



Results on Inclusive Diffraction From The ZEUS Experiment by the M_X -Method

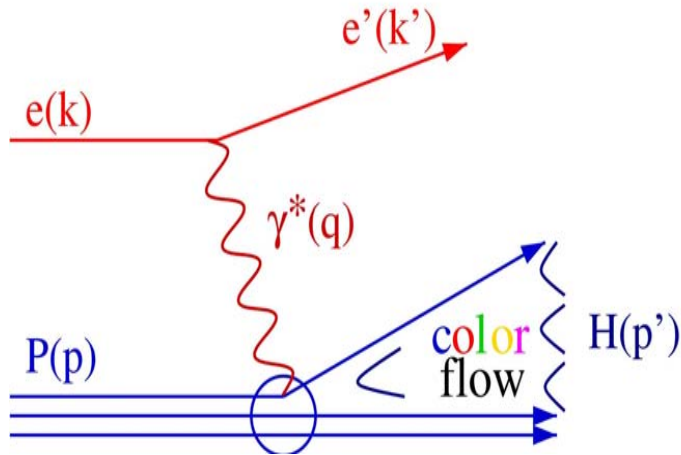


Presented by B.Löhr on behalf of ZEUS

1. DIS and diffractive DIS kinematics
2. The M_X -Method
3. $d\sigma^{\text{diff}}/dM_X$
4. Diffractive structure function $x_{\text{IP}}F_2^{\text{D}(3)}$
5. The BEKW-fit
6. The Q^2 -dependence of $x_{\text{IP}}F_2^{\text{D}(3)}$
7. Summary

(Data are from the running period 1999-2000 (prelim.) and 1998-1999 (publ.))

Inclusive DIS events :



$$s = (k+p)^2$$

center of mass energy squared

$$Q^2 = -q^2 = -(k-k')^2$$

virtuality, size of the probe

$$W^2 = M_H^2 = (p+q)^2$$

γ^* - proton cms energy squared

$$x = \frac{Q^2}{2p \cdot q}$$

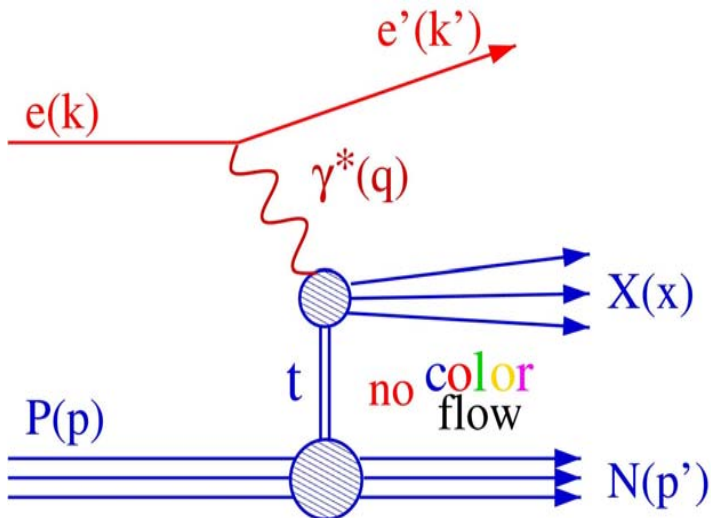
$$y = \frac{p \cdot q}{p \cdot k}$$

x: fraction of the proton carried by the struck parton

y: inelasticity, fraction of the electron momentum carried by the virtual photon

$$Q^2 = x \cdot y \cdot s$$

Diffraction DIS events :



For diffractive events in addition 2 variables

$$M_x$$

mass of the diffractive system X

$$t = (p-p')^2$$

four-momentum transfer squared at the proton vertex

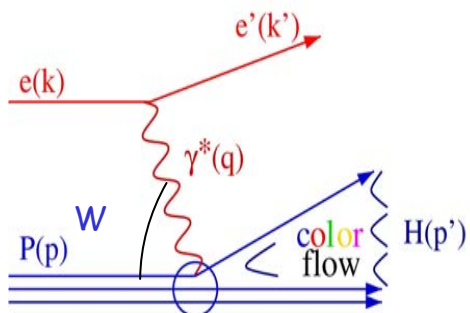
$$x_{IP} = \frac{(p-p') \cdot q}{p \cdot q} = \frac{M_x^2 + Q^2}{W^2 + Q^2}$$

momentum fraction of the proton carried by the Pomeron

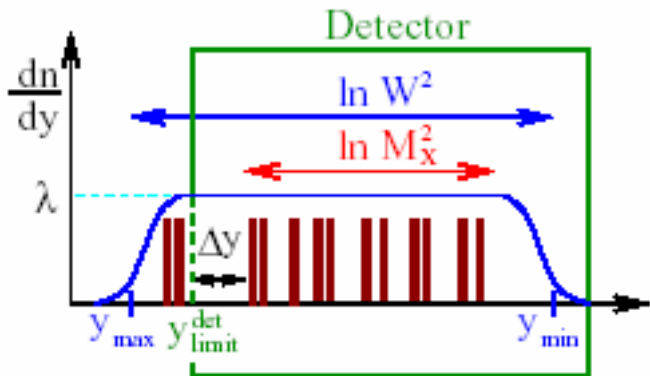
$$\beta = \frac{Q^2}{2(p-p') \cdot q} = \frac{x}{x_{IP}} = \frac{Q^2}{M_x^2 + Q^2}$$

fraction of the Pomeron momentum which enters the hard scattering

Non diffractive events



Uncorrelated particle emission between incoming p-direction and scattered quark.



$$y = \frac{1}{2} \ln \frac{E+p_z}{E-p_z}$$

Rapidity, property of a produced particle

$$\frac{dN_{\text{part}}}{dy} = \lambda = \text{const.}$$

$$W^2 = c_0 e^{y_{\text{max}} - y_{\text{min}}}$$

$$M_X^2 = c_0 e^{y_{\text{limit}} - y_{\text{min}}}$$

Poisson distr. for Δy in nondiffractive events

$$P(0) = e^{-\lambda \Delta y}$$

$$\frac{dN_{\text{nondiff}}}{d \ln M_X^2} = c \cdot e^{b \cdot \ln M_X^2}$$

Diffractive events :

$$\frac{dN_{\text{diff}}}{dM_X^2} \propto \frac{1}{(M_X^2)^n}$$

At high energies and not too low M_X
 $n \approx 1$

$$\frac{dN_{\text{diff}}}{d \ln M_X^2} \approx \text{const.}$$

Nondiffractive + diffractive contributions

$$\frac{dN}{d \ln M_x^2} = D + c \cdot e^{b \cdot \ln M_x^2}$$

D is the diffractive contribution

Two approaches for fit to the data

1.) take $D = \text{const.}$ for a limited range in $\ln M_x^2$

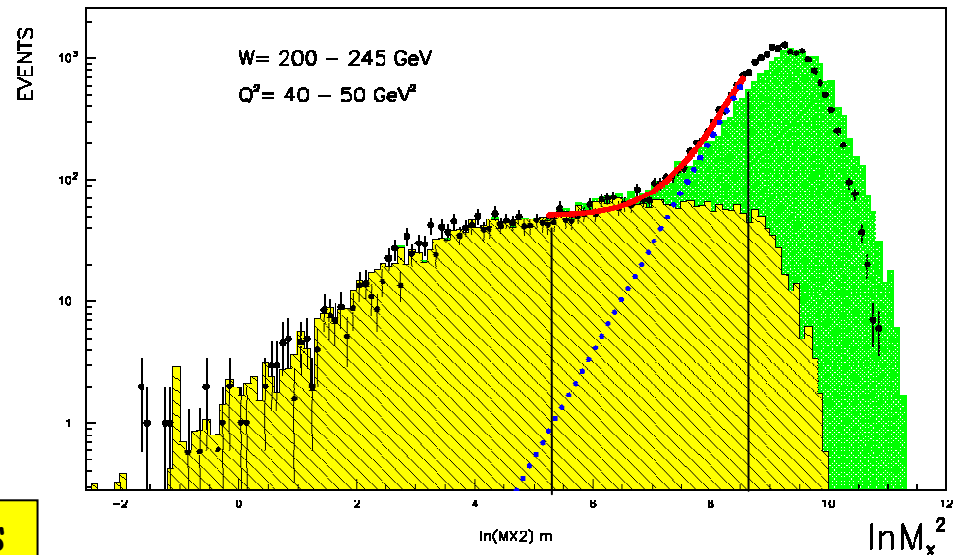
2.) take D from a BEKW-model (see later) parametrization which describes our measured data. This is an iterative procedure.

Fit slope b , c and D

for $\ln M_x^2 \leq \ln W^2 - \eta_0$

Determine diffractive events by subtracting nondiffractive events from measured data bin by bin as calculated from fitted values b and c .

Both approaches give the same results



Diffractive data selected by the M_x -method do not contain contribution from Regge exchanges. They contain, however, contributions from proton dissociative diffractive events.

MC-simulation

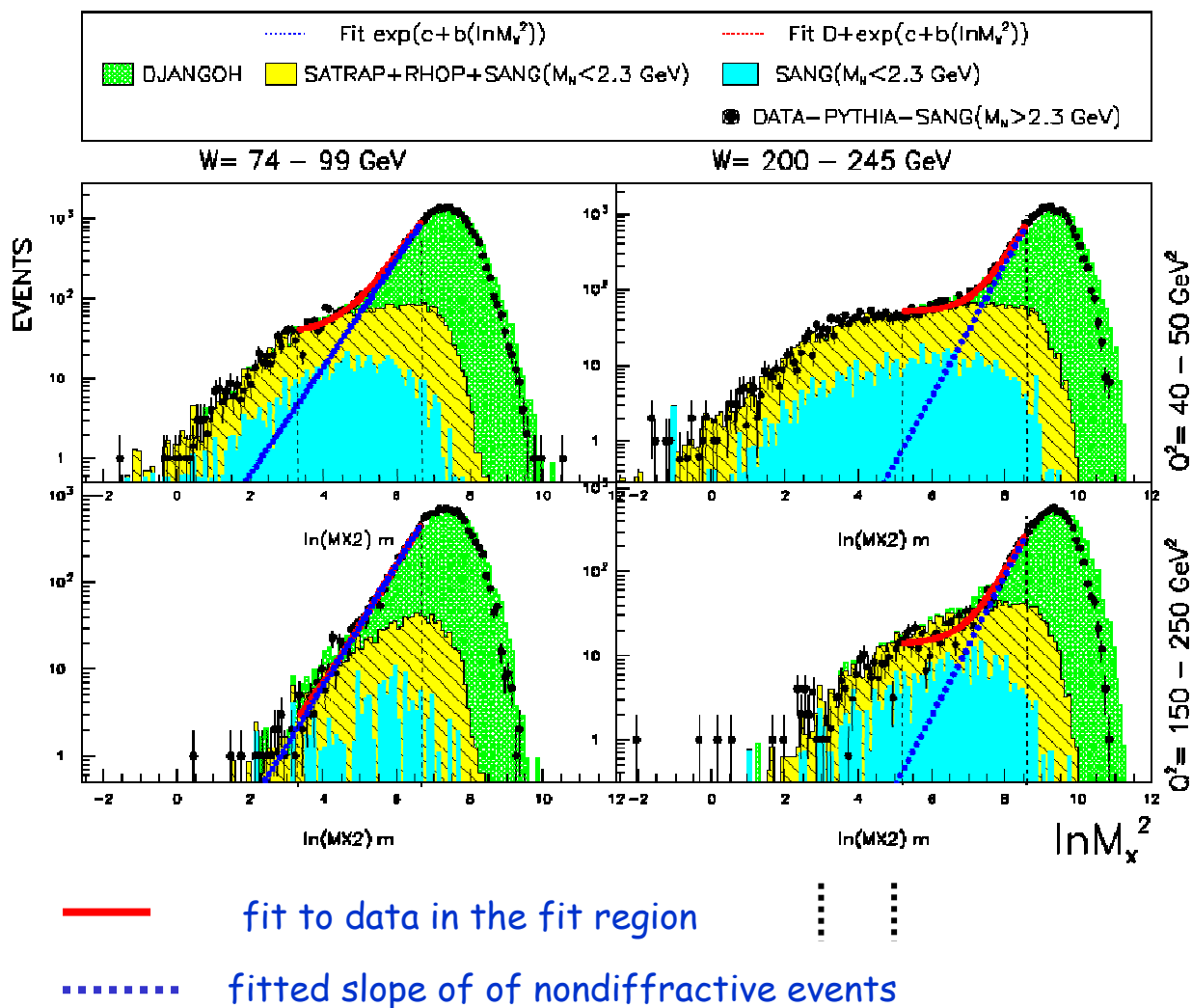
nondiffractive : DJANGO
 diffractive : SATRAP
 proton diss. : SANG

SANG adjusted to fit specially selected data which are dominated by proton dissociation



Proton dissociation can be reliably calculated for $M_N > 2.3 \text{ GeV}$ and has been subtracted from data

The ZEUS M_x -results contain contributions from proton dissociation for masses $M_N < 2.3 \text{ GeV}$.





ZEUS M_x - data from 1998 - 2000 (II)

Mx 98-99, Mx 99-00 (prel.)

Mx 98-99 : *

Published data from 1998-1999 period

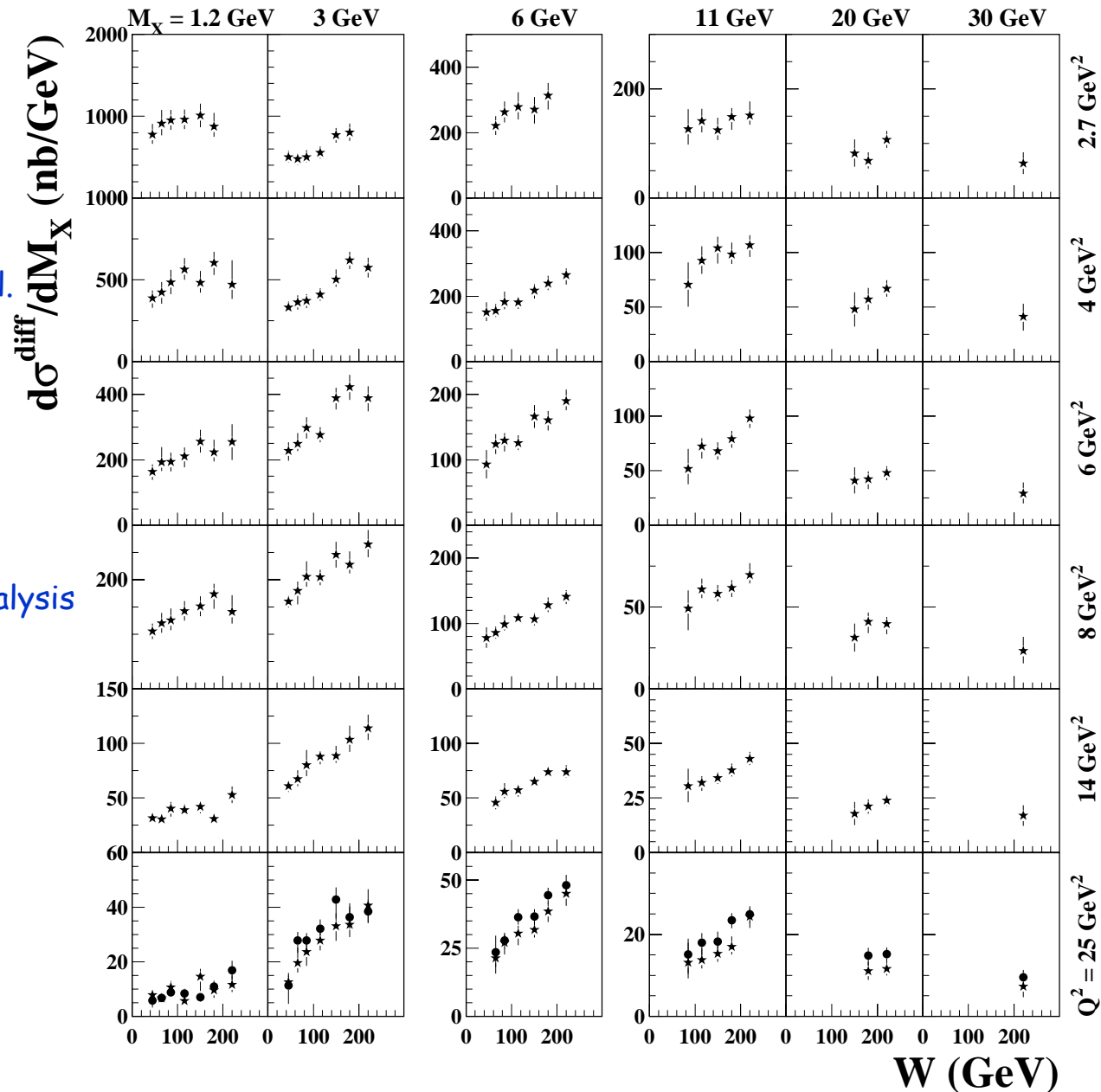
(ZEUS Coll., S.Chekanov et al. Nucl. Phys B 713, 3 (2005))

Prel. Mx 99-00: ●

Preliminary results from 1999-2000 period.

Extension of Mx 98-99 analysis to higher Q^2 .

Mx 98-99 and Mx 99-00 analyses have common bin at $Q^2 = 25 \text{ GeV}^2$





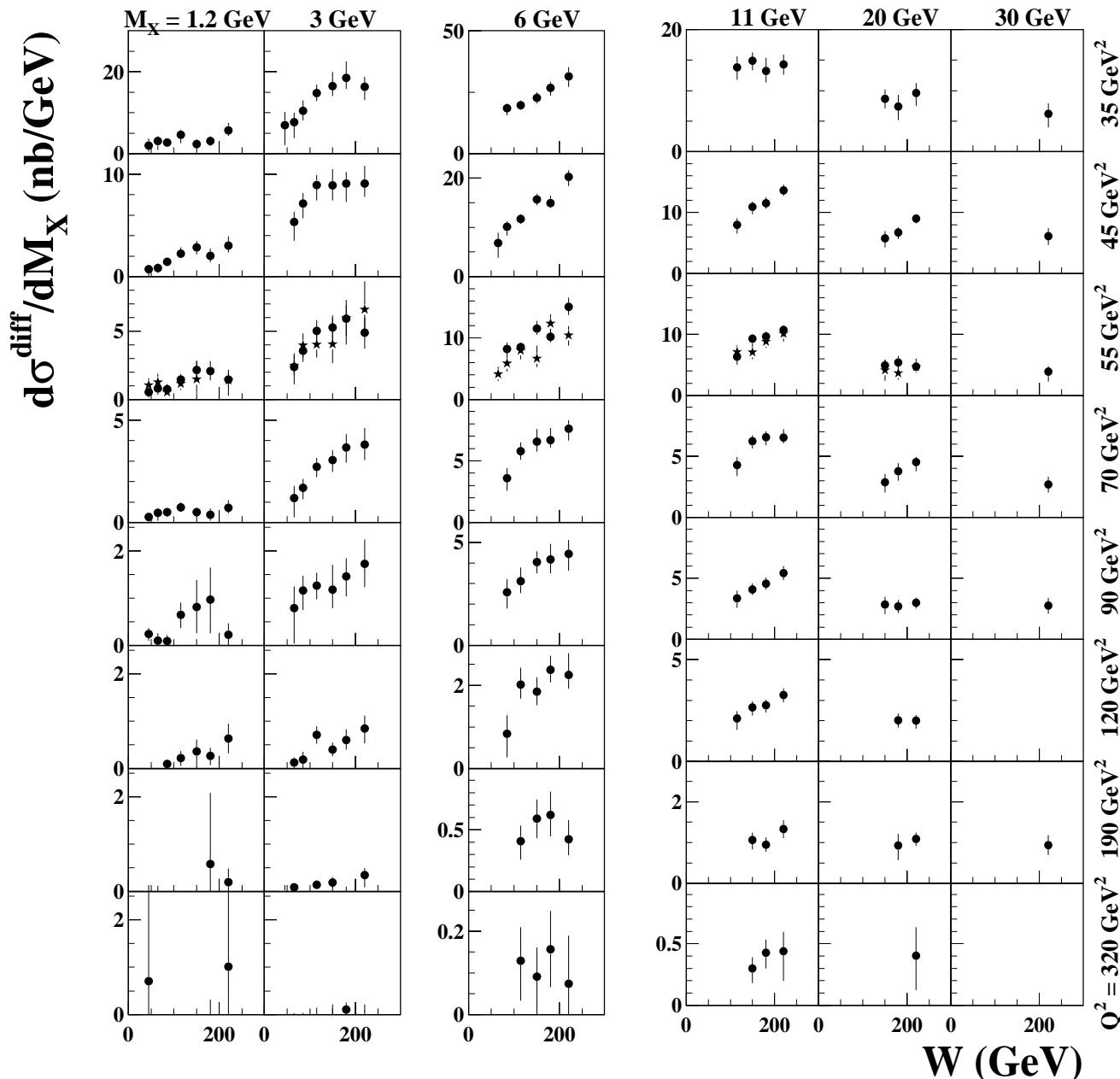
ZEUS M_x - data from 1998 - 2000 (III)



ZEUS M_x 98-99, ZEUS M_x 00 (prel.)

M_x 98-99: *

Prel.
 M_x 99-00: ●



M_x 98-99 and M_x 99-00 analyses have common bin at $Q^2 = 55 \text{ GeV}^2$

Within syst. errors good agreement between M_x 98-99 and M_x 99-00 results

Inclusive DIS:

For small x , F_2 rises rapidly as $x \rightarrow 0$

$$F_2 = c \cdot x^{-\lambda} \quad W \propto \frac{1}{x}$$

$$\lambda = \alpha_{IP}(0) - 1$$

Inclusive diffractive DIS:

$$\frac{d\sigma_{\gamma^* p \rightarrow XN}}{dM_X} = h \cdot \left(\frac{W}{W_0}\right)^{a^{diff}}$$

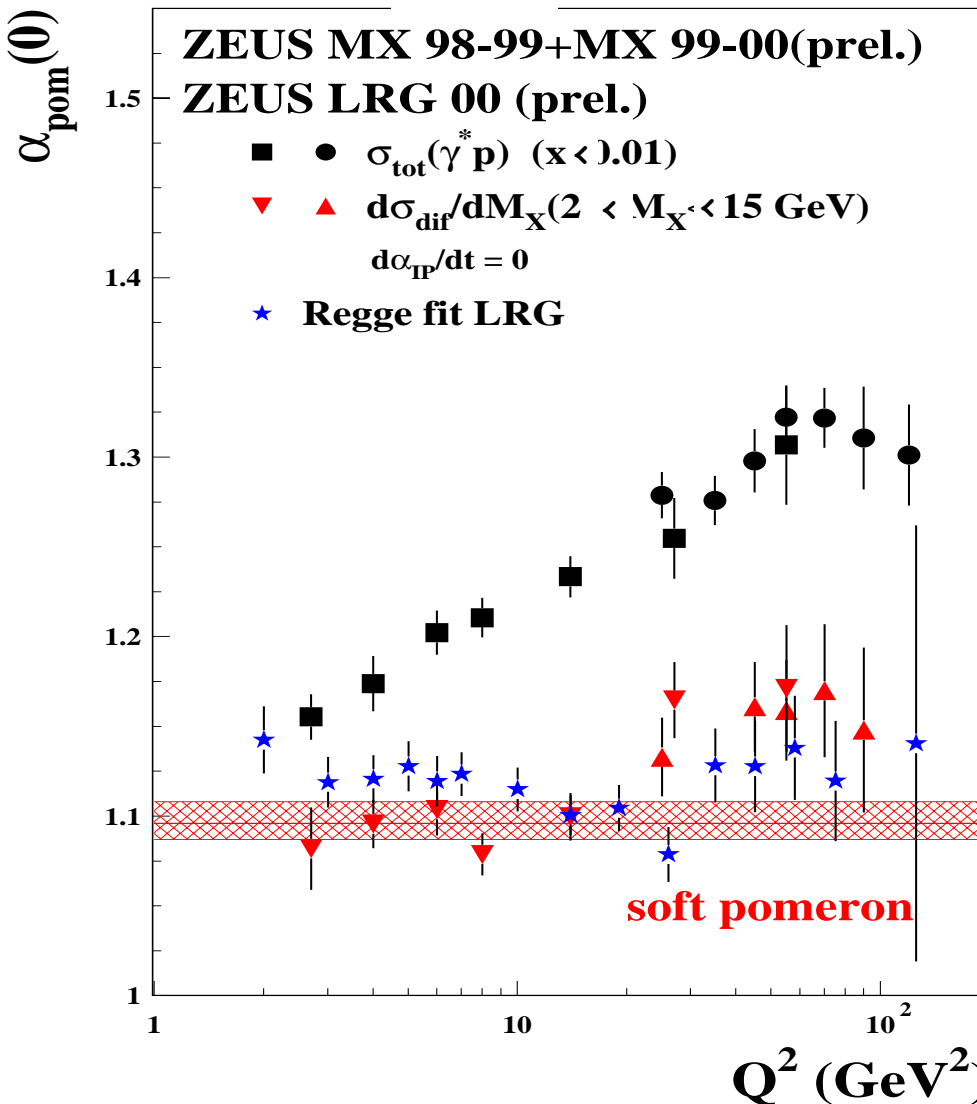
$$\bar{\alpha}_{IP} = 1 + \frac{a^{diff}}{4} \quad \text{averaged over } t$$

$$\alpha_{IP}(t) = \alpha_{IP}(0) + \alpha'_{IP} \cdot t$$

$$\frac{d\sigma}{dt} = f(t) \cdot e^{2(\alpha_{IP}(t)-1) \cdot \ln\left(\frac{W}{W_0}\right)^2}$$

$$\frac{d\sigma}{dt} \propto e^{A \cdot t} \quad \text{for small } t.$$

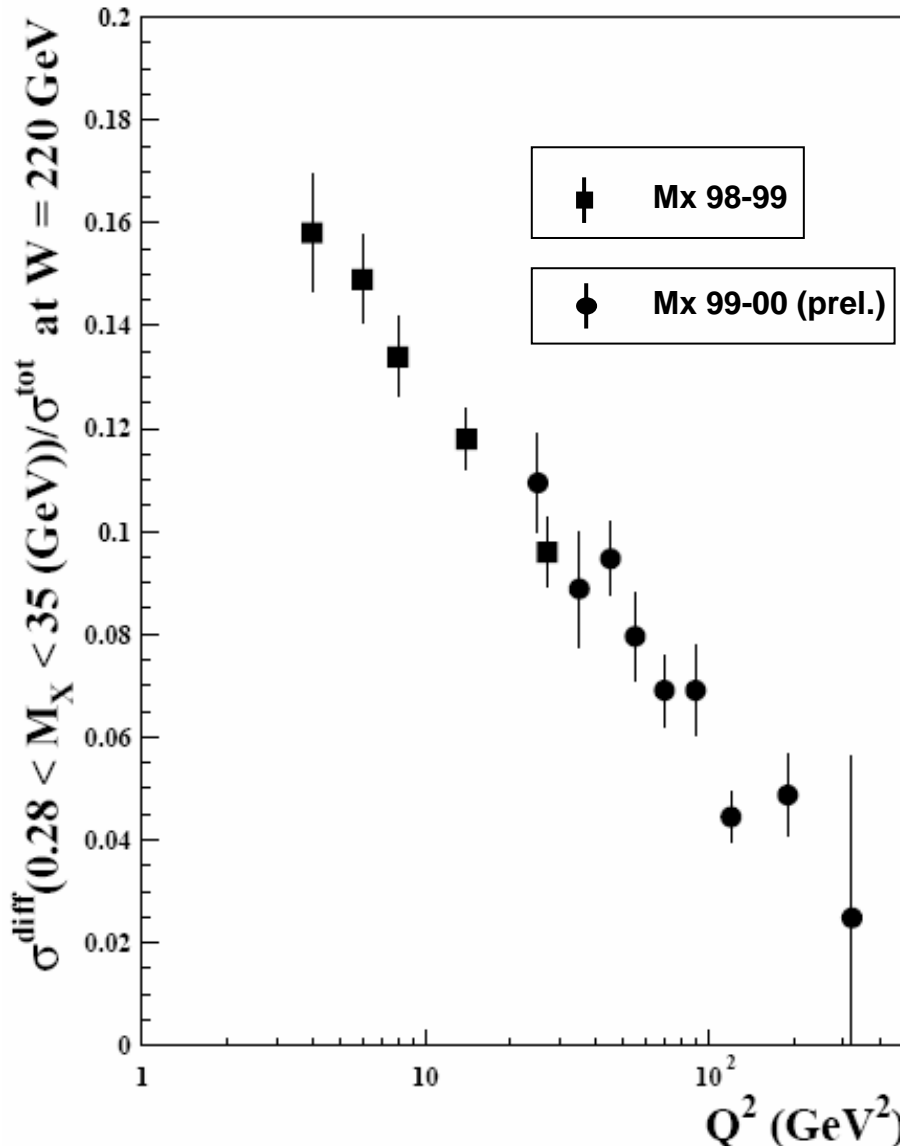
take $A = 7.9 \pm 0.5(\text{stat.}) \pm 0.5(\text{syst.}) \text{ GeV}^2$
as measured by ZEUS LPS



Inclusive DIS and inclusive diffractive DIS are not described by the same 'Pomeron'.

Ratio of total diffractive cross-section to total DIS cross-section

Ratio plotted at $W=220 \text{ GeV}$ because only there the full M_X range is covered by measurements



$$r = \sigma^{\text{diff}}(0.28 < M_X < 35 \text{ GeV}) / \sigma^{\text{tot}}$$

Within the errors of the measurements r is independent of W .

At $W=220 \text{ GeV}$, r can be fitted by

$$r = 0.22 - 0.034 \cdot \ln(1+Q^2)$$

This logarithmic dependence of the ratio of total diffractive cross-section to the total DIS cross section indicates that diffraction is a leading twist process for not too low Q^2 .

$$\frac{d^4\sigma}{dQ^2 dt dx_{IP} d\beta} = \frac{2\pi\alpha_{em}}{\beta Q^2} [1 - (1-y)^2] \cdot F_2^{D(4)}(Q^2, t, x_{IP}, \beta)$$

ZEUS neglects the contribution from longitudinal structure function

H1 defines : sizable only at high y

$$\sigma_r^D = F_2^D - \overbrace{\frac{y^2}{1+(1-y)^2} F_L^D}$$

If t is not measured :

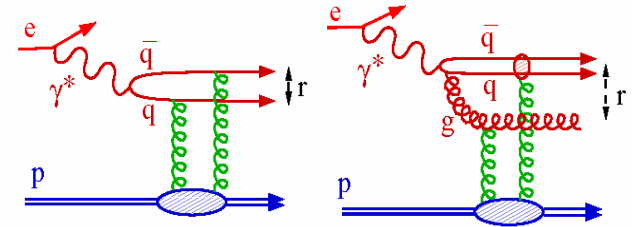
$$\frac{d^3\sigma_{\gamma^* p \rightarrow XN}^{diff}}{dQ^2 d\beta dx_{IP}} = \frac{2\pi\alpha^2}{\beta Q^4} [1 + (1-y)^2] \cdot F_2^{D(3)}(\beta, x_{IP}, Q^2)$$

$$\frac{1}{2M_X} \frac{d\sigma_{\gamma^* p \rightarrow XN}^{diff}(M_X, W, Q^2)}{dM_X} = \frac{4\pi^2\alpha}{Q^2(Q^2 + M_X^2)} x_{IP} F_2^{D(3)}(\beta, x_{IP}, Q^2)$$

If $F_2^{D(3)}(\beta, x_{IP}, Q^2)$ is interpreted in terms of quark densities, it specifies the probability to find in a proton which undergoes a diffractive interaction a quark carrying a fraction $x = \beta x_{IP}$ of the proton momentum.

Fit with BEKW model

(Bartels, Ellis, Kowalski and Wüsthoff, 1998)



- $x_{IP} F_2^{D(3)} = c_T \cdot F_{q\bar{q}}^T + c_L \cdot F_{q\bar{q}}^L + c_g \cdot F_{q\bar{q}g}^T$

$$F_{q\bar{q}}^T = \left(\frac{x_0}{x_{IP}}\right)^{n_T(Q^2)} \cdot \beta(1 - \beta),$$

$$F_{q\bar{q}}^L = \left(\frac{x_0}{x_{IP}}\right)^{n_L(Q^2)} \cdot \frac{Q_0^2}{Q^2 + Q_0^2} \cdot \left[\ln\left(\frac{7}{4} + \frac{Q^2}{4\beta Q_0^2}\right)\right]^2 \cdot \beta^3(1 - 2\beta)^2,$$

$$F_{q\bar{q}g}^T = \left(\frac{x_0}{x_{IP}}\right)^{n_g(Q^2)} \cdot \ln\left(1 + \frac{Q^2}{Q_0^2}\right) \cdot (1 - \beta)^\gamma$$

assume $n_T(Q^2) = c_4 + c_7 \ln\left(1 + \frac{Q^2}{Q_0^2}\right)$, $n_L(Q^2) = c_5 + c_8 \ln\left(1 + \frac{Q^2}{Q_0^2}\right)$,

$$n_g(Q^2) = c_6 + c_9 \ln\left(1 + \frac{Q^2}{Q_0^2}\right)$$

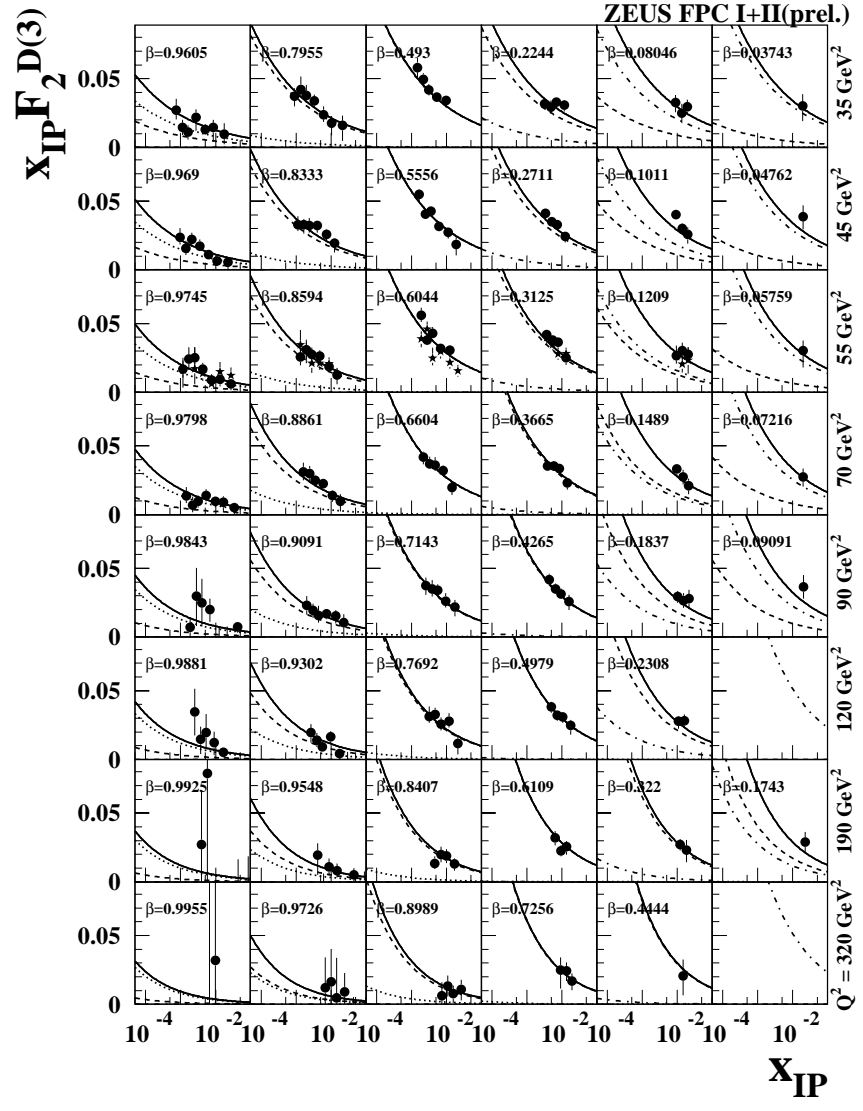
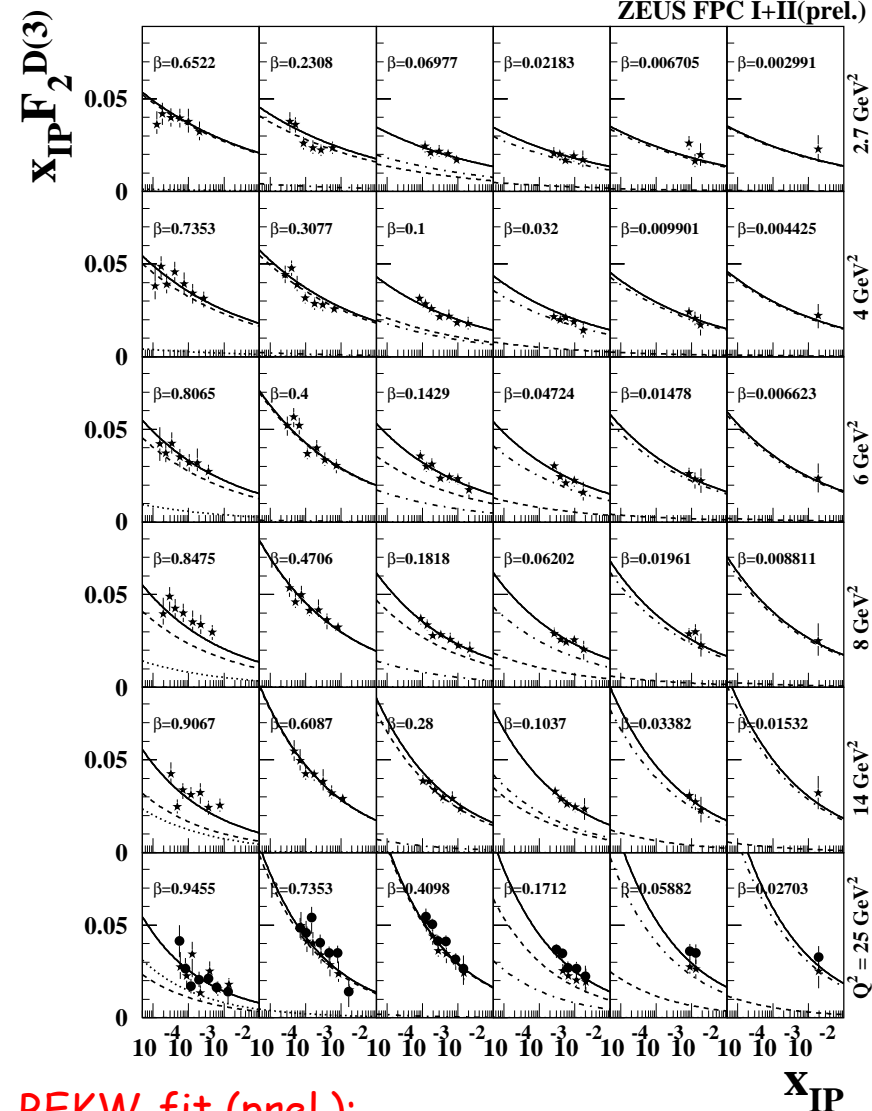
The ZEUS data support taking $n_T(Q^2) = n_g(Q^2) = n_1 \cdot \ln(1 + Q^2/Q_0^2)$ and $n_L = 0$

Taking $x_0 = 0.01$ and $Q_0^2 = 0.4 \text{ GeV}^2$ results in the **modified BEKW model (BEKW(mod))** with the 5 free parameters :

$$c_T, c_L, c_g, n_1^{T,g}, \gamma$$



$x_{IP}F_2D(3)$ Results from the Mx 98-99 and Mx 99-00 Analyses with BEKW(mod) Fit (I)

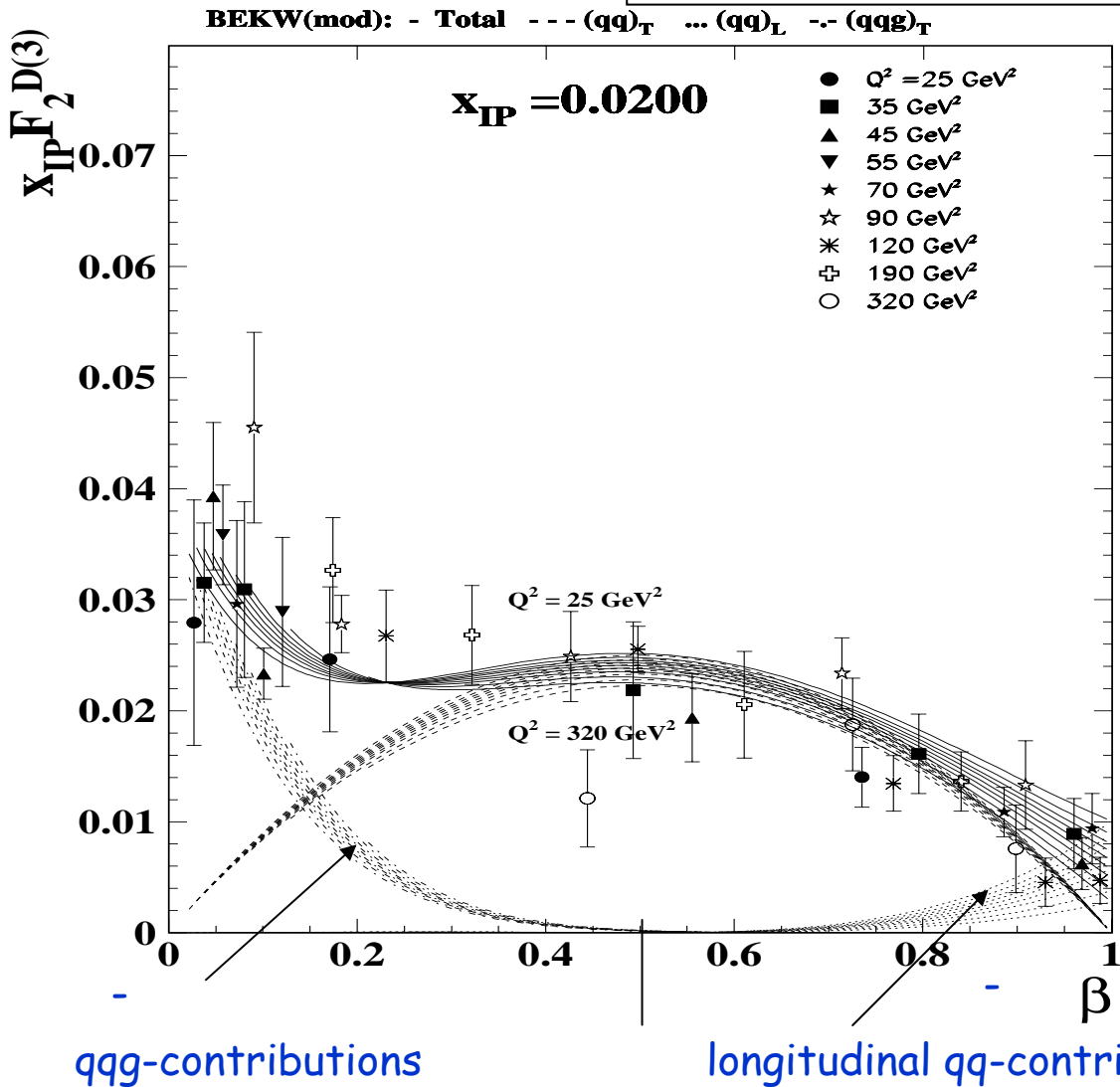


BEKW-fit (prel.):

> 400 points, 5 parameters
 $\chi^2/n_D = 0.71$, total errors

- sum of all contributions
- longitudinal qq contribution
- - - transverse qq contribution
- . - . transverse qqg contribution

ZEUS Mx 99-00 (prel.)



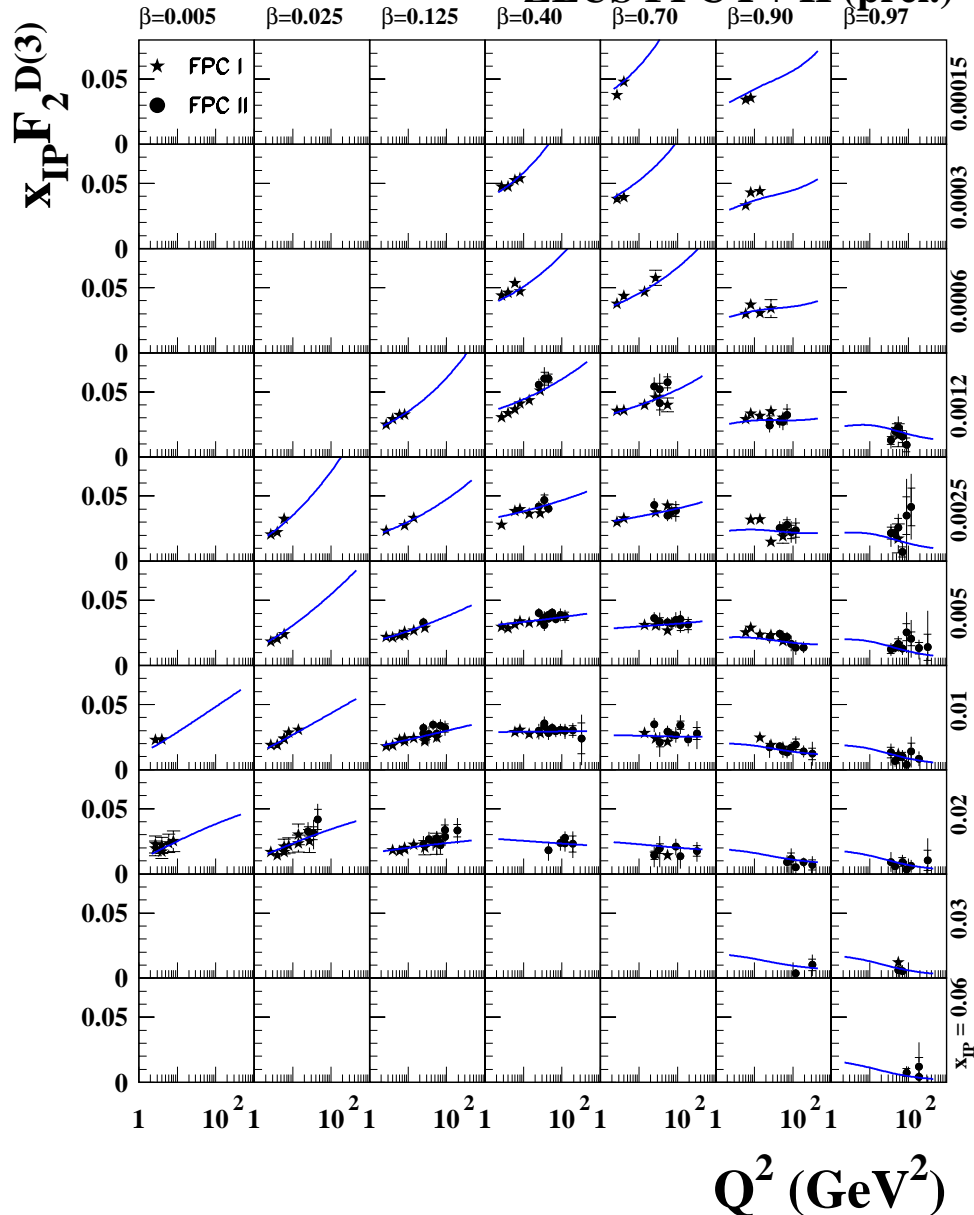
Fixed $x_{IP} = 0.02$

$25 < Q^2 < 320 \text{ GeV}^2$
in one plot

The 3 contributions
from BEKW(mod)
fit for the above
 Q^2 values plotted

The BEKW model has an effective QCD-type Q^2 -evolution incorporated.

ZEUS FPC I + II (prel.)



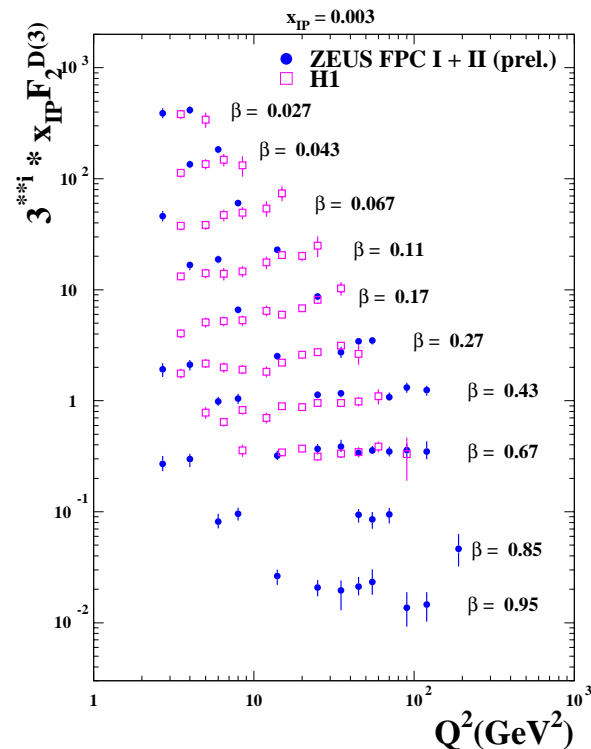
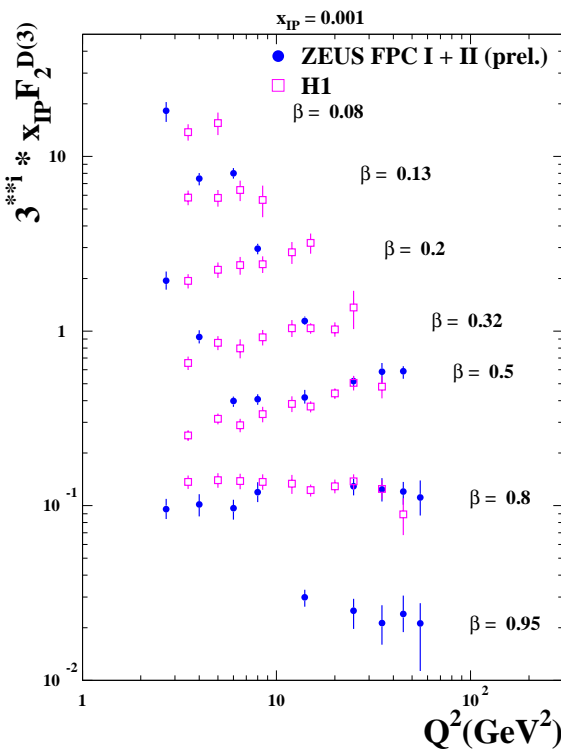
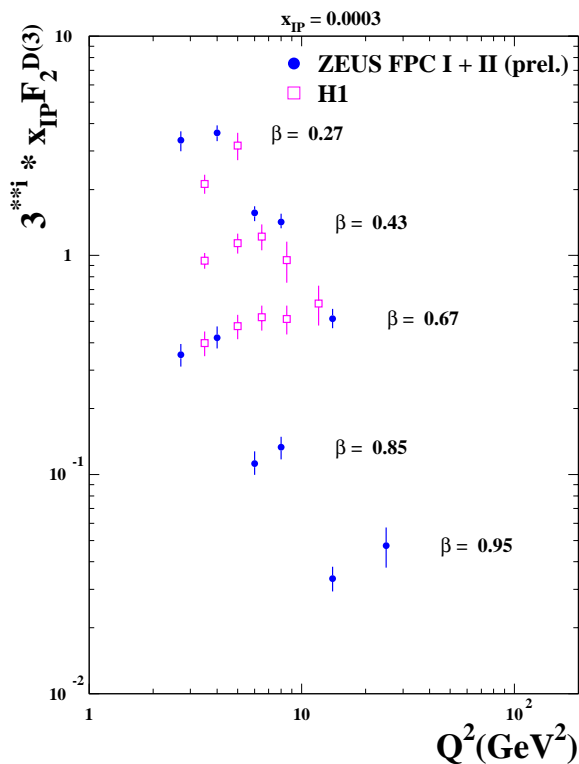
Events in fixed (x_{IP}, β) -bins

— Result of the BEKW(mod) fit

$x_{IP}F_2D(3)$ shows considerable scaling violations:

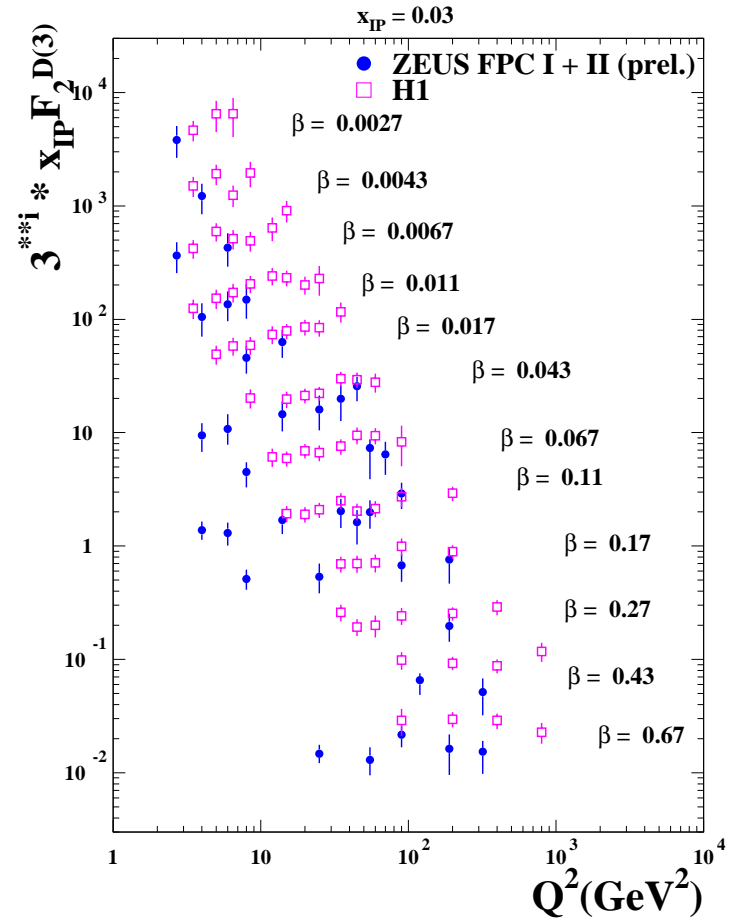
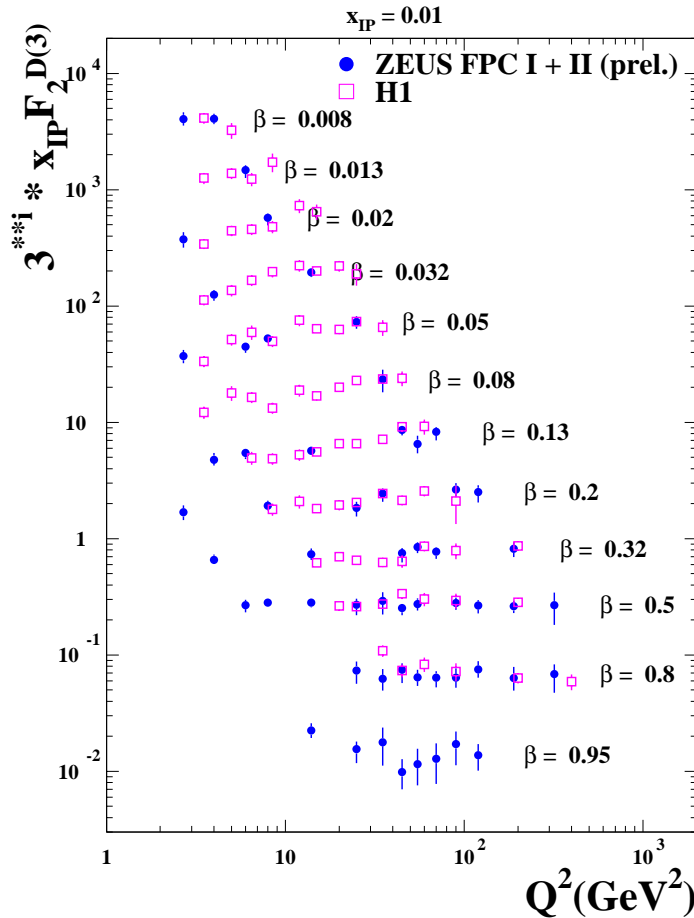
from positive scaling violations over near constancy to negative scaling violations.

$x_{IP}F_2D(3)$ Results from the Mx 98-99 and Mx 99-00 Analyses Comparison with H1 Results in H1-binning (I)



Note: ZEUS results contain contributions from p-dissociation with masses $M_{p-diss} < 2.3 \text{ GeV}$,
H1 results contain contributions with masses $M_{p-diss} < 1.6 \text{ GeV}$.

ZEUS results do not contain contributions from Reggeon-exchanges,
H1 results may contain such contributions for higher x_{IP} .



Comparison to H1 data

Fair agreement,
except maybe for a few (x_{IP}, β) bins

Note: ZEUS points are shifted to H1 bins using BEKW parametrization. Only those ZEUS point are shown for which the shift was <30%.

- ZEUS presented **preliminary results** on inclusive diffraction from **the M_x -Method** for the extraction of inclusive diffractive events from the 1999-2000 data and from **published data** taken during 1998-1999.
- The results span a wide range of the kinematic region **from low to high Q^2** .
- There is **good agreement** between the published data from 1998-1999 and the preliminary results from 1999-2000.
- The ratio of the total diffractive cross section to the total DIS cross section indicates that diffraction is a leading twist process.
- DIS and diffractive DIS cannot be describe by a single unique pomeron.
- The combined data from 1998-2000 are well described by the BEKW(mod) model.
- The diffractive structure function $x_{\text{IP}}F_2^{\text{D}(3)}$ exhibits logarithmic scaling violations.
- There is good **agreement compared to results from H1**.