

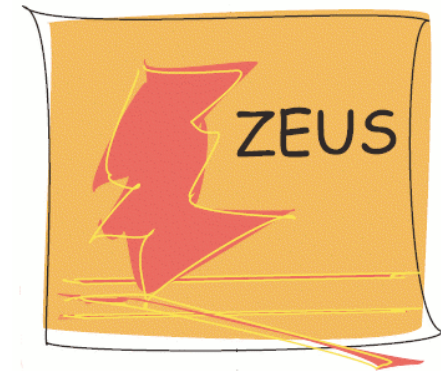
# Diffraction at HERA



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*31st July 2008*

ICHEP 08 Philadelphia USA



# Overview

- Inclusive Diffraction at HERA
  - Factorisation in Diffraction
  - Experimental techniques
  - Data to data comparisons
- QCD Fits and Diffractive PDFs
- Diffractive dijets in DIS
  - Diffractive PDFs from a Combined Fit
- Diffractive dijets in Photoproduction
  - Suppression or not of the cross section
- Summary

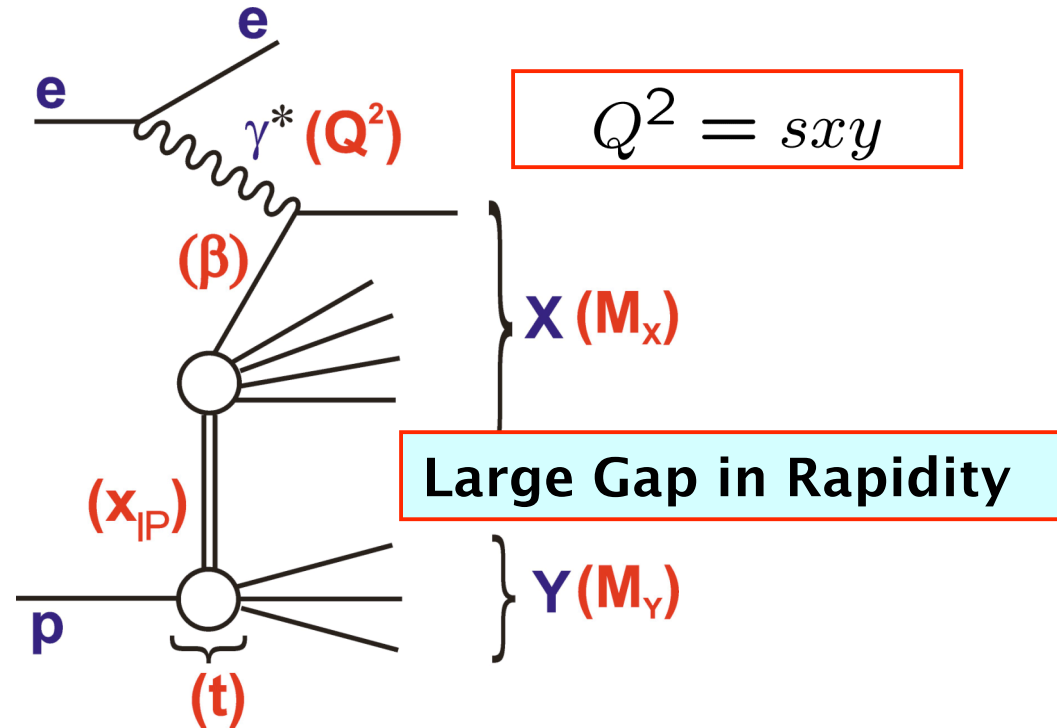
# Diffractive DIS Kinematics and Observables

$$x = x_{\mathbb{P}} \beta$$

$$\beta = \frac{Q^2}{Q^2 + M_X^2}$$

$$x_{\mathbb{P}} = \frac{Q^2 + M_X^2}{Q^2 + W^2}$$

$$Y_+ = 1 + (1 - y)^2$$

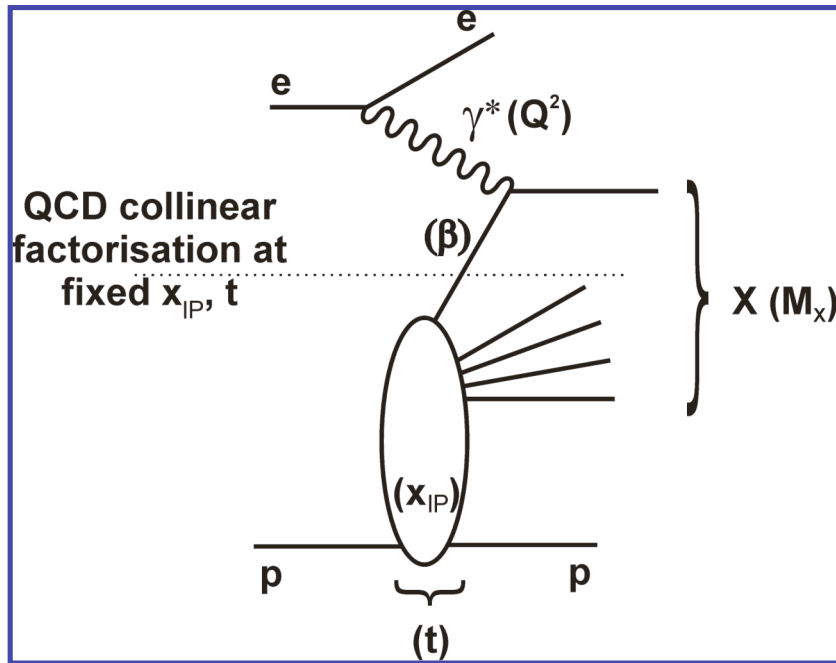


Cross section: 
$$\frac{d^4\sigma^{ep \rightarrow eXp}}{dx dQ^2 dx_{\mathbb{P}} dt} = \frac{4\pi\alpha^2}{xQ^4} Y_+ \sigma_r^{D(4)}(x, Q^2, x_{\mathbb{P}}, t)$$

$$\sigma_r^{D(4)} = F_2^{D(4)} - \frac{y^2}{Y_+} F_L^{D(4)}$$

$$\sigma_r^{D(3)} = \int_{-1}^{t_{min}} \sigma_r^{D(4)} dt$$

# Factorisation in Diffractive DIS



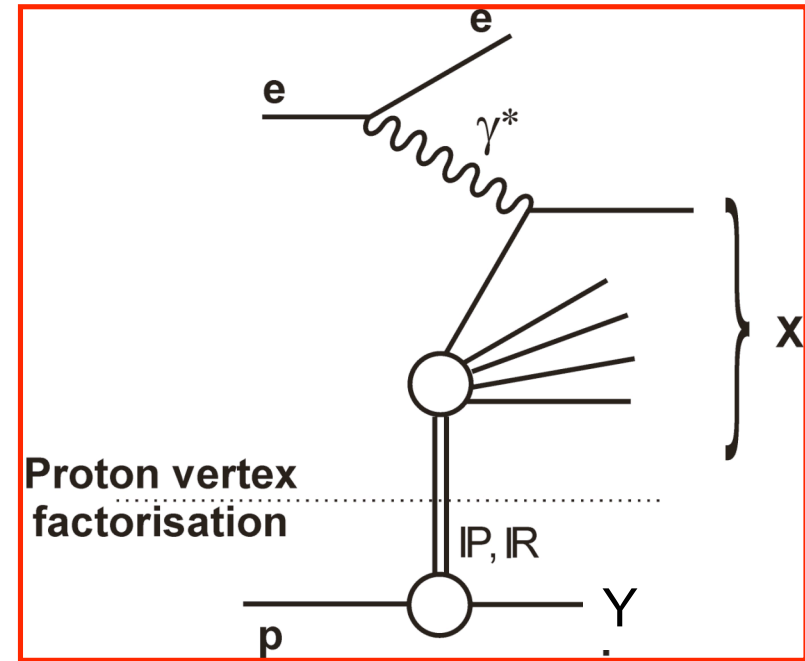
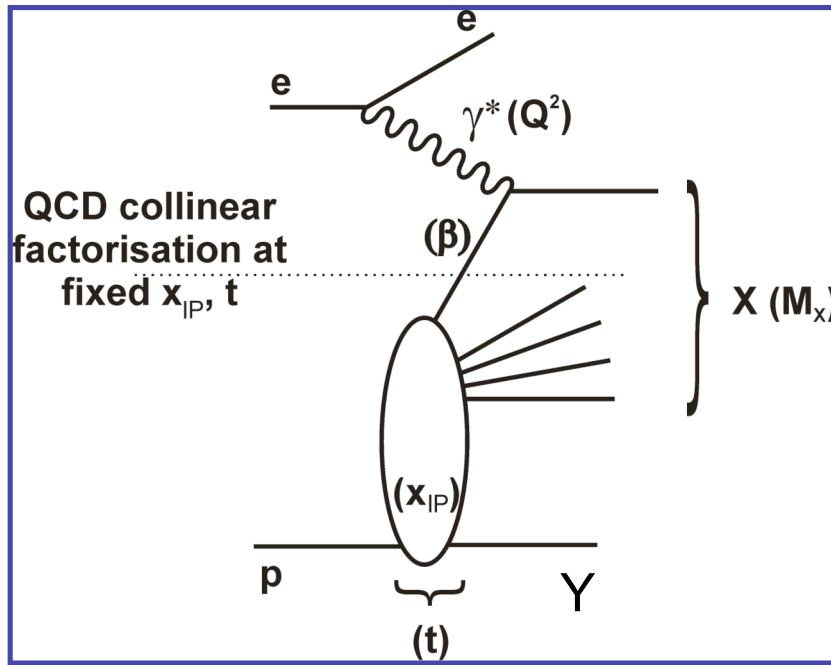
QCD hard scattering collinear factorisation (Collins) at fixed  $x_{IP}$  and  $t$

$$d\sigma_{partoni}(ep \rightarrow eXY) = f_i^D(x, Q^2, x_{IP}, t) \otimes d\sigma^{ei}(x, Q^2)$$

Applied after integration over measured  $M_Y$  and  $t$  ranges



# Factorisation in Diffractive DIS



QCD hard scattering collinear factorisation (Collins) at fixed  $x_{IP}$  and  $t$

$$d\sigma_{partoni}(ep \rightarrow eXY) = f_i^D(x, Q^2, x_{IP}, t) \otimes d\sigma^{ei}(x, Q^2)$$

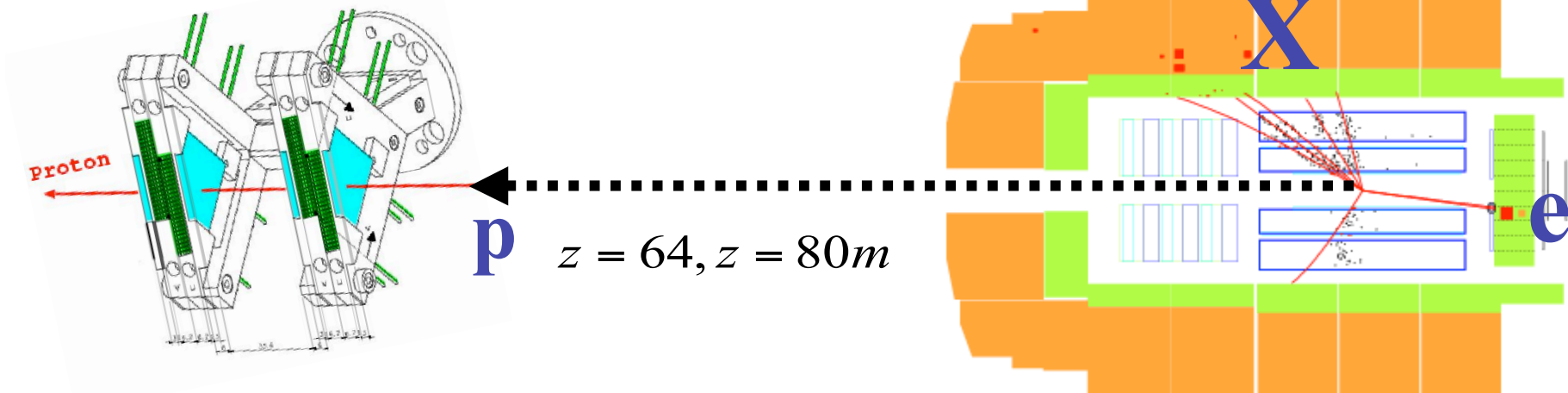
Applied after integration over measured  $M_Y$  and  $t$  ranges

'Proton vertex' factorisation of  $\beta$  and  $Q^2$  from  $x_{IP}$ ,  $t$ , and  $M_Y$  dependences

$$f_i^D(x, Q^2, x_{IP}, t) = f_{IP/p}(x_{IP}, t) \cdot f_i^{IP}(\beta = \frac{x}{x_{IP}}, Q^2)$$

# Experimentally selecting $ep \rightarrow eXp$

*I* Forward/Leading  
Proton Spectrometer



Measure Leading Proton (FPS/LPS)

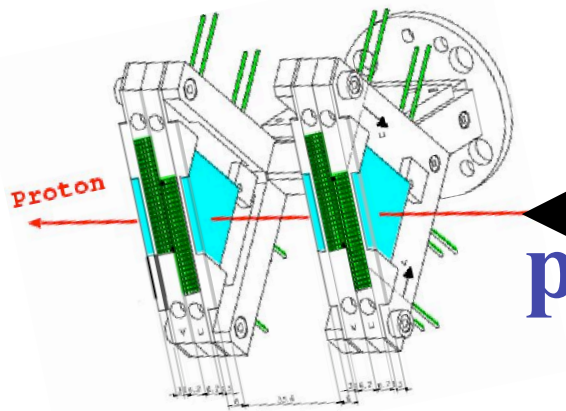
No proton dissociation

Measure the  $t$  dependence

Low detector acceptance

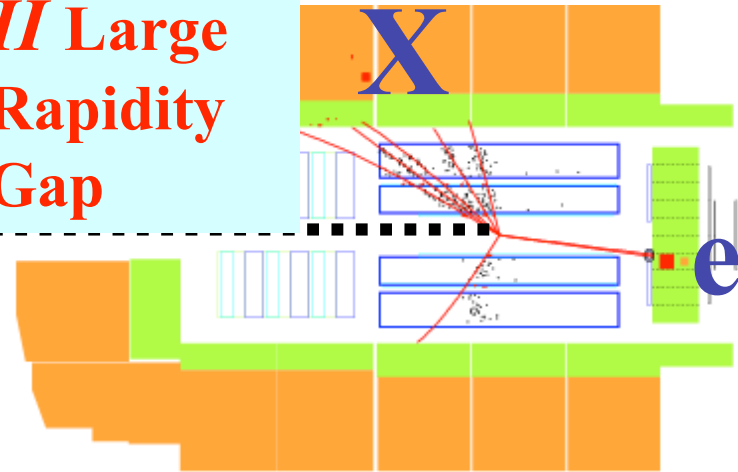
# Experimentally selecting $ep \rightarrow eXp$

## I Forward/Leading Proton Spectrometer



$$\mathbf{p}_z = 64, z = 80m$$

## II Large Rapidity Gap



Measure Leading Proton (FPS/LPS)

No proton dissociation

Measure the  $t$  dependence

Low detector acceptance

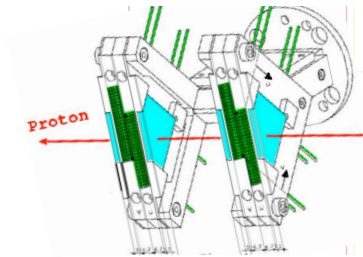
Require Large Rapidity Gap (LRG)  
spanning at least  $3.3 < \eta < \sim 7.5$

Kinematics measured from  $X$  system,  
integrate  $|t| < 1.0 \text{ GeV}^2$ ,  $M_Y < 1.6 \text{ GeV}$

High detector acceptance  $\rightarrow$  precision

# Experimentally selecting $ep \rightarrow eXp$

## I LPS/FPS



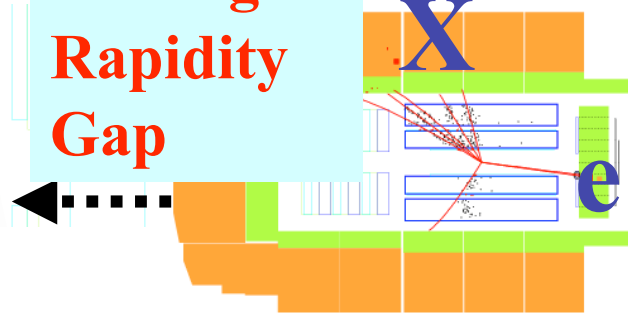
Measure Leading Proton (FPS/LPS)

No proton dissociation

Measure the  $t$  dependence

Low detector acceptance

## II Large Rapidity Gap



Require Large Rapidity Gap (LRG) spanning at least  $3.3 < \eta < \sim 7.5$

Kinematics measured from  $X$  system, integrate  $|t| < 1.0 \text{ GeV}^2$ ,  $M_Y < 1.6 \text{ GeV}$

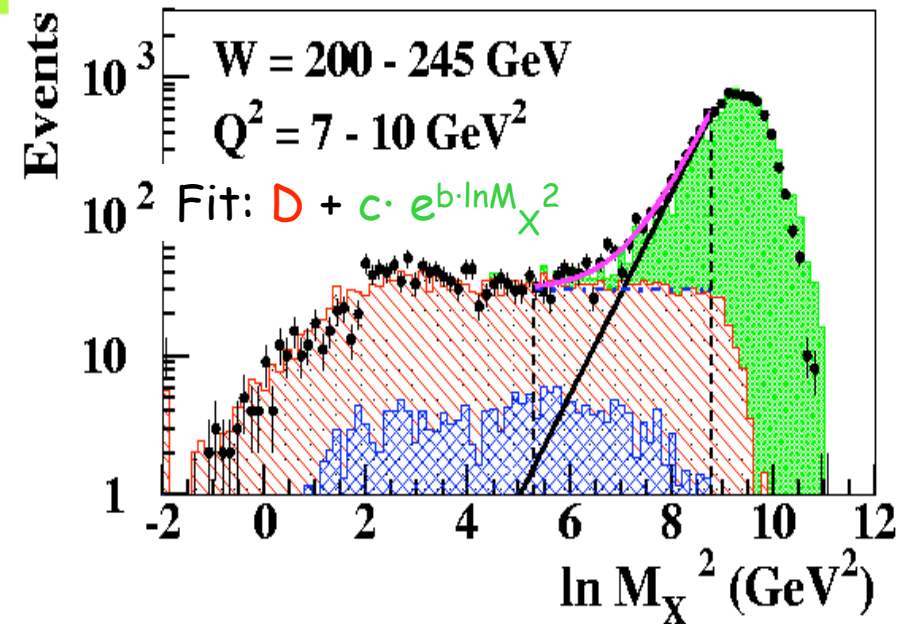
High detector acceptance  $\rightarrow$  precision

## III $M_X$ method

— Slope(nondiff) ··· Const(diff) — Fit(diff+nondiff)

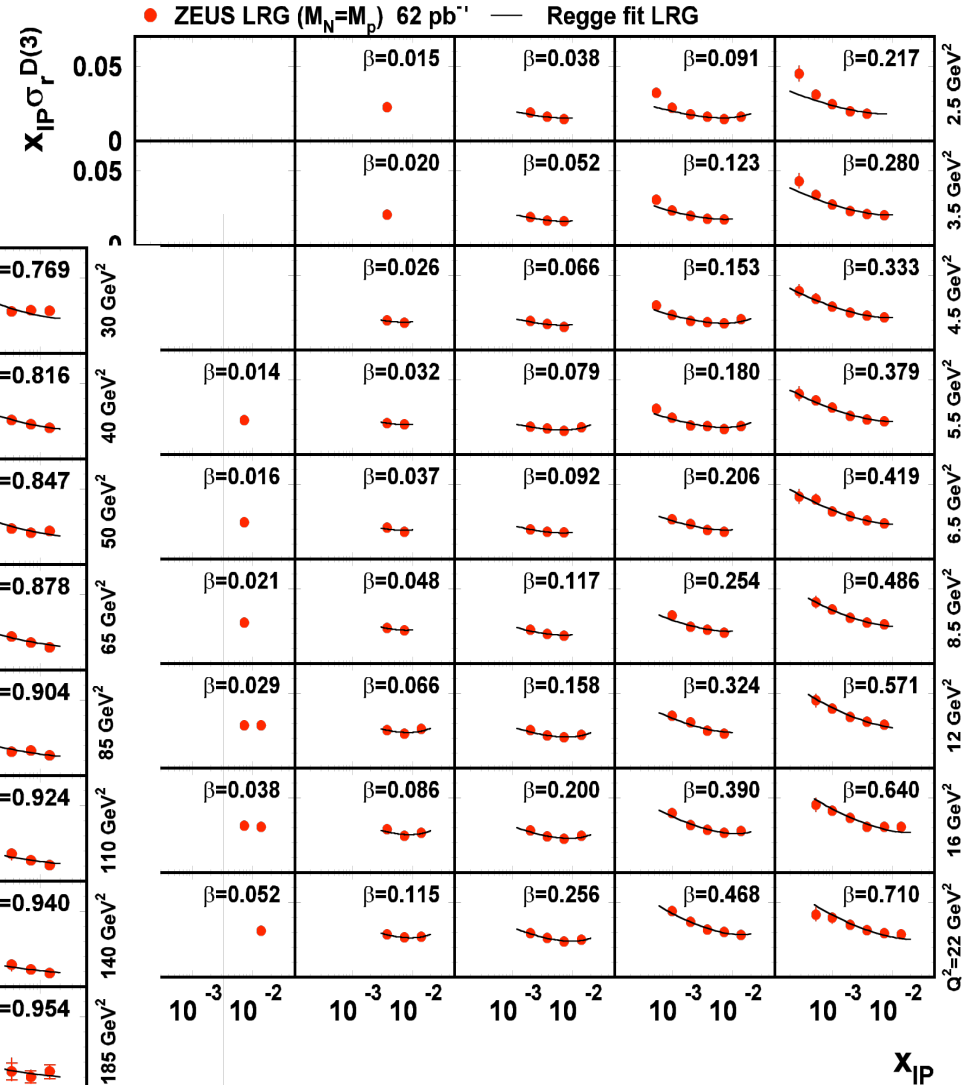
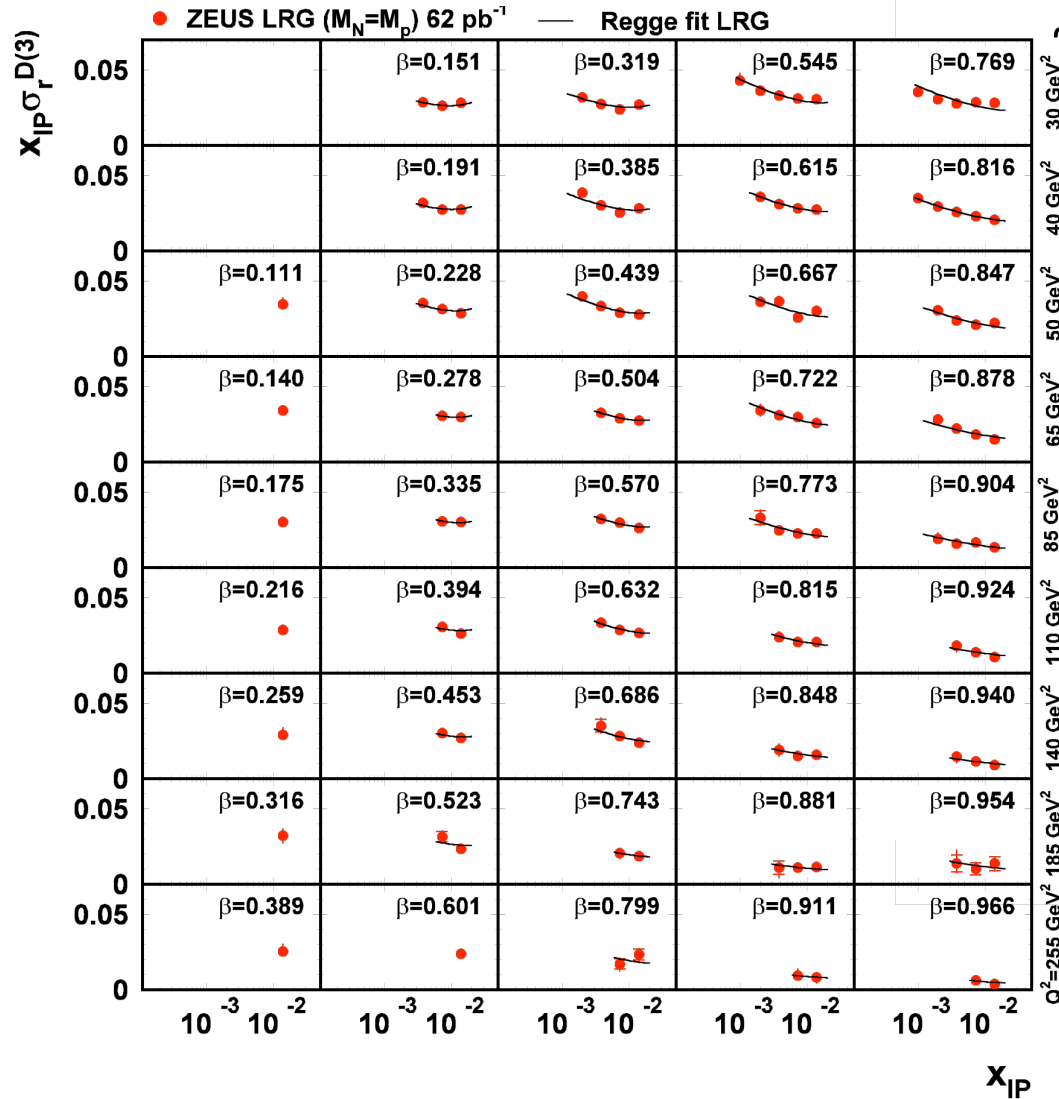
• D-PYT-Sang( $E_{\text{FPC}} > 1 \text{ GeV}$ )

■ DJG ▨ SR+Rhop ▩ Sang( $M_N < 2.3 \text{ GeV}$ )



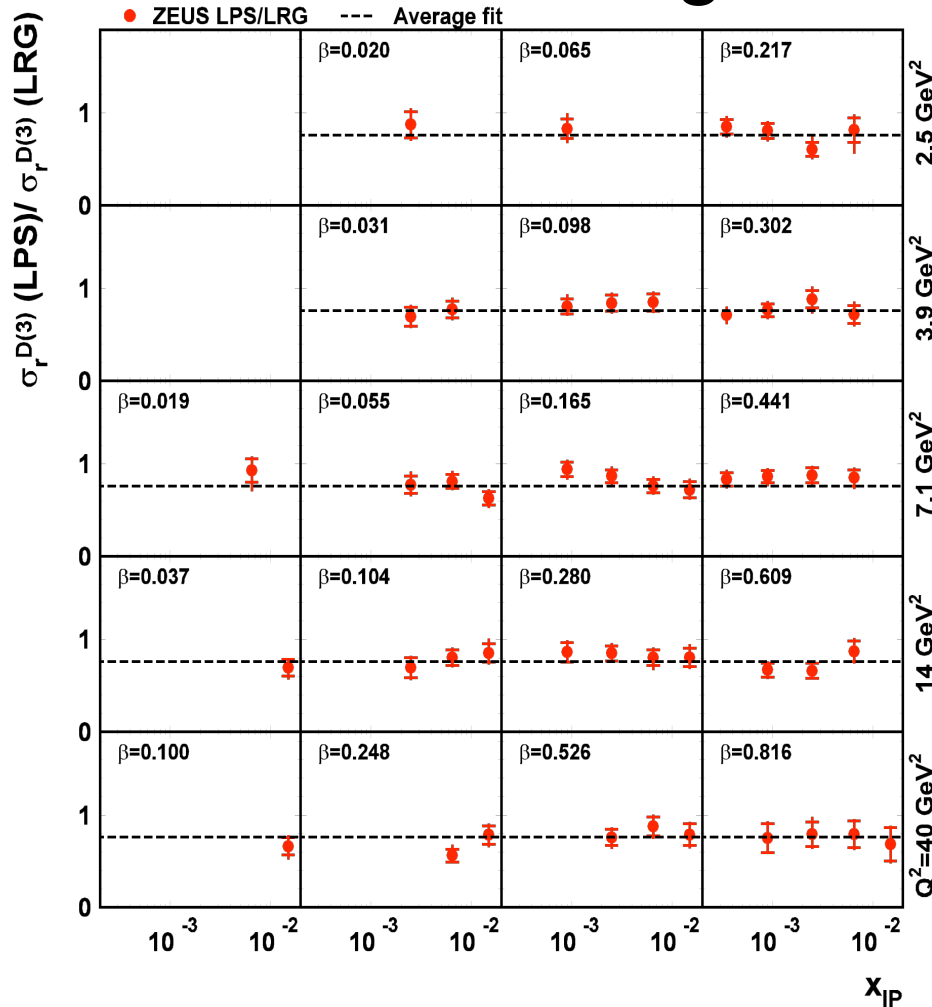
**How do the three experimental techniques compare?**

# New Large Rapidity Gap Data from ZEUS



The new LRG data from ZEUS is very precise, covering a large phase space

# Ratio of Leading Proton / Large Rapidity Gap



The LRG data contains a sizeable proton dissociation background (estimated to be 24% at ZEUS, cf. 23% at H1)

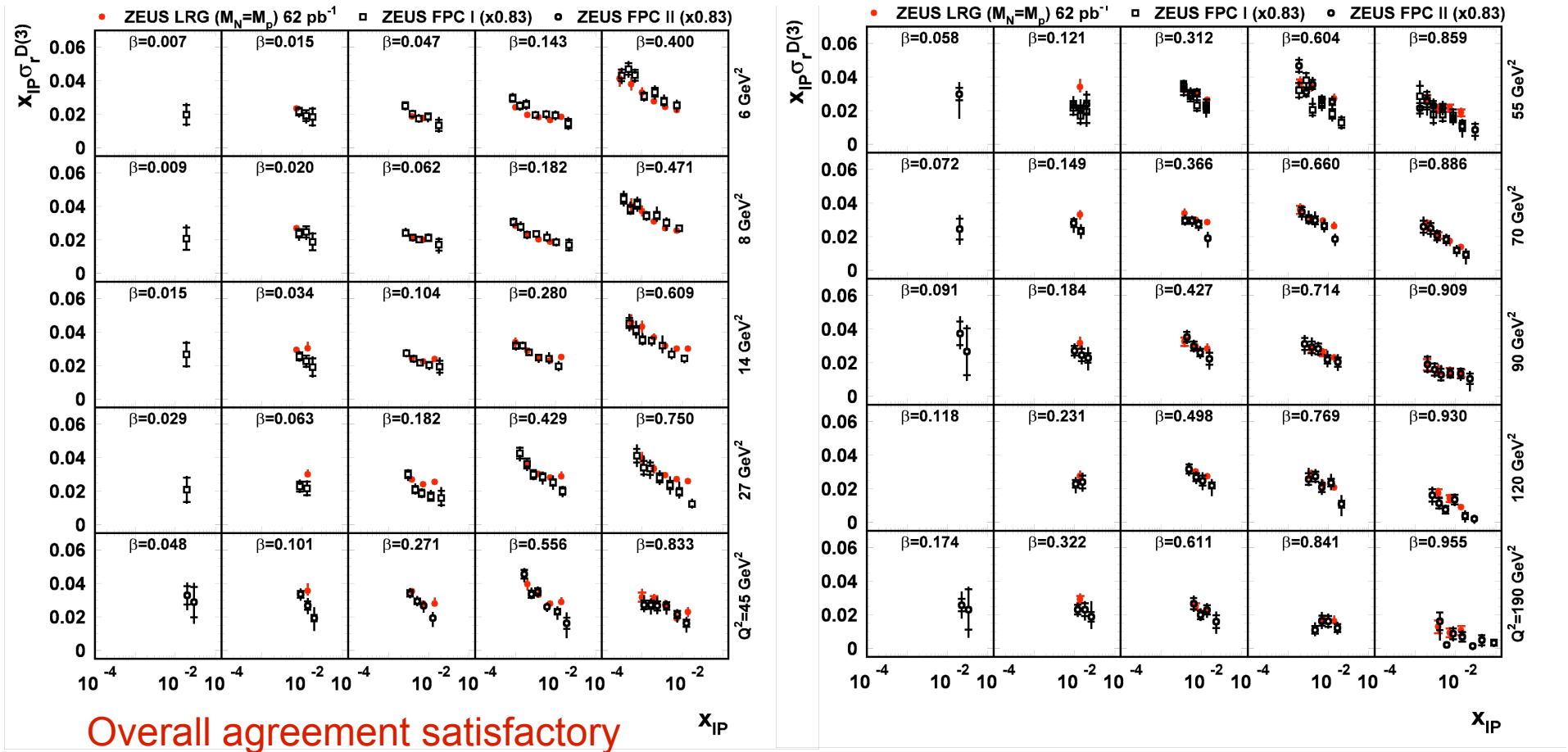
**The ratio of LPS/LRG cross sections is independent of  $Q^2$ ,  $x_{IP}$ ,  $\beta$**

ZEUS LPS / ZEUS LRG =  $0.76 \pm 0.01(\text{stat}) + 0.03 - 0.02(\text{sys}) + 0.08 - 0.05(\text{norm})$

→ p-diss. background in LRG data:  $[24 \pm 1(\text{stat}) + 2 - 3(\text{sys}) + 5 - 8(\text{norm})]\%$

# ZEUS LRG vs ZEUS $M_x$

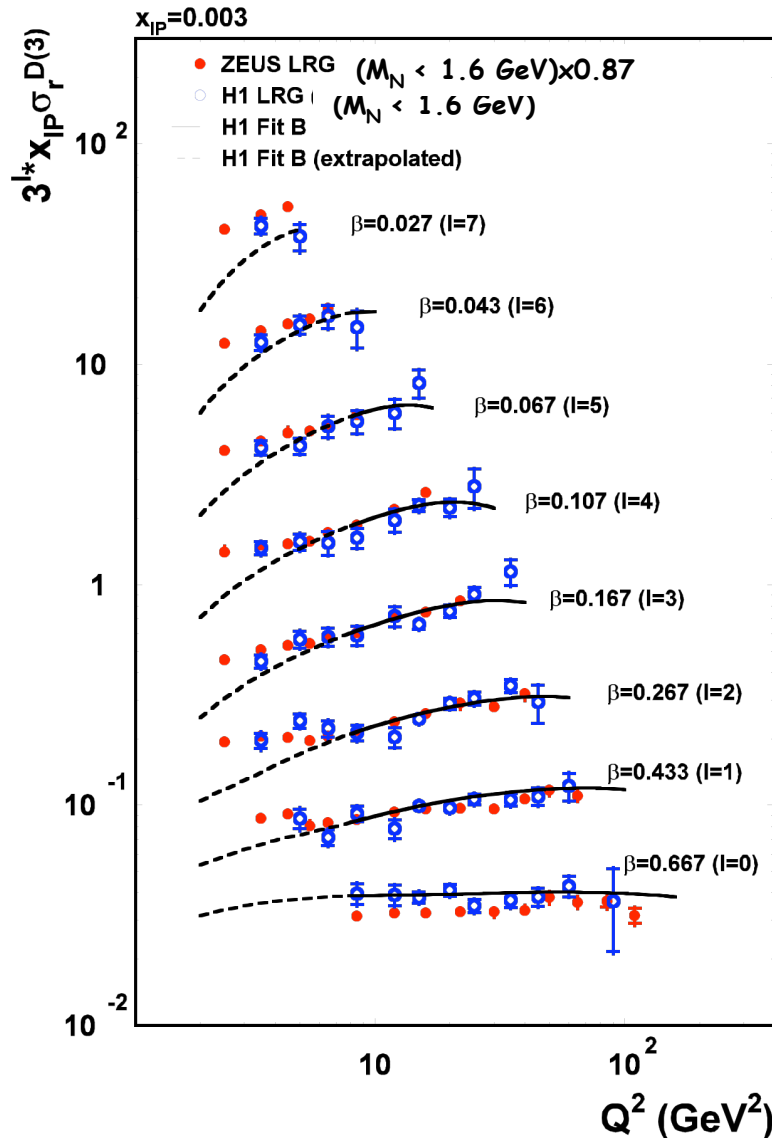
$M_x$  data ( $M_N < 2.3$  GeV) normalised to LRG ( $M_N = m_p$ ): factor  $0.83 \pm 0.04$   
(determined via a global fit) **estimates residual p-diss. background in  $M_x$  sample**



**Different  $x_{IP}$  dependence ascribed to IR suppressed in  $M_x$  data**

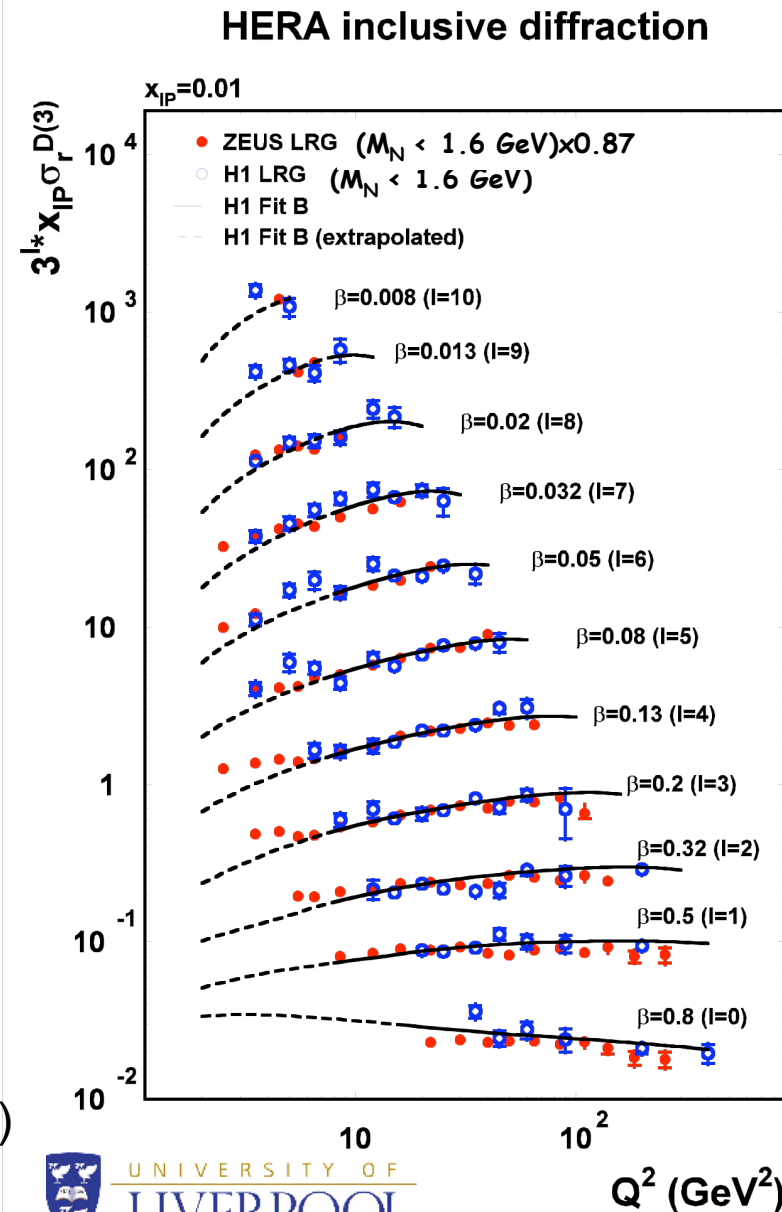


## HERA inclusive diffraction



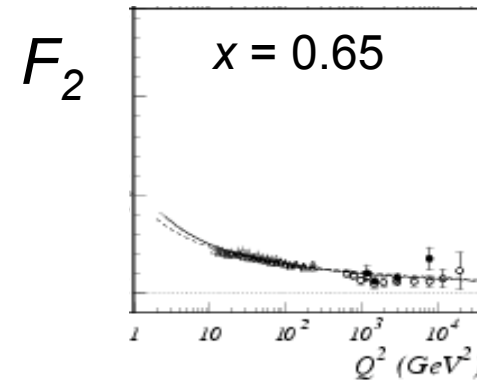
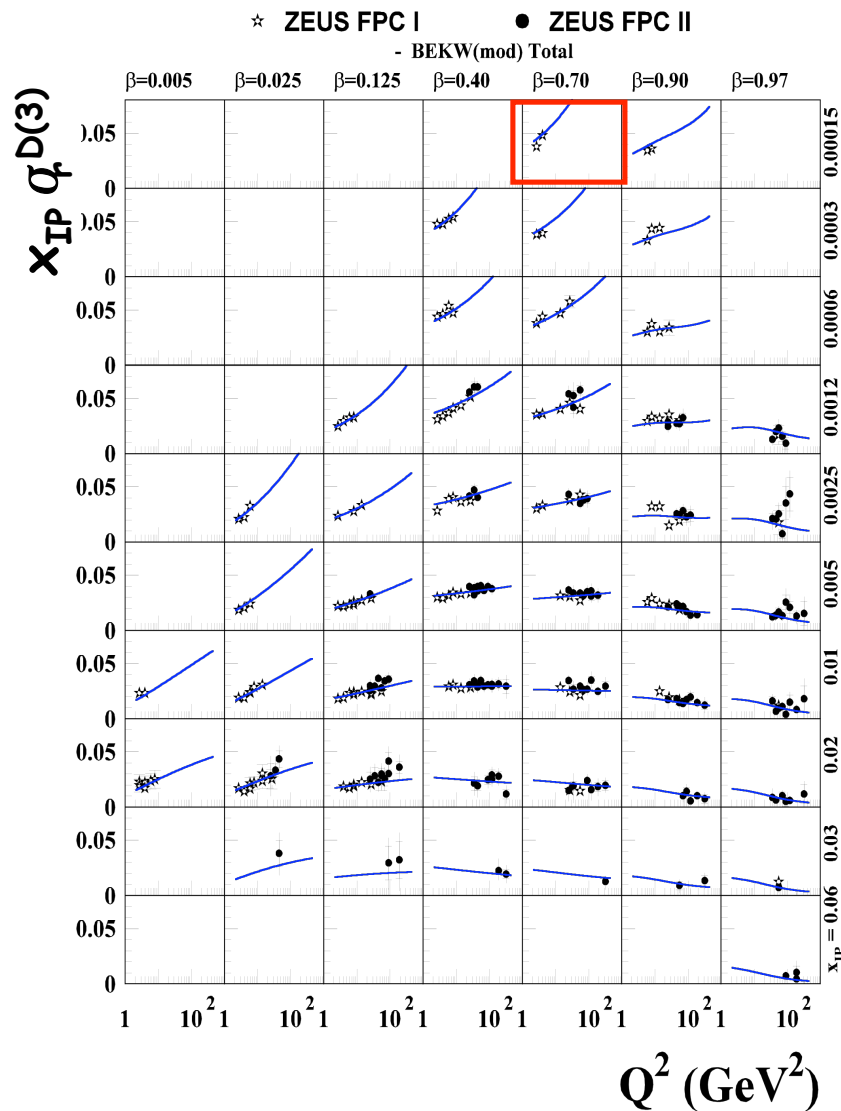
Good agreement between H1 and ZEUS  
 (ZEUS scaled by 0.87, covered by norm. unc.)

## HERA Large Rapidity Gap Data





# $Q^2$ dependence of $\sigma_r^{D(3)}$ ZEUS $M_x$ data

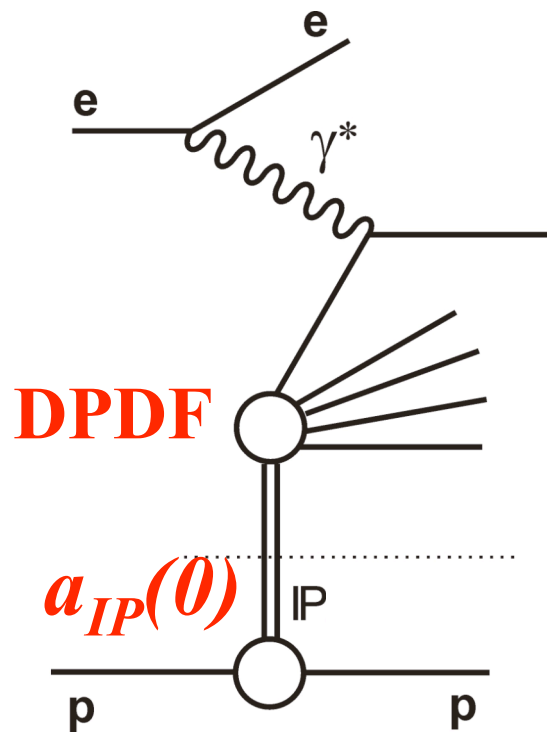


Large positive scaling violations up to high- $\beta$  values implies that the diffractive exchange is gluon-dominated

At fixed  $\beta$  the reduced cross section depends on  $x_{IP}$  - these data seem to contradict Regge factorisation

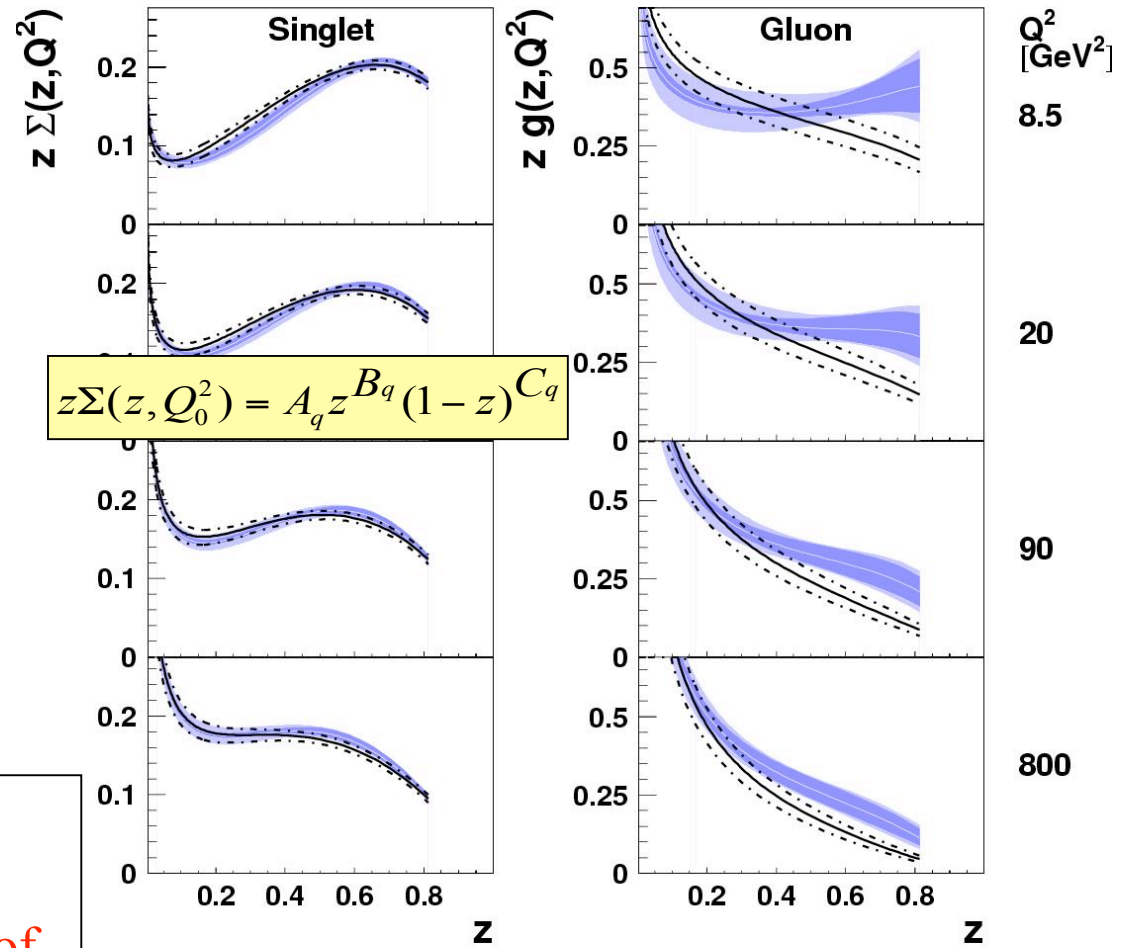
*Regge factorisation is only a useful approximation but fits made thus far are insensitive to this mild breaking*

# H1 2006 DPDF Fit



*IP* component:

- Fit  $\alpha_{IP}(0)$  ( $x_{IP}$  dependence).
- Simultaneously, fit 5 parameters of DPDFs ( $\beta$  and  $Q^2$  dependences) using NLO QCD.



$$z \Sigma(z, Q_0^2) = A_q z^{B_q} (1-z)^{C_q}$$

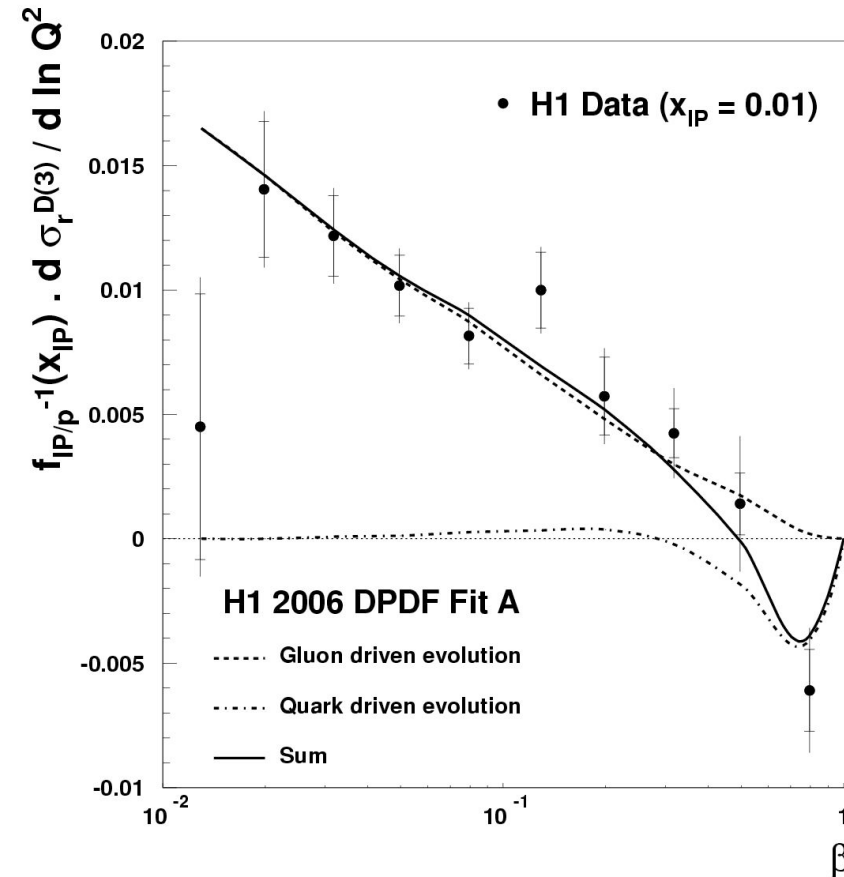
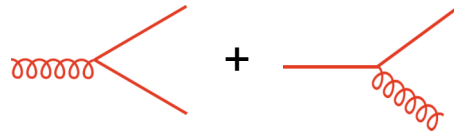
$$z_g(z, Q_0^2) = A_g (1-z)^{C_g}$$

$$z_g(z, Q_0^2) = A_g$$

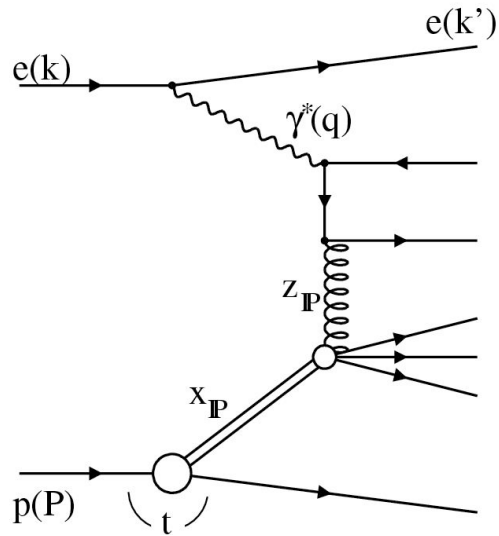
# A Closer Look at the High $z$ Region

We have only singlet quarks, so DGLAP evolution equation for  $F_2^D$  ....

$$\frac{dF_2^D}{d \ln Q^2} \sim \frac{\alpha_s}{2\pi} \left[ P_{qg} \otimes g + P_{qq} \otimes \Sigma \right]$$



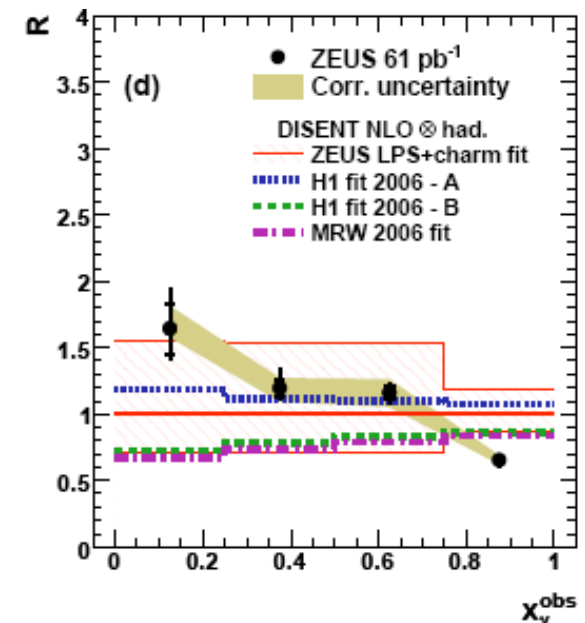
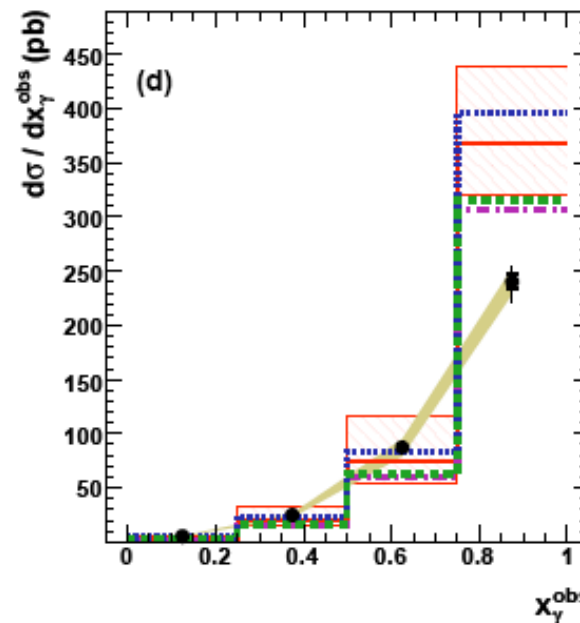
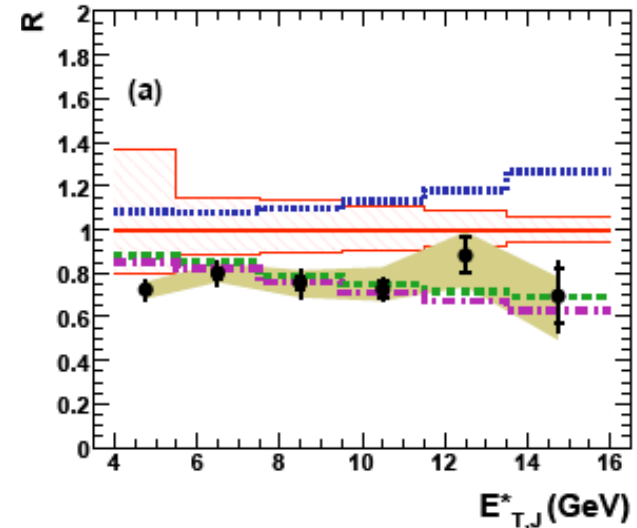
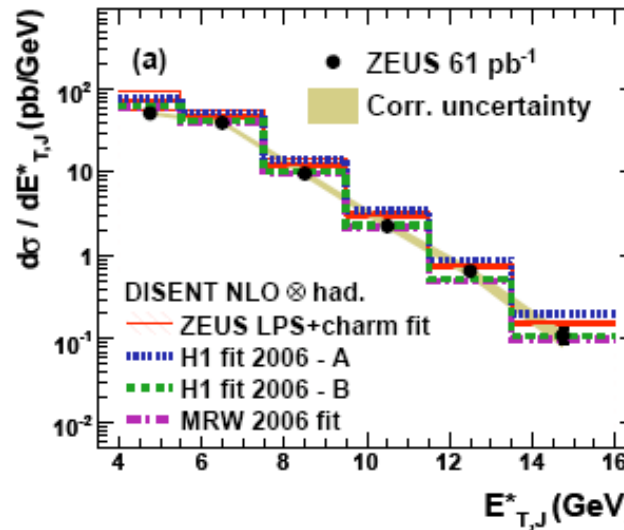
At high  $\beta$ , relative error on derivative grows,  $q \rightarrow qg$  contribution to evolution becomes important ... sensitivity to gluon is lost

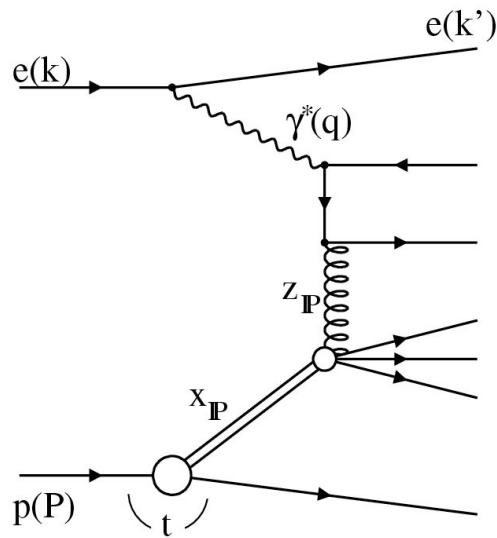


$$Z_{IP} = \frac{M_{12}^2 + Q^2}{M_X^2 + Q^2}$$

- Compare the ZEUS dijet data to H1 Fit A and Fit B
- Best agreement for H1 Fit B
- **Factorisation holds in DIS**

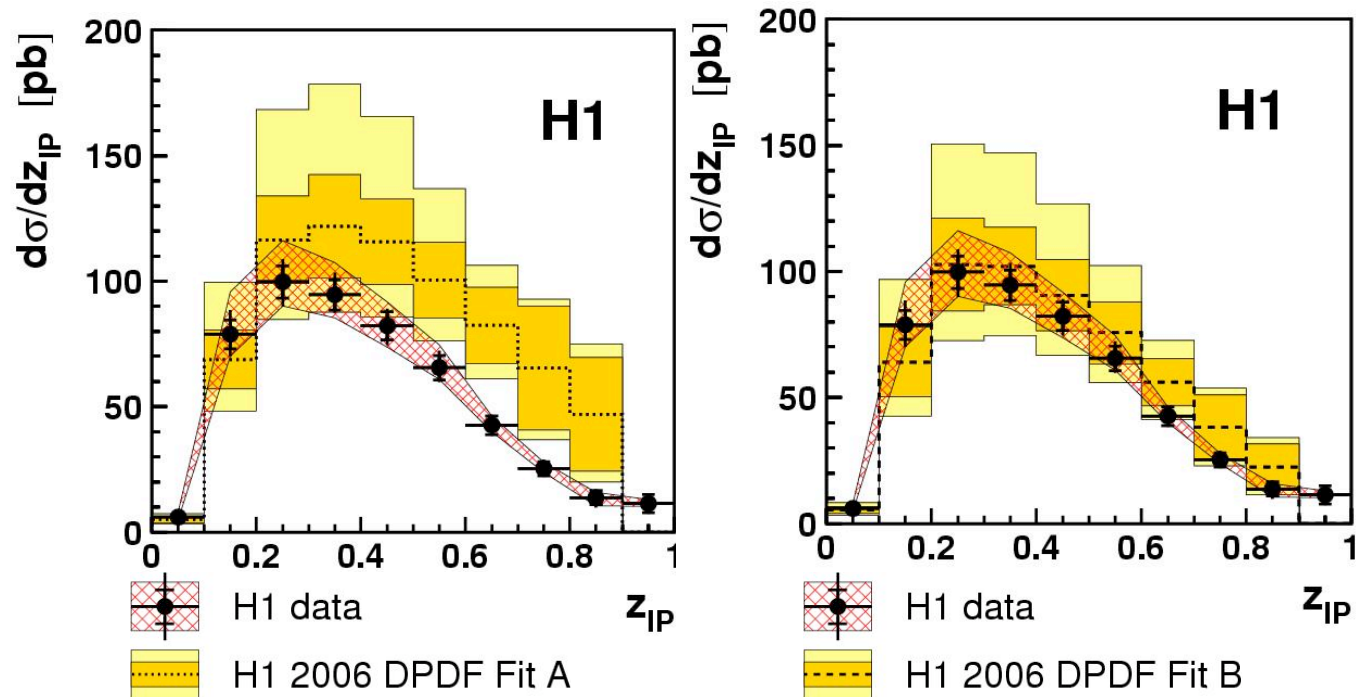
# Compare to diffractive dijets in DIS





$$z_{IP} = \frac{M_{12}^2 + Q^2}{M_X^2 + Q^2}$$

# Factorisation holds in DIS



At low  $z_{IP}$  ( $< 0.4$ ) Fit A and Fit B are similar

*The data are in good agreement with the predictions, consistent with factorisation*

At high  $z_{IP}$  the data clearly prefer Fit B

**Include the diffractive dijet in a combined fit with the inclusive H1 LRG data**

# Combined fit of dijet and inclusive data

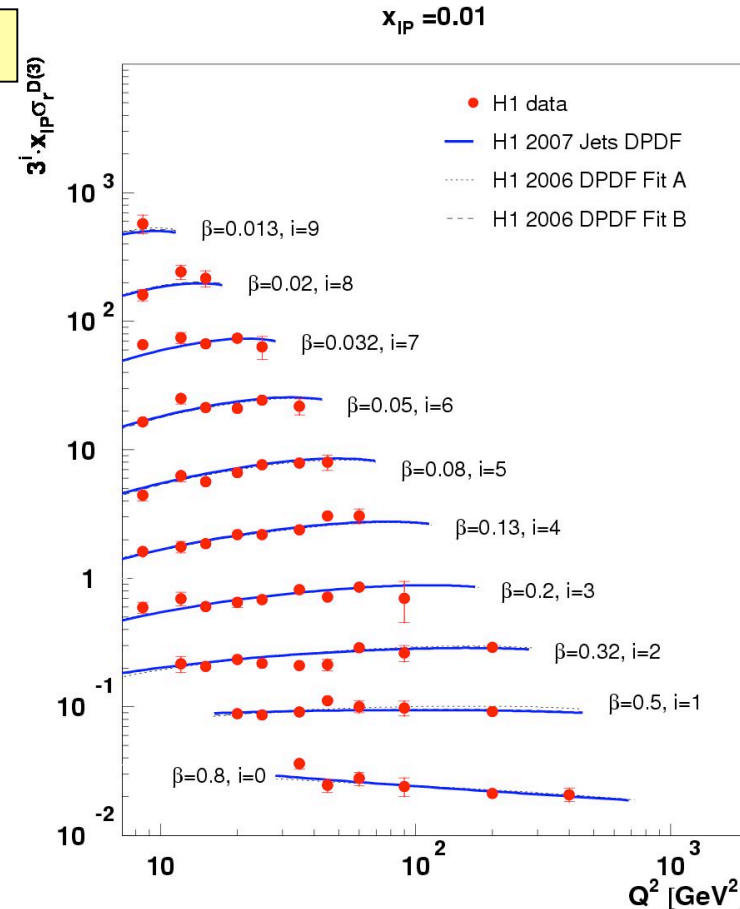
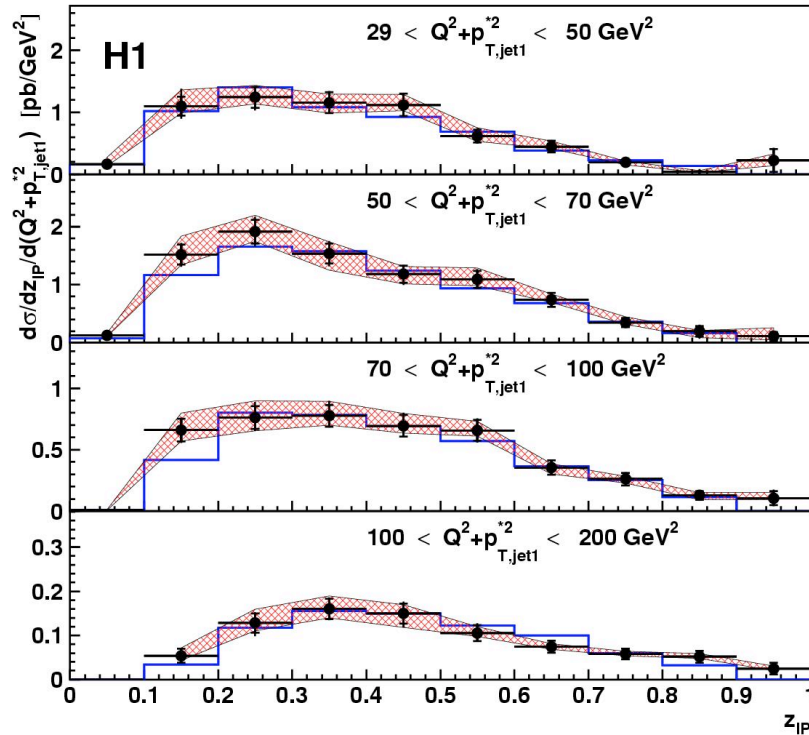


H1 data



H1 2007 Jets DPDF

$$z_g(z, Q_0^2) = A_g z^{B_g} (1-z)^{C_g}$$



- The diffractive dijet data can be used as an additional constraint in a NLO QCD fit procedure
- Details similar to the inclusive case but can now constrain 3 parameters for the gluon

**Very good simultaneous fit of both inclusive and dijet data achieved**



# Combined fit DPDFs from H1

- H1 2007 Jets DPDF
- exp. uncertainty
- exp. + theo. uncertainty
- ..... H1 2006 DPDF fit A
- H1 2006 DPDF fit B

Singlet:

$$z\Sigma(z, Q_0^2) = A_q z^{B_q} (1-z)^{C_q}$$

Gluon  
Fit A:

$$z_g(z, Q_0^2) = A_g (1-z)^{C_g}$$

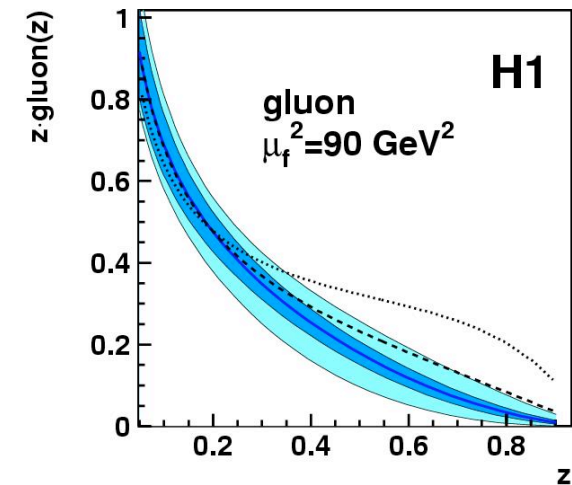
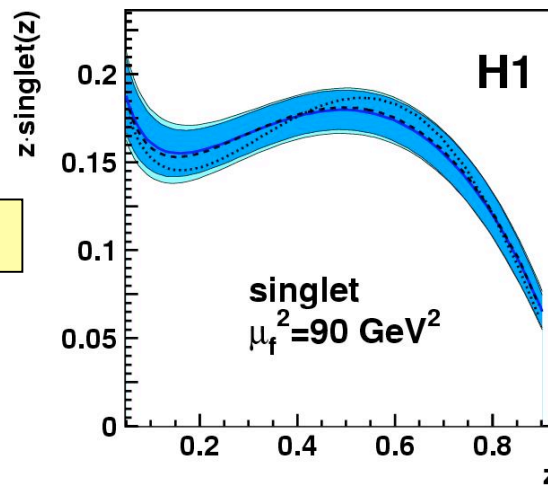
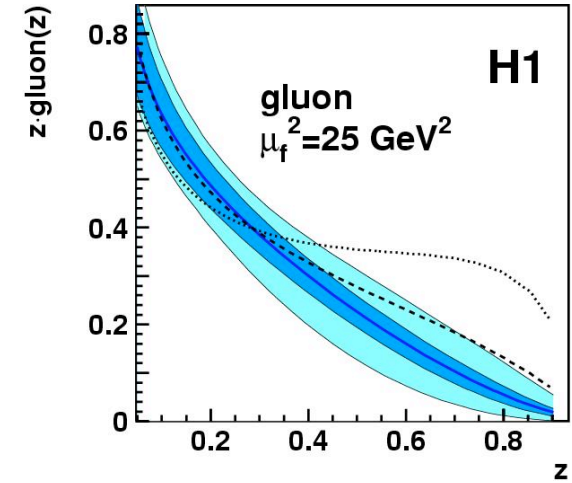
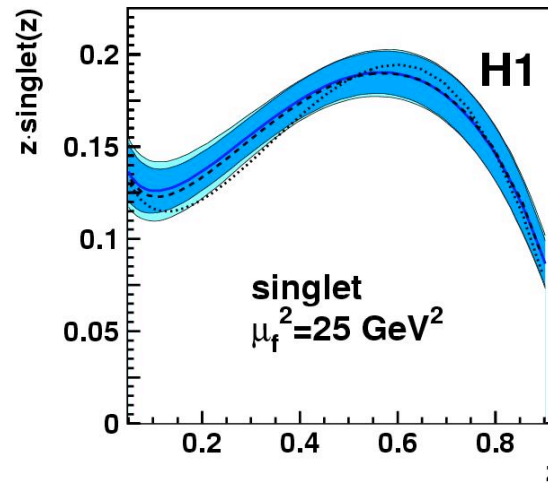
Fit B:

$$z_g(z, Q_0^2) = A_g$$

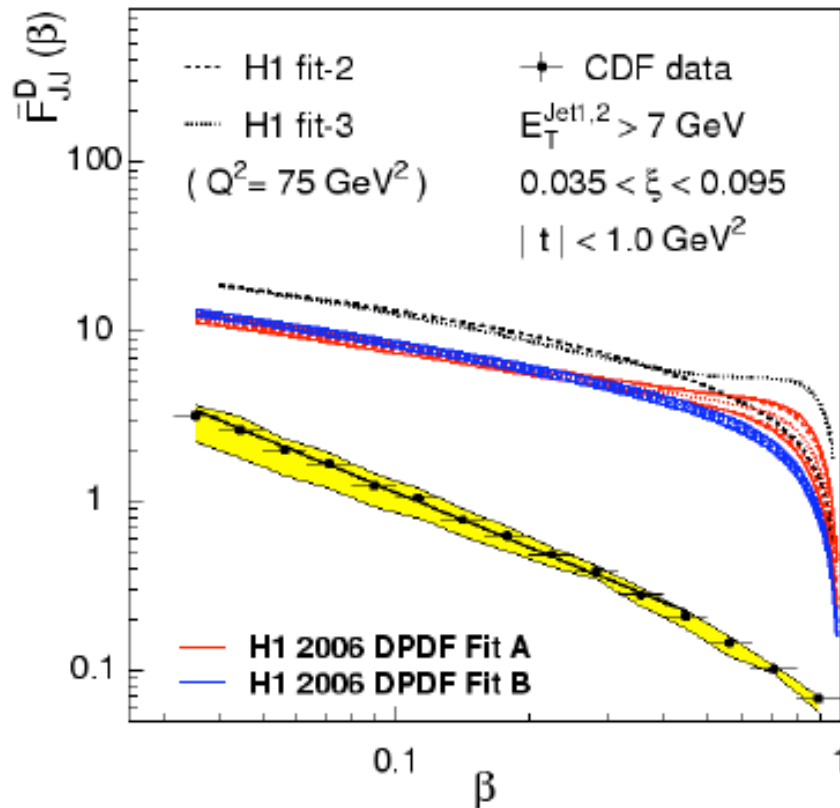
Jets:

$$z_g(z, Q_0^2) = A_g z^{B_g} (1-z)^{C_g}$$

*The singlet and gluon are constrained with similar precision across the whole kinematic range*

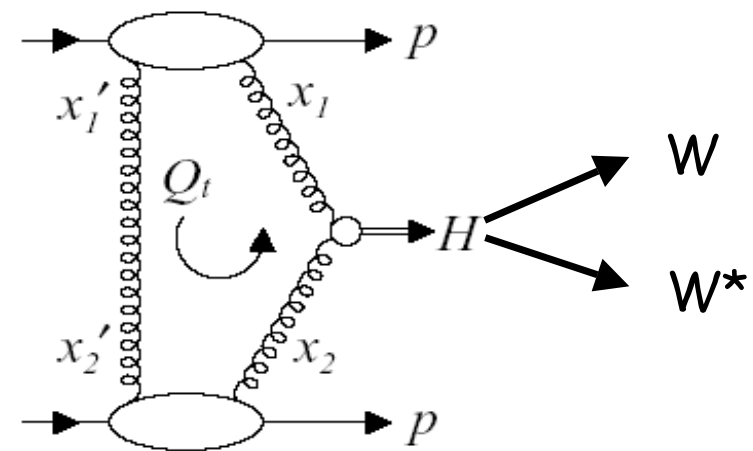


# Exporting DPDFs to Hadron-Hadron machines



When trying to use DPDFs extracted at HERA to predict diffractive dijets at CDF...

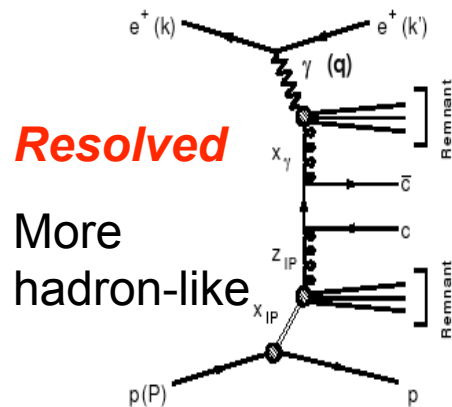
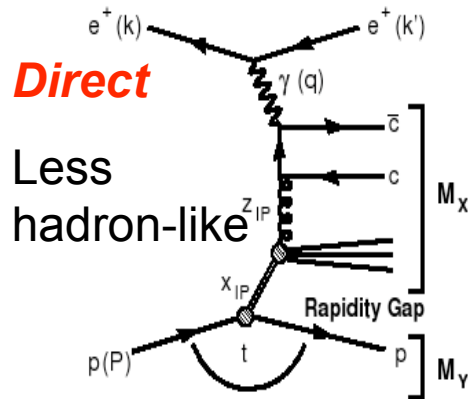
*... it simply doesn't work!*



That's a big problem when trying to make predictions for Diffractive Higgs production at the LHC



# Factorisation tests at HERA

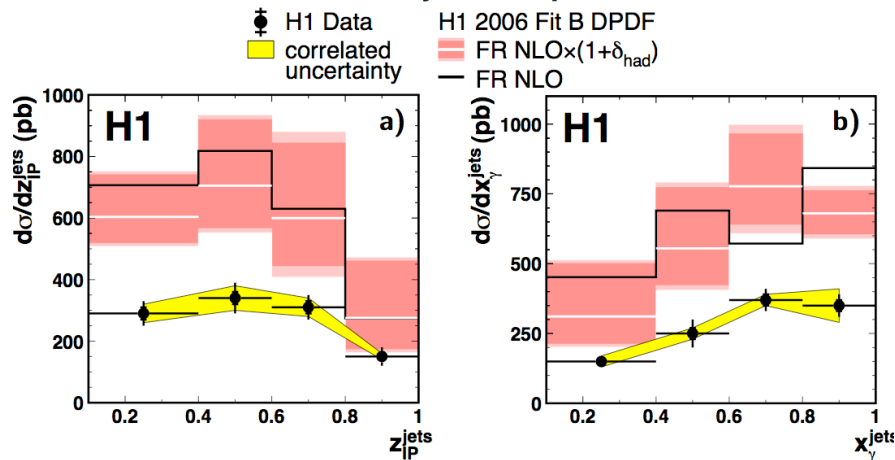


Use photoproduction at HERA as a hadron-hadron collider

How hadron-like the photon is depends on the  $x_\gamma$  variable

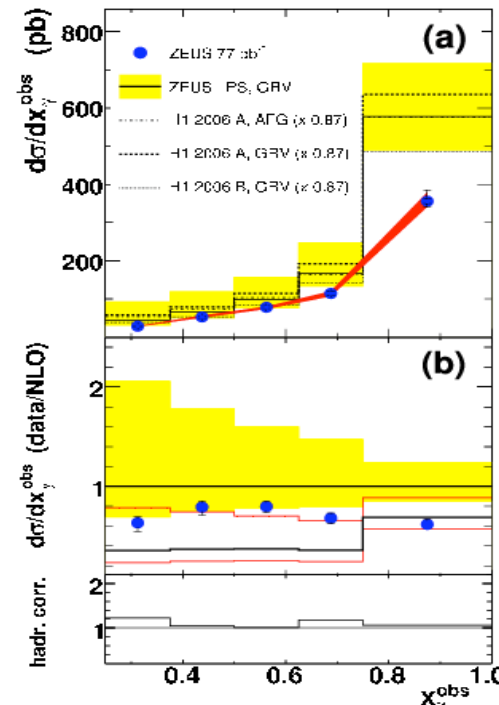
Expect *Resolved* (low  $x_\gamma$ ) to be more suppressed than *Direct* (high  $x_\gamma$ )

H1 Diffractive Dijet Photoproduction

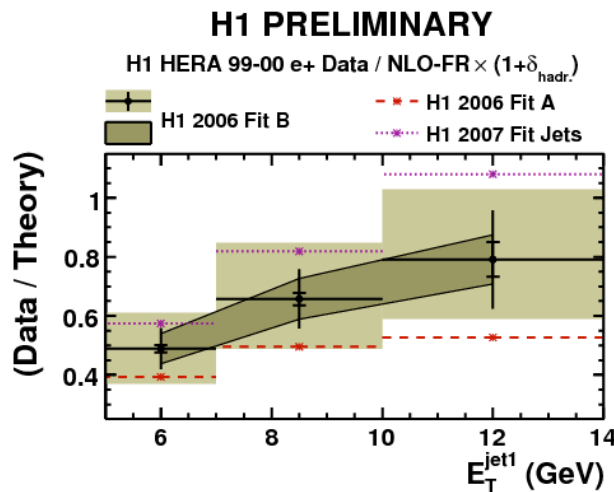
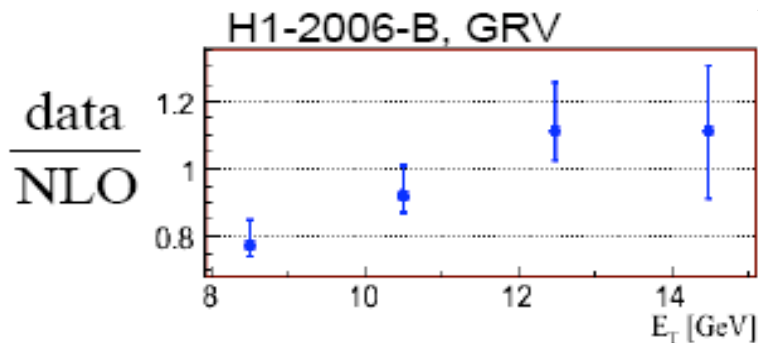


H1 saw that the cross section was suppressed in photoproduction, but independent of  $x_\gamma$

ZEUS



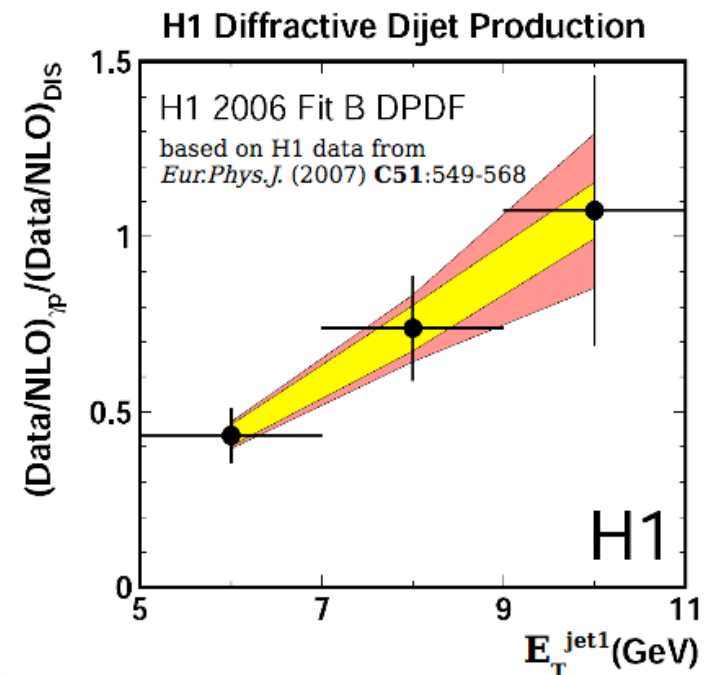
ZEUS saw consistency with no suppression but did confirm the absence of dependence of  $x_\gamma$

$E_T$  dependence of the suppression

The H1 and ZEUS dijets in photoproduction analyses have different analysis cuts on jet  $E_T$  with ZEUS being at higher  $E_T$  than H1

Looking at the Data/Theory ratio as a function of jet  $E_T$  suggests that there is an  *$E_T$  dependence of the suppression*

The suggestion of an  $E_T$  dependence is even stronger when looking at the double ratio of Data/Theory  $\gamma p$  / DIS where some systematic uncertainties cancel

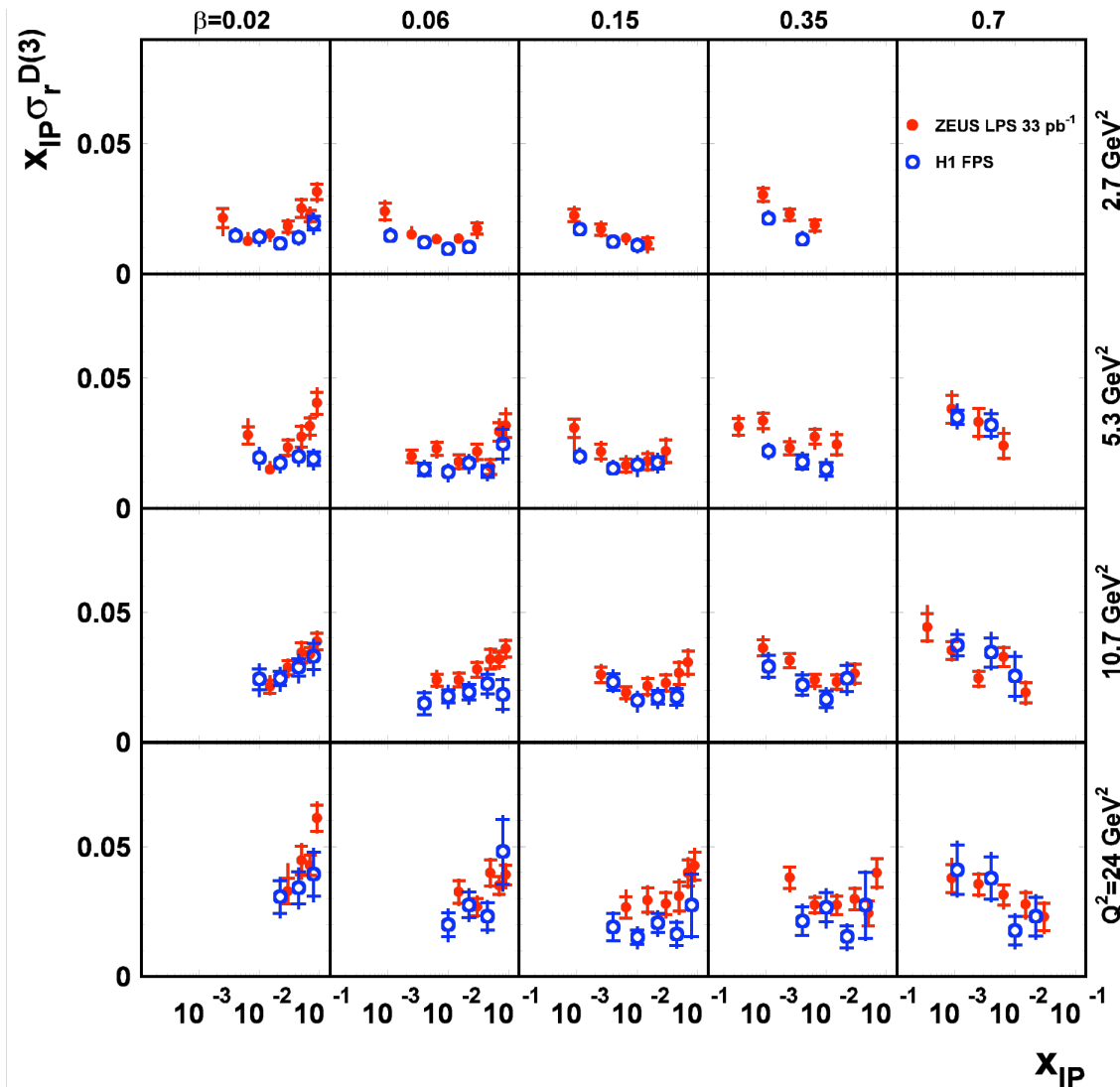


# Summary

- A wealth of data from both the H1 and ZEUS collaborations using the Leading Proton, LRG and  $M_X$  methods
  - the experimental techniques are largely consistent
- Proton vertex factorisation is a good enough approximation of the data to allow extraction of DPDFs from NLO QCD fits to  $\beta$ ,  $Q^2$  dependences of inclusive data
- Diffractive dijet data in DIS agree well with predictions of fits to inclusive data
- Diffractive dijet data in photoproduction show evidence of a suppression wrt predictions that is consistent with
  - No  $x_\gamma$  dependence
  - An  $E_T$  dependence
- Combined fit to inclusive and dijet data constrains both the quark and gluon PDFs to similar good precision2

BACK-UP SLIDES FOLLOW

# LPS vs FPS



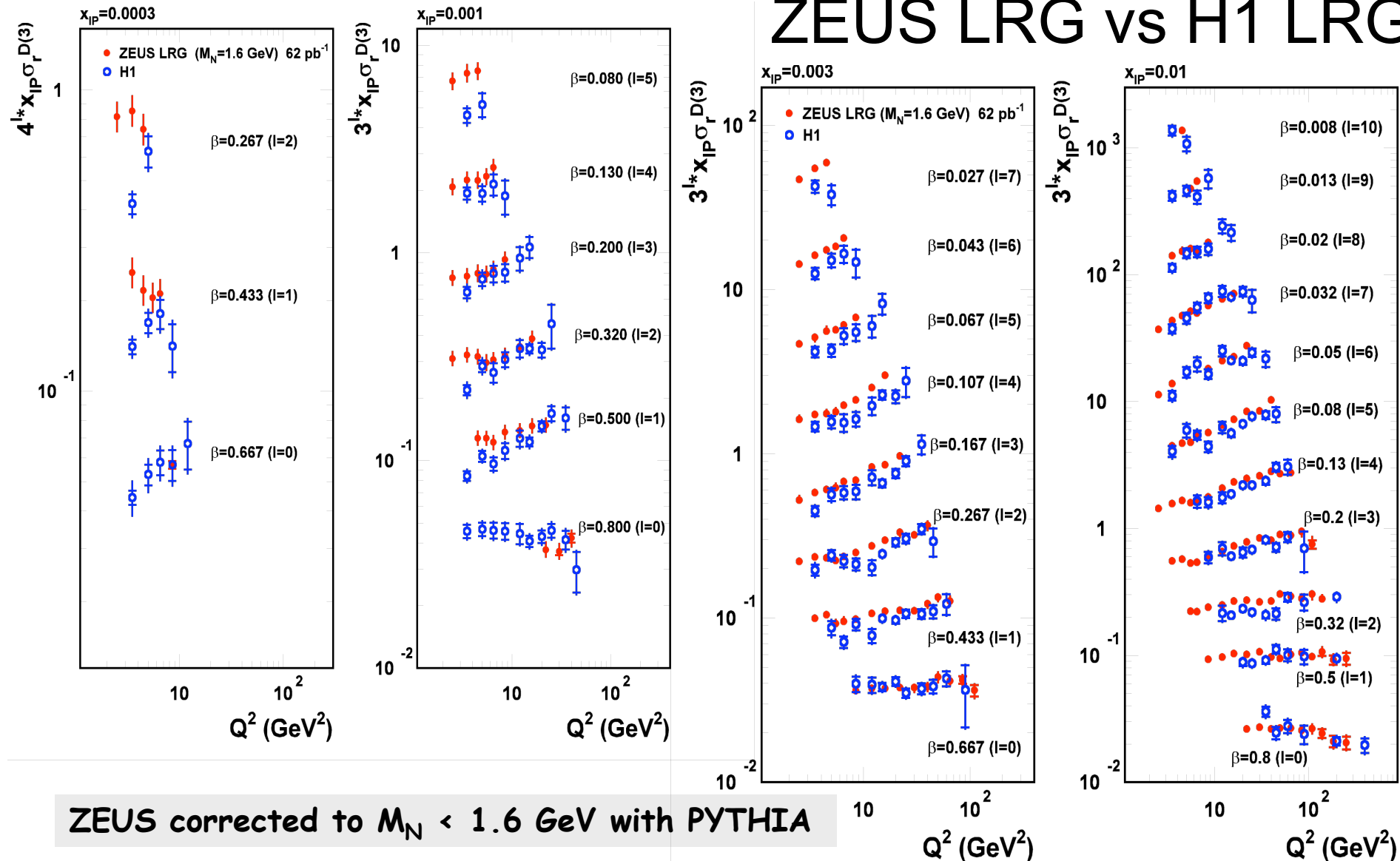
The cleanest possible comparison in principle...

...but large normalisation uncertainties:

(LPS: +11-7%, FPS: +/-10%)

**ZEUS and H1 proton-tagged data agree within normalisation uncertainties**

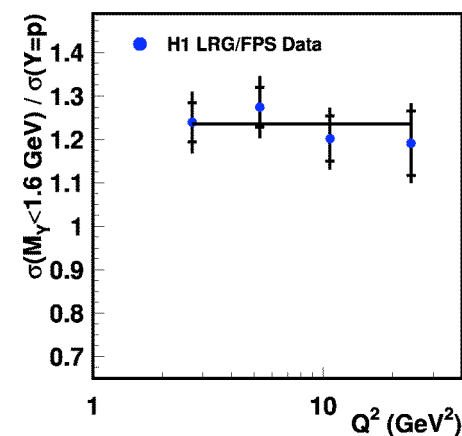
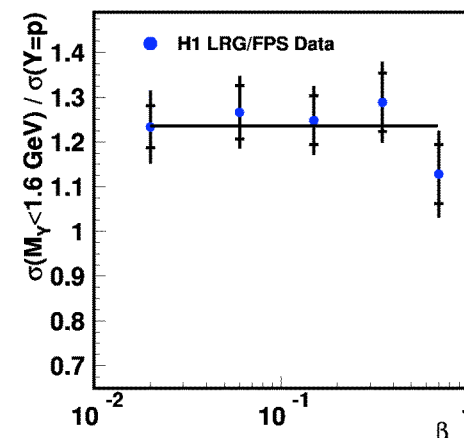
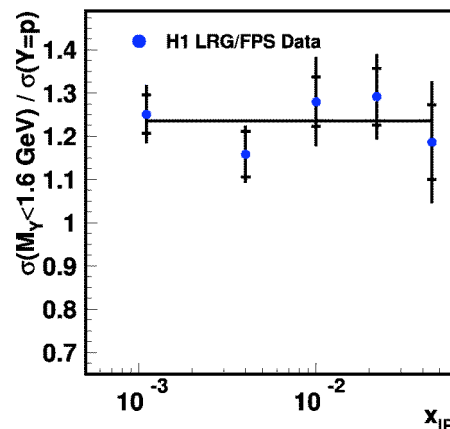
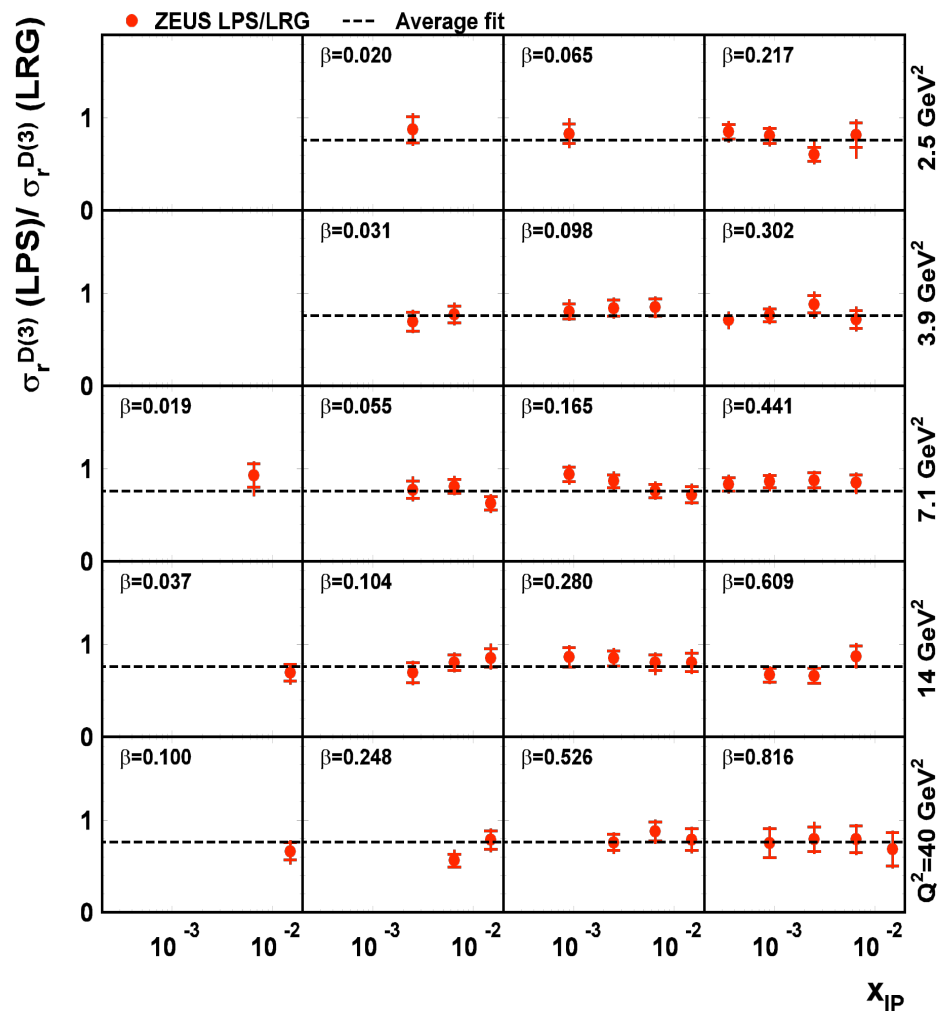
# ZEUS LRG vs H1 LRG



→ Remaining normalisation difference of 13% (global fit) covered by uncertainty on p-diss. correction (8%) and relative normalisation uncertainty (7%)

→ Shape agreement ok except low  $Q^2$

# Ratio of LPS / LRG



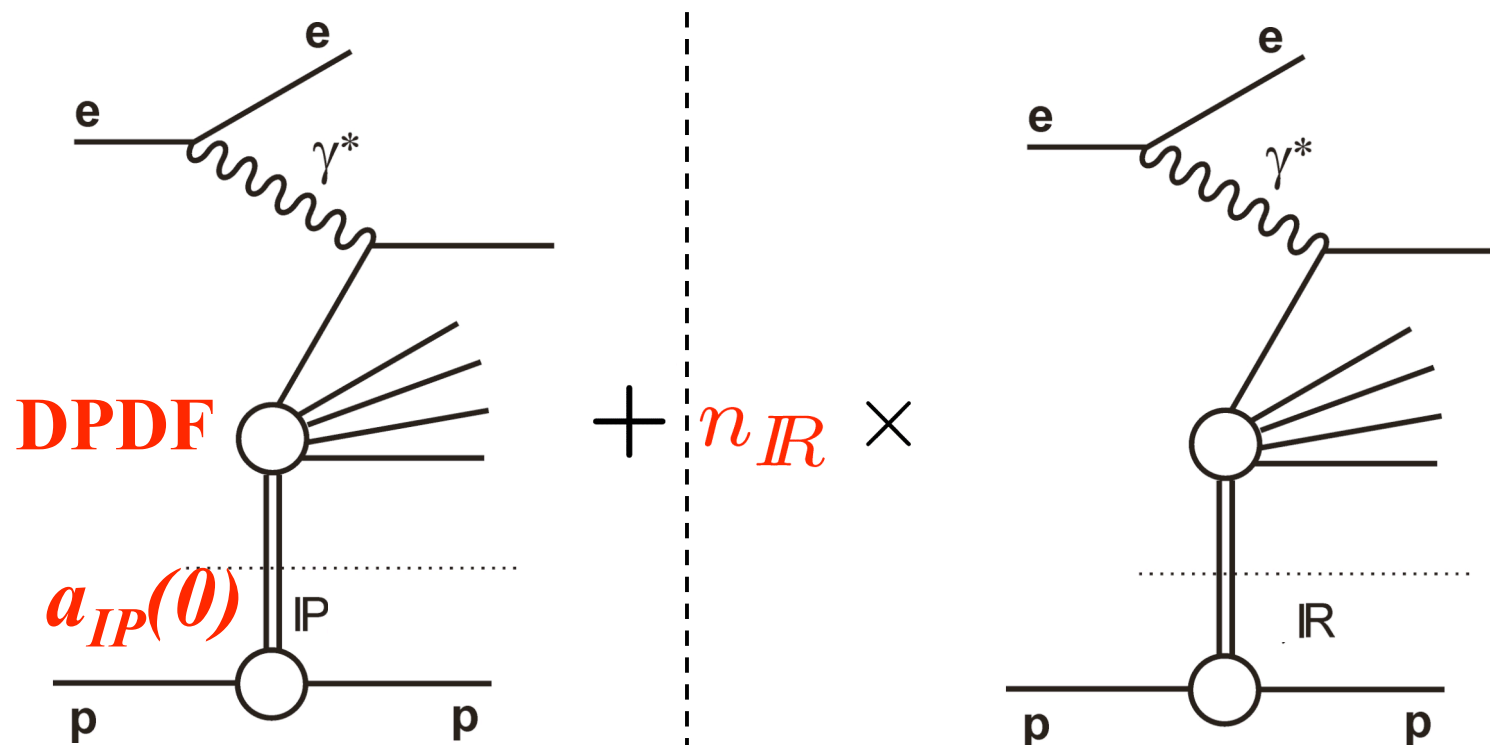
The ratio LPS/LRG cross sections  
is independent of  $Q^2$ ,  $x_{IP}$ ,  $\beta$

ZEUS LPS / ZEUS LRG =  $0.76 \pm 0.01(\text{stat}) \pm 0.03-0.02(\text{sys}) \pm 0.08-0.05(\text{norm})$

→ p-diss. background in LRG data:  $[24 \pm 1(\text{stat}) \pm 2-3(\text{sys}) \pm 5-8(\text{norm})]\%$

H1 LRG / H1 FPS =  $1.23 \pm 0.03(\text{stat}) \pm 0.16(\text{sys})$

# H1 2006 DPDF Fit - Overview



**IP component:**

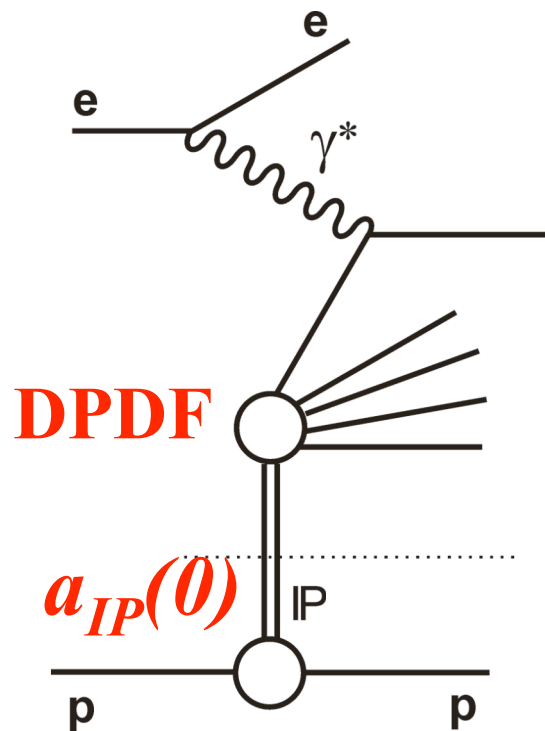
- Fit  $\alpha_{IP}(0)$  ( $x_{IP}$  dependence).
- Simultaneously, fit 5 parameters of DPDFs ( $\beta$  and  $Q^2$  dependences) using NLO QCD.

**IR component:**

- Fit  $n_{IR}$  one parameter for normalisation.
- All flux parameters taken from previous H1 data. PDFs taken from Owens-pion



# H1 2006 DPDF Fit - Details



*IP* component:

- Fit  $\alpha_{IP}(0)$  ( $x_{IP}$  dependence).
- Simultaneously, fit 5 parameters of DPDFs ( $\beta$  and  $Q^2$  dependences) using NLO QCD.

- Parameterise quark singlet  $z\Sigma(z, Q_0^2)$  and gluon  $zg(z, Q_0^2)$  densities, where  $z$  is parton momentum fraction ( $= \beta$  for QPM).

- Parameterisation used is  $z\Sigma(z, Q_0^2) = A_q z^{B_q} (1-z)^{C_q}$  and  $zg(z, Q_0^2) = A_g (1-z)^{C_g}$  (gluon insensitive to  $B_g$ )

- Results reproducible with Chebyshev polynomials.

- Fit is stable with variations of, e.g.  $\beta_{max}$  – the maximum value of  $\beta$  allowed in the fit.

- Fit stable for  $Q_{min}^2 > 8.5 \text{ GeV}^2$ .

- Fit all data with:  
 $Q^2 \geq 8.5 \text{ GeV}^2$  (and  $M_X > 2 \text{ GeV}$ ,  $\beta \leq 0.8$ )

# H1 2006 DPDF Fit - Results

## • Fit A

$$Q_0^2 = 1.75 \text{ GeV}^2$$

$$\chi^2 \sim 158 / 183 \text{ d.o.f.}$$

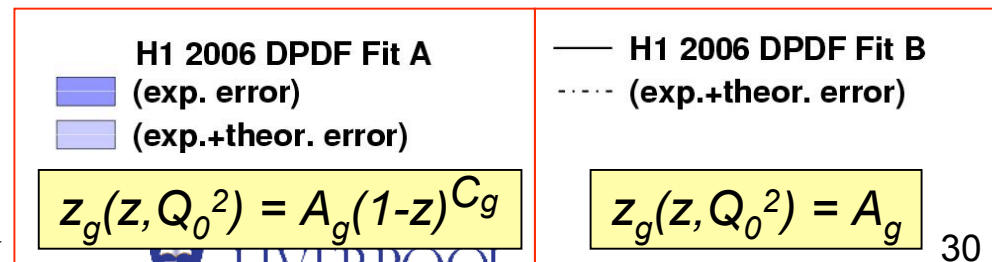
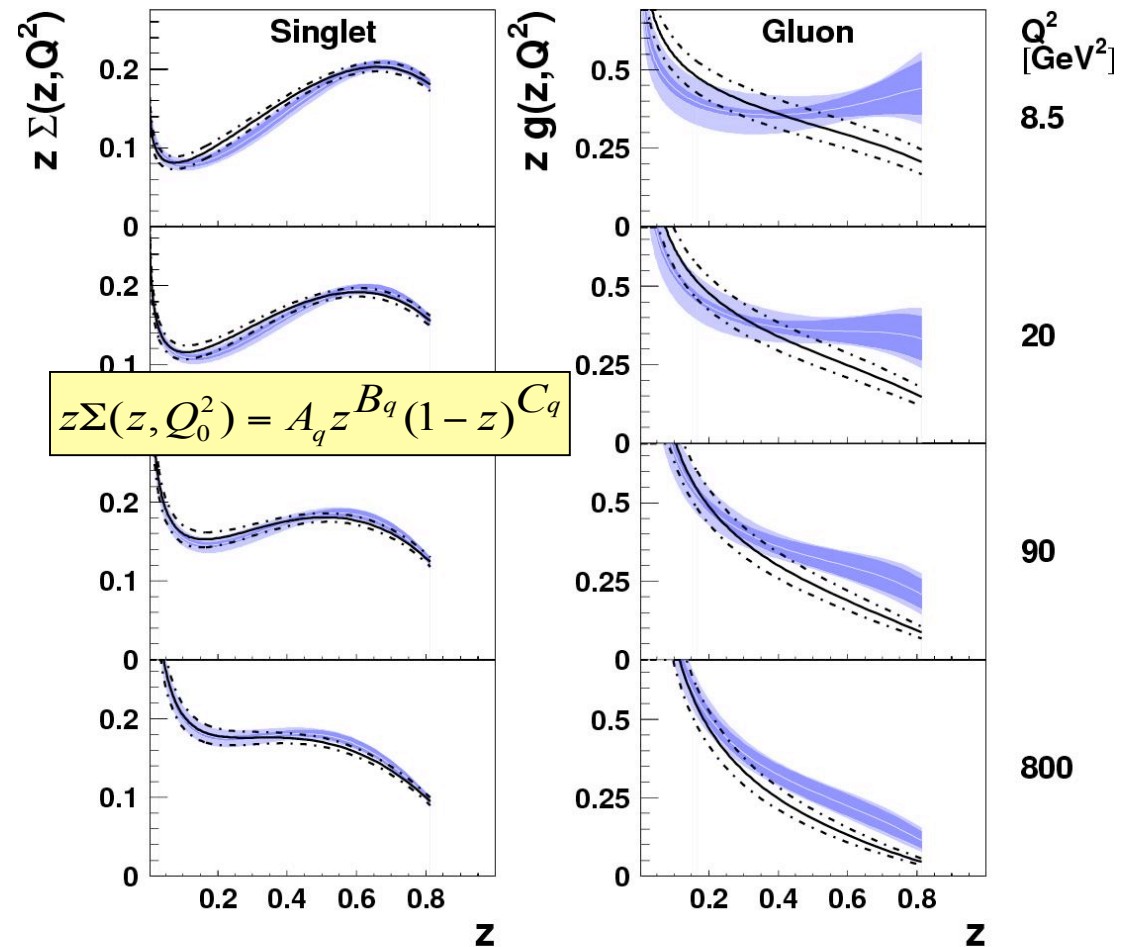
## • Fit B

Drop  $C_g$  - gluon is parameterised as a constant at the starting scale!

$$\chi^2 \sim 164 / 184 \text{ d.o.f.}$$

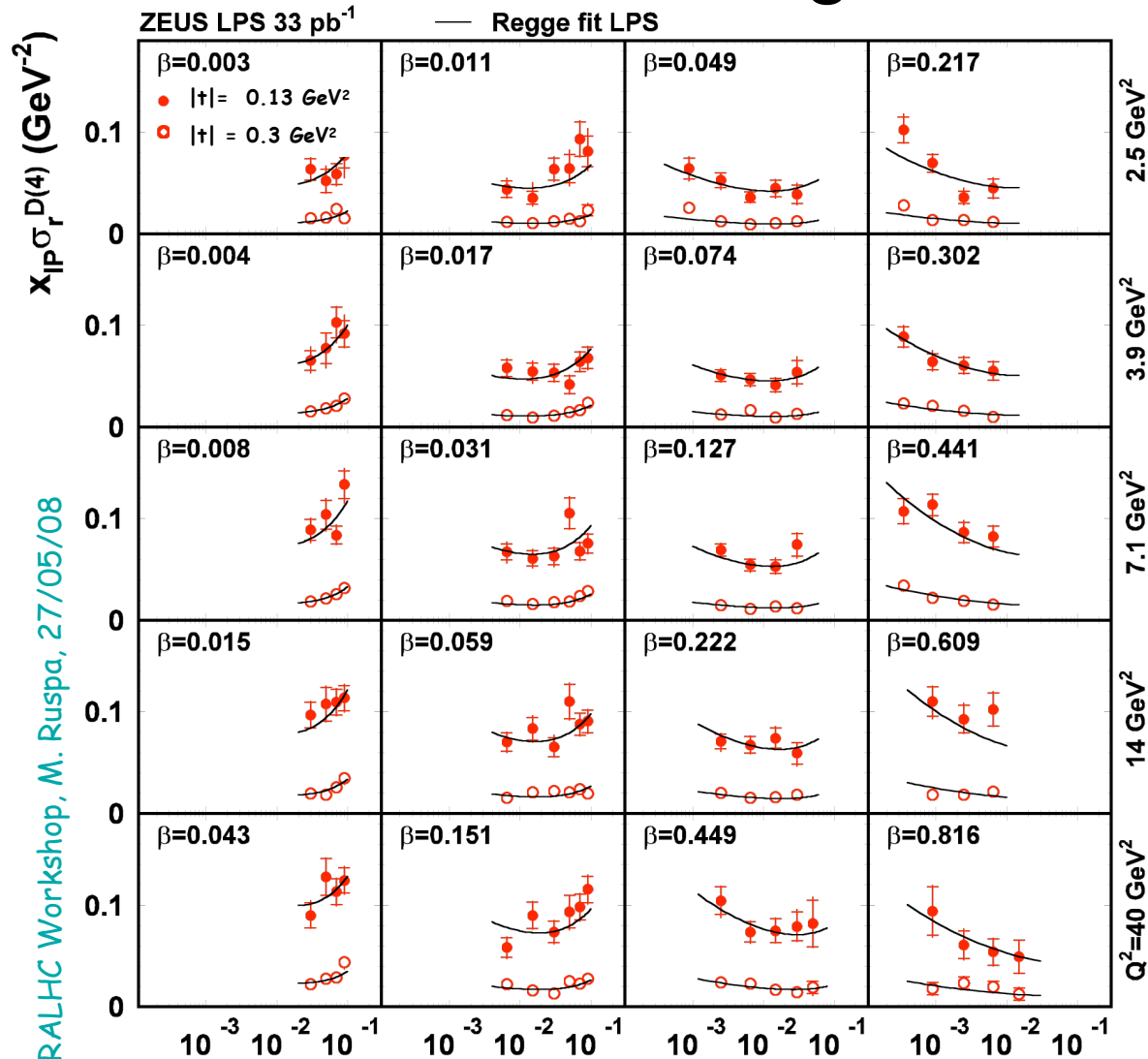
$$Q_0^2 = 2.5 \text{ GeV}^2$$

- Quarks very stable
- Gluon similar at low  $z$
- No sensitivity to gluon at high  $z$



# $x_{IP}$ dependence of $\sigma_r^{D(4)}$

## Leading Proton data



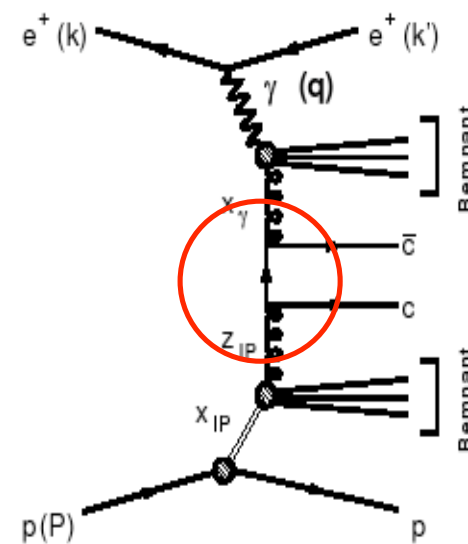
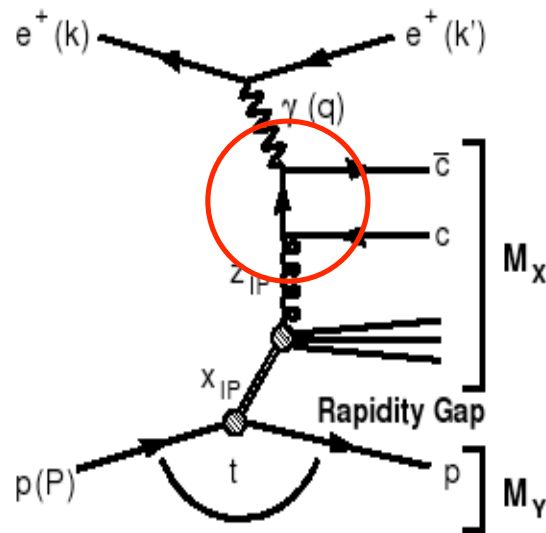
First measurement in two  $t$  bins

Same  $x_{IP}$  dependence in two  $t$  bins

Low  $x_{IP}$ :  $\sigma_r^{D(4)}$  falls with  $x_{IP}$  faster than  $1/x_{IP}$

High  $x_{IP}$ :  $x_{IP} \sigma_r^{D(4)}$  flattens or increases with  $x_{IP}$   
(Reggeon and  $\pi$ )

# Suppression



- It was predicted that the suppression would depend on the configuration of the photon in photoproduction - it doesn't
- It seems we only need a hard scale to resolve the scatter
- That can be provided by the  $Q^2$  of the photon in DIS or by the  $E_T$  of the jets in photoproduction
- The latter seems to require playing the hard interaction backwards, but there is no time-ordering of that hard interaction