

# *Leading Neutrons at HERA*

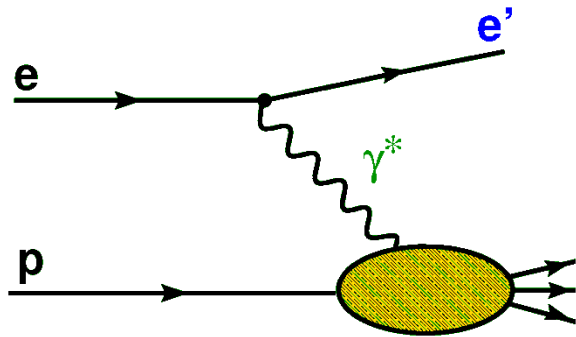
**Vitaliy Dodonov, MPI for Nuclear Physics, Heidelberg**  
**On behalf of the H1 and ZEUS collaborations**



**International Workshop on Diffraction in High-Energy Physics**  
**La Londe-les-Maures, France, September 9-14 2008**

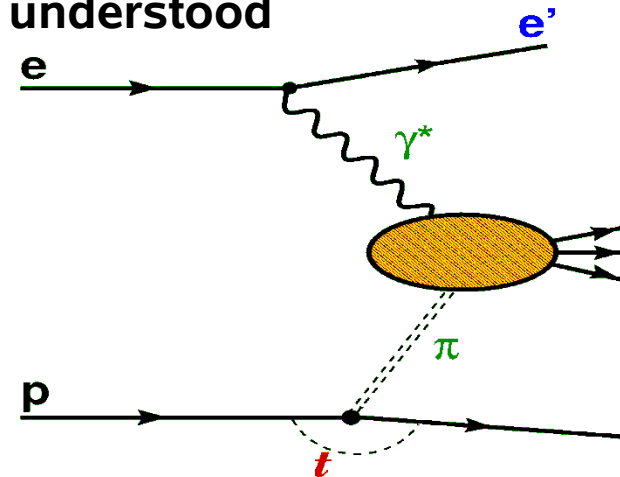
# Motivation

- Significant fraction of  $ep$  scattering events contain a leading neutron in the final state carrying a substantial portion of the energy of the incoming proton:  
 $e+p \rightarrow e+n+X$
- Production mechanism is not yet completely understood



Leading neutron can come from “standard fragmentation”

- From hadronization of p remnant
- Implemented in MC models (e.g. Lund)



Leading neutron can be produced via exchange of virtual particle

- charged iso-vector ( $\pi^+$ ,  $\rho^+$ )

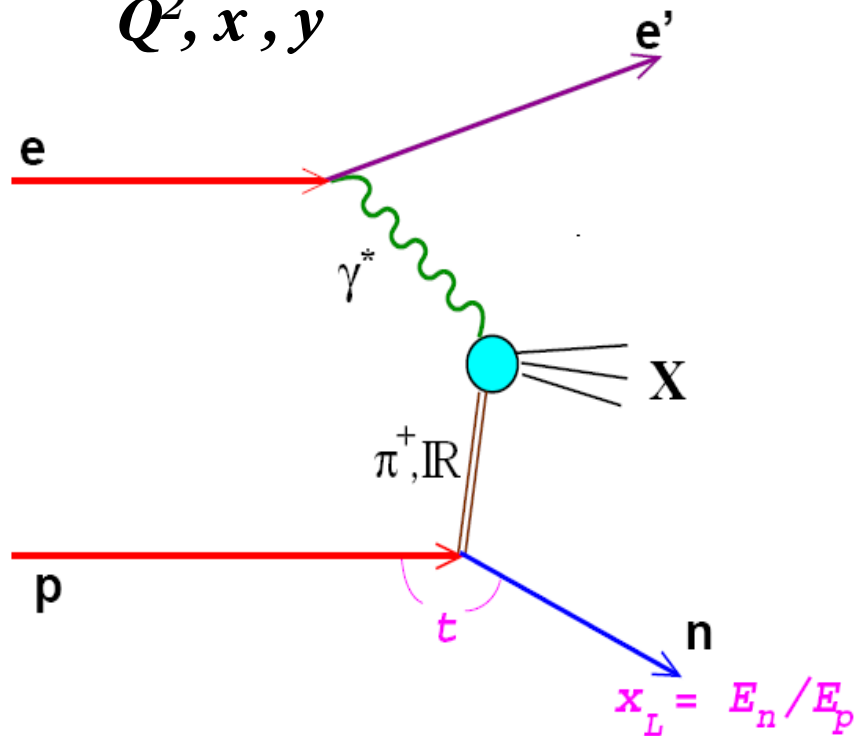
Results discussed in this talk:

- Leading Neutron spectra in DIS and photoproduction
- Leading Neutron production cross section
- Comparison with models

## Photon vertex:

lepton variables:

$$Q^2, x, y$$



## Proton vertex:

leading baryon variables:

$$x_L = E_n / E_p$$

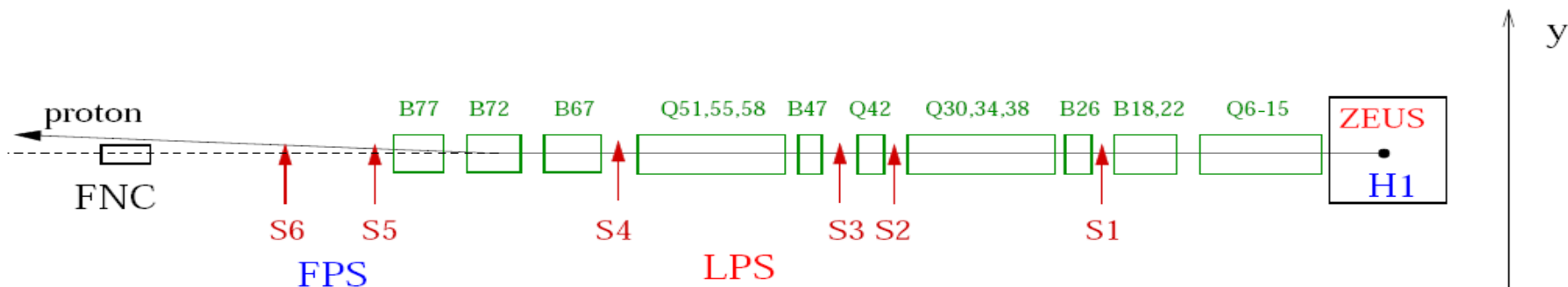
$$t = (\mathbf{p} - \mathbf{p}')^2 \approx -\frac{\mathbf{p}_T^2}{x_L} - \frac{(1 - x_L)^2}{x_L} m_n^2$$

In the  $\pi$ -exchange model cross section dependence on leading baryon variables is independent of kinematics at photon vertex

$$\sigma_{ep \rightarrow enX}(x, Q^2, x_L, t) = f_{\pi^+/p}(x_L, \mathbf{p}_T) \cdot \sigma_{e\pi \rightarrow eX}(x/(1-x_L), Q^2)$$

Leading Neutron production affected by absorption and rescattering effects: evidences of vertex factorization violation

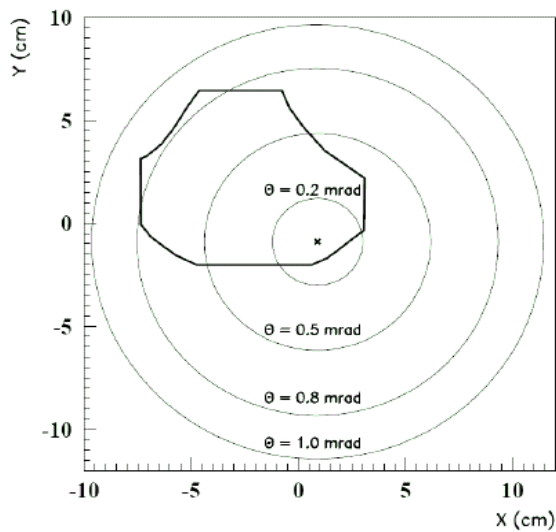
# H1 and ZEUS Forward Neutron Calorimeters



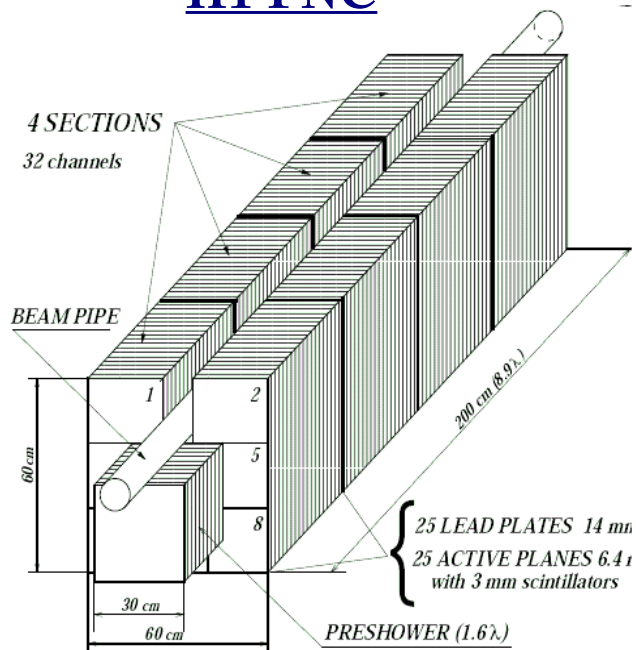
Position – 105m from the interaction point.

Geometrical acceptance is limited by beam-line elements:  $< 0.8$  mrad.

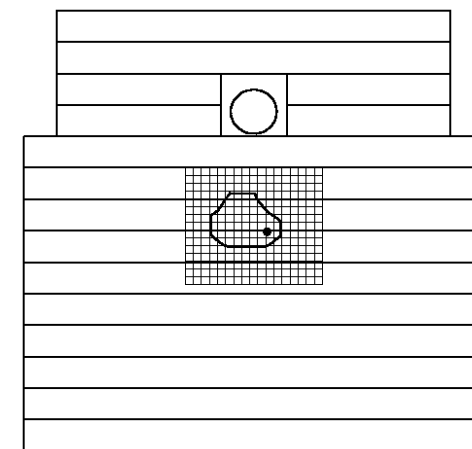
## Geometrical Acceptance



## H1 FNC



## ZEUS FNC



14 towers, 17x15 grid of the FNT hodoscopes,  
 $\sigma_E/E \approx 0.7/\sqrt{E}$

$\sim 10 \lambda$  lead-scintillator sandwich calorimeters

## H1: $F_2^{\text{LN}(3)}(Q^2, x, x_L)$ - triple diff. cross section of LN production in DIS

- HERA-II data, much higher statistics compared to old result H1prelim-08-111
- Upgraded Forward Neutron Calorimeter (better resolution and photon identification)
- Kinematic range:  $6 \text{ GeV}^2 < Q^2 < 100 \text{ GeV}^2$ ,  $1.5 \times 10^{-4} < x < 3 \times 10^{-2}$ ,  $p_{T,n} < 0.2 \text{ GeV}$

## ZEUS: relative to inclusive cross section vs $x_L$ and $p_T$ in DIS and $\gamma p$

Leading Neutrons are selected from inclusive data sets (i.e. no LN tag): Nucl.Phys.B776 (2007) 1

- DIS:  $Q^2 > 2\text{-}3 \text{ GeV}^2$ , 3 subsets  $Q^2 \approx 2.7, 8.9, 40 \text{ GeV}^2$
- $\gamma p$ :  $Q^2 < 0.02 \text{ GeV}^2$ ,  $e^+$  tagged  $\Rightarrow 150 < W_{\gamma p} < 270 \text{ GeV}$

Leading Neutron yields:

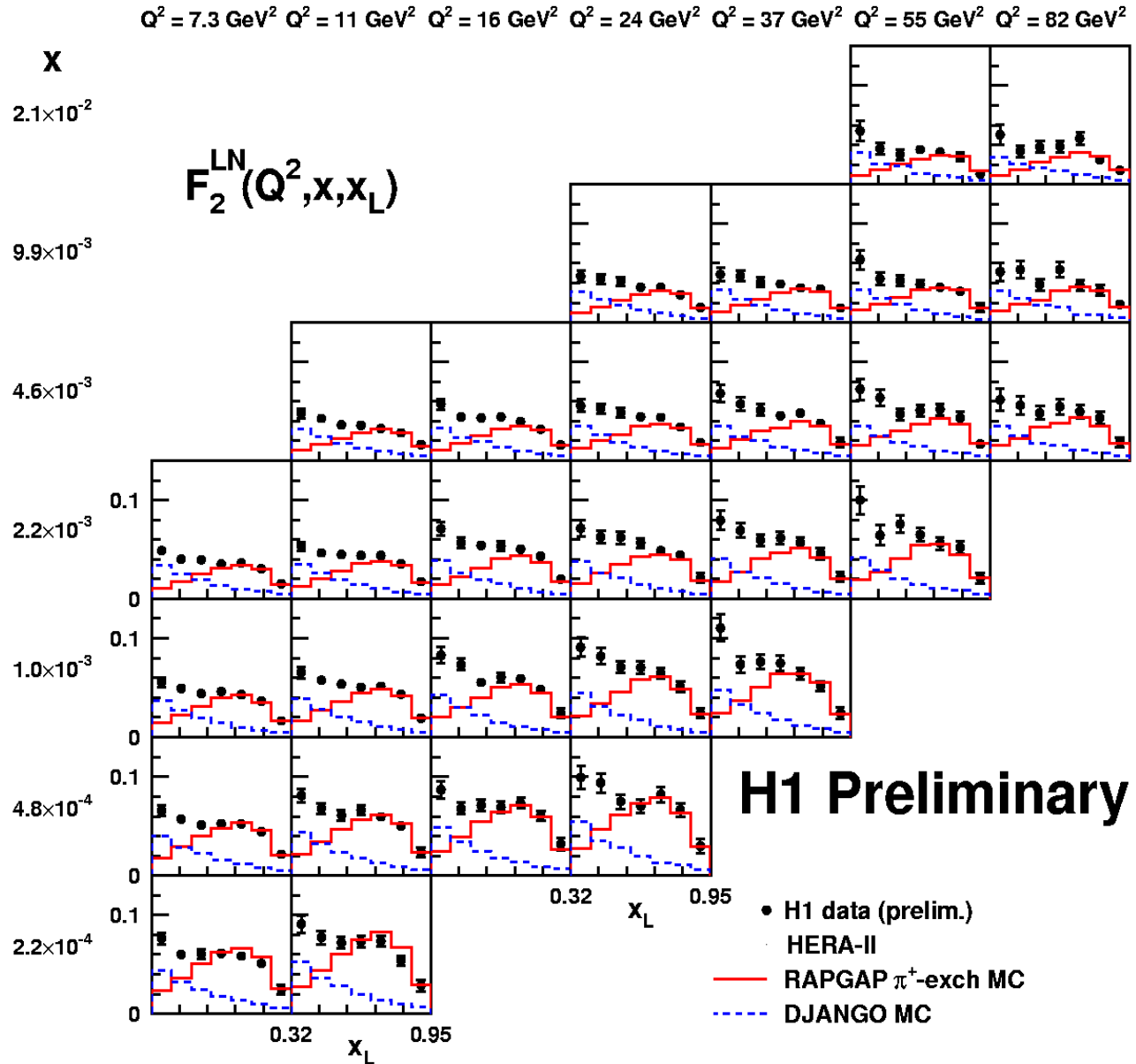
- DIS,  $\gamma p$  have very different inclusive cross sections  $\sigma_{\text{inc}}$
- For sensible comparisons look at LN yields:  $r_{\text{LB}} = \sigma_{\text{LB}}/\sigma_{\text{inc}}$
- Additional benefit: systematic uncertainties of central detector cancel

# Triple differential reduced cross section

$$\frac{d^3 \sigma(ep \rightarrow eNX)}{dQ^r dx dx_L} =$$

$$= \frac{4\pi\alpha^2}{xQ^\epsilon} \left[ 1 - y + \frac{y^2}{\Upsilon} \right] F_2^{\text{LN}}(Q^r, x, x_L)$$

- DJANGO (standard fragmentation) predicts too low cross section +  $x_L$  spectrum shape is too different
- RAPGAP  $\pi^+$ -exch. describes data well for  $x_L > 0.7$



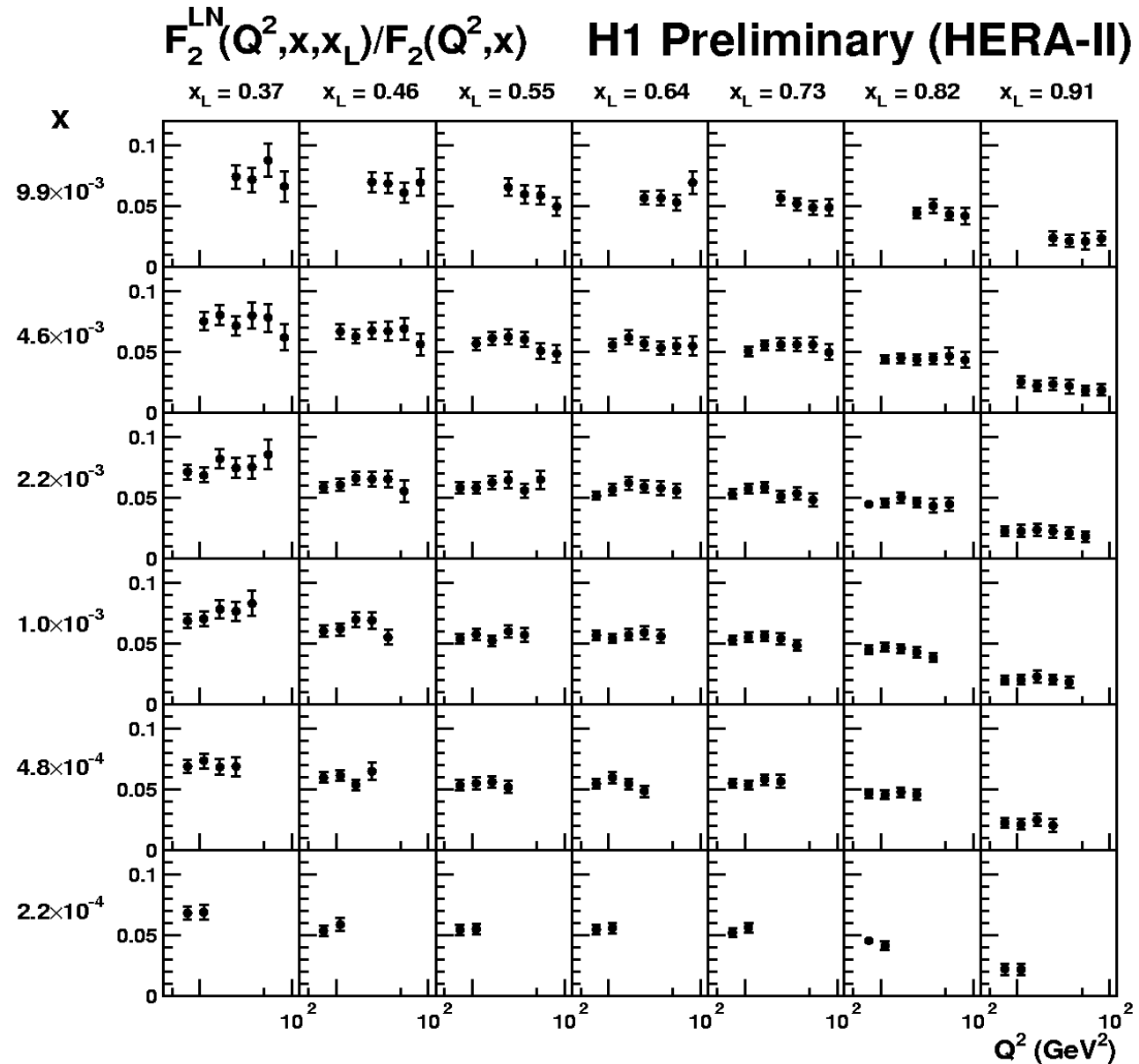
# $F_2^{\text{LN}}(Q^2, x, x_L)$ to $F_2(Q^2, x)$ ratio

$F_2(Q^2, x)$  from the H1 parameterization  
(Eur.Phys.J.C21 (2001) 33)

$F_2^{\text{LN}}(Q^2, x, x_L)/F_2(Q^2, x)$   
is mostly flat in  $Q^2$  and  $x$

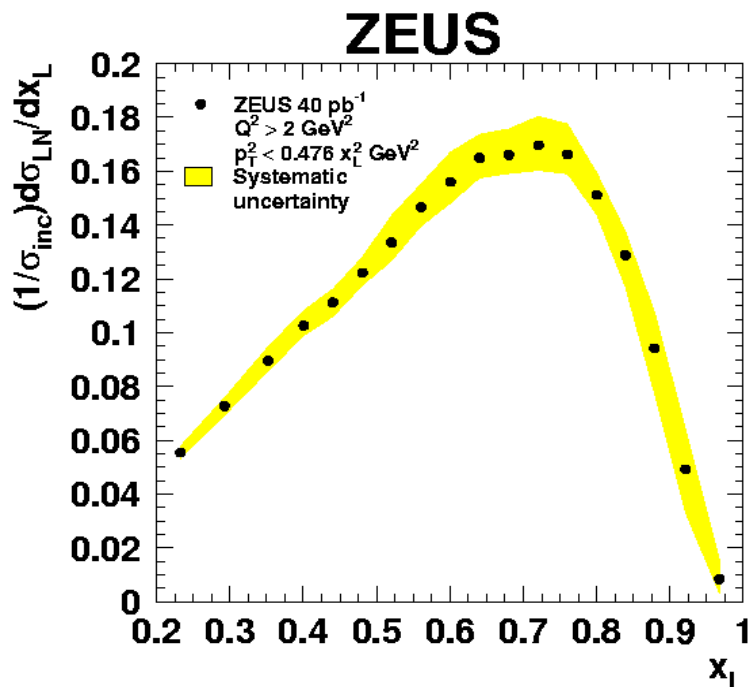
=> consistent with vertex factorization  
(or common suppression of LN cross section)

**LN production rate,  
kinematics is approx.  
independent of  $(Q^2, x)$**



# DIS: $x_L$ and $p_T^2$ distributions

Leading Neutron:  $p_T^2 < 0.476 x_L^2 \text{ GeV}^2$

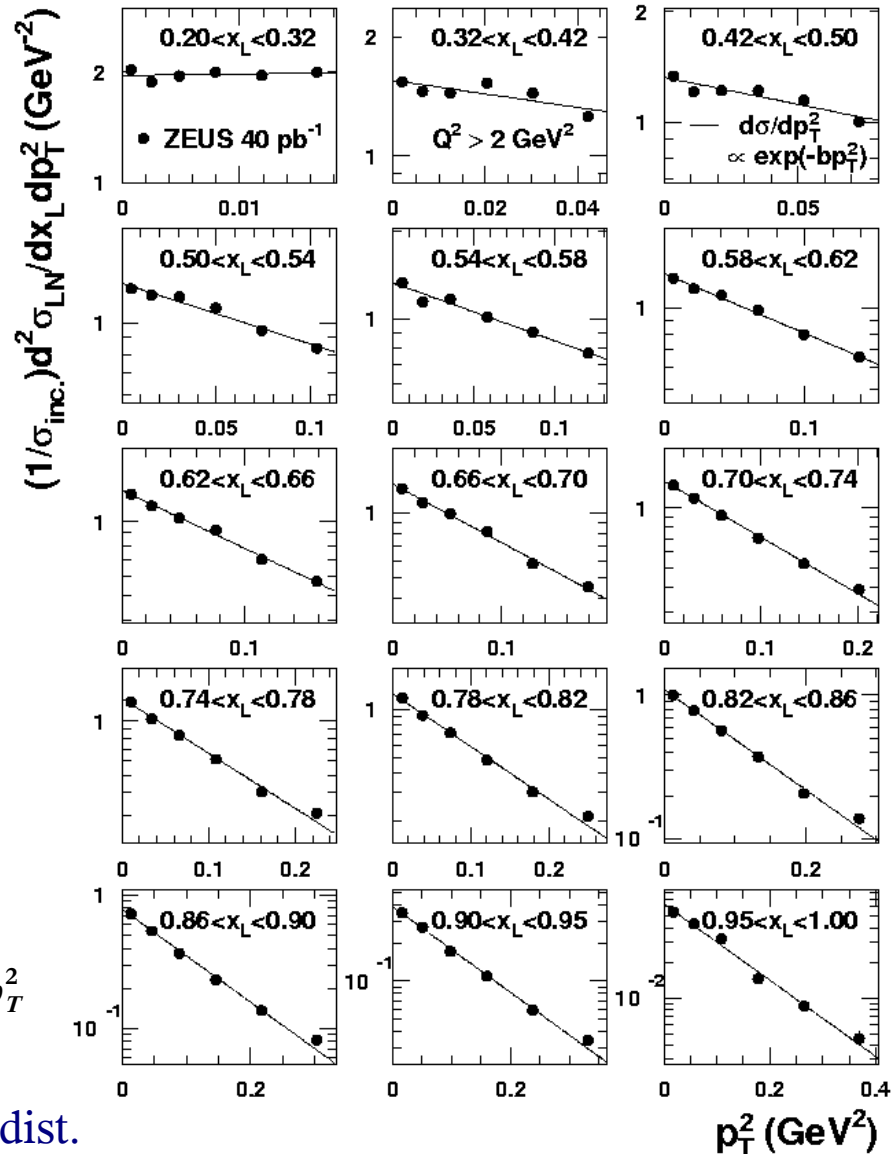


- LN yield  $\rightarrow 0$  at kin. limit  $x_L \rightarrow 1$
- Below  $x_L \approx 0.7$  yield drops due to decreasing  $p_T^2$  range

Fit by exp. in  $p_T^2$ : 
$$\frac{1}{\sigma_{\text{inc}}} \frac{d^2 \sigma_{\text{LN}}}{dx_L dp_T^2} = a(x_L) \cdot e^{-b(x_L) p_T^2}$$

Intercept  $a(x_L)$  and slope  $b(x_L)$  fully describe  $(x_L, p_T^2)$  dist.

Log Scale



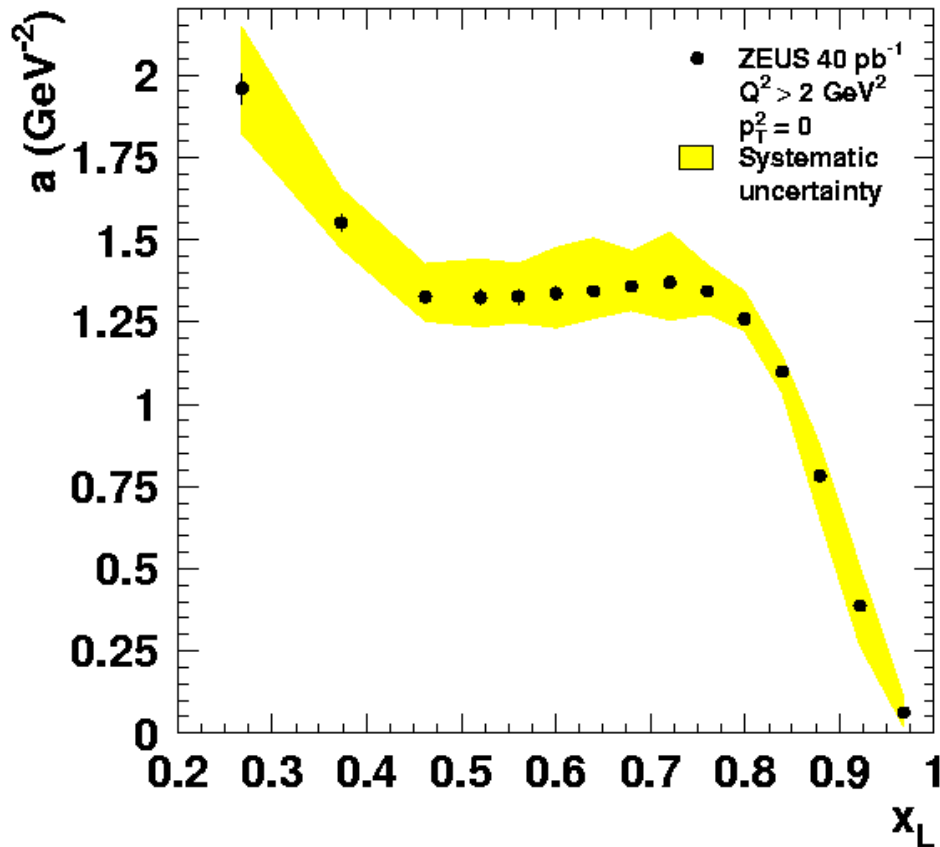


# DIS $p_T^2$ distributions: intercepts and slopes

$$\frac{1}{\sigma_{\text{inc}}} \frac{d^2 \sigma_{\text{LN}}}{dx_L dp_T^2} = a(x_L) \cdot e^{-b(x_L) p_T^r}$$

Intercepts  $a(x_L)$

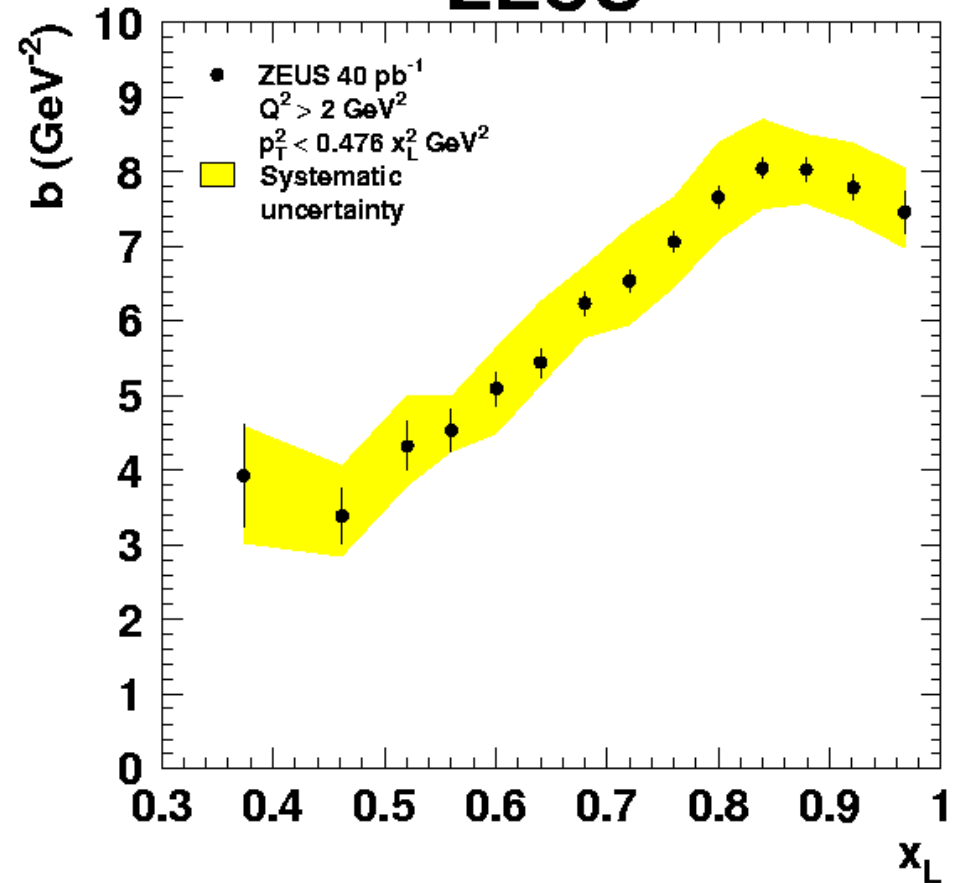
**ZEUS**



- LN intercepts fall with  $x_L$
- bump/plateau/shoulder  $0.4 < x < 0.8$

Slopes  $b(x_L)$

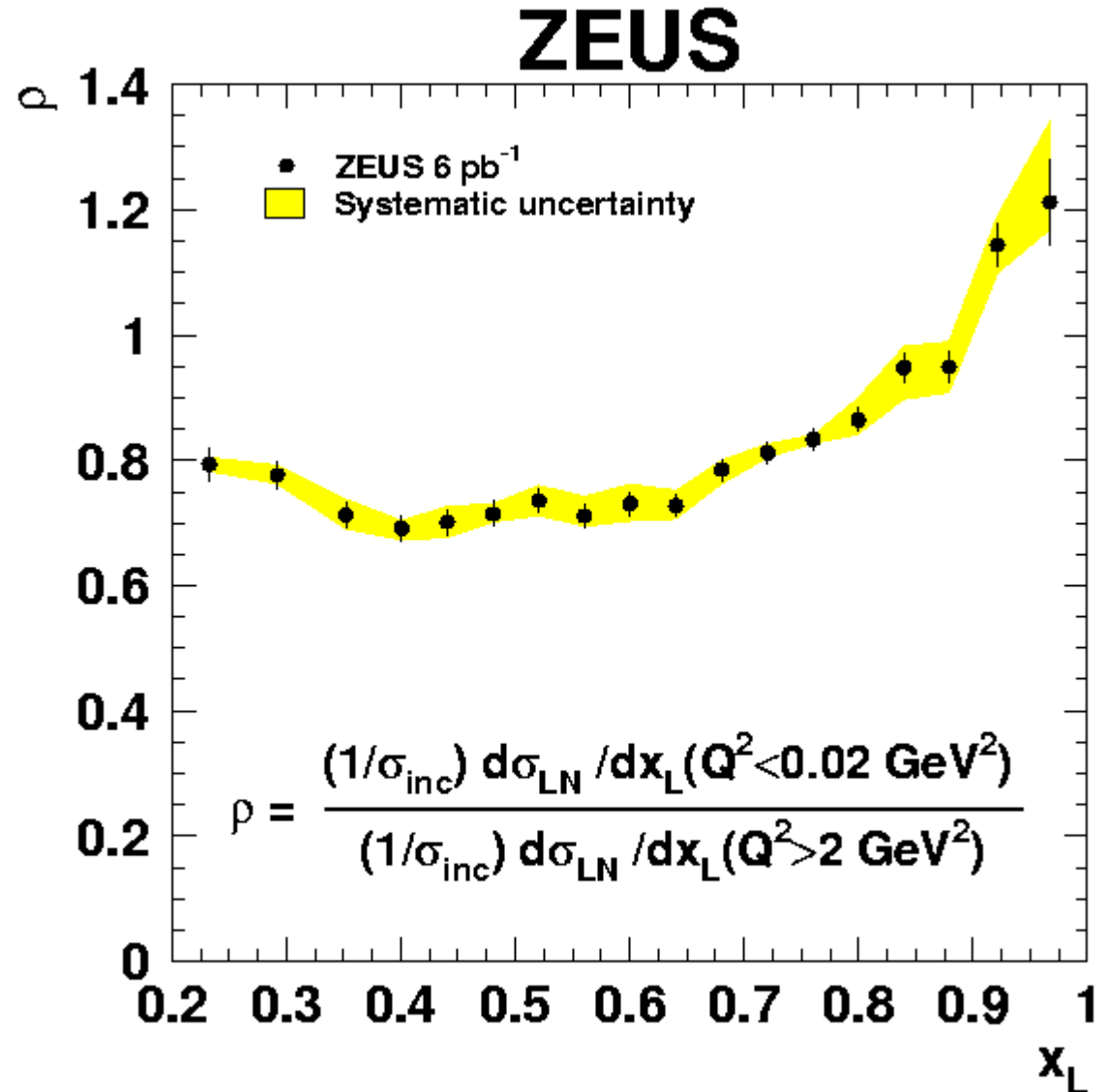
**ZEUS**



- LN slopes sharp rise with  $x_L$

## Compare $\gamma p$ /DIS: $x_L$ distributions

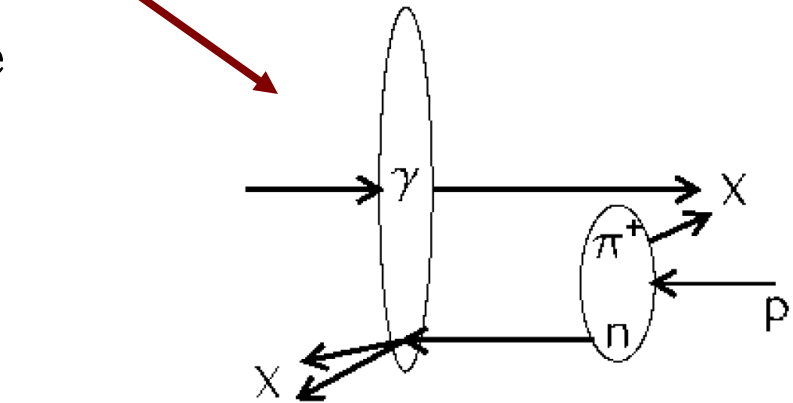
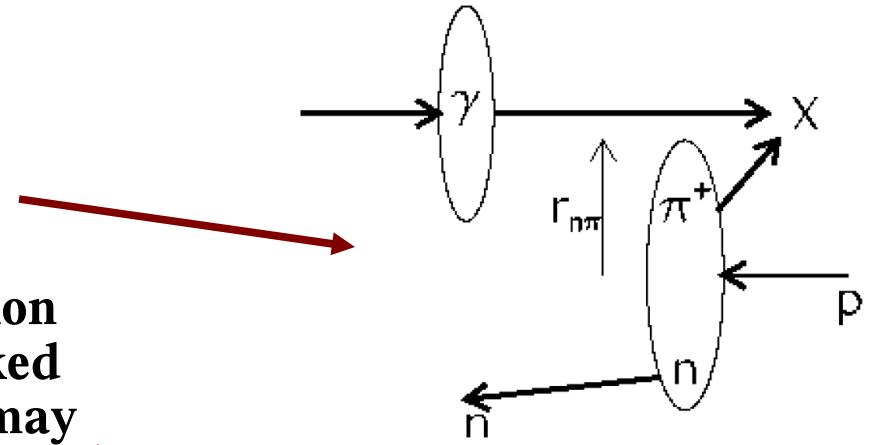
- Combine all DIS  $Q^2 > 2 \text{ GeV}^2$ , compare to  $\gamma p$   $x_L$  dist.
- Ratio  $\sim 70\%$  mid- $x_L$ , rising above 1 as  $x_L \rightarrow 0.9$



# Exchange model refinements: Absorptions (Rescattering)

For e.g. LN production via  $\pi$ -exchange:

- In DIS  $\gamma^*$  is small; small chance both  $n, \pi$  scatter on  $\gamma^*$ :  $n$  reaches detector
- In photoproduction  $\gamma^*$  large; if  $n$ - $\pi$  separation smaller rescattering of  $n$  may occur:  $n$  kicked to lower  $x_L$  and higher  $p_T$  (migration) and may escape detection (absorption loss)
- In another language: multi-Pomeron exchange
- Compare photoproduction and DIS:
  - $x_L, p_T$  distributions
  - effects of absorption?
- Effects of absorption/migration estimated:  
D' Alesio, Pirner;  
Nikolaev, Speth, Zakharov;  
Kaidalov, Khoze, Martin, Ryskin;  
Kopeliovich, Potashnikova, Schmidt, Soffer

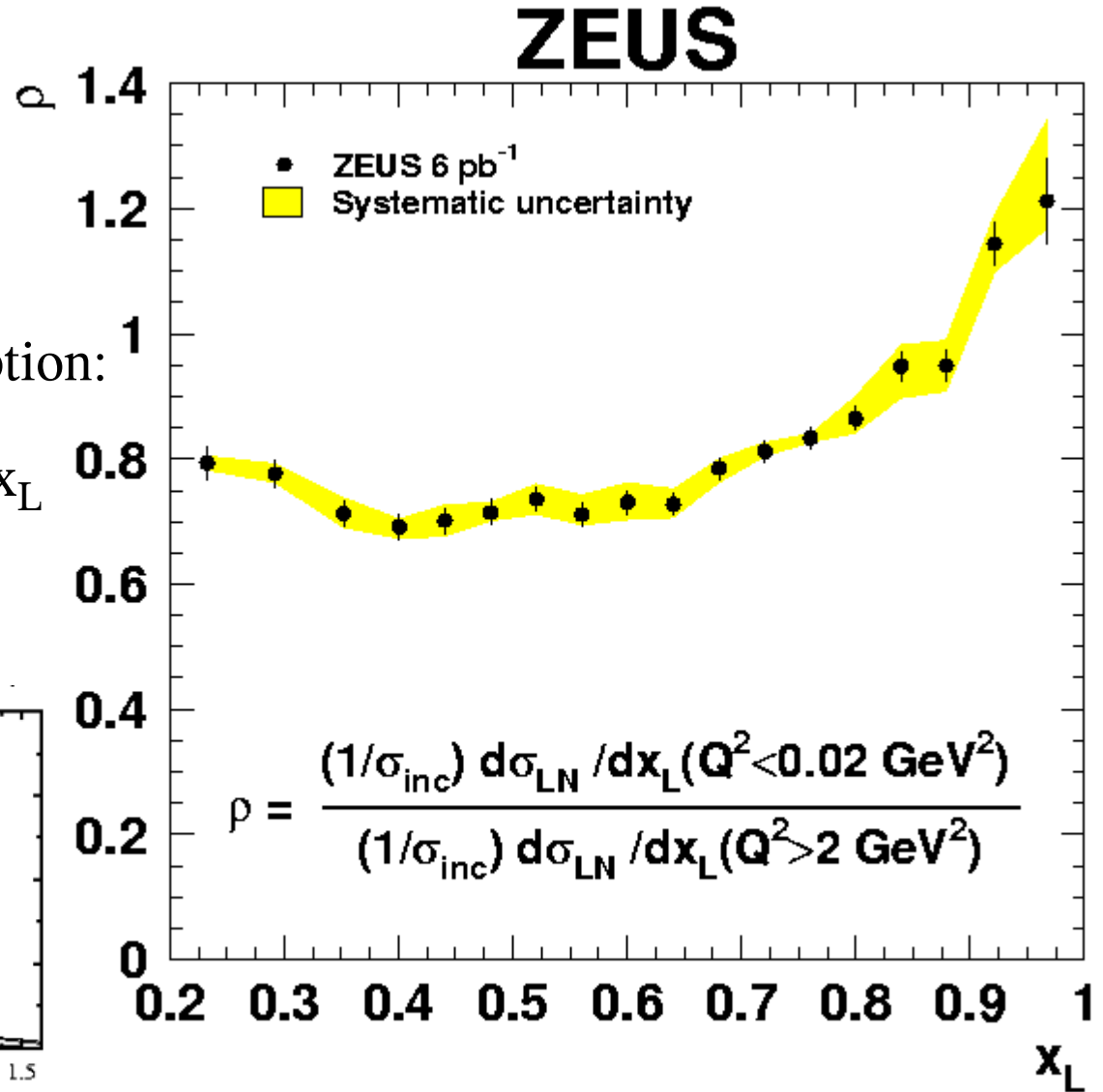
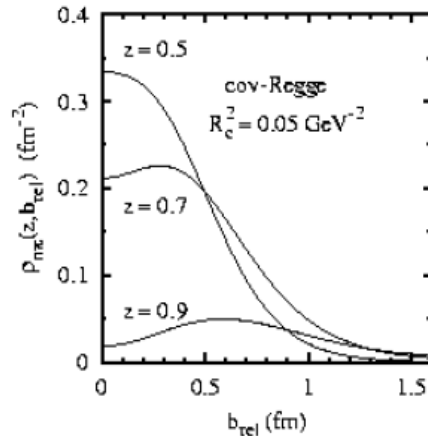
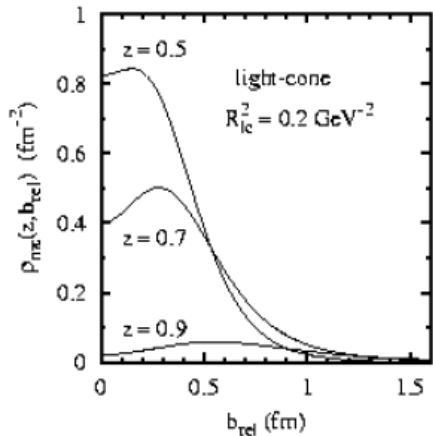


# Compare $\gamma p$ /DIS: $x_L$ distributions

- Combine all DIS  $Q^2 > 2 \text{ GeV}^2$ , compare to  $\gamma p$   $x_L$  dist.
- Ratio  $\sim 70\%$  mid- $x_L$ , rising above 1 as  $x_L \rightarrow 0.9$

Qualitatively consistent with absorption:

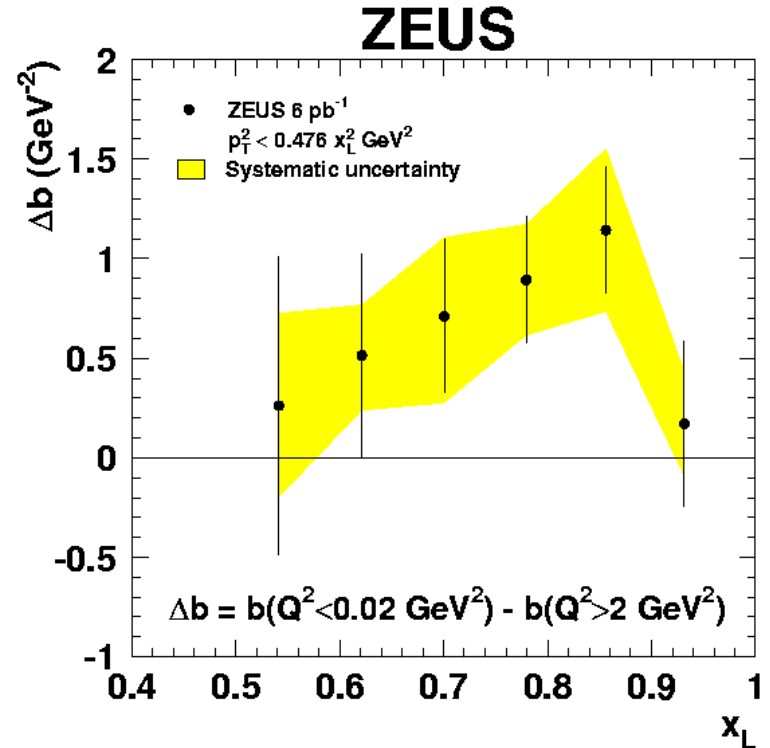
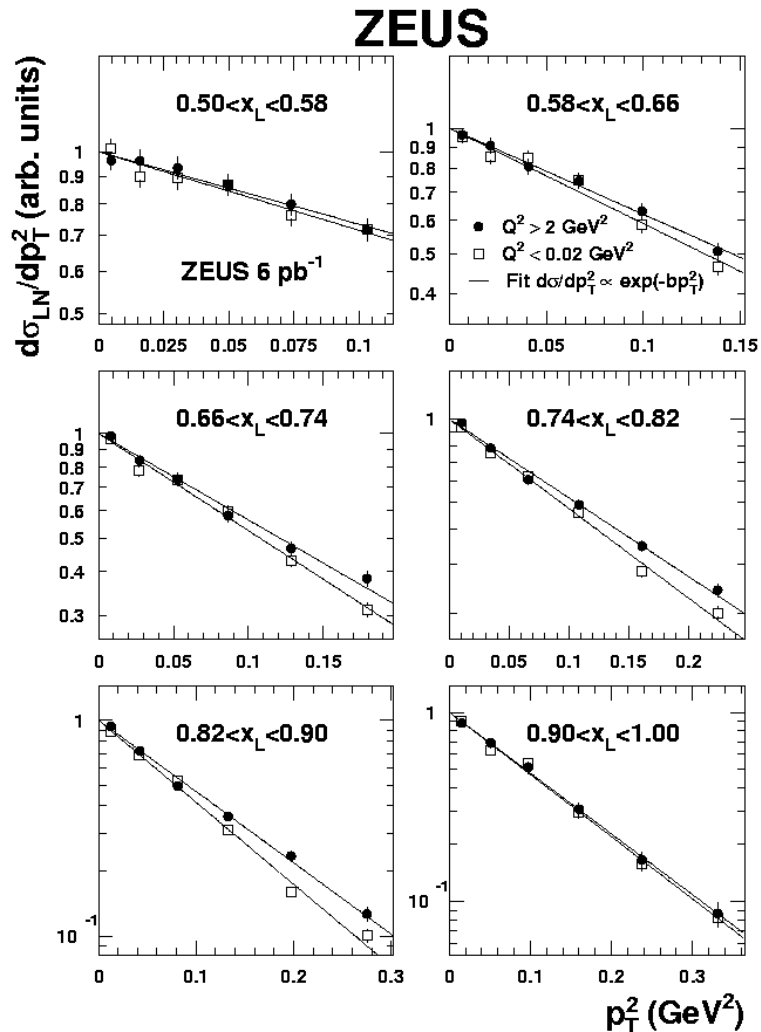
- Exchange model: mean  $n$ - $\pi$  separation  $r_{n\pi}$  decreases at lower  $x_L$
- Smaller  $r_{n\pi} \Rightarrow$  more absorption at lower  $x_L$



# Compare $\gamma p$ /DIS: $p_T^2$ distributions

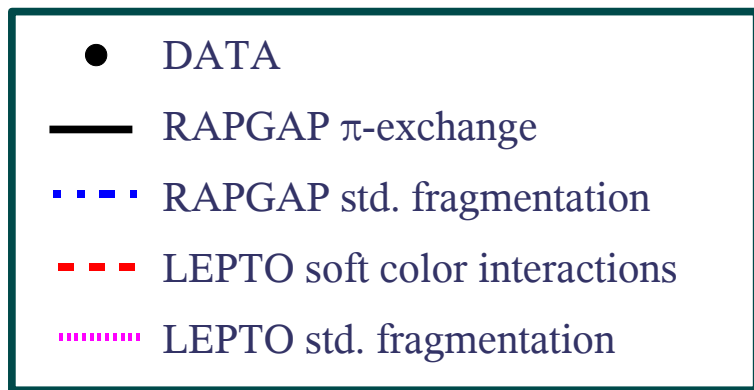
normalized at  $p_T^2 = 0$

$$\Delta b = b(\gamma p) - b(\text{DIS})$$

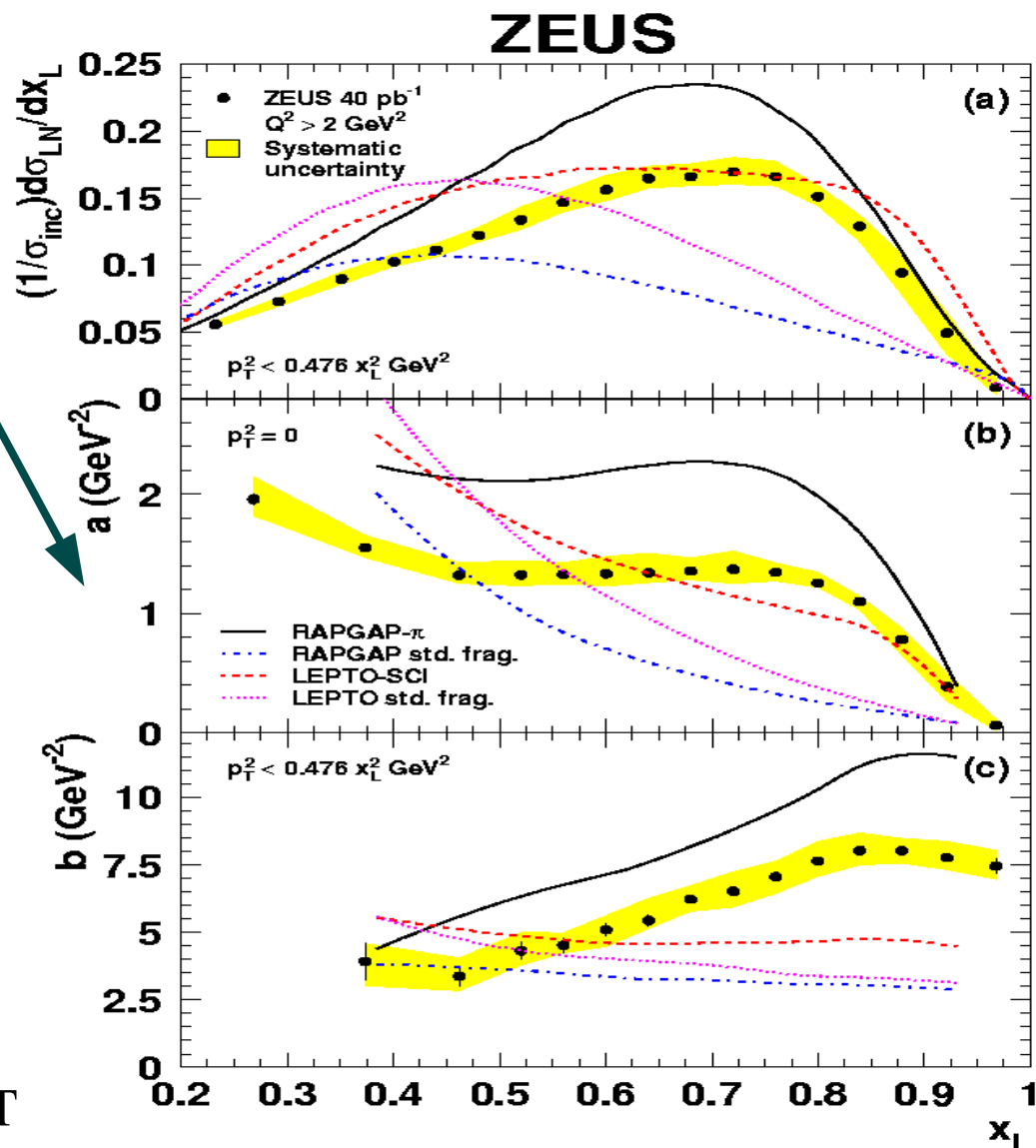


- Small but clear difference:  
 $b(\gamma p) > b(\text{DIS})$  for  $0.6 < x_L < 0.9$
- Qualitatively consistent with absorption:  
 more absorp. at small  $r_{n\pi} \sim$  large  $p_T$   
 fewer LN at high  $p_T \Rightarrow$  larger slope

# Model comparisons: DIS

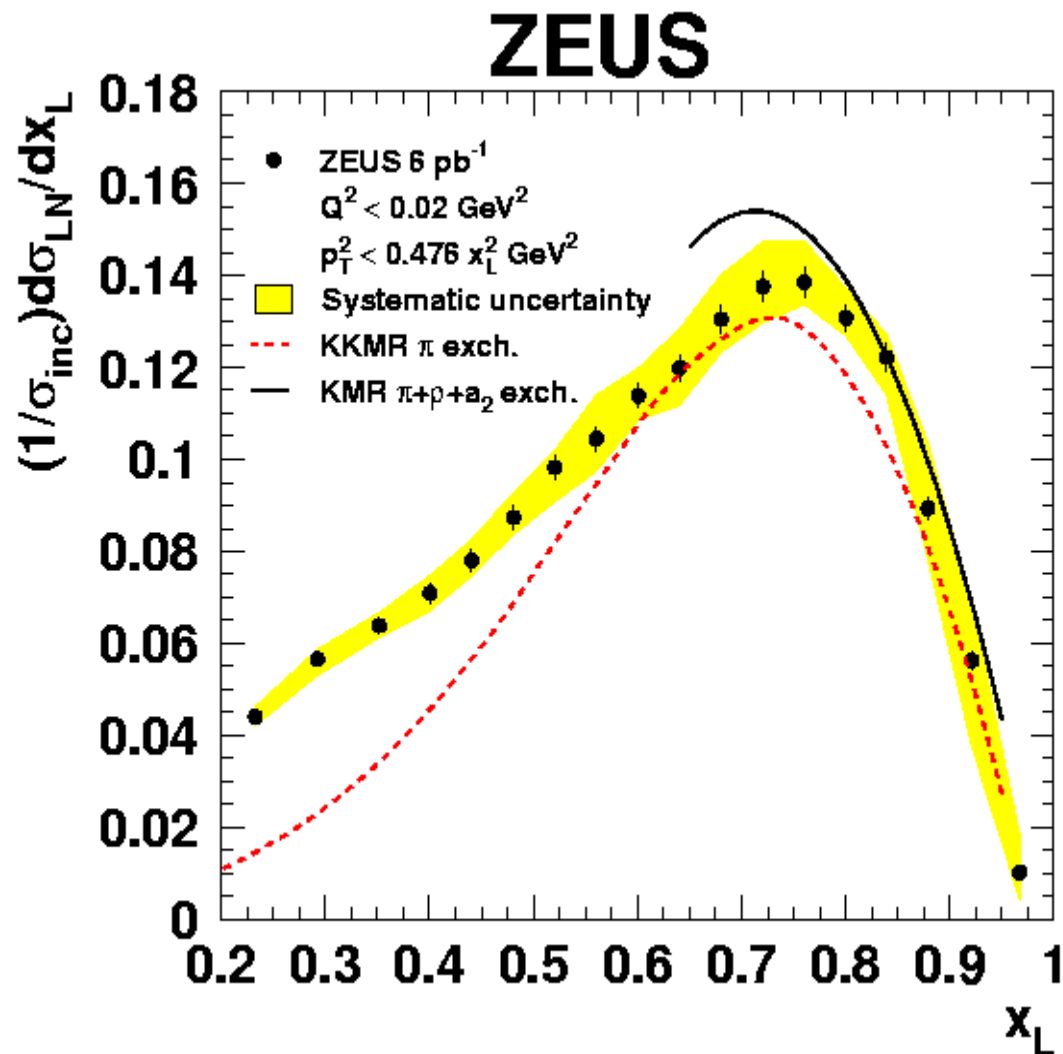


- None of models describe data well
- Both std. frag. too few n, too low  $x_L$
- LEPTO-SCI ~ OK in shape, magnitude, but slopes too small, not  $x_L$  dependent
- RAPGAP  $\pi$ -exch. closest to data (but slopes too high)
- Other DIS,  $\gamma p$  std. frag. models also fail: ARIADNE, CASCADE, PYTHIA, PHOJET

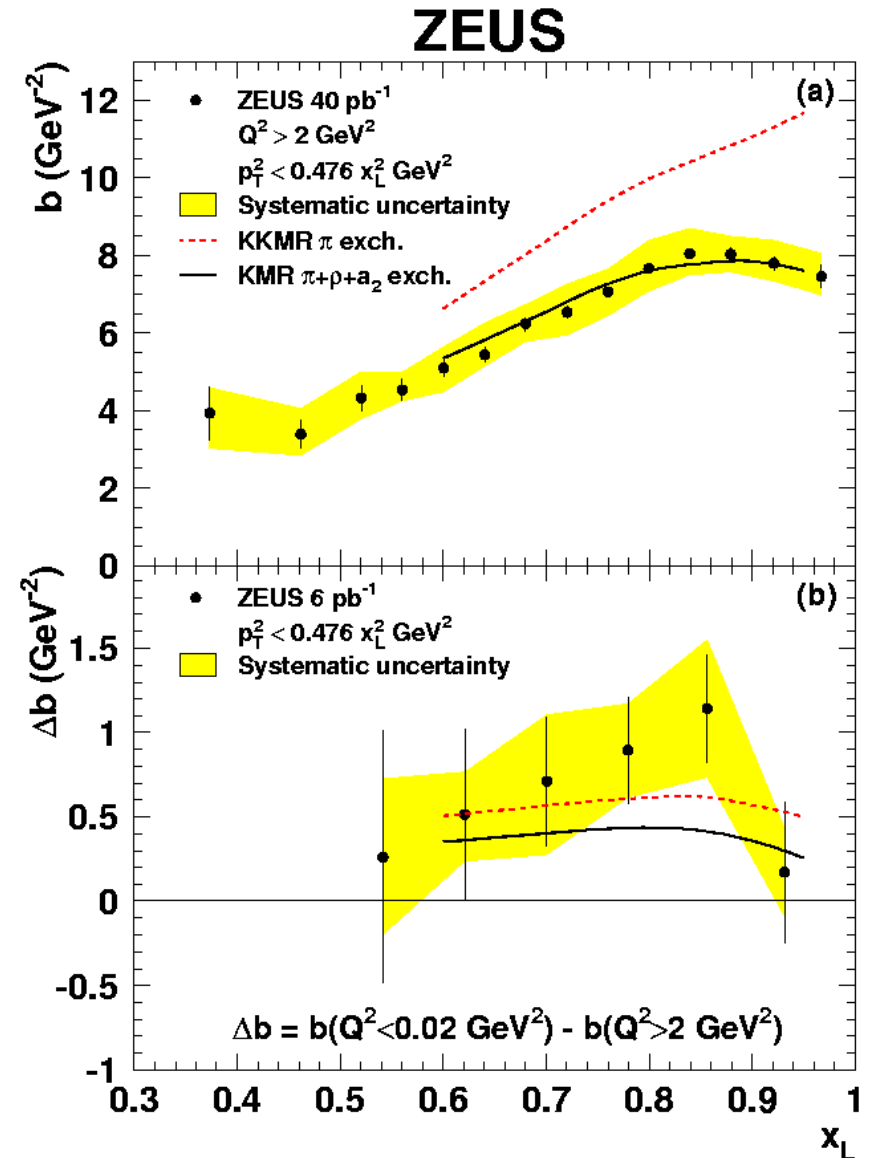


# Compare: $\pi$ -exch. with absorption and $(\rho, a_2)$ exchanges

- Compare with calculations of Kaidalov, Khoze, Martin and Ryskin:
  - start with pure  $\pi$ -exch.
  - some  $n$  rescatter on  $\gamma$
  - rescattered  $n$  migrate in  $(x_L, p_T)$
  - add  $(\rho, a_2)$  exchanges
- Overall loss  $\sim 50\%$  from pure OPE
- Reasonable agreement with LN in  $\gamma p$



- Absorption+migration with pion exchange alone doesn't describe slopes; too high in magnitude, no turnover at high  $x_L$ ;  $\Delta b$  is  $\sim$  OK
- Addition of  $(\rho, a_2)$  exchanges gives good description of both slopes magnitude and  $x_L$  dependence,  $\Delta b$  is still OK





# Pion Structure Function from $F_2^{\text{LN}}$

According to  $\pi$ -exchange model we can extract  $F_2^\pi$  from measured  $F_2^{\text{LN}}$ :

$$F_2^{\text{LN}(3)}(\beta, Q^2, x_L) = \Gamma_\pi(x_L) \cdot F_2^\pi(\beta, Q^2)$$

where

$$\beta = x/(1-x_L)$$

$\Gamma_\pi(x_L)$  is integrated over  $t$  pion flux

For pion flux

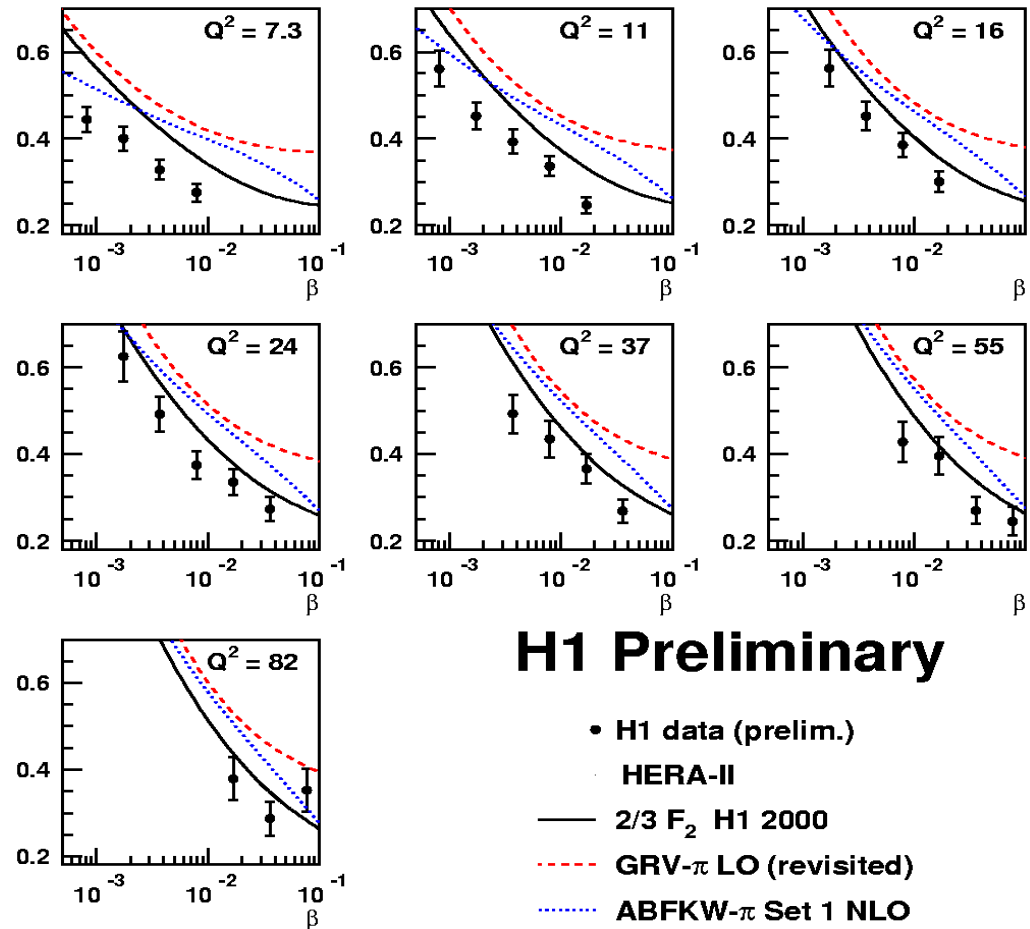
$$f_{\pi^+/p} = \frac{1}{2\pi} \frac{g_{p\pi n}^2}{4\pi} (1-x_L) \frac{-t}{(m_\pi^2 - t)^2} \exp\left(-R_{\pi n}^2 \frac{m_\pi^2 - t}{1-x_L}\right)$$

at  $x_L = 0.73$  we obtained  $\Gamma_\pi = 0.131$

Data compared to parameterizations:

- 2/3 of  $F_2 \Rightarrow$  additive quark model
- GRV- $\pi$  LO (revisited)
- ABFKW- $\pi$  Set 1 NLO

$$F_2^{\text{LN}(3)}(x_L = 0.73)/\Gamma_\pi, \Gamma_\pi = 0.131$$



# Pion Structure Function from $F_2^{LN}$

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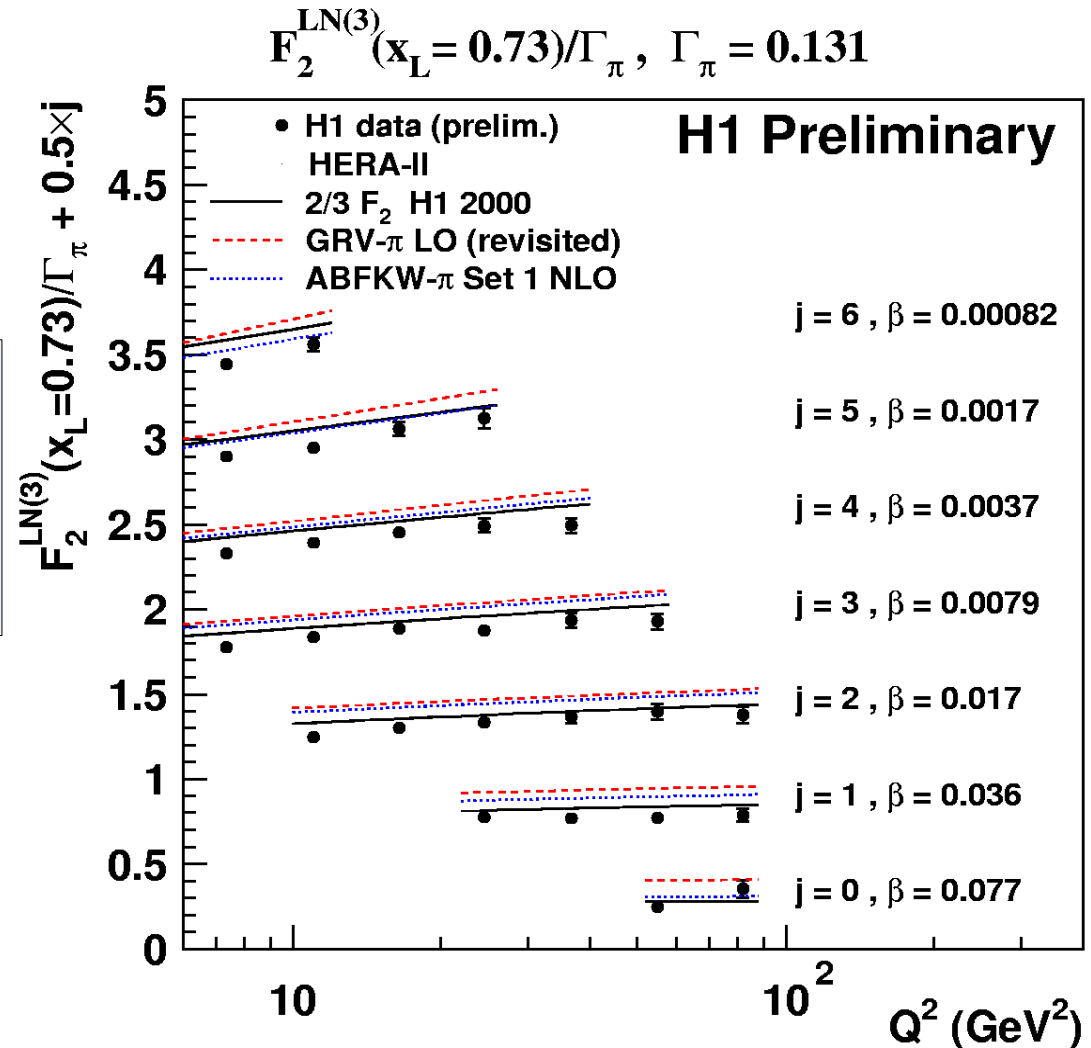
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## Summary

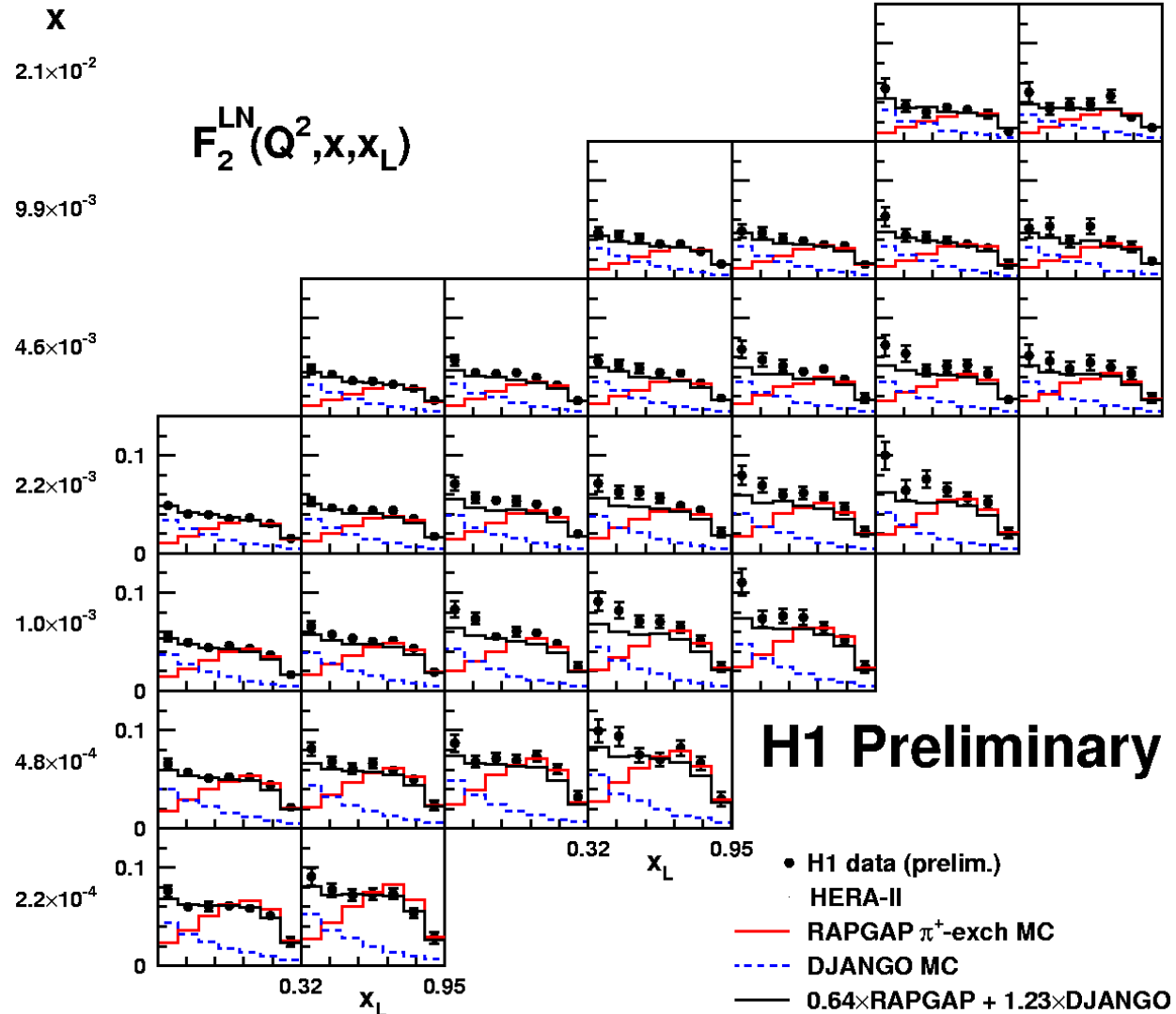
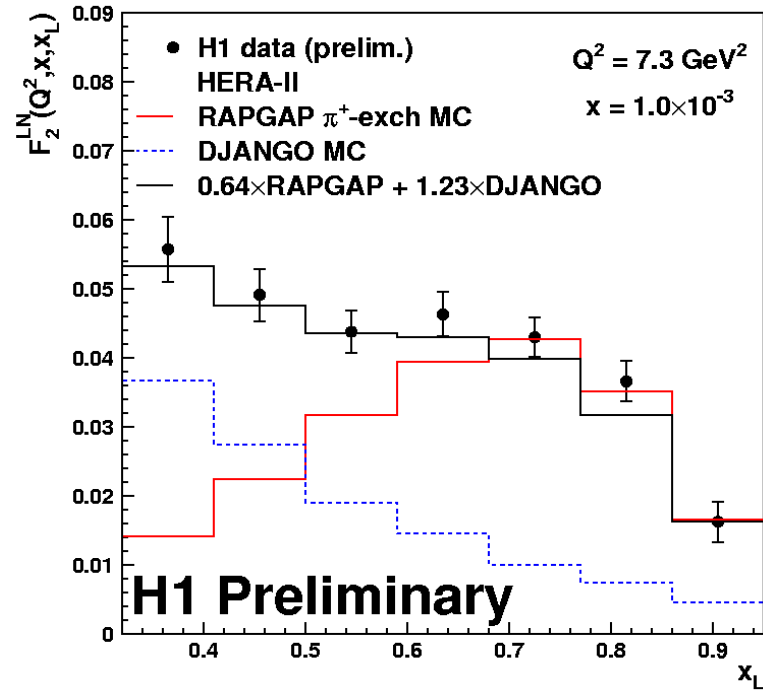
- **New measurement of  $F_2^{\text{LN}(3)}(Q^2, x, x_L)$  structure function from HERA-II data.**
  - ⇒  $F_2^{\text{LN}}/F_2$  ratio is consistent with vertex factorisation
  - ⇒  $F_2^{\text{LN}(3)}$  is interpreted in terms of pion structure function  $F_2^\pi$  and compared to different parameterizations of  $F_2^\pi$
- **Standard fragmentation models do not describe leading neutron production**
- **Models with virtual particle exchange describe data much better:  $x_L$  shape and magnitude are OK but  $p_T$  slopes are still off**
- **Account for absorption and  $(\rho, a_2)$  exchanges ⇒ promising agreement with data**

# Backup

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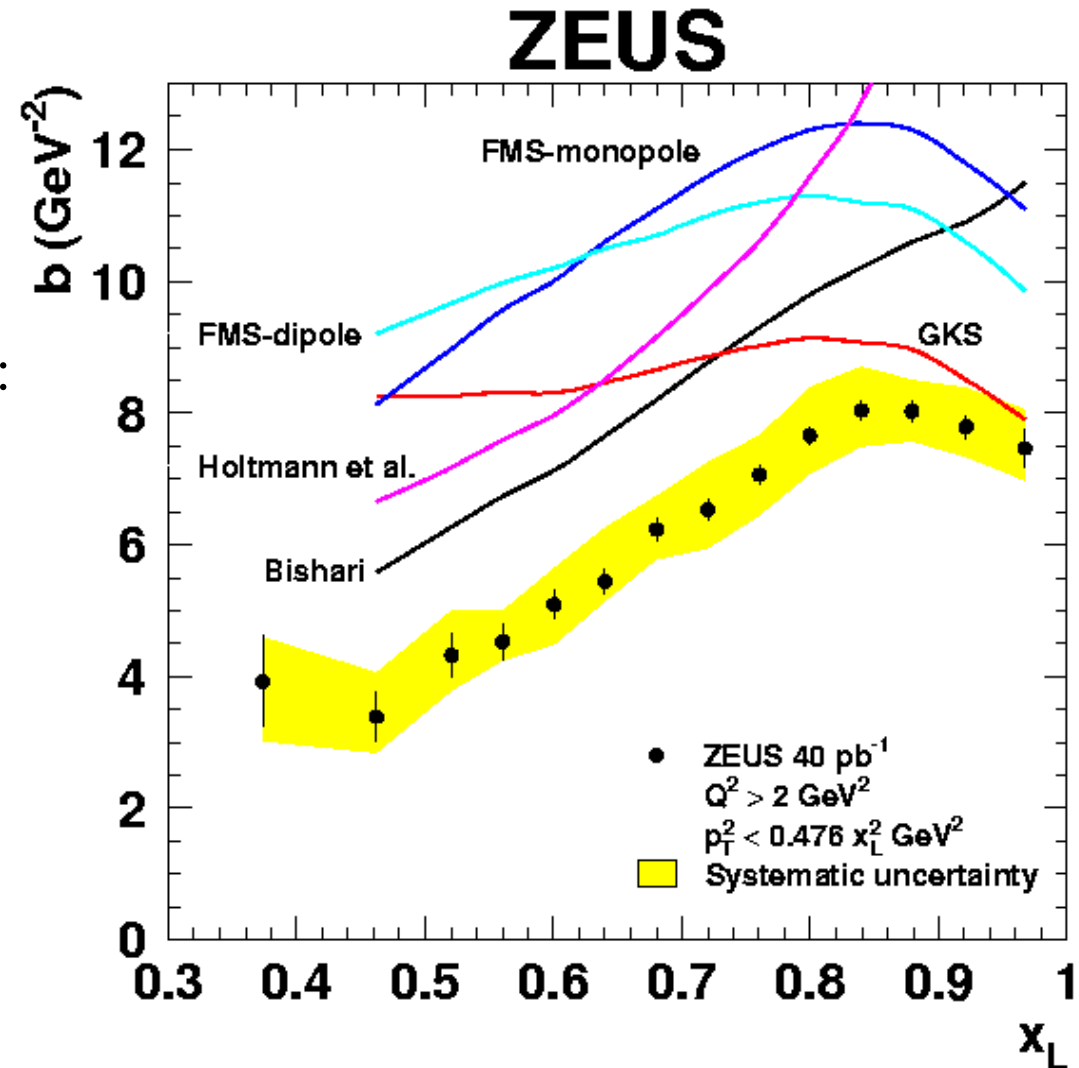
# Backup

$Q^2 = 7.3 \text{ GeV}^2$   $Q^2 = 11 \text{ GeV}^2$   $Q^2 = 16 \text{ GeV}^2$   $Q^2 = 24 \text{ GeV}^2$   $Q^2 = 37 \text{ GeV}^2$   $Q^2 = 55 \text{ GeV}^2$   $Q^2 = 82 \text{ GeV}^2$



# Model comparisons: DIS slopes

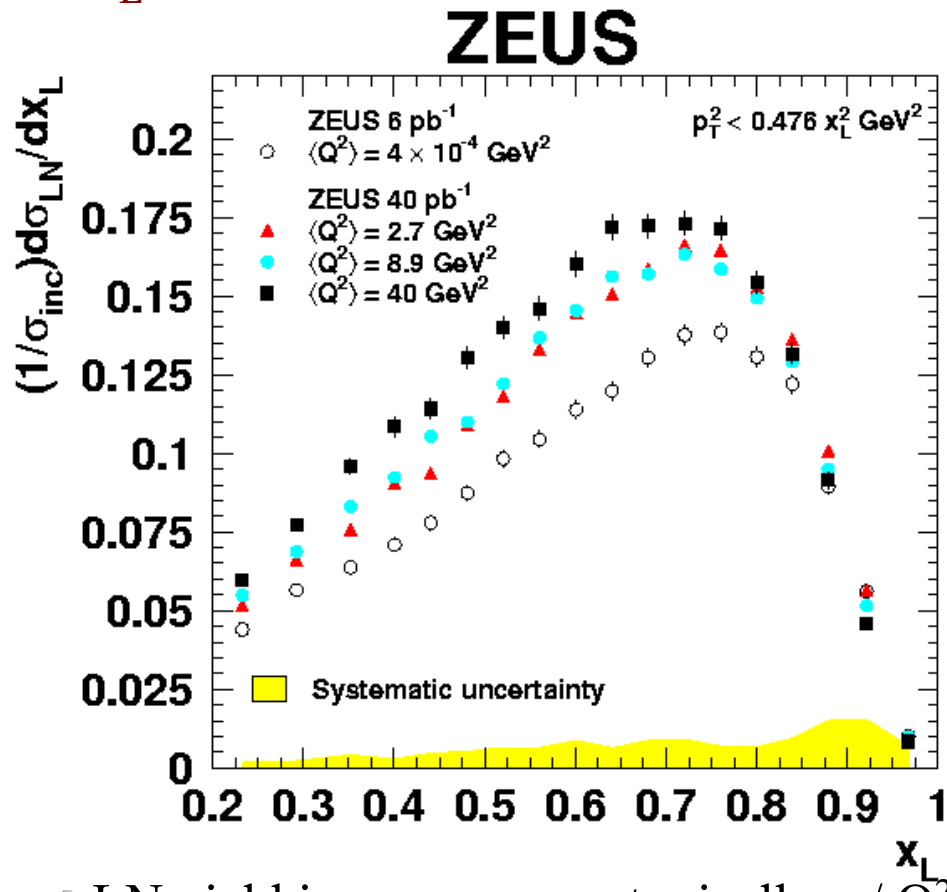
- Numerous parameterizations of pion flux  $f_{\pi/p}(x_L, p_T)$  in literature
- Here compare to measured DIS  $b(x_L)$ :
- Best agreeing models shown here; others wildly off
- All give too large  $b(x_L)$
- More refinement needed: absorption/migration



# $Q^2$ dependence of LN production

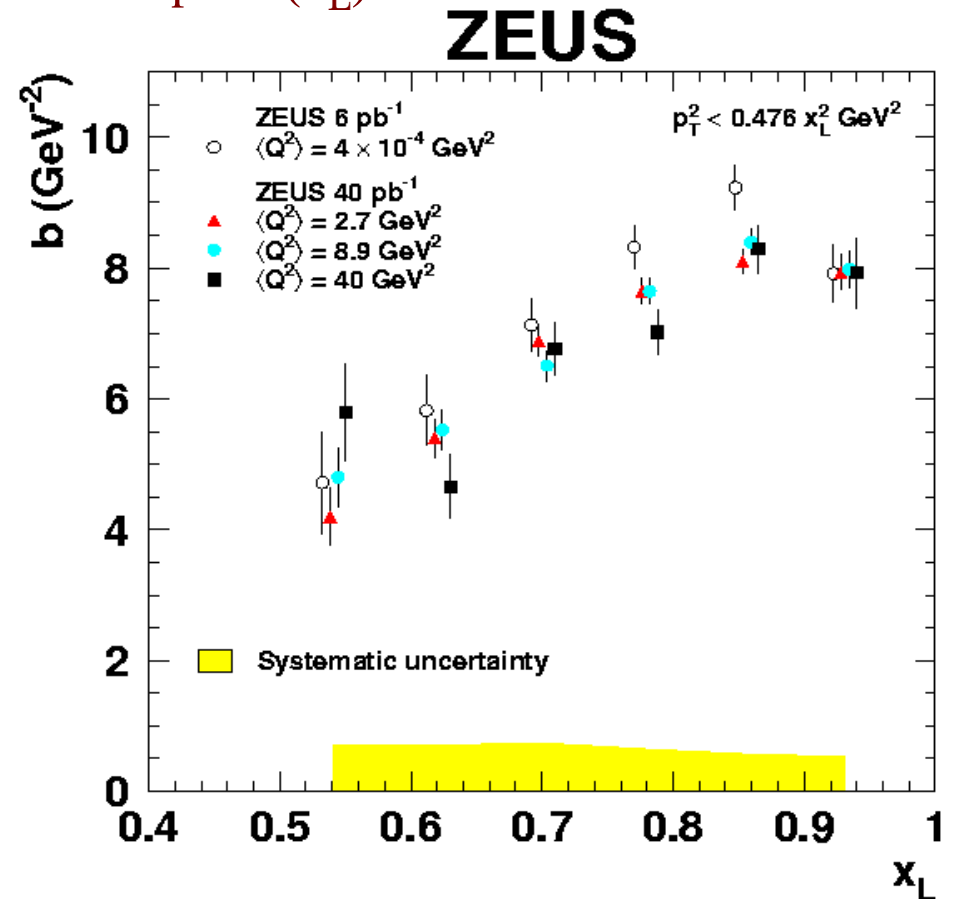
## 3 $Q^2$ bins in DIS + 1 $Q^2$ bin in $\gamma p$

- $x_L$  distributions:



- LN yield increases monotonically w/  $Q^2$
- Consistent with absorption:  
larger  $Q^2 \Rightarrow$  smaller  $\gamma^*$  less absorption

- slopes  $b(x_L)$  :



- Slopes for 3  $Q^2$  approx. same
- Slope for  $\gamma p$  significantly larger

