

# High $E_T$ Jet Physics at HERA

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On behalf of the H1 and ZEUS collaborations



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# Introduction: Jets, Algorithms...

## Definition of a generic hadronic jet:

*Group of particles which are 'close' to each other*

Partons in QCD calculations  
Final state hadrons in data and MC

Many definition of 'closeness'

### Cone algorithms

use geometrical information (no E)

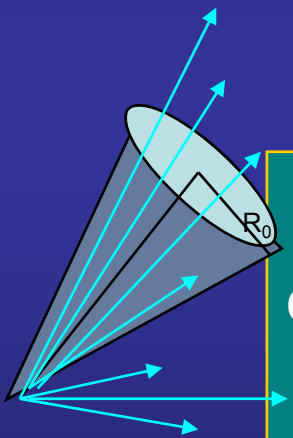
$$\Delta R = \sqrt{\Delta^2\eta + \Delta^2\phi} < R_0$$

### Clustering algorithms

use geometrical information + E

$$\text{e.g. } d_{ij} = \min(E_{T,i}^2, E_{T,j}^2)(\Delta^2\eta + \Delta^2\phi)/R_0^2 > \text{any } E_{T,k}^2$$

Theoretical problems with overlapping cones  
→ large uncertainty  
Not used at HERA since early days



# Clustering Algorithms

Group particles together if

**Mass**

$$d_{ij} = 2 E_i E_j (1 - \cos\theta_{ij}) < d_{\text{cut}}$$

**JADE** algorithm  
used extensively in  $e^+e^-$   
problems with ghost jets

**$k_{\perp}$**

$$d_{ij} = 2 \min(E_i^2, E_j^2) (1 - \cos\theta_{ij}) / E_0^2 < d_{\text{cut}} \quad E_0 \dots \text{hard scale}$$

**Durham** algorithm  
allows to vary resolution scale  $d_{\text{cut}}$   
→ subjets

**Longitudinally invariant  $k_{\perp}$**

$$d_{ij} = \min(E_{T,i}^2, E_{T,j}^2) (\Delta\eta^2 + \Delta\phi^2) / R_0^2 > \text{any } E_{T,k}^2 \quad R_0 \dots \text{radius, chosen } = 1$$

**Longitudinally Invariant  $k_{\perp}$**  algorithm  
combines features of cone and durham-

The remaining objects are called **Jets**



# Hadronisation corrections

**Needed** to compare data and QCD calculations

Hadronisation corrections applied to NLO

**Obtained** from LO+PS Monte Carlos

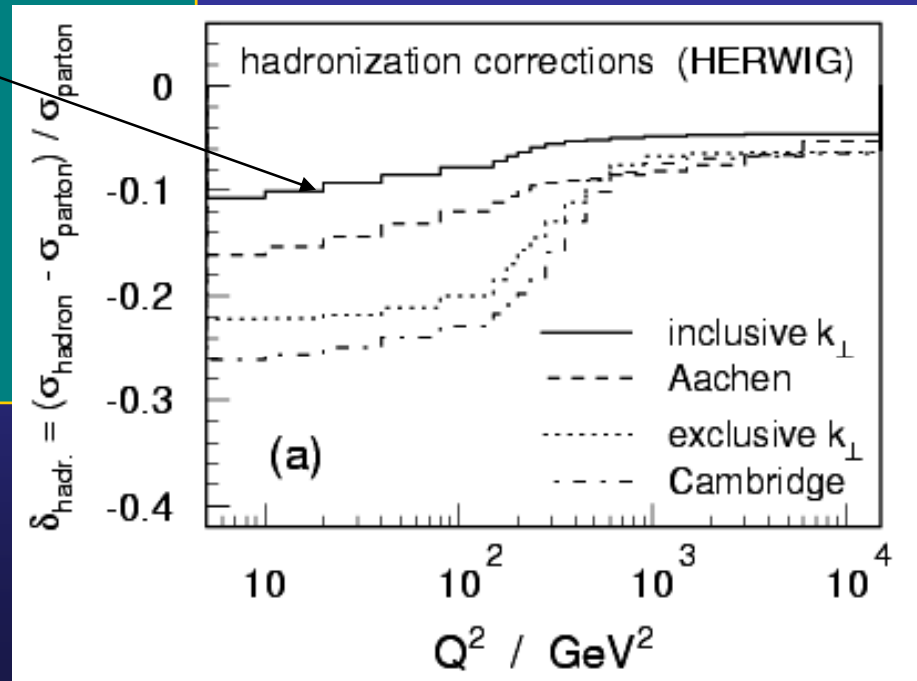
Describe jet production at HERA (apart from normalization)

**Longitudinally invariant  $k_T$  algorithm**

Smallest hadronisation correction  
(smallest uncertainty?)  
Preferred algorithm at HERA

**Durham (exclusive)  $k_T$  algorithm**

Allows to vary scale  
Only used for subjet studies



# Major experimental uncertainties

- Energy scale of scattered electron ( $\pm 1\%$ )
- Uncertainty in detection efficiency for scattered electron ( $\pm 2\%$ )
- Uncertainties in trigger efficiencies and event selections
- Uncertainties in correction for detector acceptances

## Error on Jet cross sections

- $< \pm 1\%$
- $\pm 2\%$
- $< \pm 3\%$
- $< \pm 3\%$

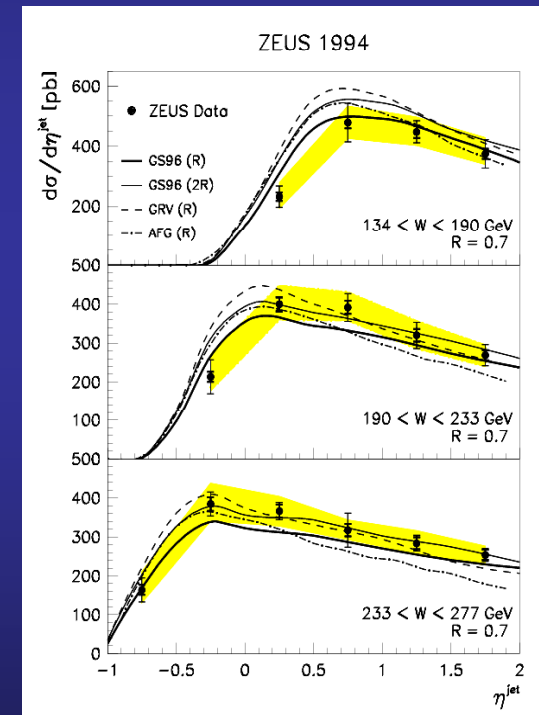
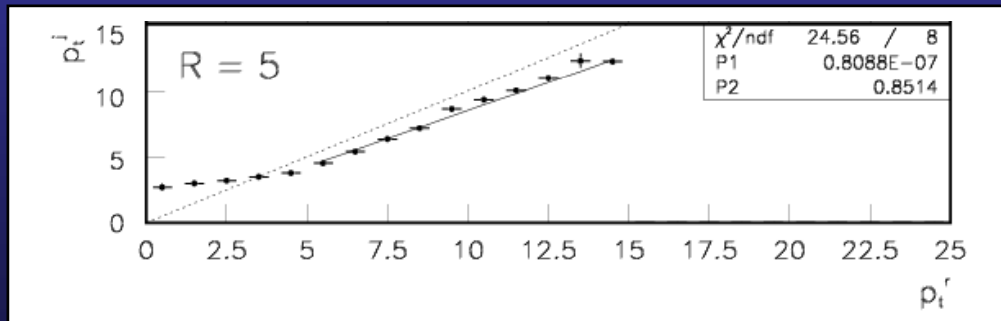
## Jet Energy Scale

$E_T(\text{jet})$  spectrum exponentially falling  $\rightarrow$  **Dominant** error  
 First ZEUS jet publications  $\pm 10\%$  error in cross section



Select Jets with **high(er)**  $E_T$   
 Dedicated **effort** to understand hadronic energy scale

Use of fact that  $\mathbf{p}_T$  (scattered electron) =  $\mathbf{p}_T$  (hadronic system)  
 Study events with single jets and small remaining hadronic energy



**Current Jet Energy Scale Uncertainty  $\pm 1\%$  (for  $E_T(\text{jet}) > 10$  GeV)**

**$\pm 3\%$**

$$d\sigma_{\text{jet}} = \sum_{a=q,\bar{q},g} \int dx \underbrace{f_a(x, \alpha_S, \mu_F)}_{\text{PDFs}} \underbrace{d\hat{\sigma}_a(x, \alpha_S(\mu_R), \mu_R, \mu_F)}_{\text{hard scattering cross section}} \underbrace{D(z, \mu_{F'}, \alpha_S)}_{\text{Fragmentation function}}$$

PDFs

hard scattering  
cross section

Fragmentation  
function

## Major theoretical uncertainties

### Hadronisation corrections

Typical uncertainty  $\pm 1\%$

Corrections in general small  $<10\%$

Uncertainty taken as difference between estimations based on different MCs  
(unsatisfactory)

### Proton PDFs

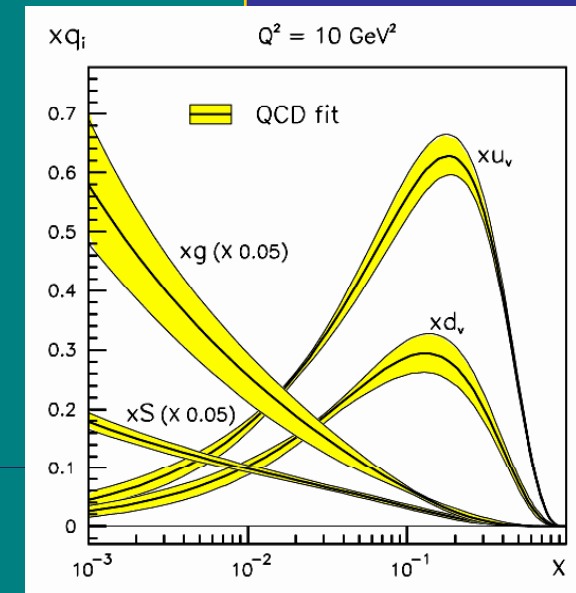
Typical uncertainty  $\pm 1-2\%$

Traditionally taken from difference obtained with  
various sets of PDFs (unsatisfactory)

Covariance matrix  $V_{p_\mu, p_\lambda}$  of the fitted parameters  $\{p_\lambda\}$  available  
correct evaluation of error on cross section

$$(\Delta\sigma_{\text{jet}})^2 = \sum_{\lambda, \mu} \frac{\partial\sigma_{\text{jet}}}{\partial p_\mu} V_{p_\mu, p_\lambda} \frac{\partial\sigma_{\text{jet}}}{\partial p_\lambda}$$

In  $\gamma P$ : additional uncertainty due to PDFs of photon



$$d\sigma_{\text{jet}} = \sum_{a=q,\bar{q},g} \int dx \underbrace{f_a(x, \alpha_S, \mu_F)}_{\text{PDFs}} \underbrace{d\hat{\sigma}_a(x, \alpha_S(\mu_R), \mu_R, \mu_F)}_{\text{hard scattering cross section}} \underbrace{D(z, \mu_{F'}, \alpha_S)}_{\text{Fragmentation function}}$$

## Strong coupling $\alpha_s$

Current uncertainty  $\pm 1.7\%$

Enters PDFs value assumed to evolve to different scales in fit to inclusive DIS data  
 Enters  $d\hat{\sigma}_a$  governs strength of interaction

Current world average value  $\alpha_s(M_Z) = 0.1176 \pm 0.0020$

Used consistently in PDFs and  $d\hat{\sigma}_a$

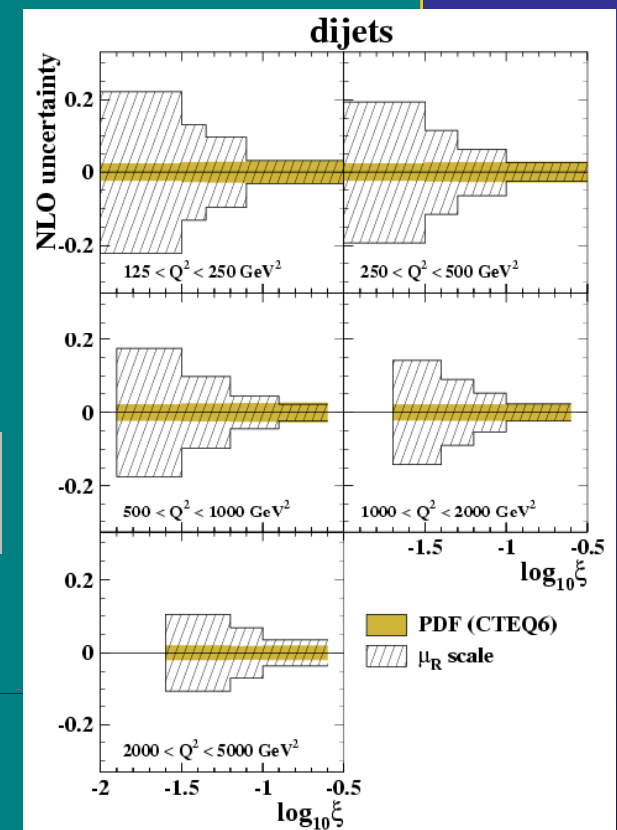
Only free parameter of pQCD: jet cross sections sensitive to it

Uncertainty can be large  
 Dominating theoretical error

## Terms beyond $\alpha_s^2$

Corresponding uncertainty NOT known  
 Estimated through residual dependence on renormalization  $\mu_R$   
 and factorization  $\mu_F$  scales

→ Choice of scales:  $Q$ ,  $E_{T,\text{jet}}$  or linear combination  
 Customary, but arbitrary to vary scales by factor 2



# Jet Production Processes at HERA

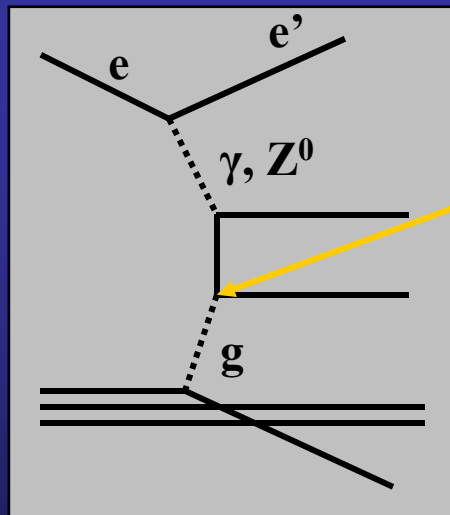
## Event classes

Photoproduction:  $Q^2 \sim 0$  (real photons)

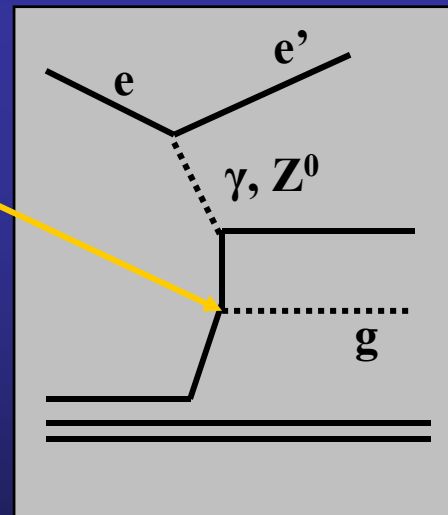
Deep inelastic scattering:  $Q^2 > \text{few GeV}^2$  (virtual photons)

## Jet production mechanisms (LO in $\alpha_S$ )

Boson-Gluon Fusion



QCD Compton Scattering



$\alpha_S$

Di-jets

Multi-jet

Higher order processes ( $\alpha_S^n, n > 1$ )  
Multi-parton interactions ( $\rightarrow$  L. Stanco)



# Inclusive-Jet Cross Sections in DIS

## Data sample

$$Q^2 > 125 \text{ GeV}^2$$

$$L = 82 \text{ pb}^{-1} \text{ of } e^\pm p \text{ collisions}$$

## Jet reconstruction

With  $k_T$  algorithm in the longitudinally invariant inclusive mode  
And in the Breit frame

## Jet selection

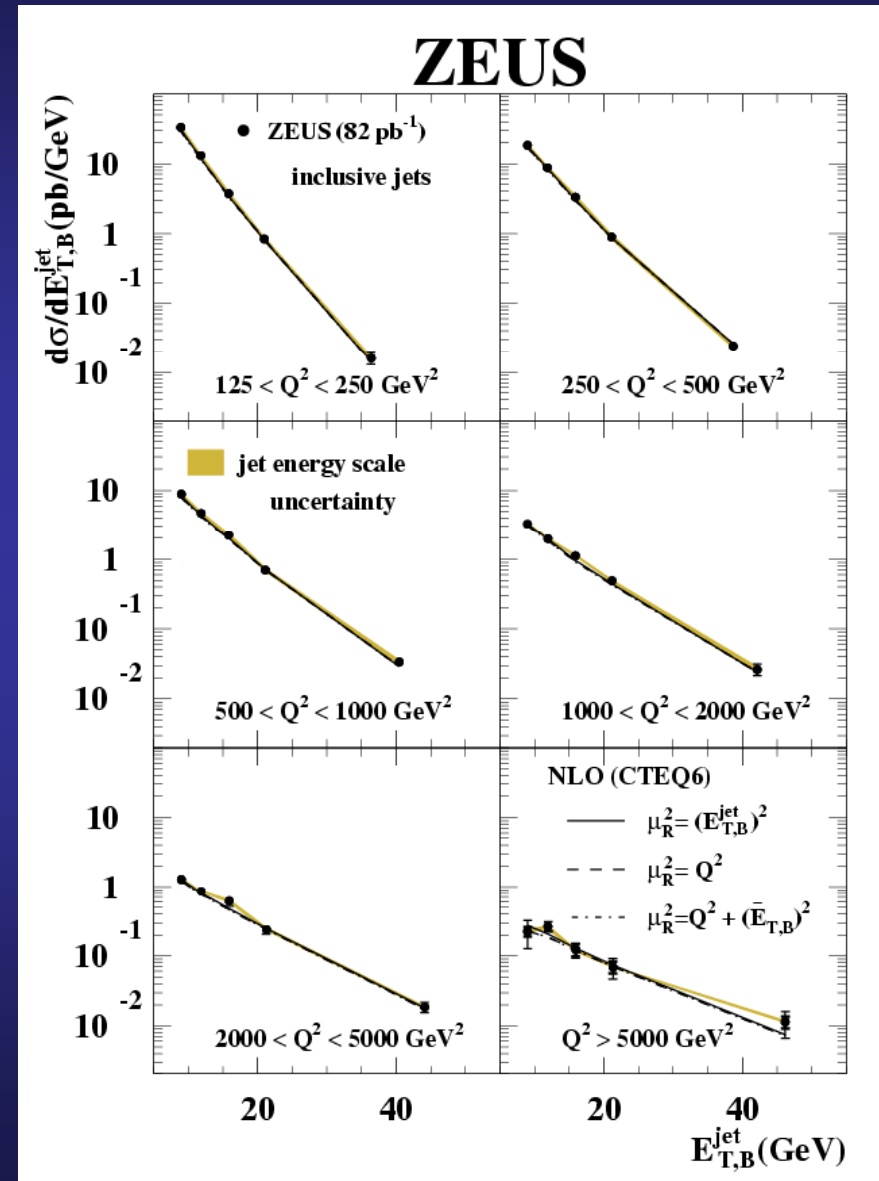
$$E_T^{\text{Jet}} (\text{Breit}) > 8.0 \text{ GeV}$$

$$E_T^{\text{Jet}} (\text{lab}) > 2.5 \text{ GeV}$$

$$-2 < \eta^{\text{Jet}} (\text{Breit}) < 1.5$$

## Results

$d\sigma/dE_T^{\text{jet}}(\text{Breit})$  in large range of  $Q^2$   
Nice description by NLO QCD

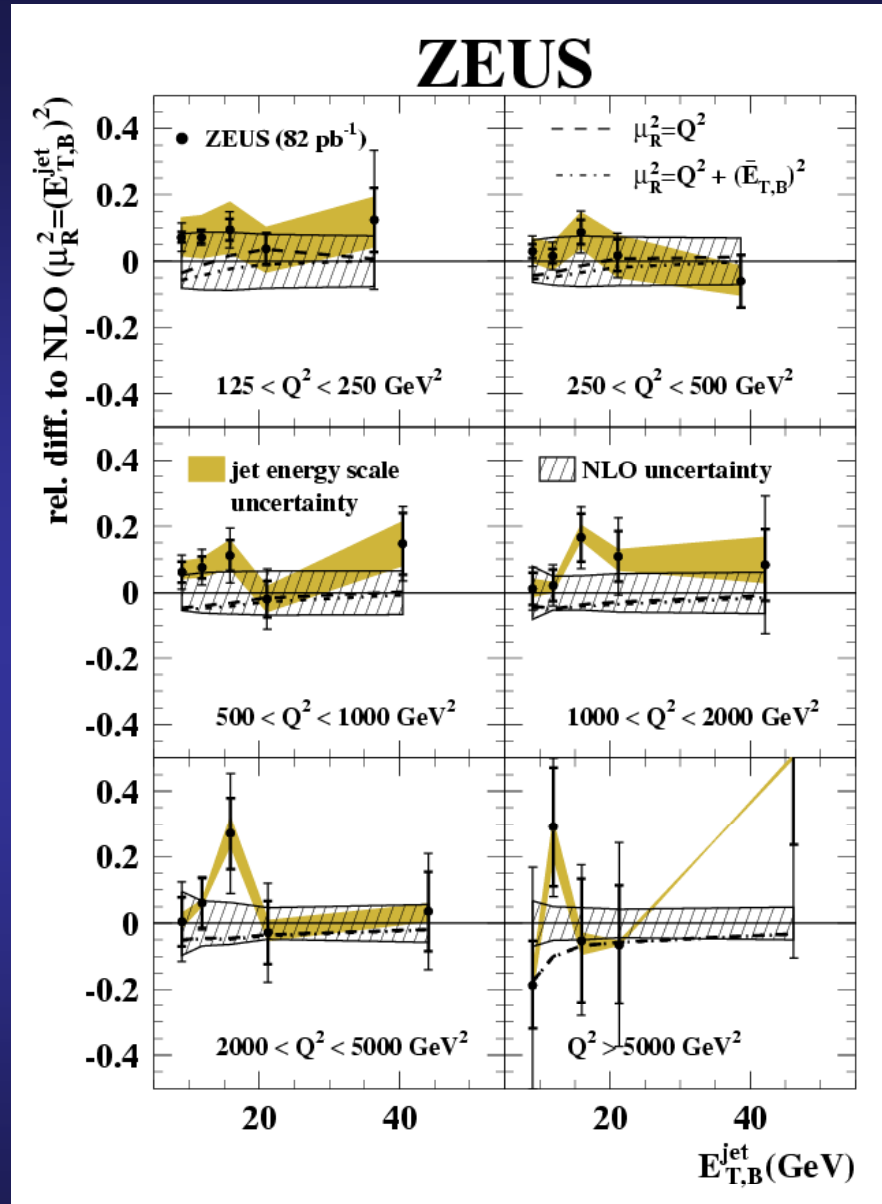


## Ratio to NLO QCD

Dominant experimental uncertainty: energy scale

Theoretical uncertainty  $\sim$  experimental uncertainty

With HERA II data statistical uncertainty at high  $Q^2$  will be significantly reduced



## Determination of strong coupling constant $\alpha_s(M_Z)$

$Q^2 > 500 \text{ GeV}^2$

Smaller experimental (E scale) uncertainties  
Smaller theoretical (PDFs and scale) uncertainties

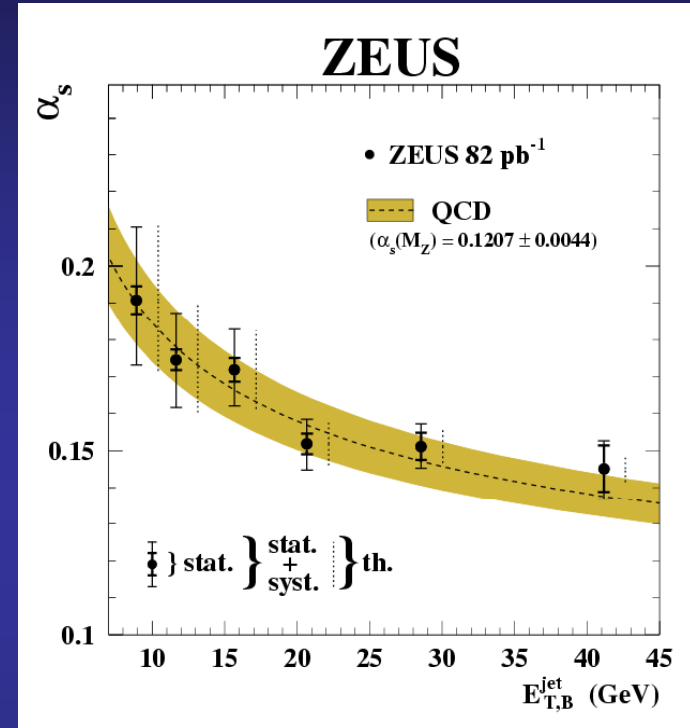
Parameterize theoretical cross section as

$$d\sigma/dQ^2(\alpha_s(M_Z)) = C_1\alpha_s(M_Z) + C_2\alpha_s^2(M_Z)$$

(same value of  $\alpha_s$  in PDF and calculation)

Determine  $C_1$  and  $C_2$  from  $\chi^2$  fits

Fit parameterized theoretical cross  $d\sigma/dQ^2(\alpha_s(M_Z))$   
to measurement



## Result

$$\alpha_s(M_Z) = 0.1207 \pm 0.0014 \text{ (stat.) } \begin{matrix} +0.0035 \\ -0.0033 \end{matrix} \text{ (syst.) } \begin{matrix} +0.0022 \\ -0.0023 \end{matrix} \text{ (theo.)}$$

One of world's most precise...

PDG:  $\alpha_s(M_Z) = 1.176 \pm 0.002$

# Multi-Jets in DIS

## Data sample

$150 \text{ GeV}^2 < Q^2 < 15,000 \text{ GeV}^2$   
 $L = 65 \text{ pb}^{-1}$  of  $e^+p$  collisions

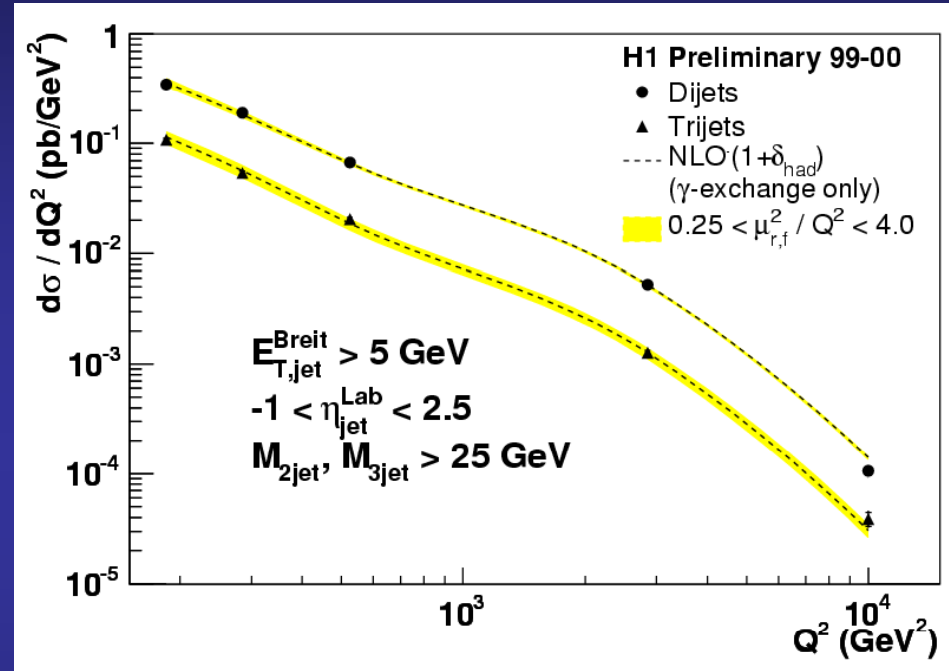
## Jet reconstruction

With  $k_T$  algorithm in the  
 longitudinally invariant inclusive mode  
 And in the Breit frame

## Jet selection

$E_T^{\text{Jet}} (\text{Breit}) > 5.0 \text{ GeV}$   
 $-1 < \eta^{\text{Jet}} (\text{lab}) < 2.5$   
 Dijets:  $m_{2\text{jet}} > 25 \text{ GeV}$   
 Trijets:  $m_{3\text{jet}} > 25 \text{ GeV}$

To reduce infrared regions ( $E_T^1 \sim E_T^2$ )



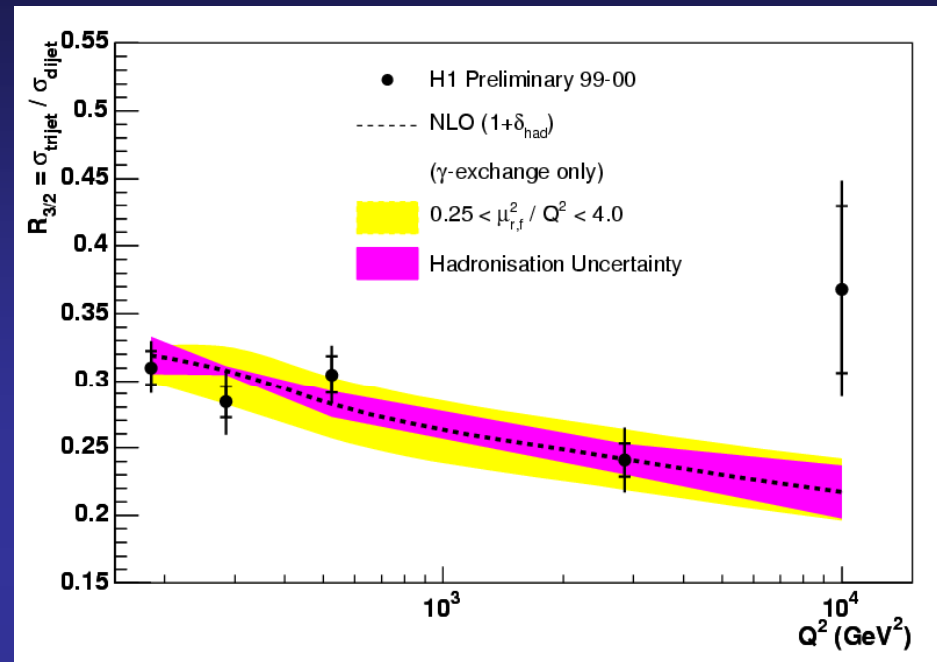
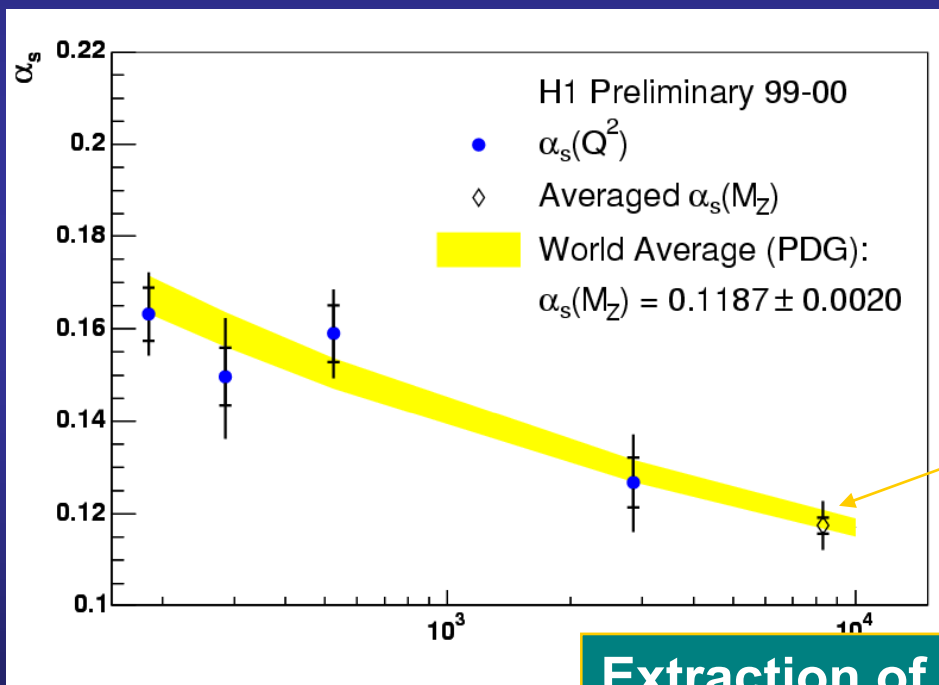
## Results

$d\sigma/dQ^2$  in large range of  $Q^2$   
 Nice description by NLO QCD

## Ratio of 3-jet/2-jet

Many experimental uncertainties cancel

Theoretical uncertainty ~ experimental uncertainty



Average  $\alpha_s(M_Z)$

## Extraction of $\alpha_s(M_Z)$

Similar procedure as with incl. jets

$$\alpha_s(M_Z) = 0.1175 \pm 0.0017 \text{ (stat.)} \pm 0.005 \text{ (syst.)} \begin{matrix} +0.0054 \\ -0.0068 \end{matrix} \text{ (theo.)}$$

# Di-Jets in Photoproduction

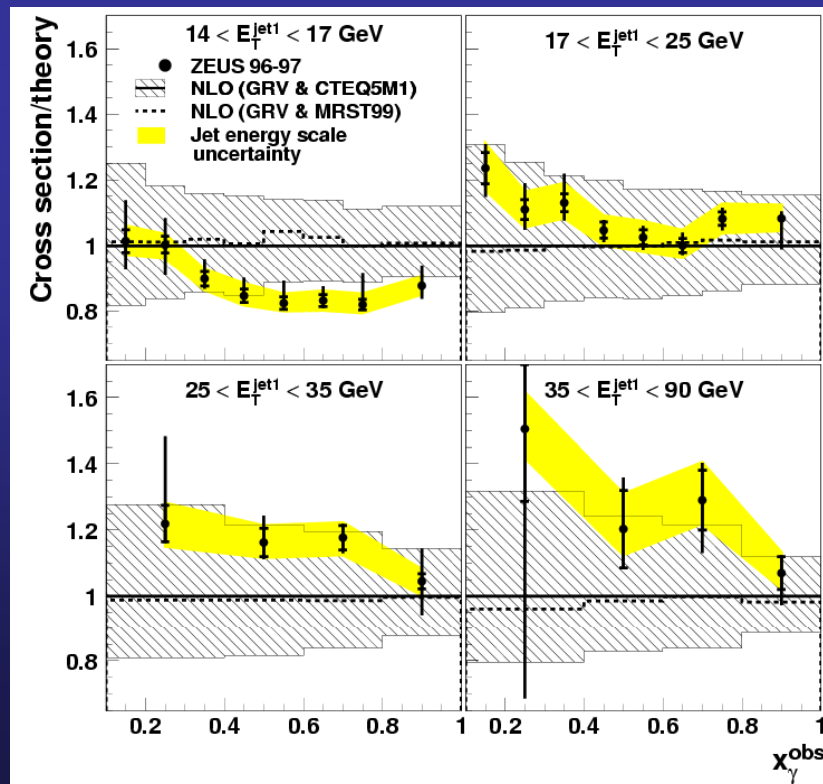
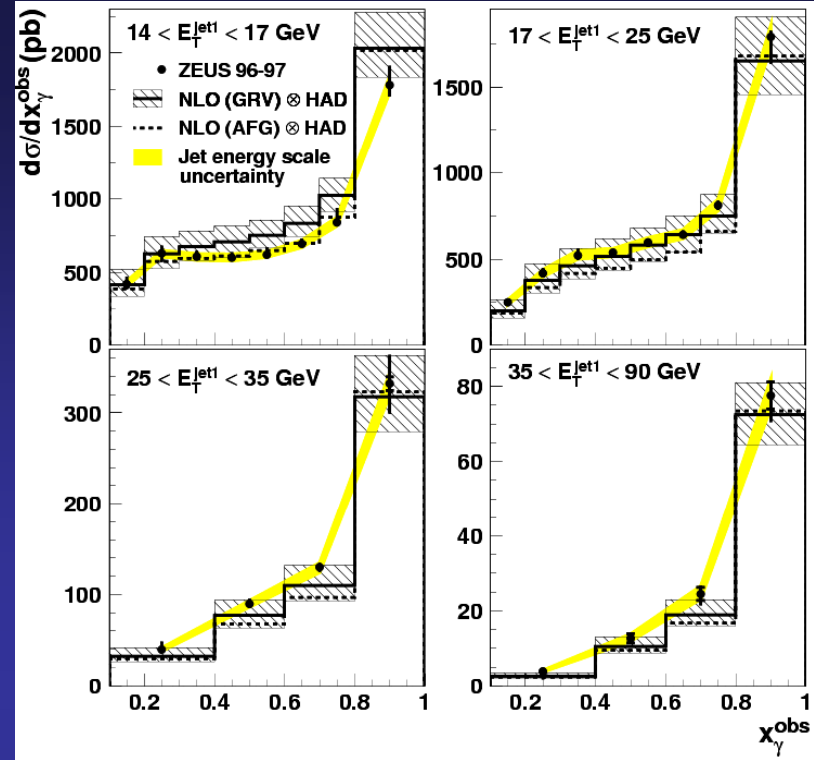
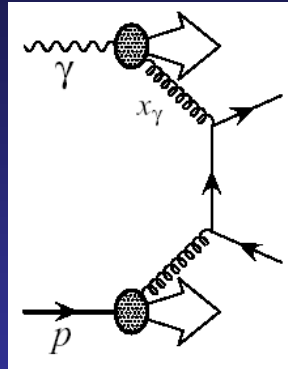
## Data sample

$$Q^2 < 1 \text{ GeV}^2$$

$$134 < W < 277 \text{ GeV}$$

## Jet reconstruction

With  $k_T$  algorithm in the longitudinally invariant inclusive mode



## Jet selection

$$E_T^{\text{Jet}} > 14.0 \text{ and } 11.0 \text{ GeV}$$

$$-1 < \eta^{\text{Jet}} (\text{lab}) < 2.4$$

$d\sigma/dx_\gamma$  in bins of  $E_T^{\text{jet}}$

Large discrepancy with theory



Photon PDF inadequate?  
NLO pQCD calculation inadequate?

Photon PDFs determined from  $\gamma\gamma$  interactions at low scales

At HERA photon PDFs being probed at high scales (jet  $E_T$ )

Similar study of H1 shows 'perfect' agreement with NLO pQCD



H1 uses slightly higher  $E_{T,2}$  cut at 15 GeV

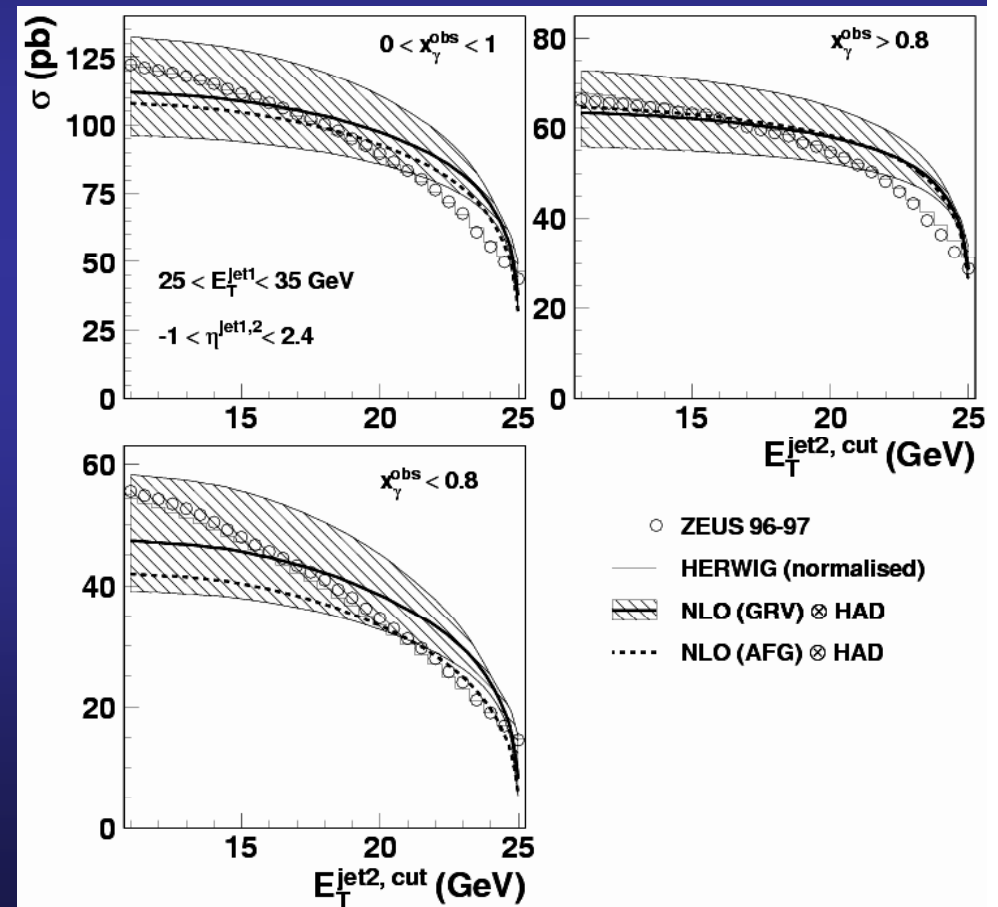
Study of dependence on  $E_{T,2}$  cut

Dependence **NOT** reproduced by NLO

H1 cut more fortunate

**NNLO calculations needed**

**Until then, no meaningful constraint on photon PDFs**

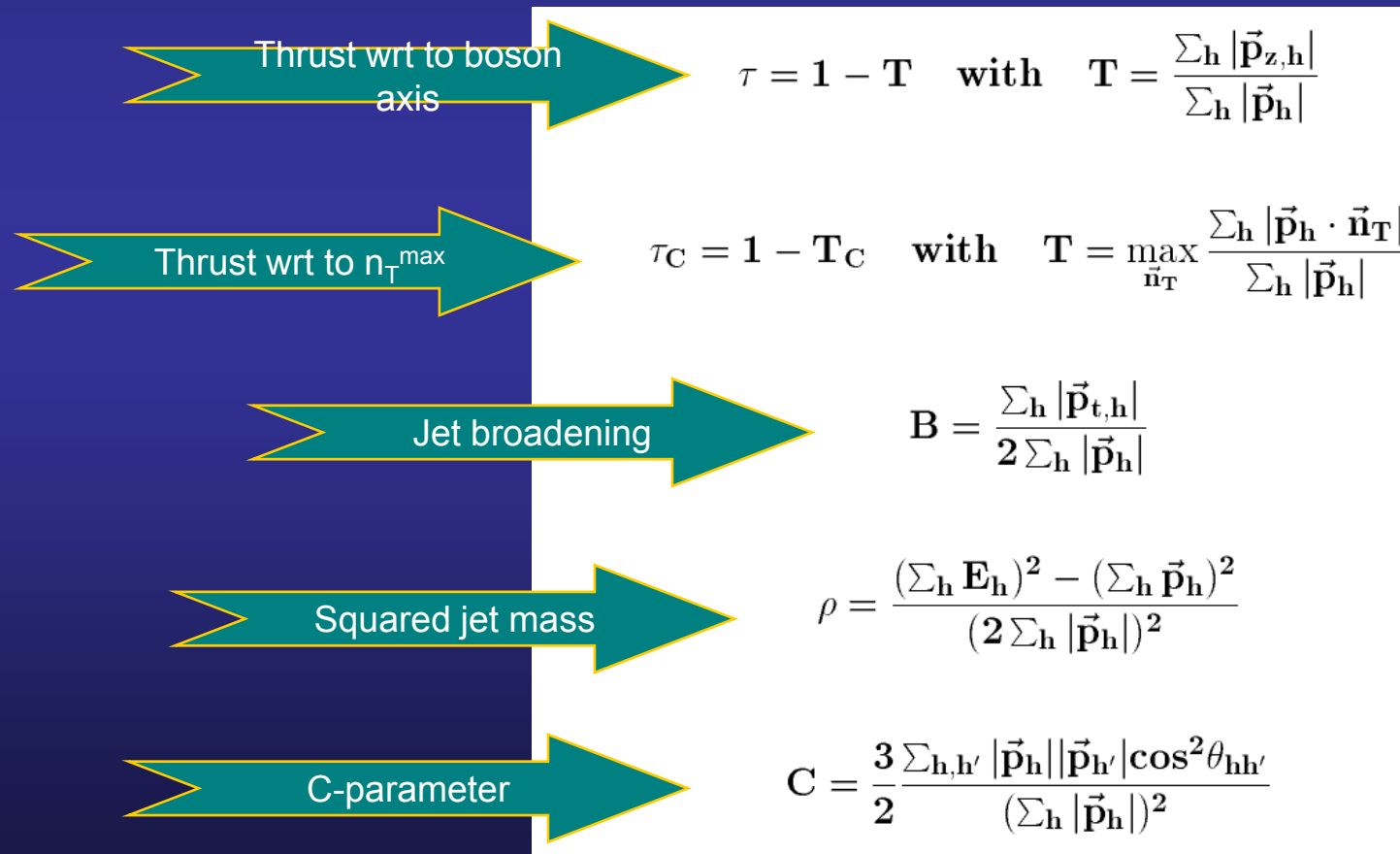


# Event Shapes

## Data sample

$196 < Q^2 < 40,000 \text{ GeV}^2$   
 $L = 106 \text{ pb}^{-1}$  of  $e^\pm p$  data

## Event shape of hadronic final state





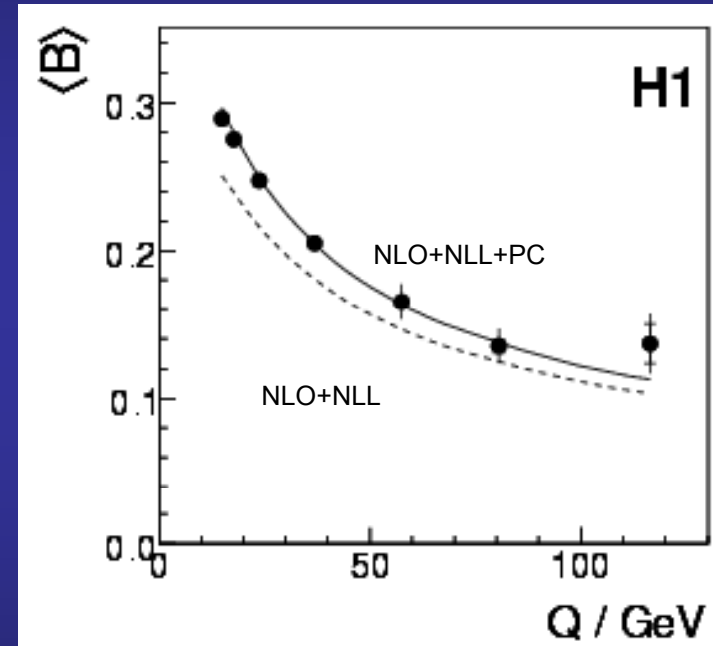
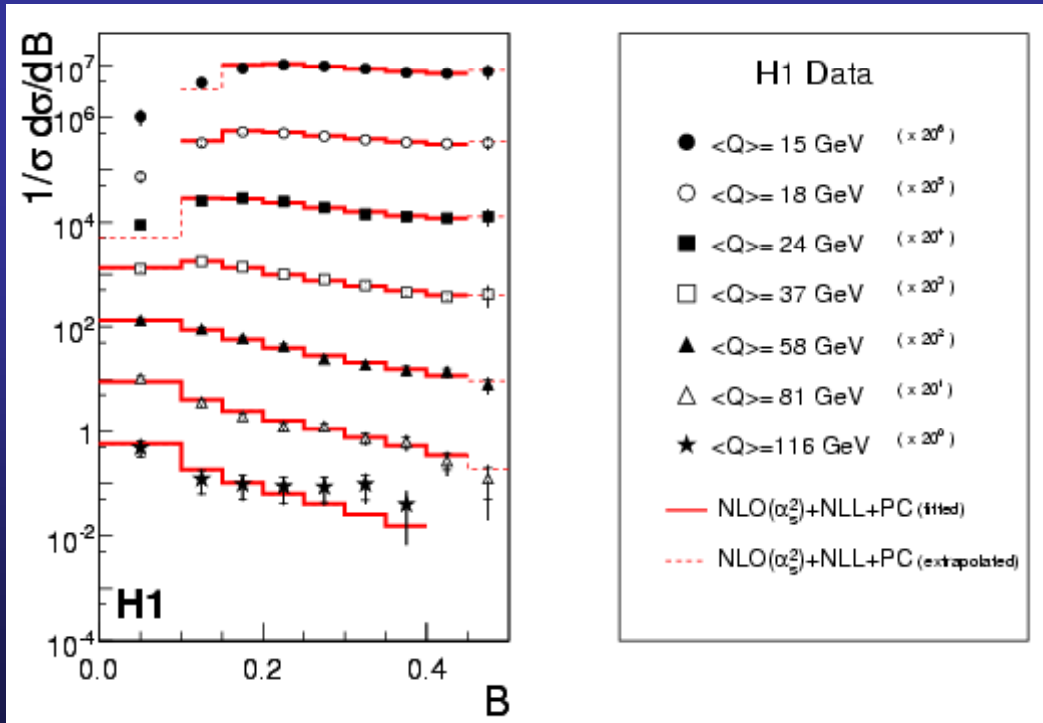
# Theoretical calculations

Available at NLO level including resummed next-to-leading logarithms (NLO+NLL)  
 Hadronisation corrections taken care of by calculable **power corrections** ( $\sim 1/Q$ )

$$P_V \sim \alpha_0 P_V^{\text{calc}} \quad \text{where } \alpha_0 \text{ is a universal parameter (independent of ES variable } V)$$

## Differential distribution

Nicely reproduced by NLO+NLL+PC

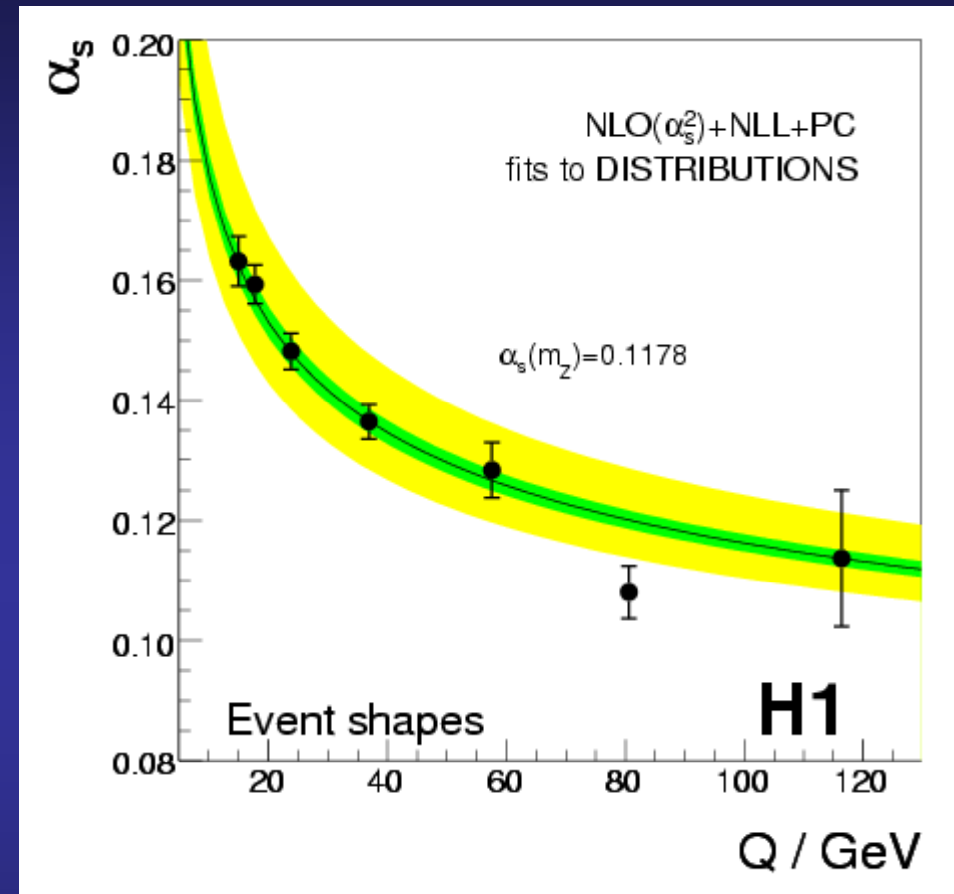
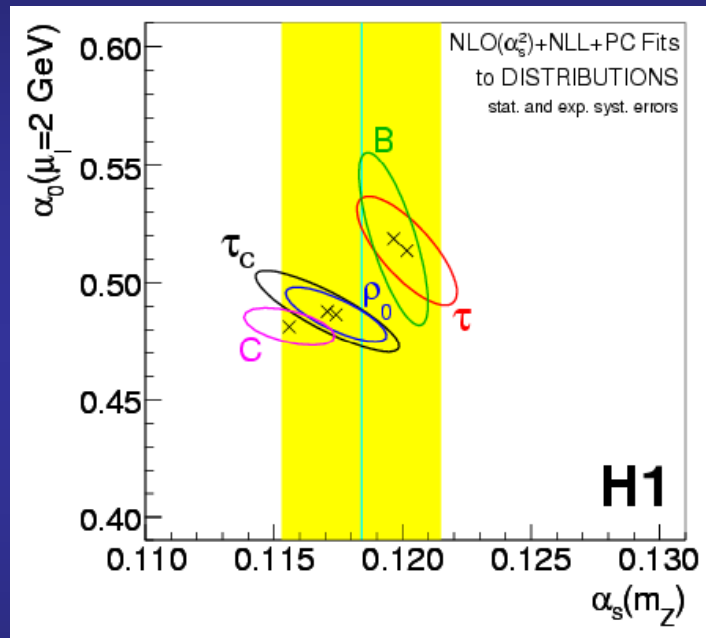


## Mean of ES Variable versus Q

Nicely reproduced by NLO+NLL+PC  
 (2 parameter fit:  $\alpha_s$  and  $\alpha_0$ )

## Fit to differential distributions

→ ~Consistent values of  $\alpha_s$  and  $\alpha_0$   
 ( $\alpha_0$  indeed universal within 10%)

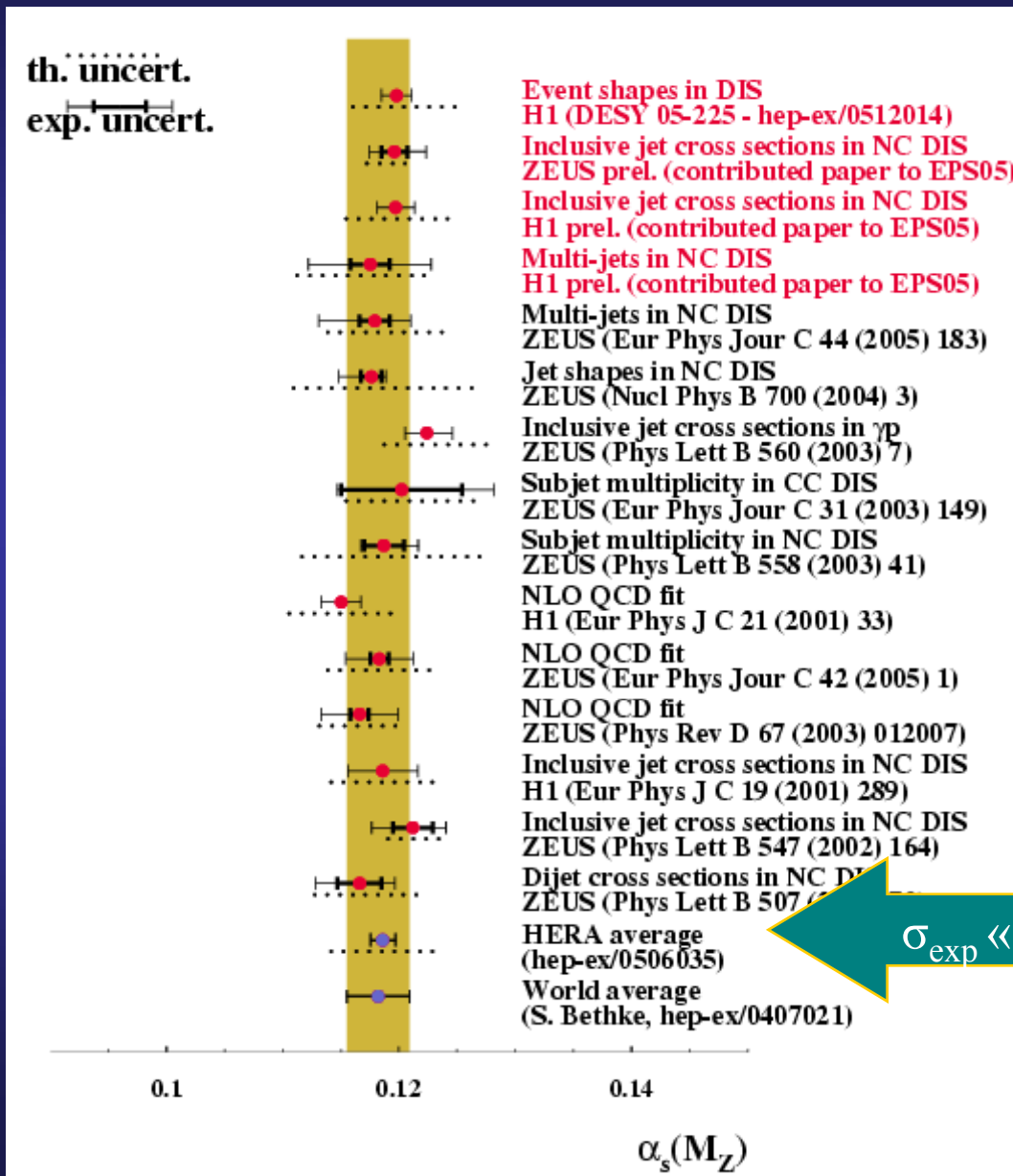


## Combined fit over all ES Variables

$$\rightarrow \alpha_s(M_Z) = 0.1198 \pm 0.0013 \text{ (exp)} \quad \begin{matrix} +0.0056 \\ -0.0043 \end{matrix} \text{ (theo)}$$

$$\rightarrow \alpha_0 = 0.476 \pm 0.008 \text{ (exp)} \quad \begin{matrix} +0.0018 \\ -0.0059 \end{matrix} \text{ (theo)}$$

# HERA Measurements of $\alpha_S$



See

C.Glassman  
hep-ex/0506035

$\sigma_{\text{exp}} \ll \sigma_{\text{theo}}$

# Conclusions

## Precision jet physics at HERA

Experimental uncertainties often  $< 3\%$   
Uncertainties dominated by jet energy scale uncertainty

## Performed large number of measurements

Photoproduction (inclusive jets, dijets, multijets...)  
DIS NC (inclusive, dijets, multijets, subjets...)  
DIS CC (inclusive...)

## Results provide

Constraints on proton PDF (included in NLO QCD fits of  $F_2$ )  
Constraints on photon PDF (needs better calculations)  
High precision measurements of the strong coupling constant  $\alpha_s(M_Z)$

$$\alpha_s(M_Z)_{\text{HERA}} = 0.1193 \pm 0.0005 (\text{exp}) \pm 0.0025 (\text{theo})$$

**Backup Slides**

# Frames for Jet Finding

## $e^+e^-$ annihilation

Laboratory frame = center-of-mass frame  
(unless there is significant initial state radiation)

Events  $p_T$  balanced  
Cone algorithm invariant to longitudinal boost  
→ Analysis in **laboratory** frame

## $p\bar{p}$ colliders

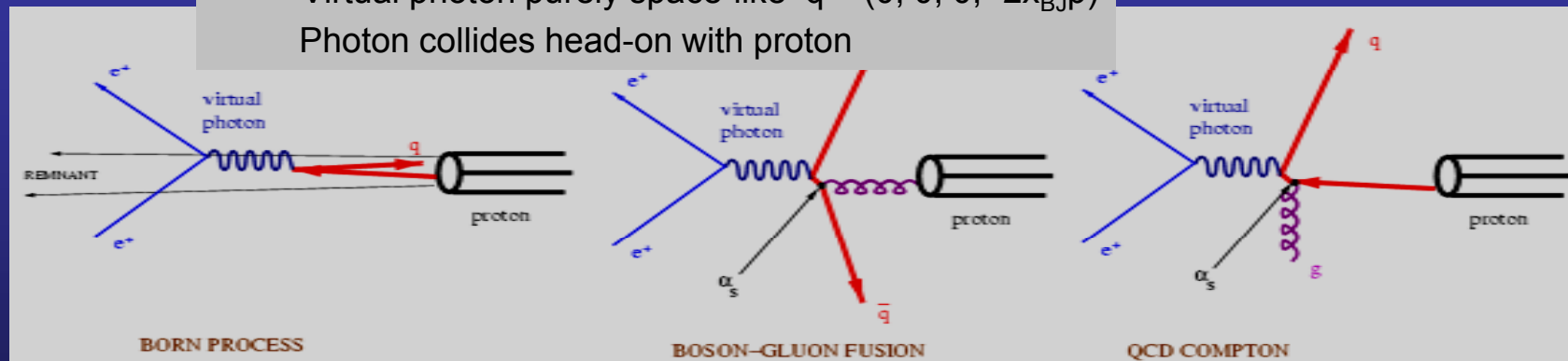
## Deep Inelastic Scattering

Hadronic system recoils against scattered electron  
→ jets have  $p_T$  in laboratory frame

### BREIT FRAME

$$2x_{BJ} + \vec{p} = 0$$

Virtual photon purely space-like  $\vec{q} = (0, 0, 0, -2x_{BJ}p)$   
Photon collides head-on with proton

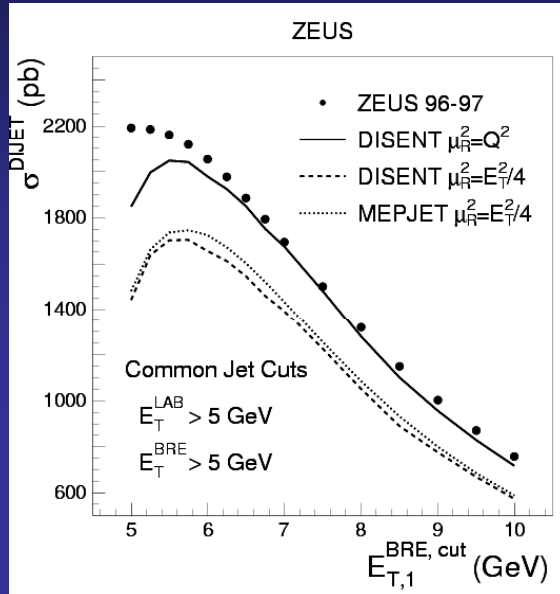


### High- $E_T$ Jet Production in Breit Frame



suppression of Born contribution  
suppression of Beam remnant jet(s)

# $E_T$ – Cuts for Dijet Selection



## Symmetric cuts

For instance  $E_{T,1}, E_{T,2} > 5 \text{ GeV}$

**Reduced phase space** for real emission of soft gluons close to cut

**Complete** cancellation of the soft and collinear singularities with the corresponding singularities in the virtual contributions **not possible**

Unphysical behavior of calculated cross section

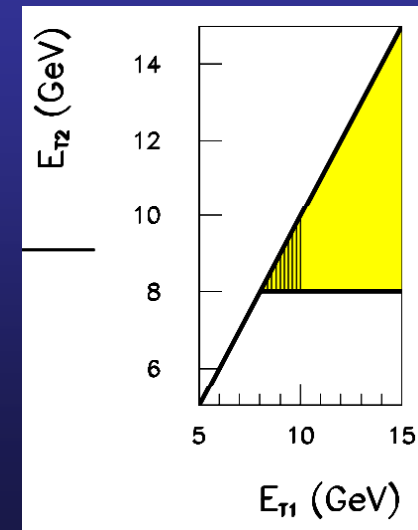
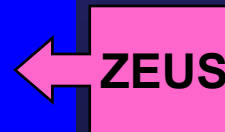
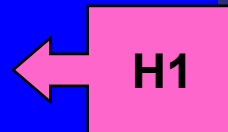
## Solutions

Additional cut on **Sum of  $E_T$**

for instance  $E_{T,1} + E_{T,2} > 13 \text{ GeV}$

**Asymmetric cuts**

for instance  $E_{T,1} > 8 \text{ GeV}, E_{T,2} > 5 \text{ GeV}$



# Jets in Photo-production

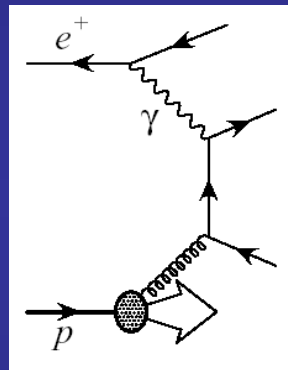
Photon has very low virtuality  $Q^2 \sim 0 \text{ GeV}^2$

Only one inclusive variable  $W_{\gamma p}$  ... photon – proton center of mass

To  $\mathcal{O}(\alpha_s)$   $\rightarrow$  2 types of processes contribute to jet production

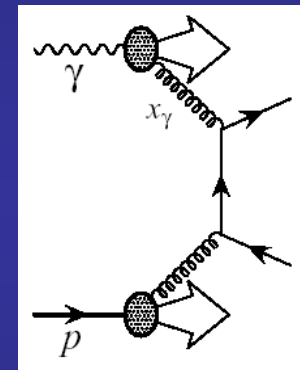
## Direct photo-production

Photon interacts as an entity



## Resolved photo-production

A parton with momentum fraction  $x_\gamma$  in the photon enters the hard scattering process



## Momentum fraction $x_\gamma$

Can be reconstructed as

$$x_\gamma = \frac{\sum_{\text{jets}} (\mathbf{E} - \mathbf{p}_z)}{\sum_{\text{all hadrons}} (\mathbf{E} - \mathbf{p}_z)}$$

Direct  $x_\gamma = 1$

Resolved  $x_\gamma < 1$

