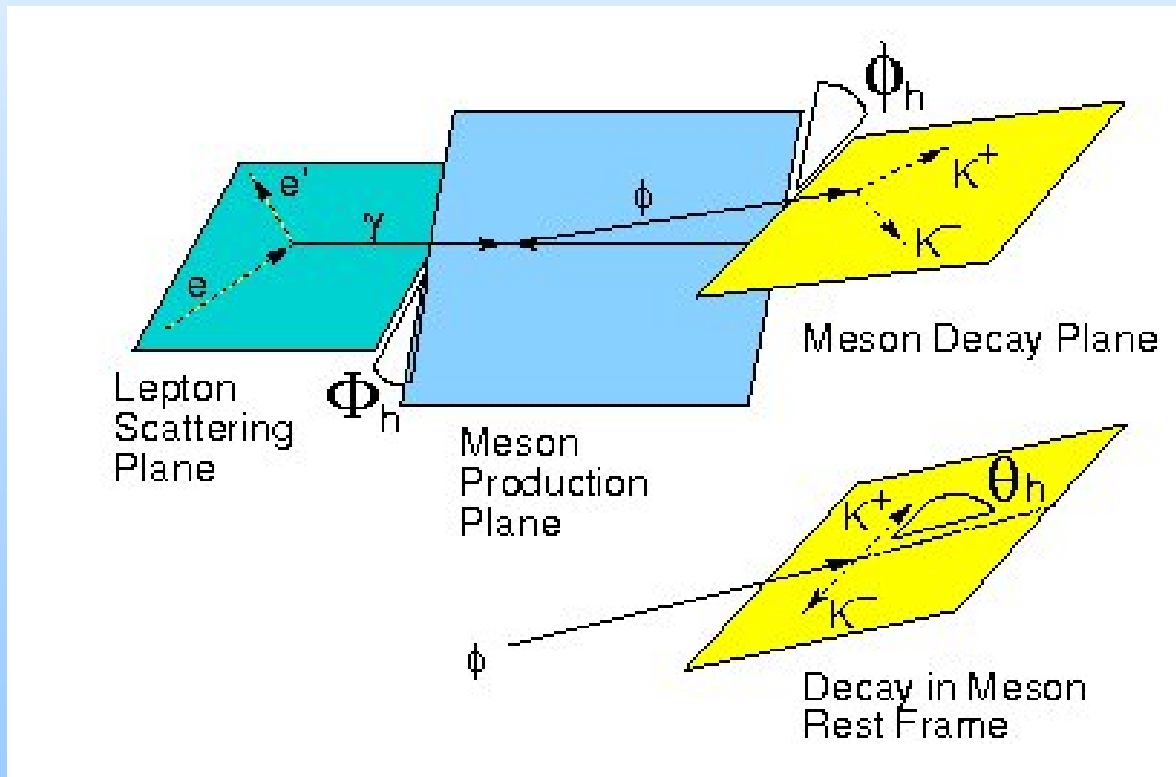
M_V, Q^2, t hard scales \rightarrow pQCD applicable

Helicity amplitudes $T_{\lambda\rho,\lambda\gamma}$:

3 angles, 15 spin density matrix elements, 6 helicity amplitudes:

SCHC: T_{00}, T_{11} single flip: T_{01}, T_{10} double flip T_{1-1}, T_{-11}

pQCD ($|t| < Q^2$) prediction: $T_{-11}, T_{10} < T_{01} < T_{11} < T_{00}$



Spin Density Matrix Elements

$$r_{kl}^{ij}(Q^2)$$

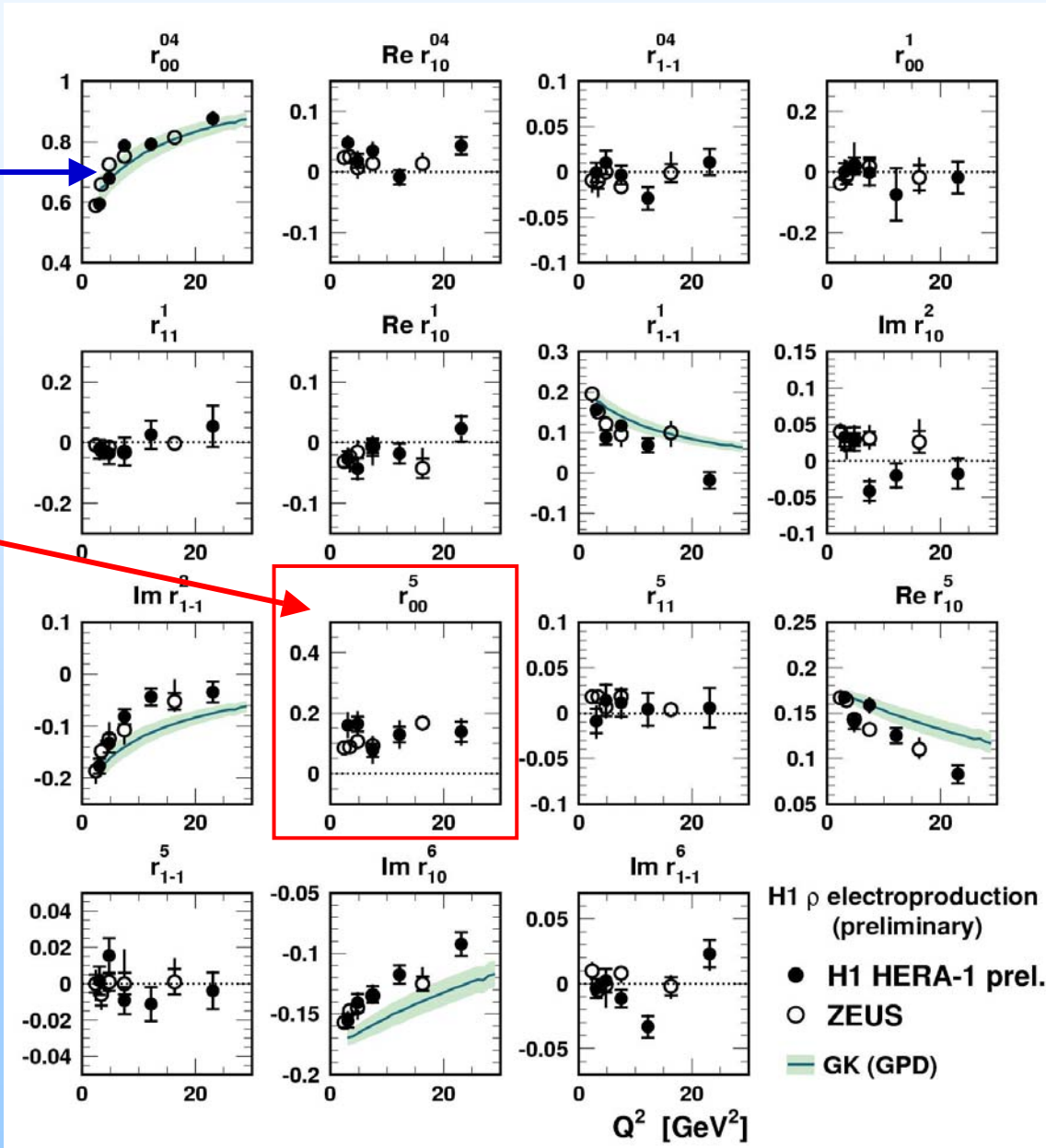
$$t < 3 \text{ GeV}^2$$

- 5 SCHC ME $\neq 0$ agree with GPD (GK) calculation
- other agree with SCHC prediction (---)
- except $r_{00}^5 \sim T_{10}T_{00}^*$

Goloskokov - Kroll (GK) consider skewed GPD

$$\gamma^* p \rightarrow \rho^0 Y$$

Vertex factorization



Helicity amplitudes – t dependence $r_{kl}^{ij}(t)$

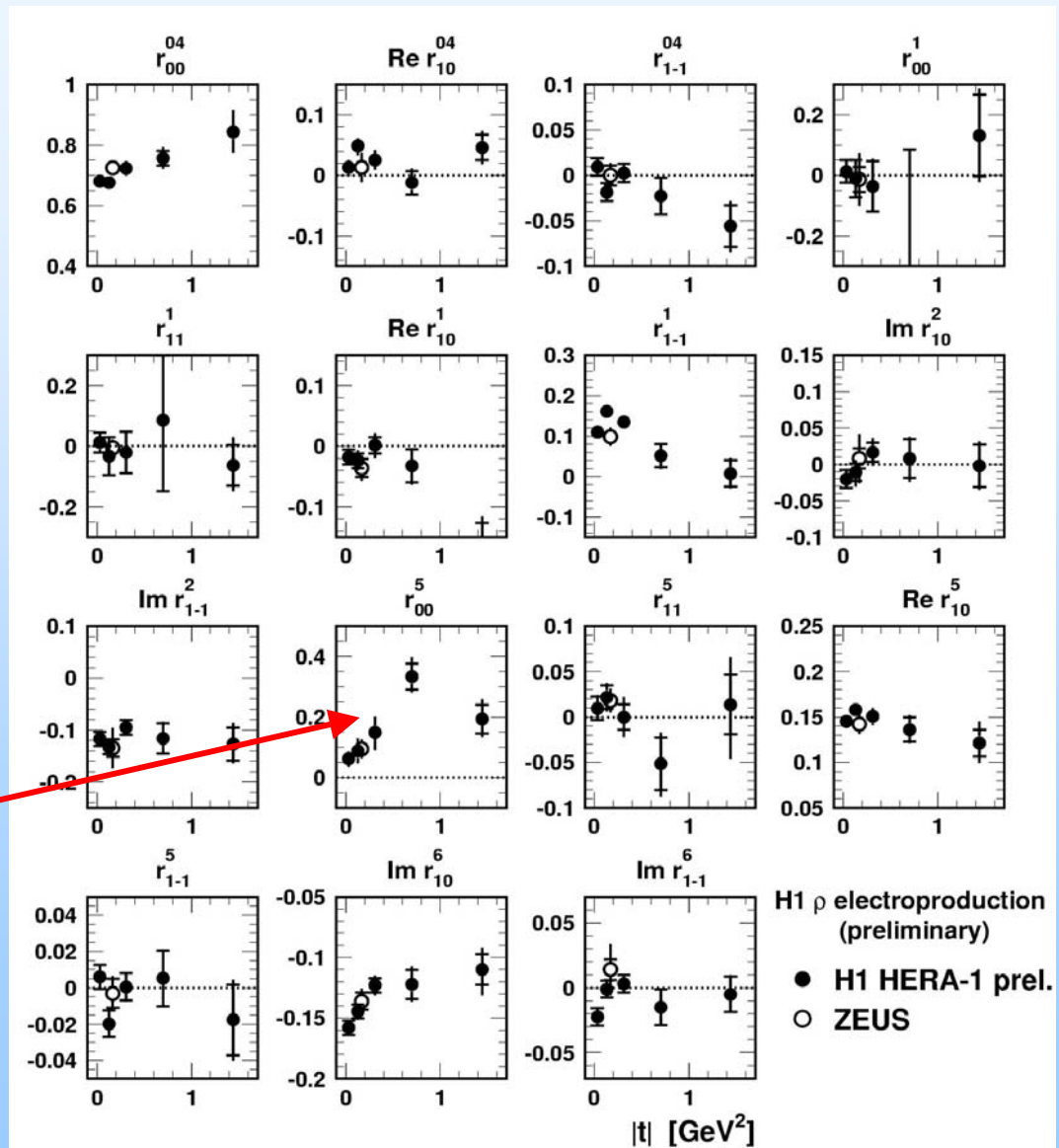
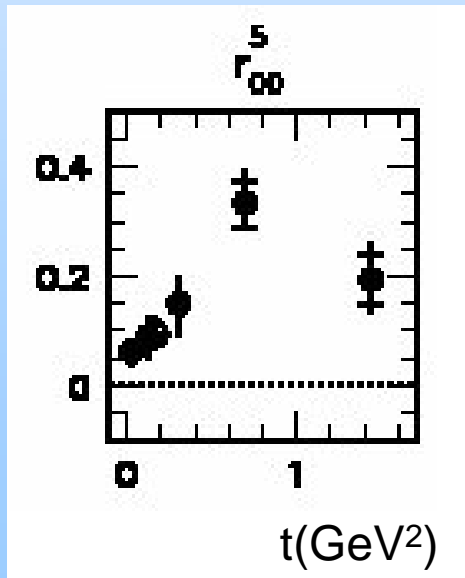


Vertex factorization

SCHC (.....)

single helicity flip

$$r_{00}^5 \neq 0$$



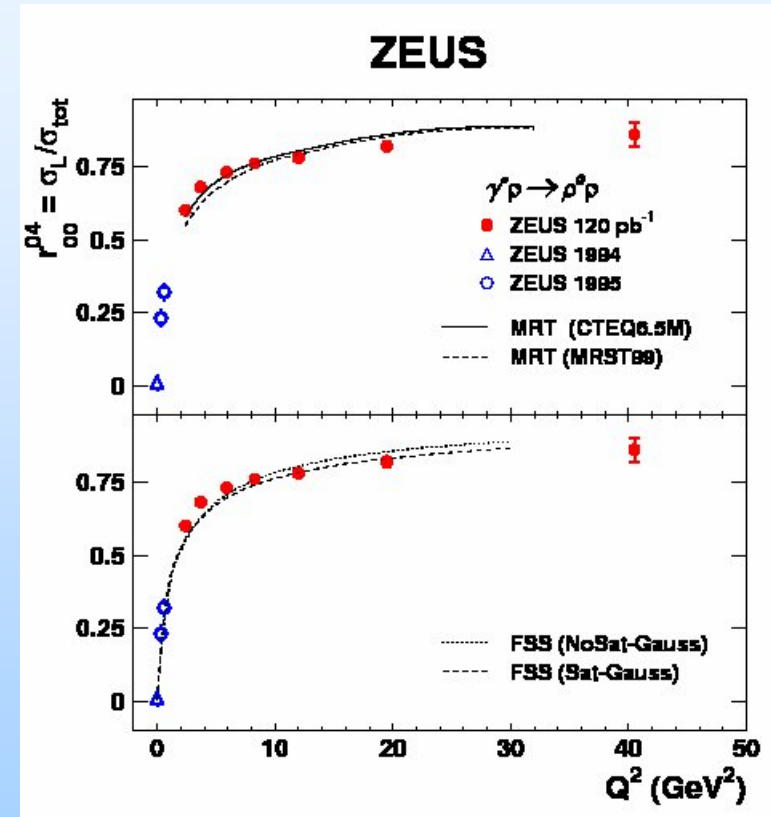
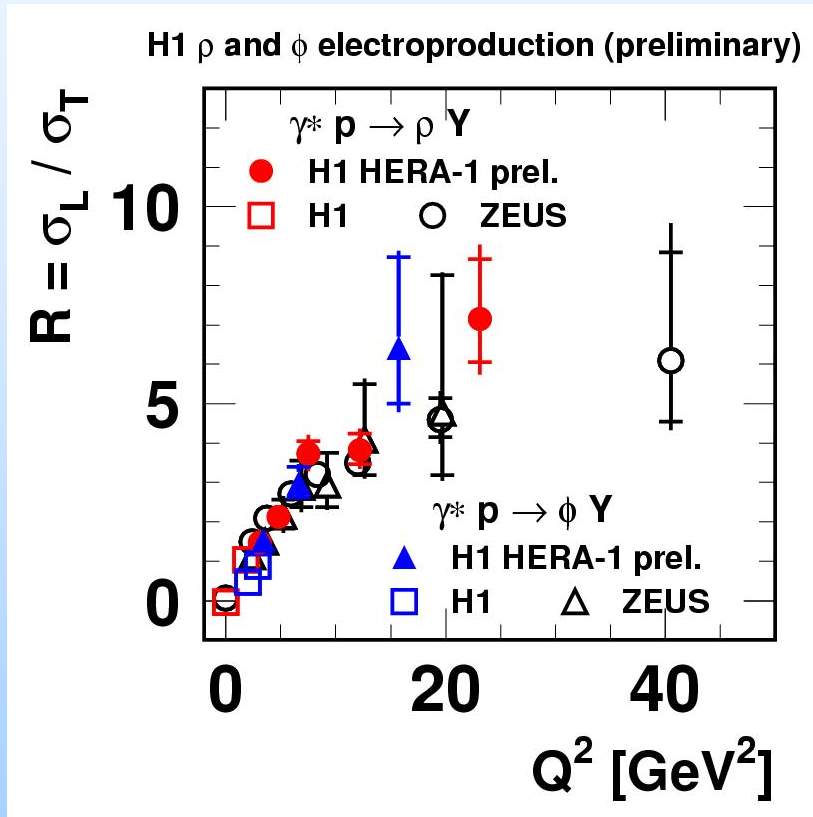
$$\sigma_{tot} = \sigma_T + \varepsilon \cdot \sigma_L$$

$$\langle \varepsilon \rangle = 0.98$$

$$R = \sigma_L / \sigma_T = \varepsilon^{-1} \cdot r_{00}^4 / (1 - r_{00}^4)$$

for $r_{00}^4 \rightarrow 1$

error of R large and asymmetric

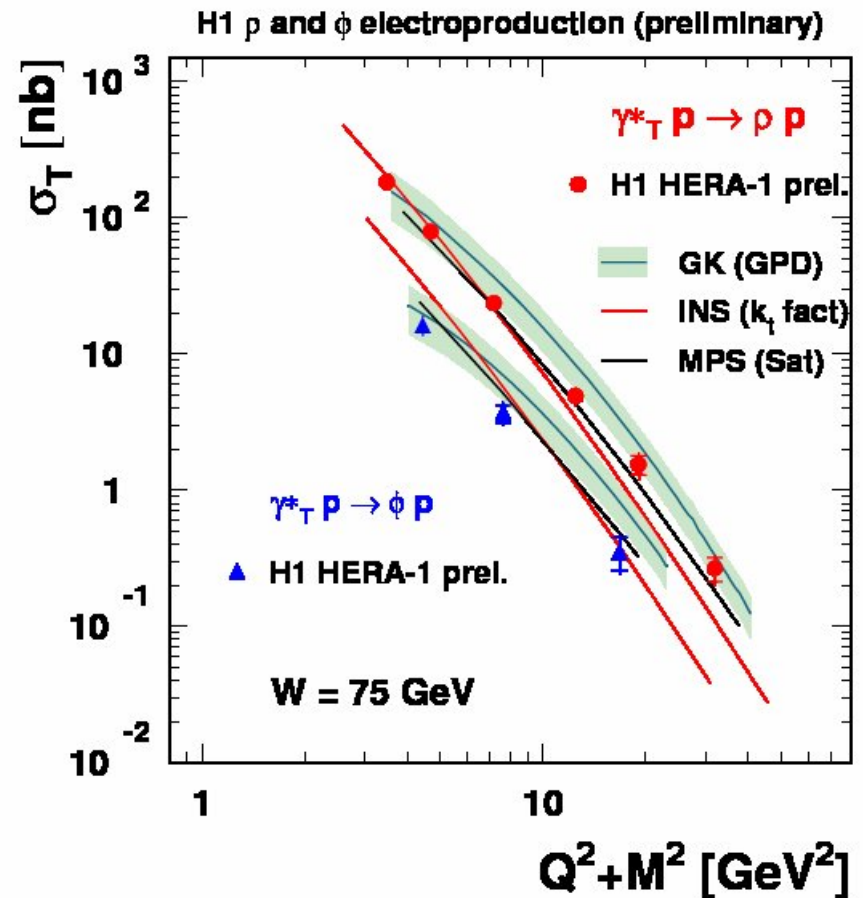
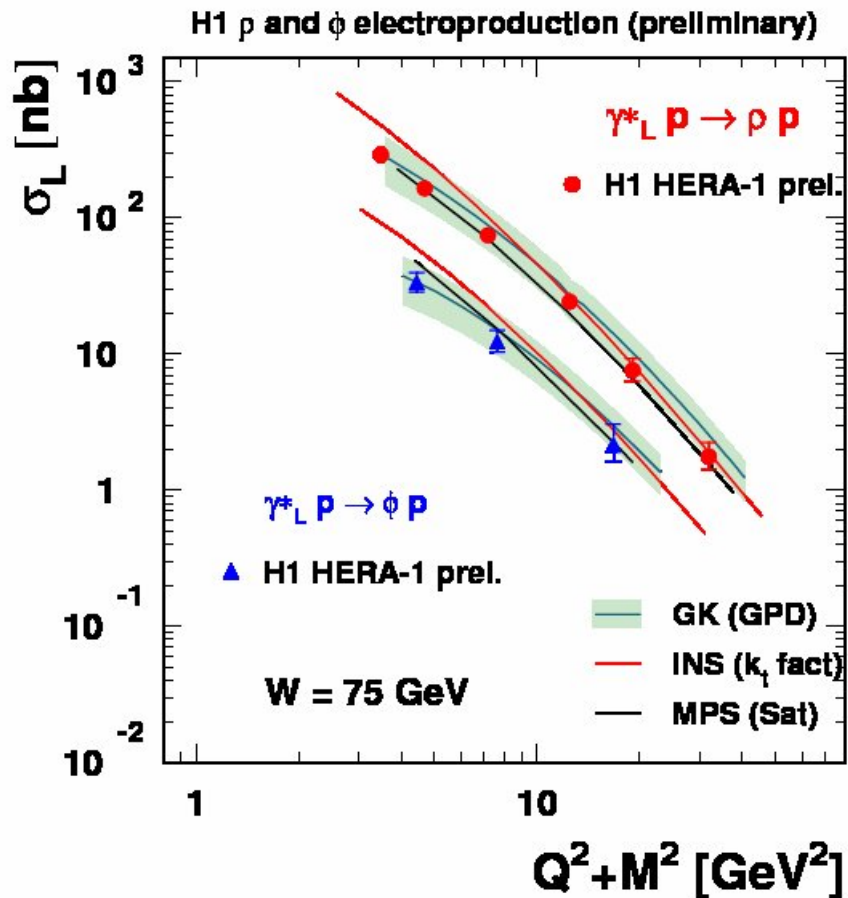


- $R \sim Q^2$
- Leveling off for $Q^2 > 10 \text{ GeV}^2$
- σ_L dominates at large Q^2



- σ_L and σ_T different Q^2 dependence
- $\sigma_L \rightarrow 0$ for $Q^2 \rightarrow 0$ gauge invariance
- σ_L dominates at large Q^2
- GK describes σ_L better than σ_T

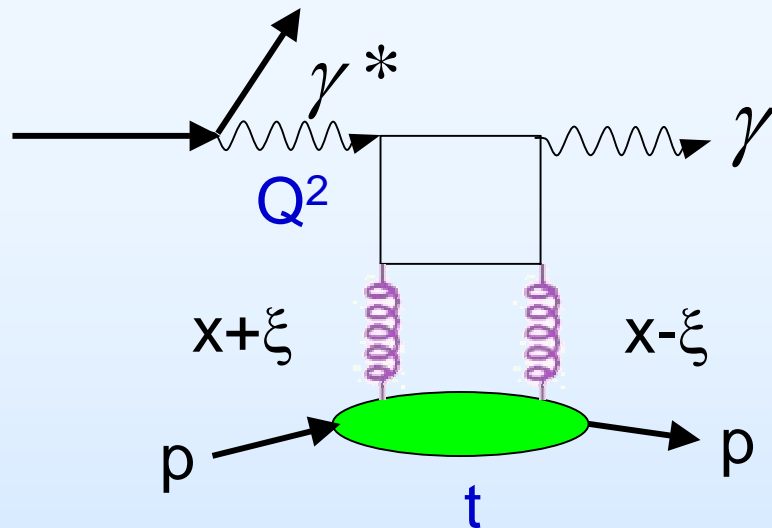
$$\gamma^* p \rightarrow \rho^0 p$$



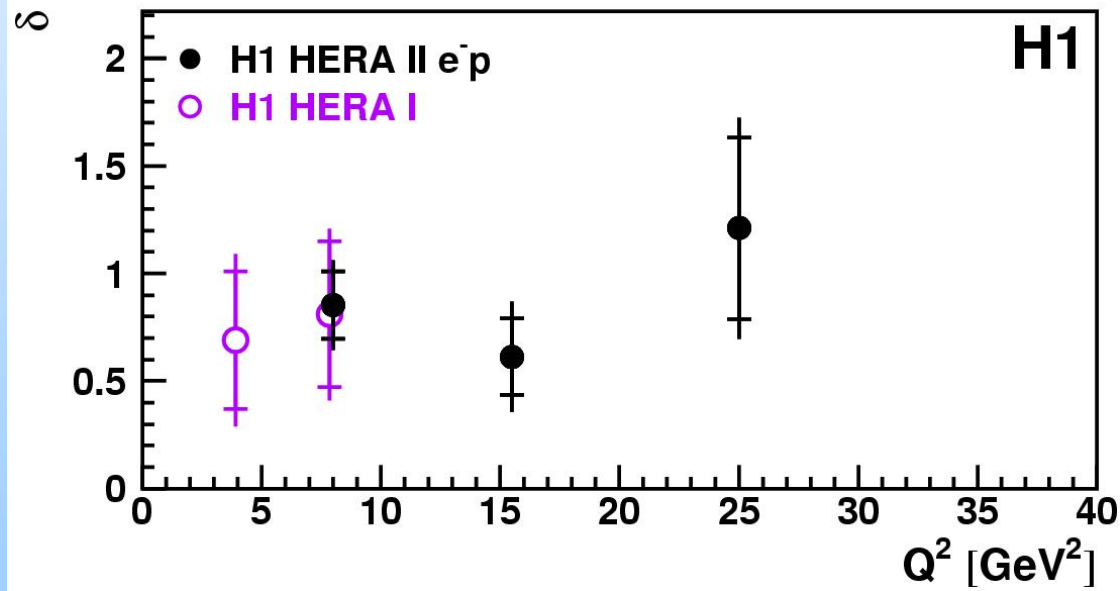
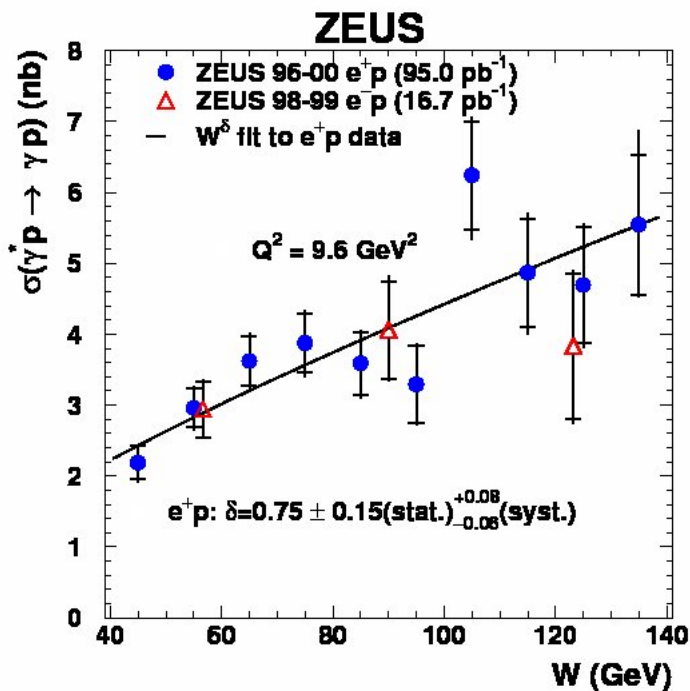
DVCS

$$\gamma^* p \rightarrow \gamma p$$

$$\sigma \sim W^\delta$$



Skewing



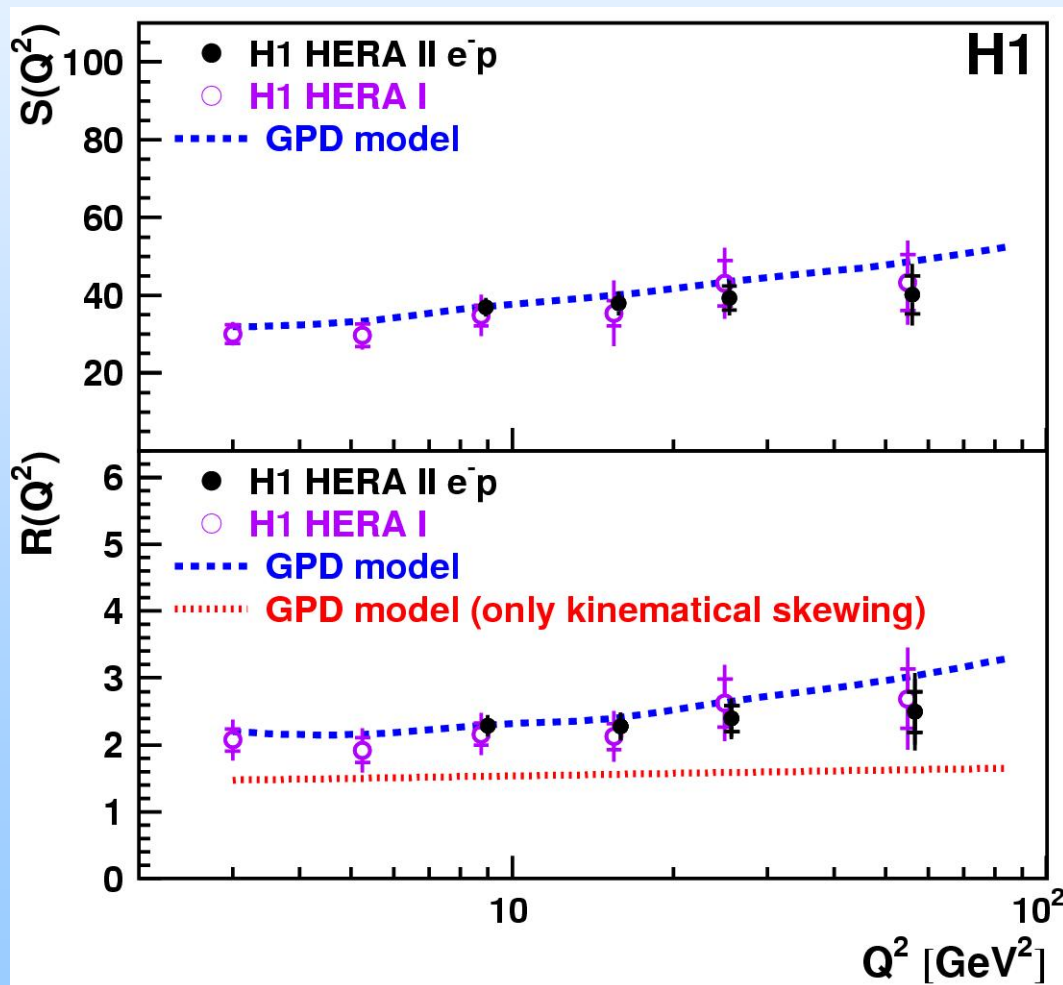
hard process

Dimensionless observables:

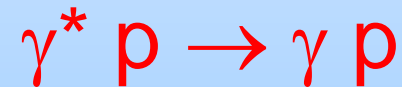
$$S = \sqrt{\frac{\sigma_{DVCS} \cdot Q^4 \cdot b(Q^2)}{1 + \rho^2}}$$

Q^2 -dependence of GPD

$$R(Q^2) = \frac{\text{Im} A(\gamma^* p \rightarrow \gamma p)_{t=0}}{\text{Im} A(\gamma^* p \rightarrow \gamma^* p)_{t=0}} = \frac{\sqrt{\pi \sigma_{DVCS} b(Q^2)}}{\sigma_T(\gamma^* p \rightarrow X) \sqrt{1 + \rho^2}} \sim \frac{GPD}{pdf}$$



- GPD model (A. Freund)
P.R. D68 (2003) 096006

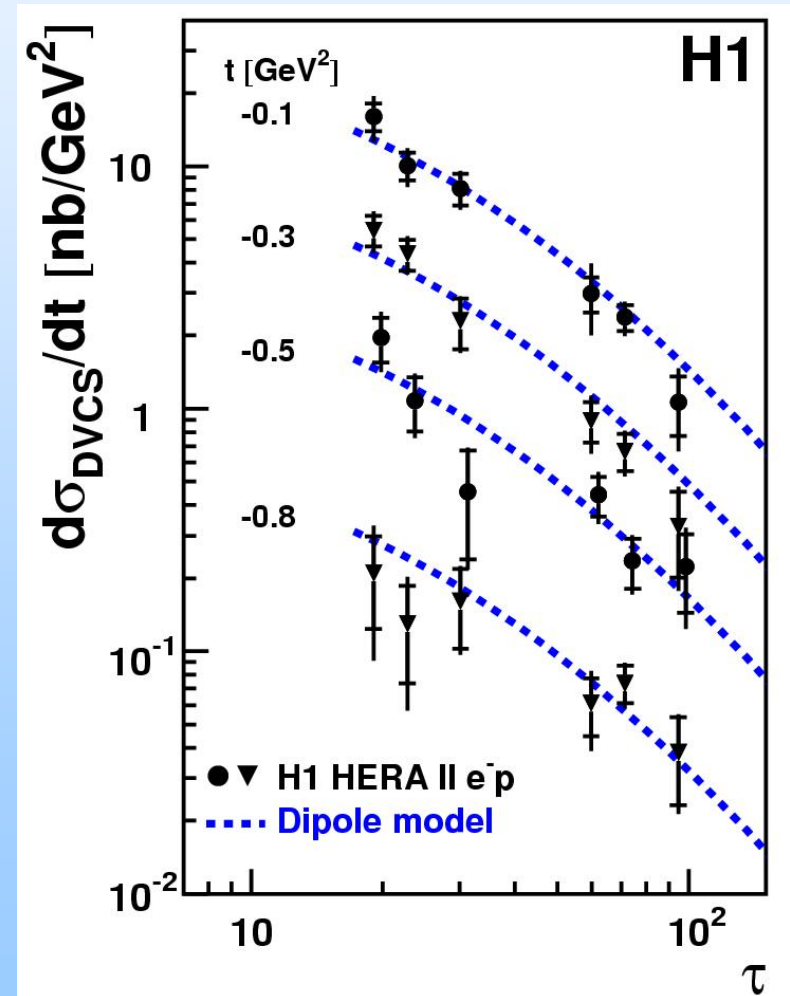
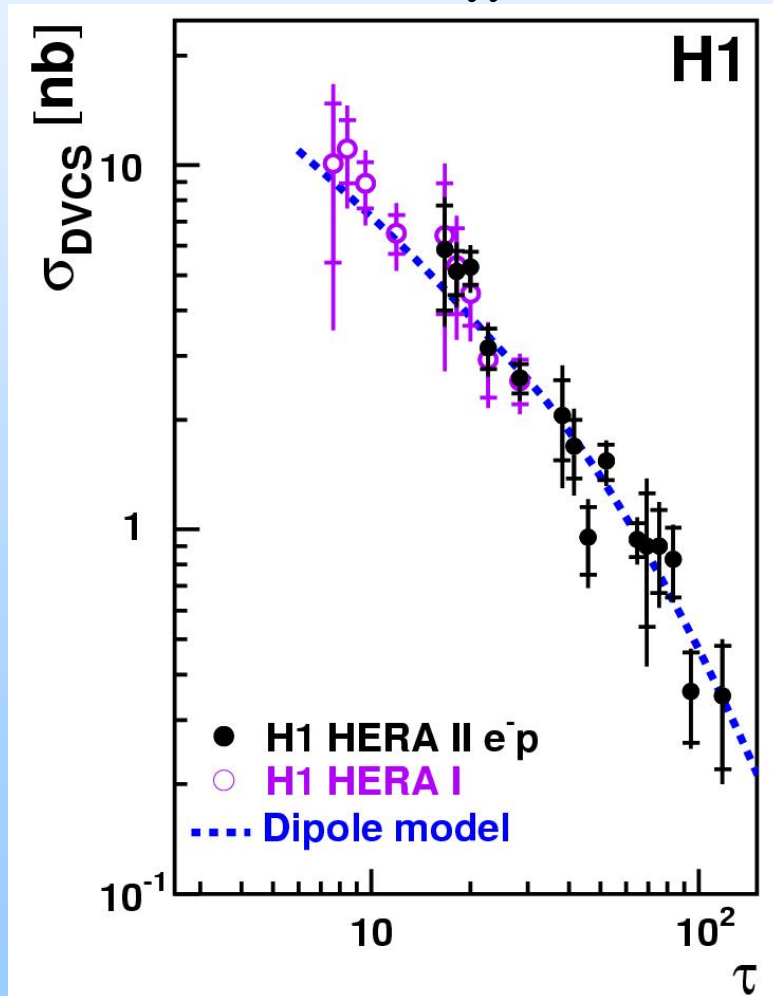


← kinematical skewing not sufficient

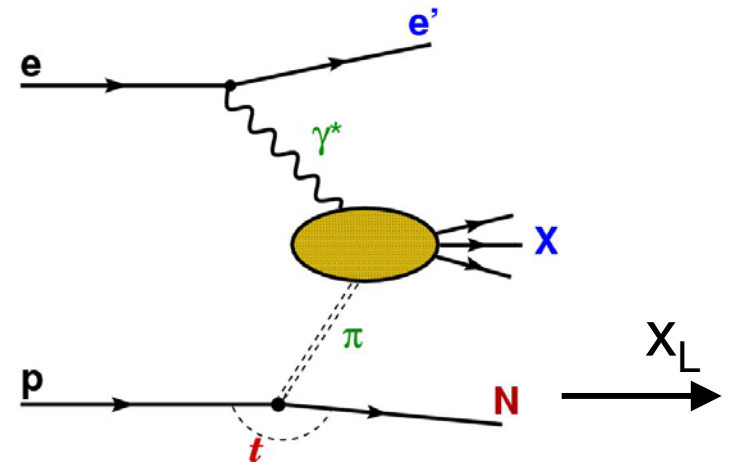
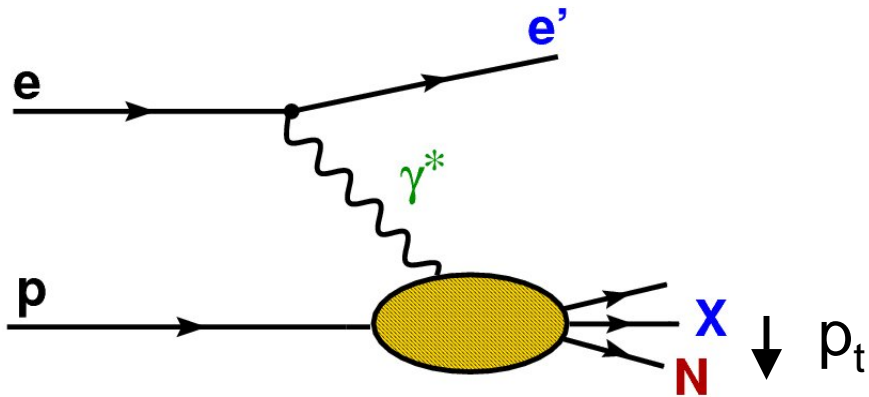
Dipole model predicts geometrical scaling

$$\sigma_{tot}^{\gamma^* p}(x, Q^2) = \sigma_{tot}^{\lambda^* p}(\tau = Q^2 / Q_s^2)$$

$$Q_s(x) = Q_0 \left(\frac{x_0}{x}\right)^{\lambda/2} \quad \text{saturation scale}$$



Leading baryons in $e p \rightarrow e N X$ reactions



- Comparison with standard fragmentation models

- Limiting fragmentation

$$\frac{d^2\sigma}{dx_L dp_t^2}(W^2, Q^2, x_L, p_t^2) = g(x_L, p_t^2)G(W^2, Q^2)$$

compare $\gamma p \rightarrow NX$ with

$$\gamma^* p \rightarrow NX$$

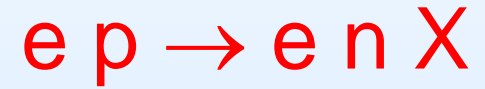
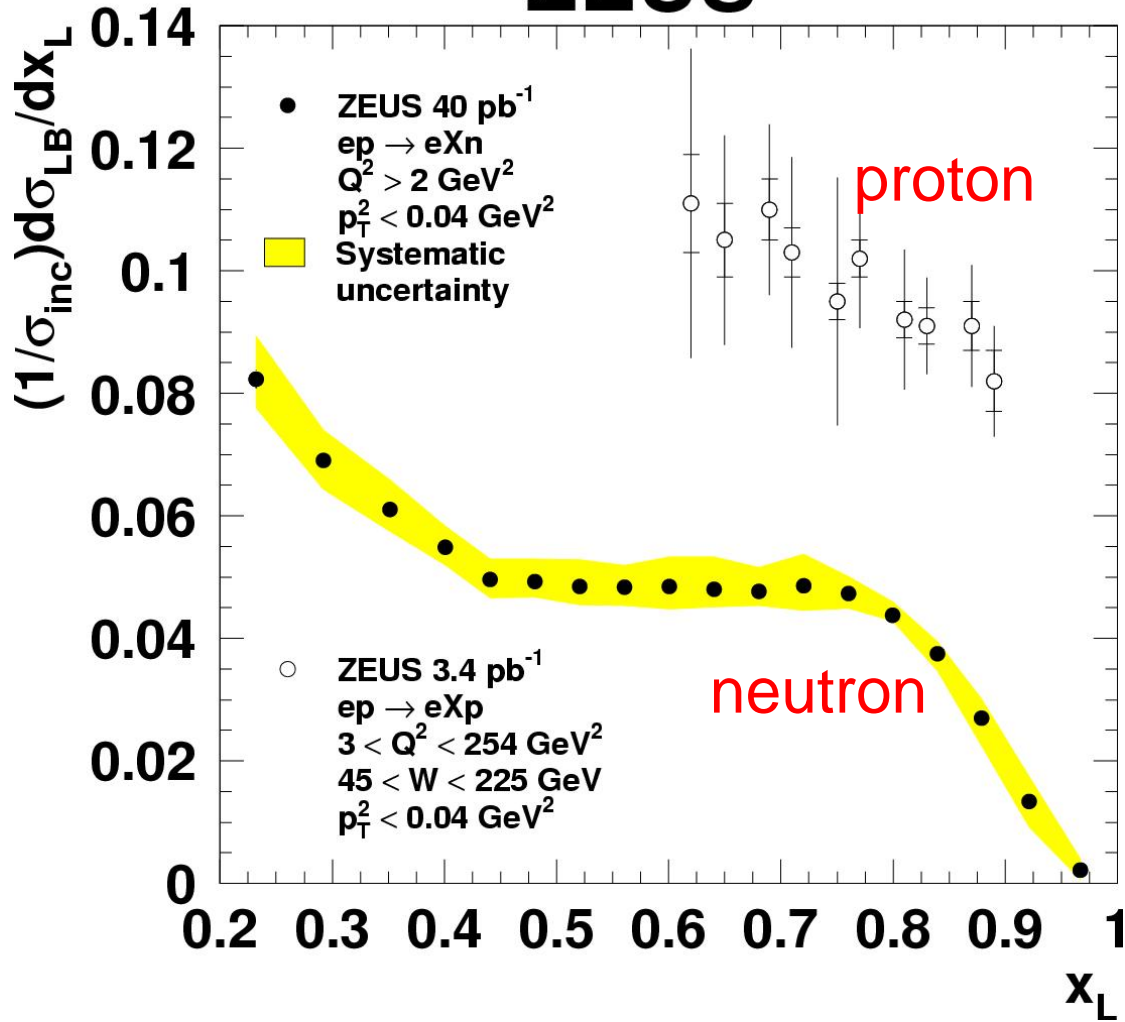
- π exchange, factorization

$$\frac{d^2\sigma(W^2, Q^2, x_L, t)}{dx_L dt} = f_{\pi/p}(x_L, t) \cdot \sigma_{\gamma^* \pi}((1-x_L)W^2, Q^2)$$

- $F_2^\pi(x, Q^2)$?

- absorption / migration

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$$r = \frac{\sigma(ep \rightarrow epX)}{\sigma(ep \rightarrow enX)} \approx 2$$

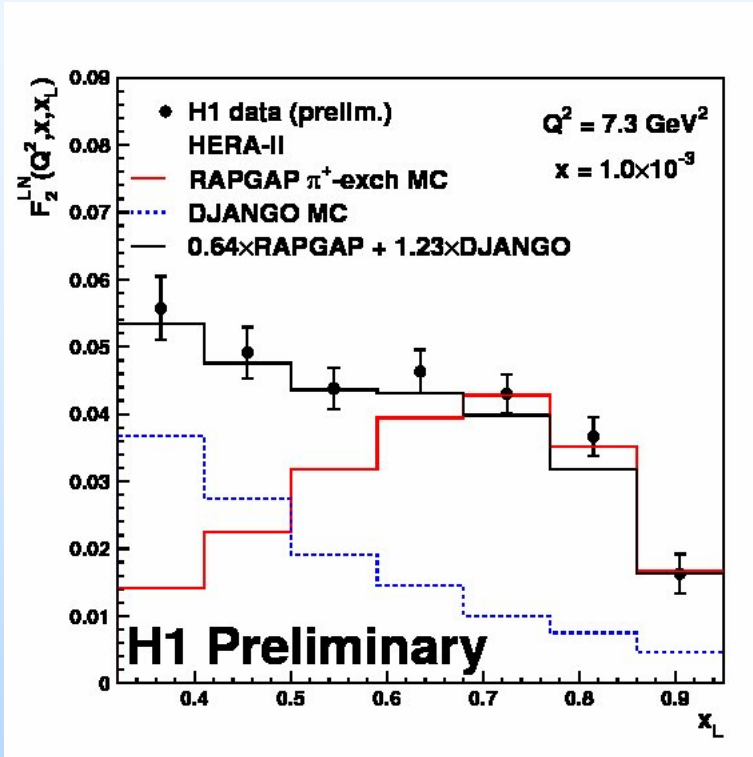
isospin 1 exchange:

$$r = \frac{1}{2}$$

$$\frac{1}{\sigma_{incl}} \cdot \frac{d^2\sigma_{LN}}{dx_L \cdot dp_T^2} = a(x_L) \cdot \exp(-b(x_L)p_T^2)$$

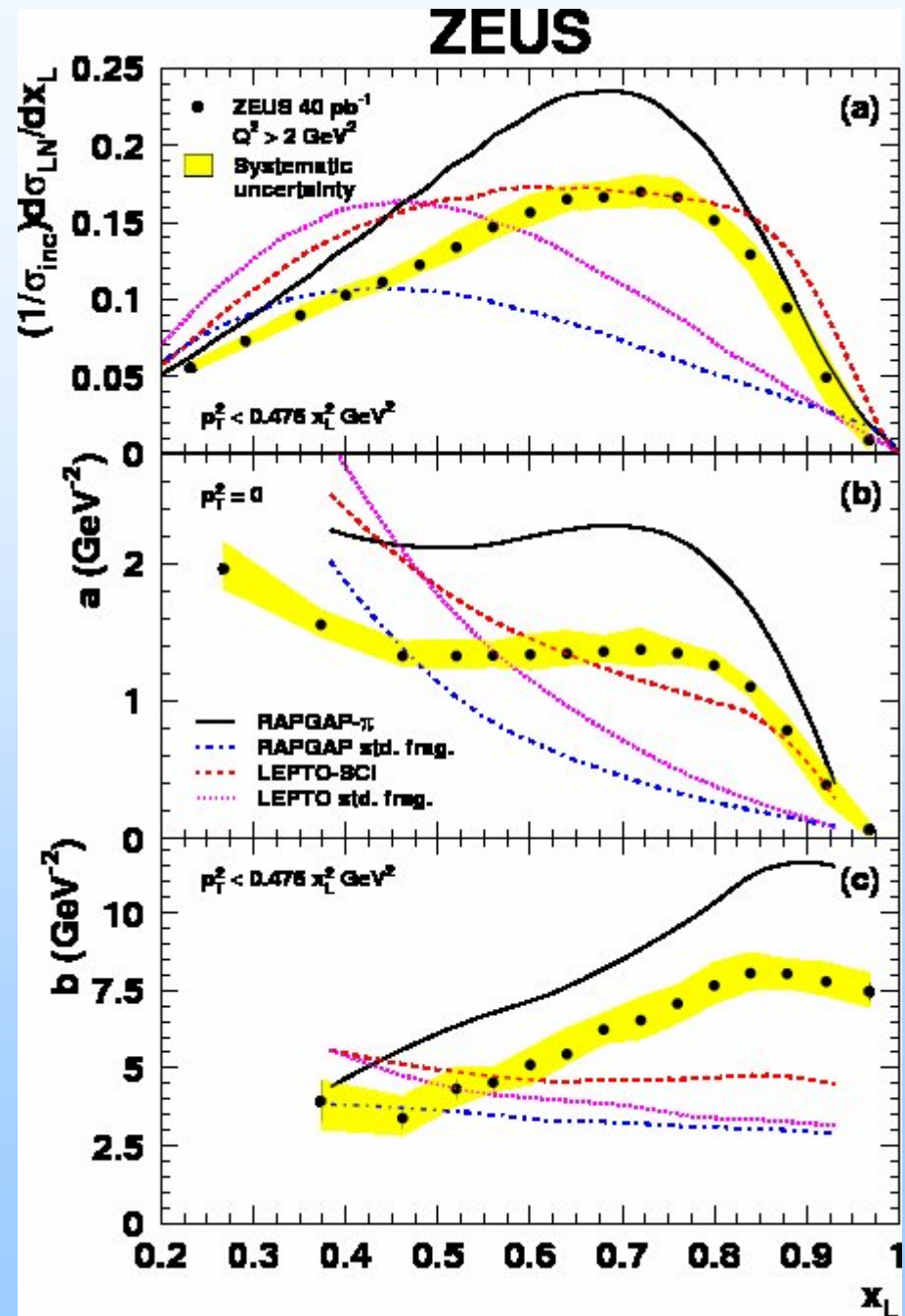
$e p \rightarrow e n X$

Comparison with models:



- all fragmentation models fail
- best mixture of Django + RAPGAP π -exchange

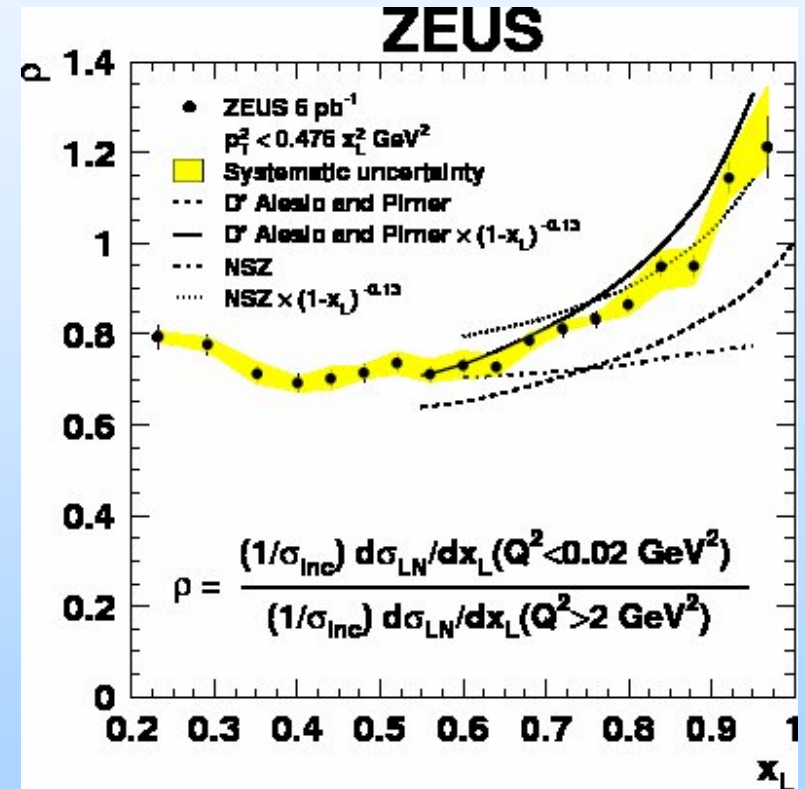
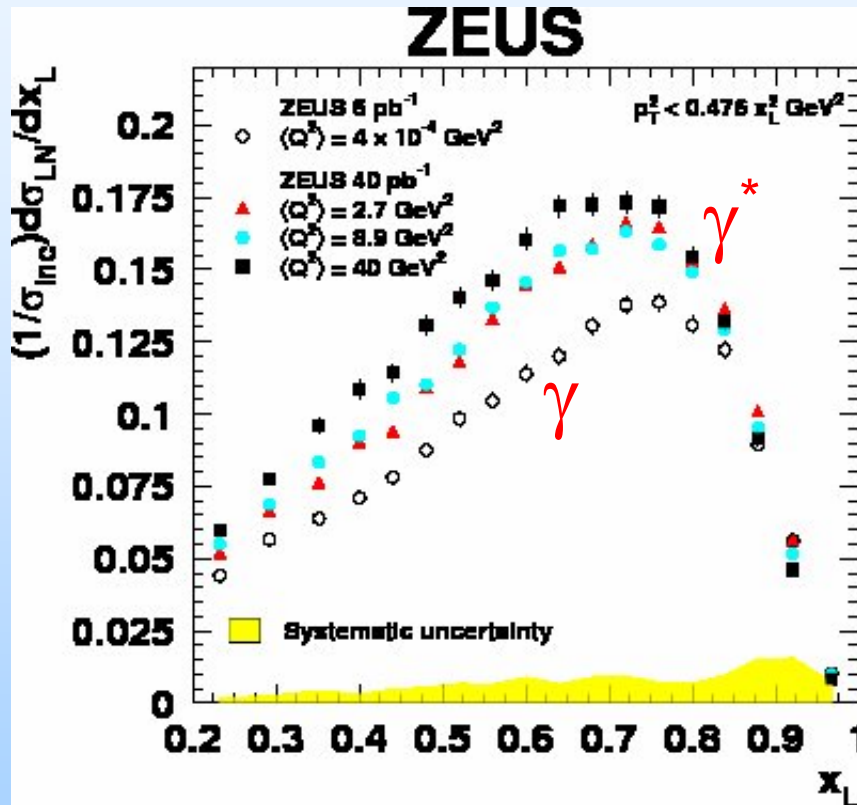
$$\Rightarrow F_2^\pi(x, Q^2)$$



Absorption / migration effects

$$e p \rightarrow e n X$$

Compare photoproduction / DIS



Absorption large

- large photon size $Q^2 \approx 0$
depletion for photoproduction

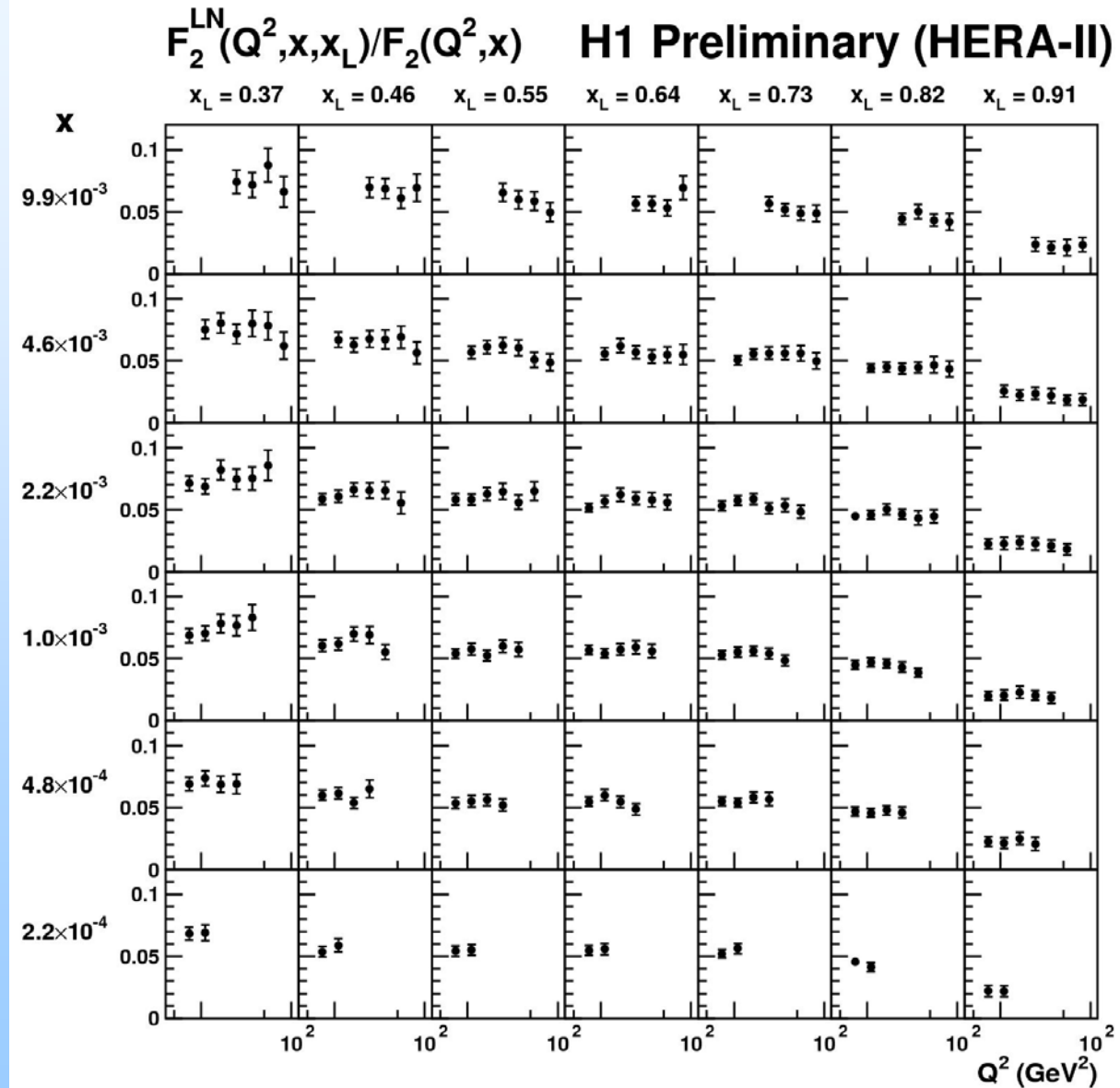
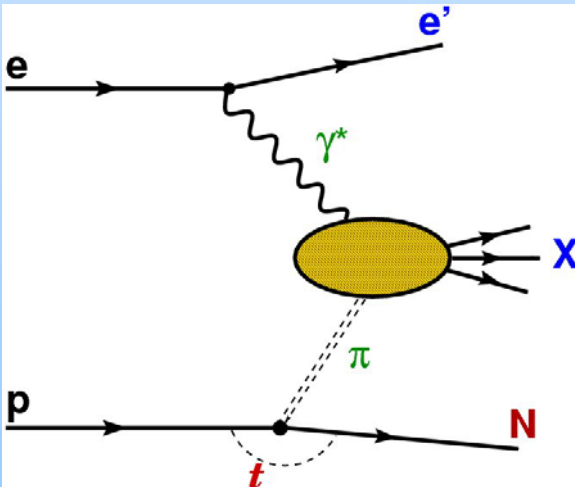
- absorption models
- migration $x_L < 0.5$

e p → e n X

- $$\frac{F_2^{LN(3)}(x, Q^2, x_L)}{F_2(x, Q^2)}$$

independent of x, Q^2

- factorization of vertices

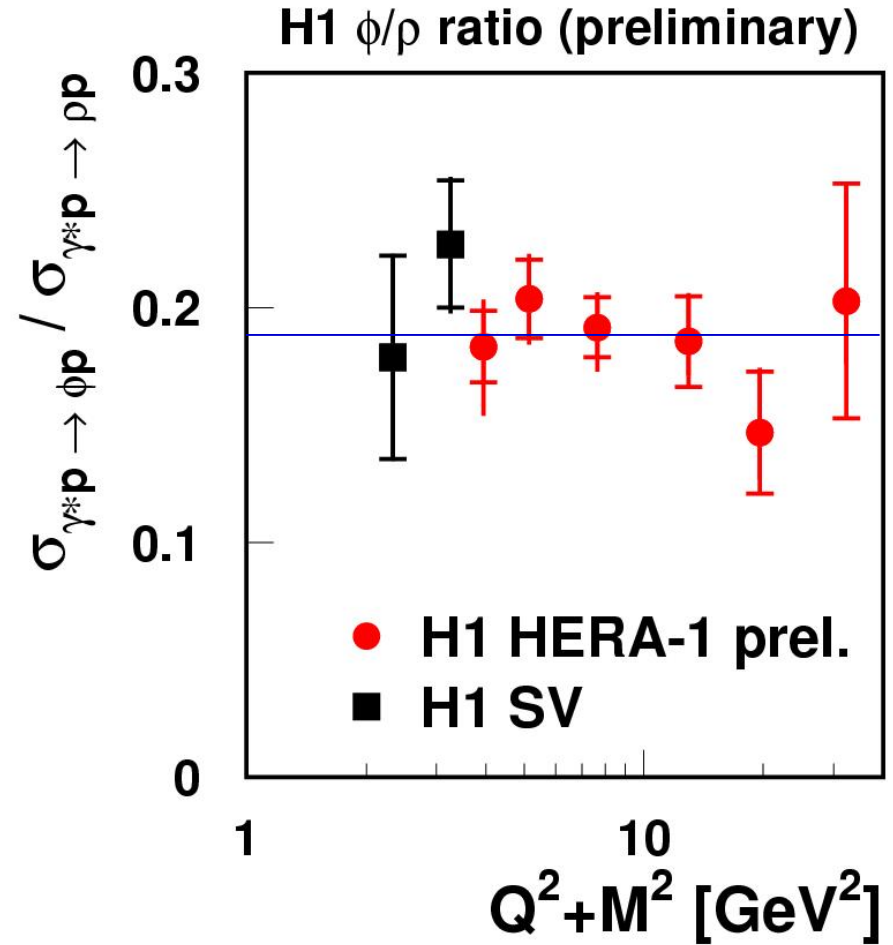
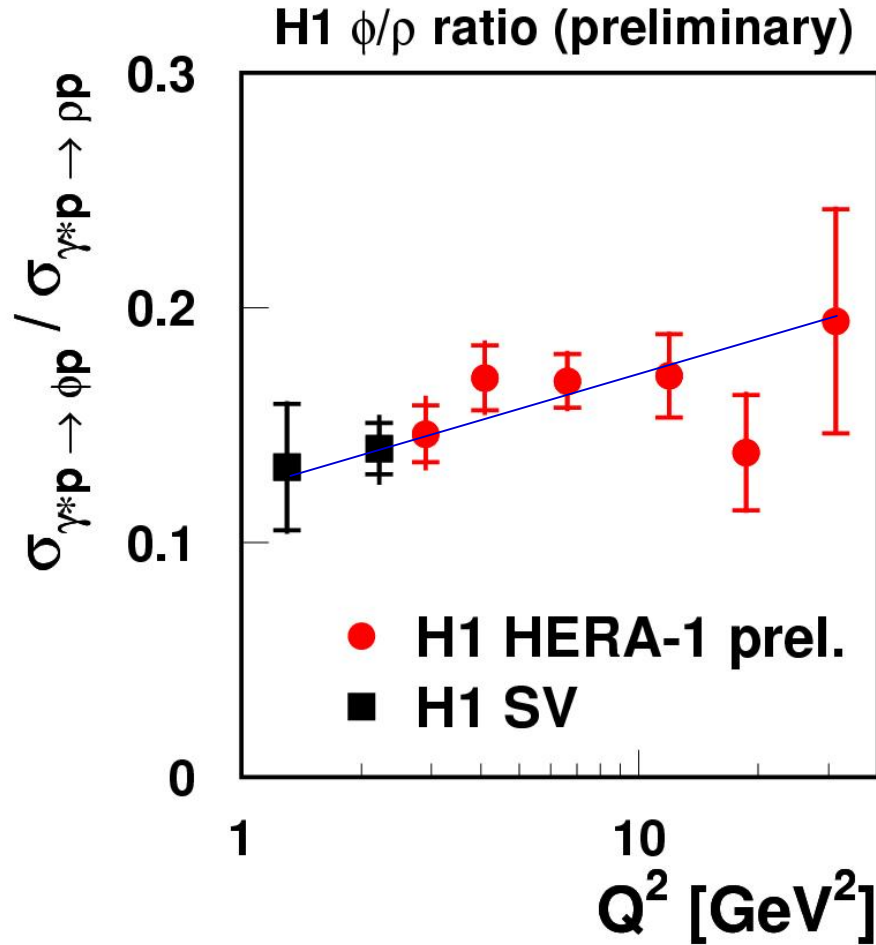


Summary

- $ep \rightarrow eVp$ hard process: $M_V = M_\Psi, M_\Upsilon; Q^2, t$ large
- described by
 - dipole
 - 2-gluon exchange
 - GPD} models
- constrain gluon structure function at small x
- improved theoretical calculations needed
- **leading particles** : $ep \rightarrow eNX$ observed
- standard fragmentation models fail
- violated: vertex factorization / limited fragmentation
- absorption/migration effects observed
- π structure function estimated

BACK UP :

$$\gamma^* p \rightarrow V p$$



$Q^2 + M_V^2$ hard scale

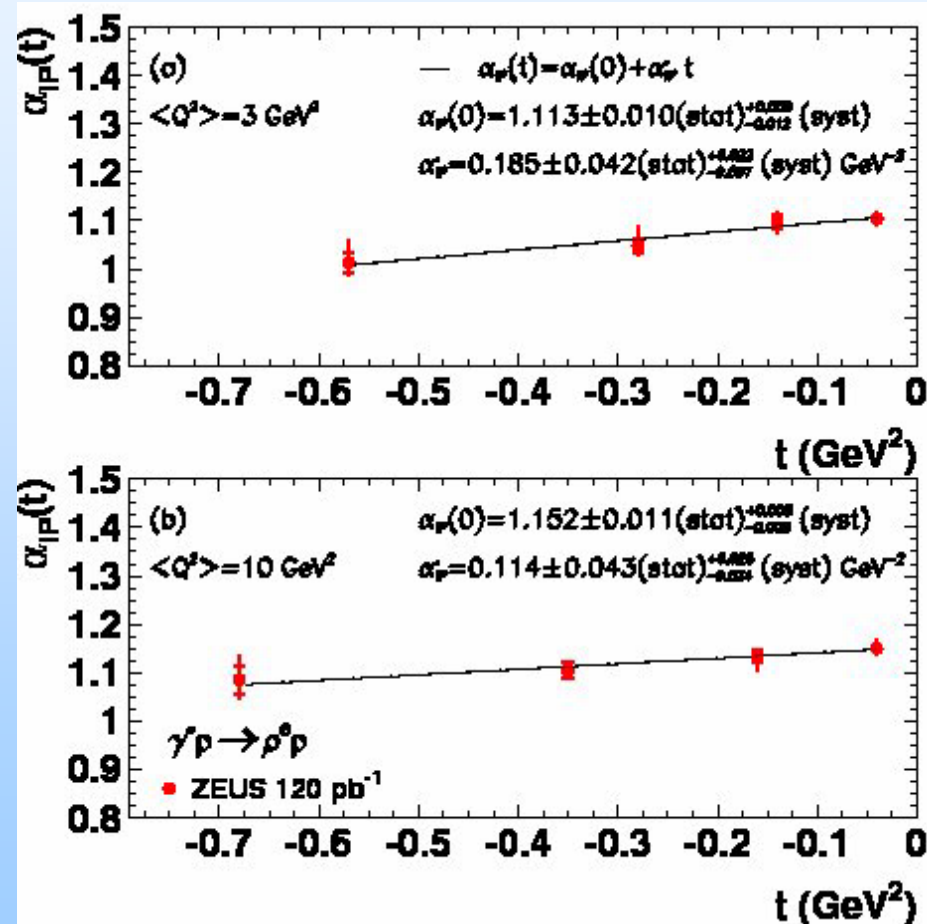
Regge inspired description



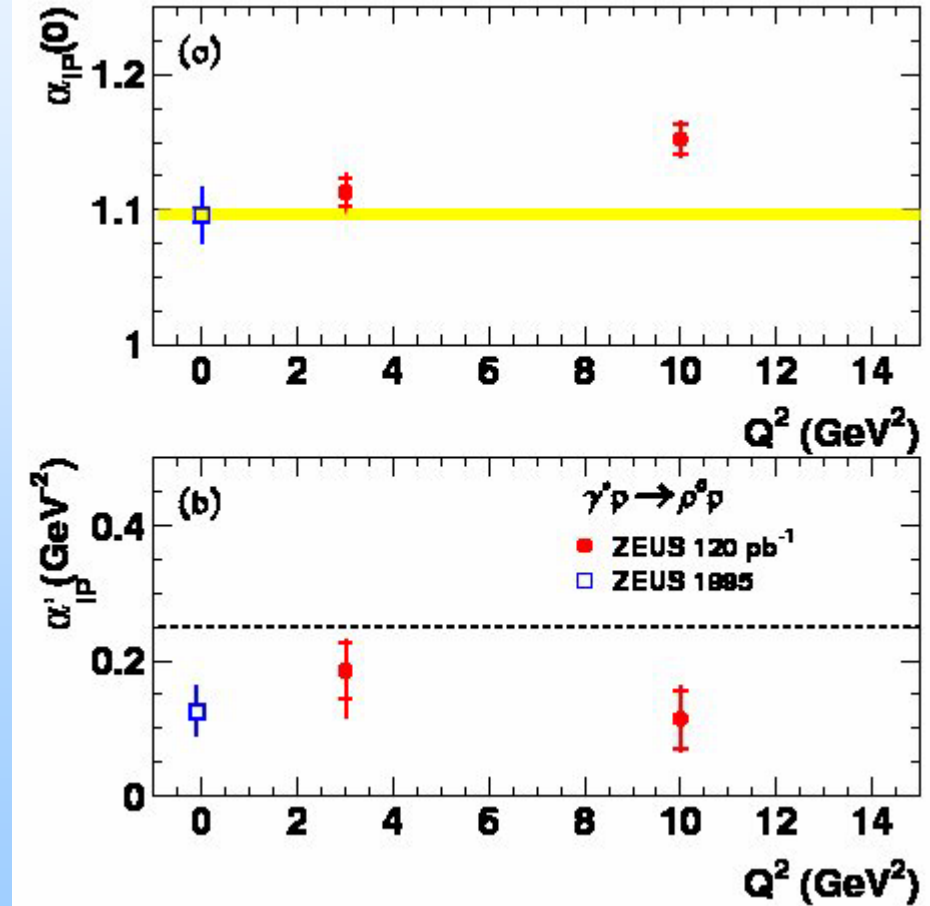
$$\frac{d\sigma}{dt} \sim F(t) \cdot W^{4 \cdot (\alpha(t) - 1)}$$

$$\alpha_p(t) = \alpha_p(0) + \alpha'_p \cdot t$$

W dependence analyzed for $t = \text{const}$



$\alpha_p(0)$ grows with Q^2

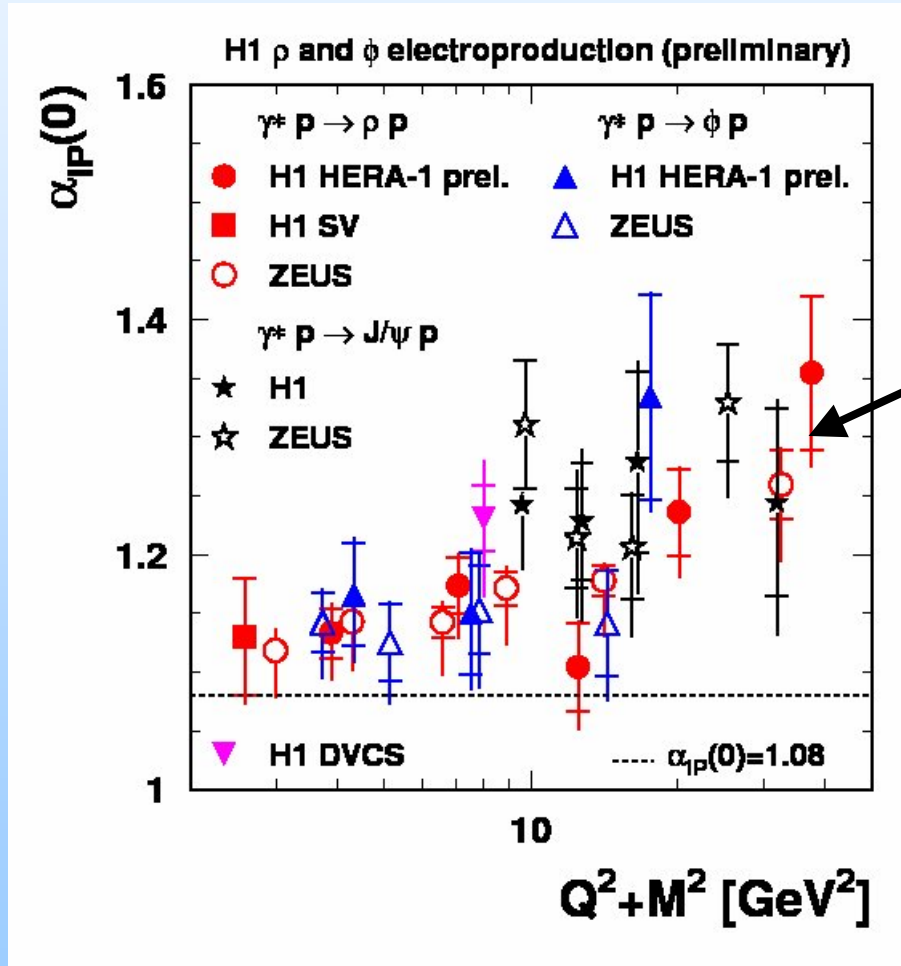


α'_p independent of Q^2 ?

Regge inspired description:

$$d\sigma/dt \sim F(t) W^{4(\alpha(t) - 1)}, \text{ analyzed for } t = \text{const}$$

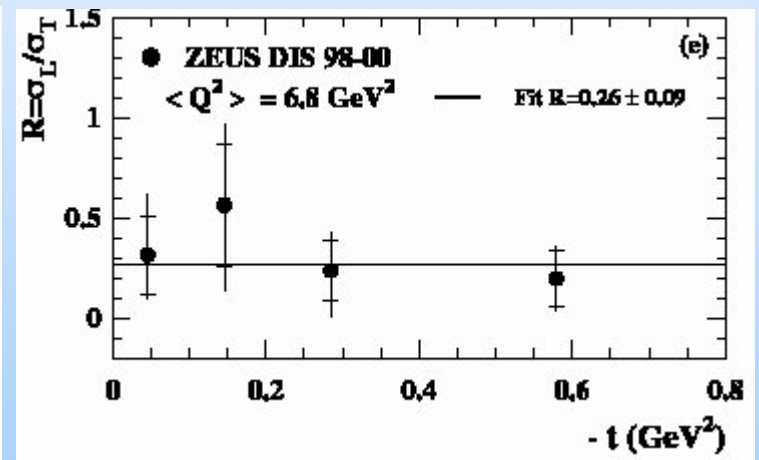
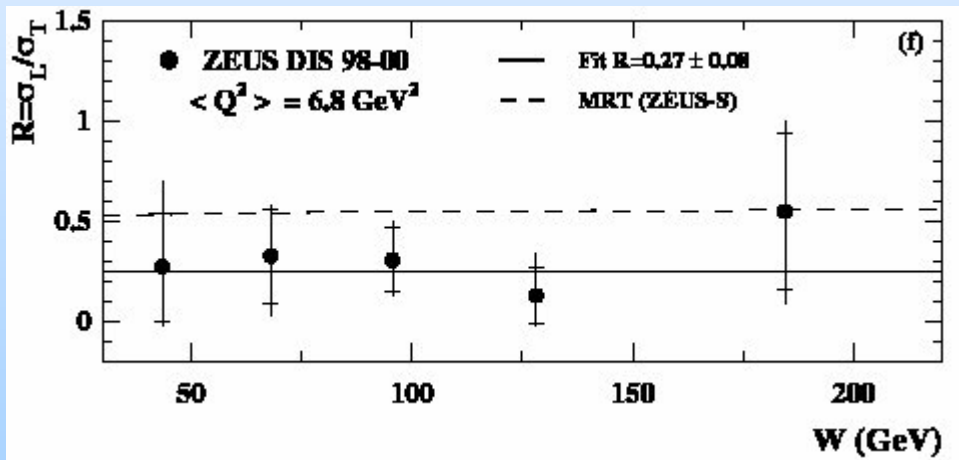
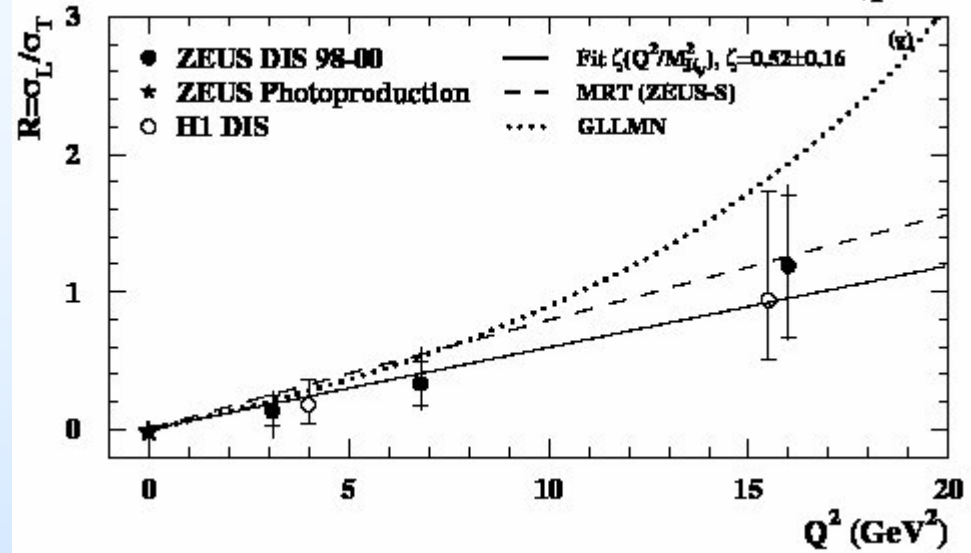
$$\alpha_{\text{IP}}(t) = \alpha_{\text{IP}}(0) + \alpha'_{\text{IP}} t$$



hard Pomeron

soft Pomeron

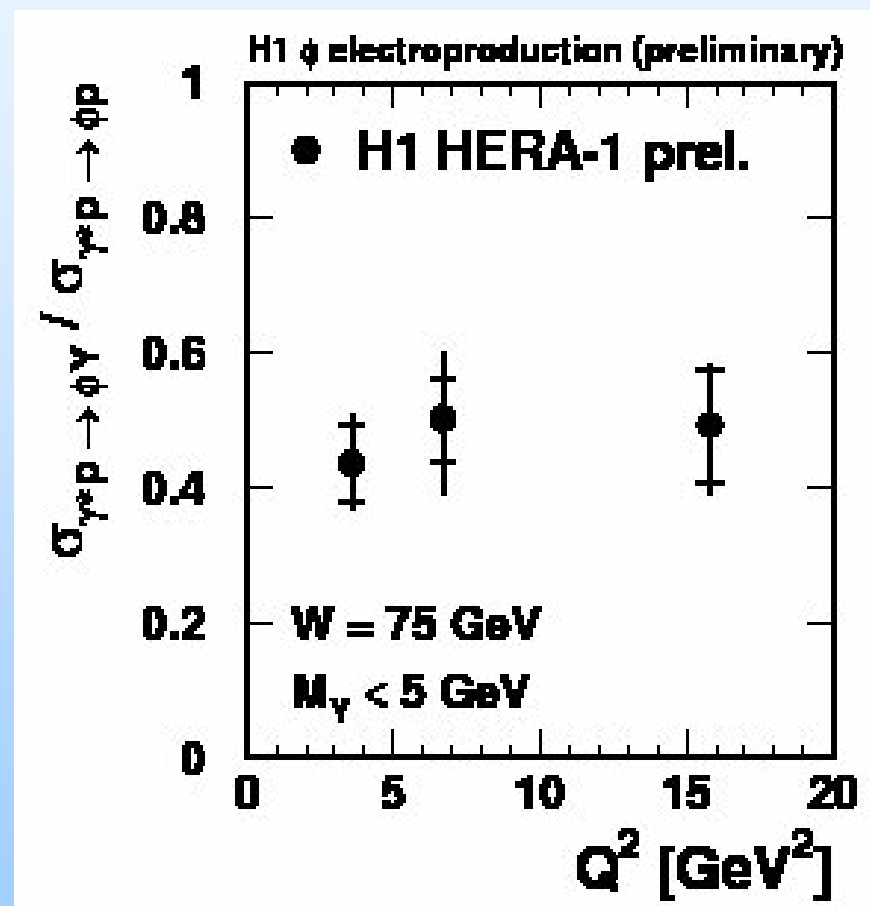
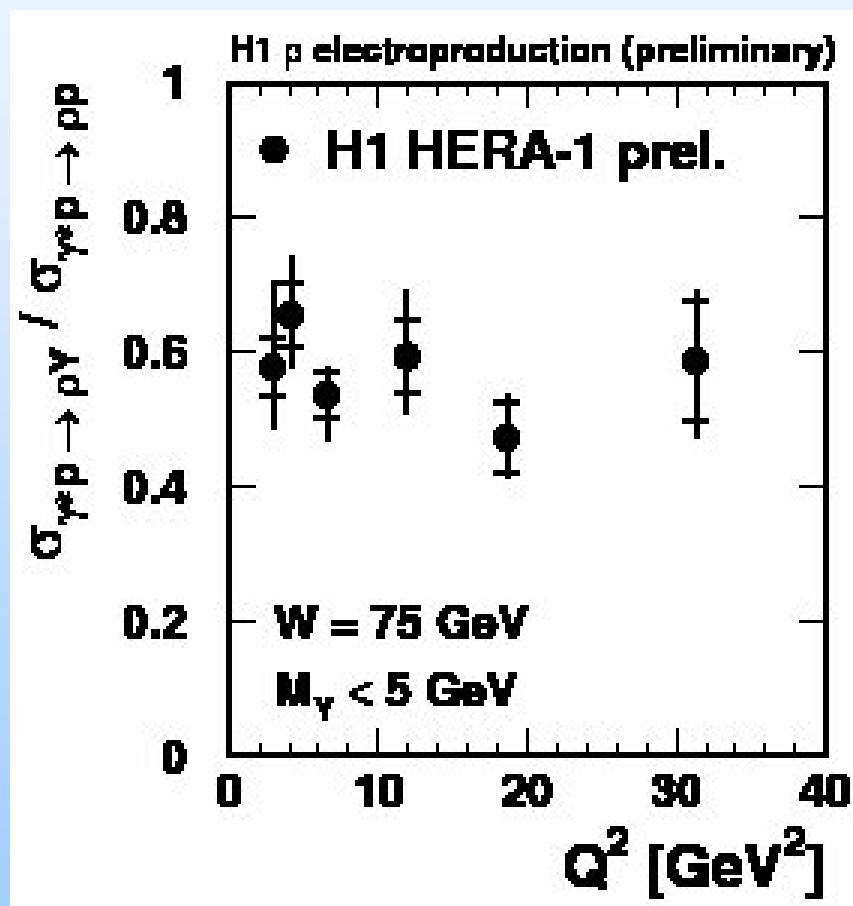
Scaling in $Q^2 + M_V^2$



- $R \sim Q^2/M^2 \rightarrow$ slower increase for Ψ than ρ , Φ
- no W , t dependence of R

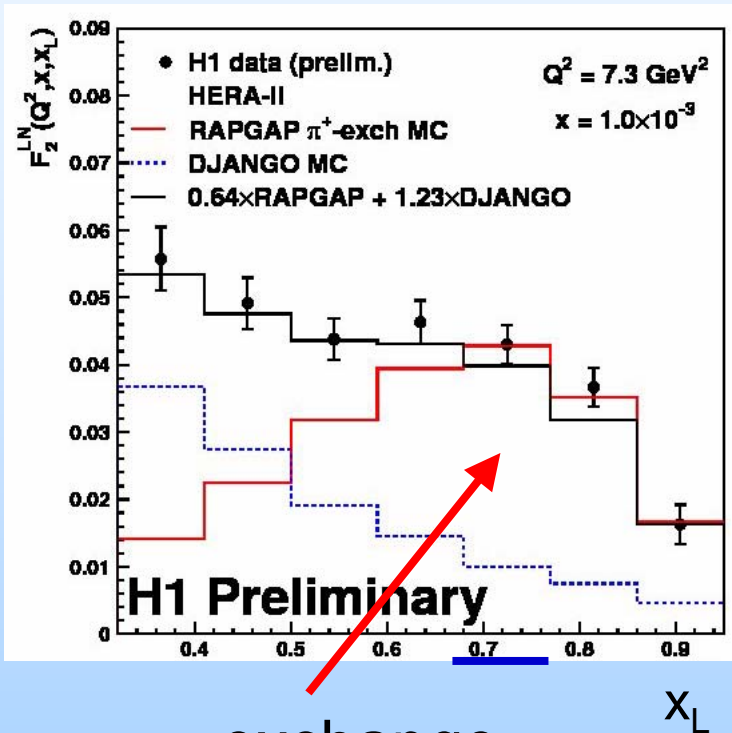
Vertex factorization:

elastic / proton dissociation: universality of Q^2 , W dependence, helicity amplitudes



π structure function:

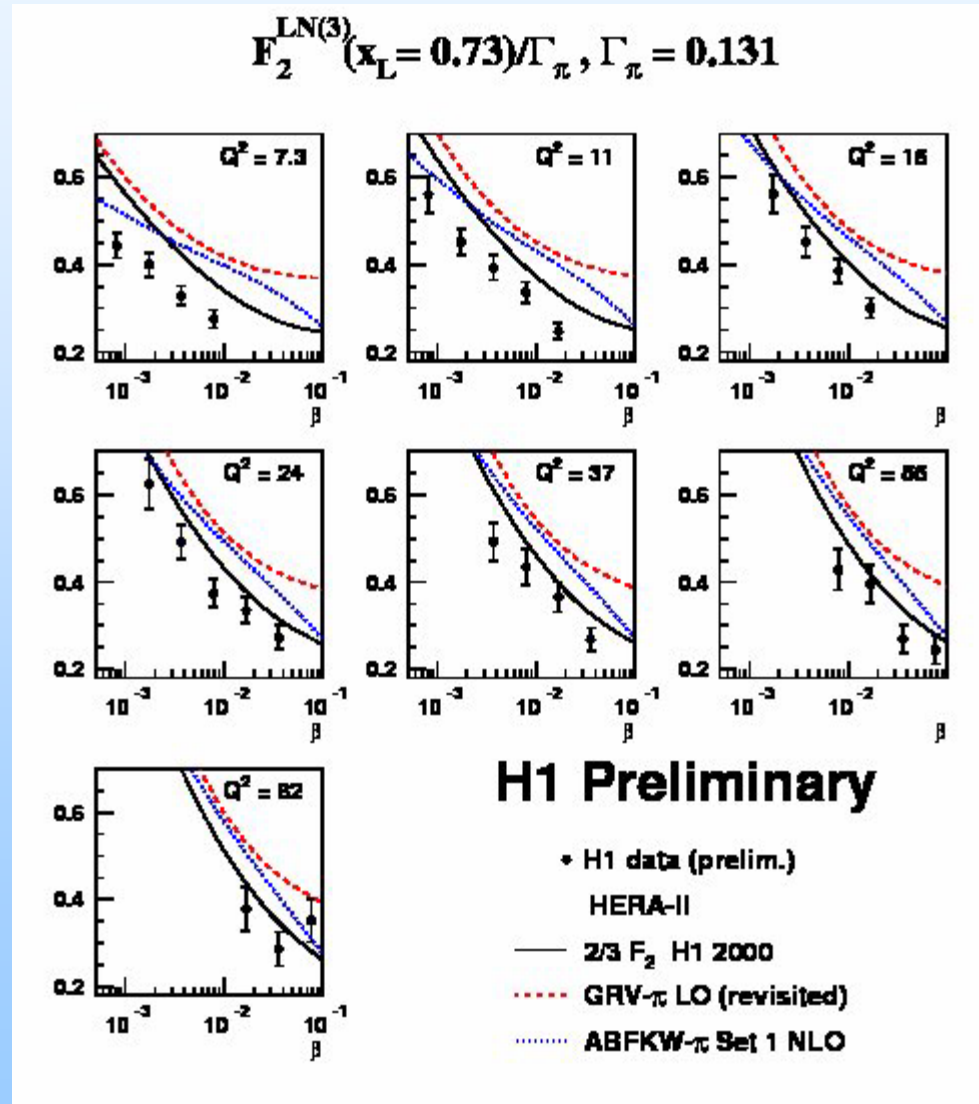
$$F_2^{LN(3)}(\beta, Q^2, x_L) = f_{\pi/p}(x_L) \cdot F_2^\pi(\beta, Q^2) \quad \beta = x/(1-x_L)$$



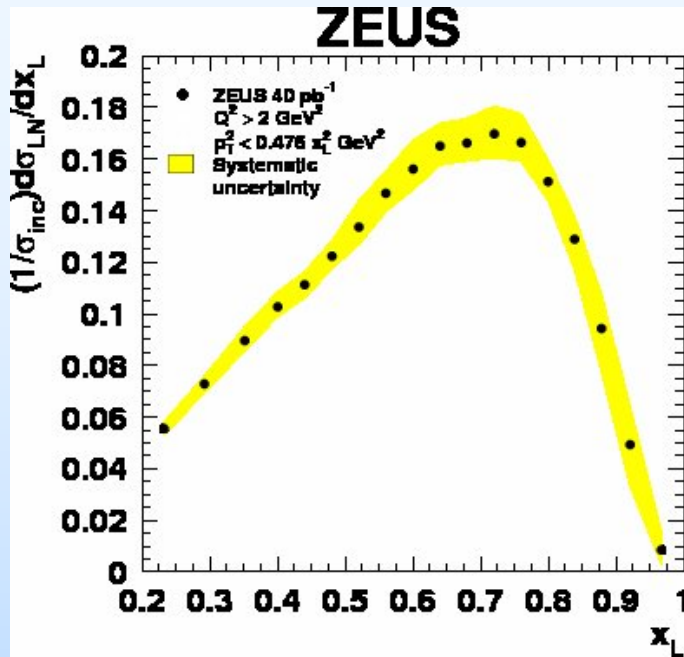
π exchange

— additive quark model

$$F_2^\pi = 2/3 \cdot F_2^p$$

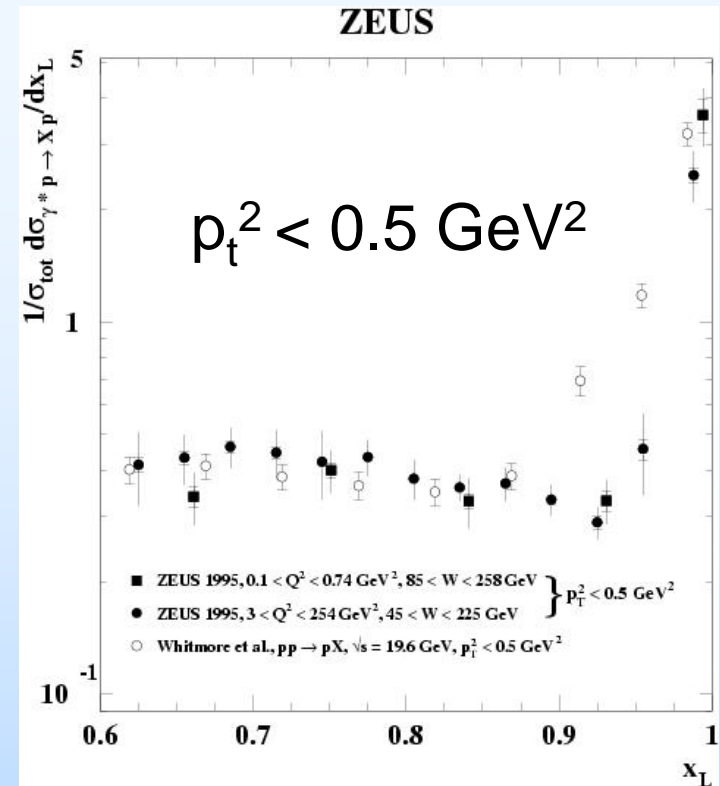


$$p_t^2 < 0.48 \cdot x_L^2$$



$$e p \rightarrow e n X$$

- $x_L \rightarrow 1$ n yield $\rightarrow 0$
- similarity to $p p \rightarrow n X$



$$e p \rightarrow e p X$$

- $x_L \rightarrow 1$ diffractive peak
- similarity to $p p \rightarrow p X$
- flat for $x_L < 0.95$