

Experimental results on heavy quark fragmentation



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OUTLINE :

Introduction

b fragmentation functions

b fragmentation branchings

b fragmentation ratios

c fragmentation in e^+e^-

c fragmentation at HERA

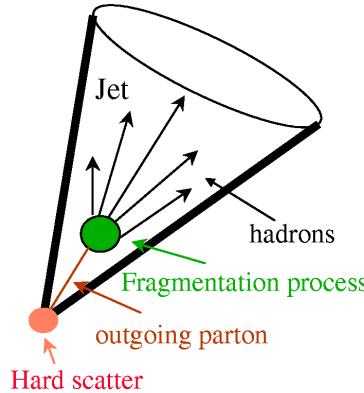
Summary

B A C K U P :

estimates of extrapolation factors

fragm. branchings for excited D mesons

Heavy Quark fragmentation issues



Important to measure

HQ fragmentation to find :

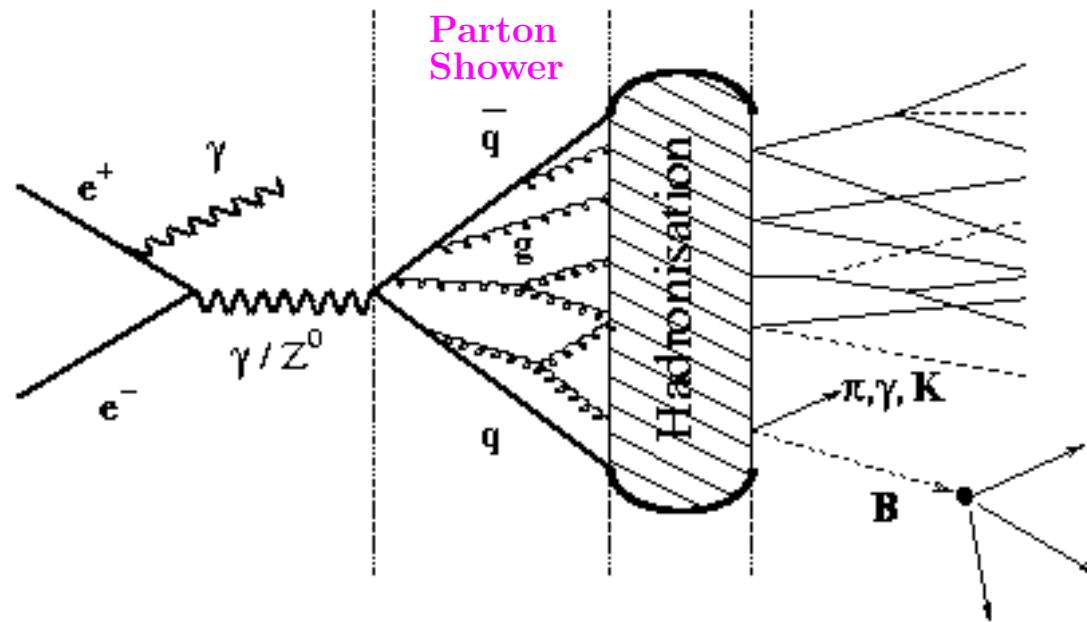
- 1) What is the proper parameterisation for the fractional transfer of b/c -quark energy/momentum to a given B/D -meson (z) ?
fragmentation function (FF), $f(z)$ or $D(z)$
- 2) What are the relative fragm. branchings (FB's) of B/D -hadrons ?
$$f(c \rightarrow D) = \frac{N(D)}{N(c)} = \frac{\sigma(D)}{\sum_{\text{all}} \sigma(D)}$$
 - a) Are u and d quarks produced equally ? $R_{u/d} = \frac{c\bar{u}}{cd}$
 - b) What is the s -quark production suppression ? $\gamma_s = \frac{2c\bar{s}}{cd+c\bar{u}}$
 - c) Are vector (B^*/D^*) and pseudoscalar (B/D) mesons produced as predicted by spin counting ? $P_v = \frac{V}{V+PS}$ ($= 0.75$?)
- 3) Are these functions, branchings and ratios universal ?
compare results in e^+e^- annihilations with those at HERA

Heavy Quark fragmentation in e^+e^- annihilations

Production

Fragmentation

Decays



HQ fragmentation is hard

harder for larger m_Q

e.g., for Peterson param.:

$$f(z) \propto \frac{1}{z(1-1/z-\epsilon/(1-z))^2}$$

$$\epsilon(b) \sim \frac{m_c^2}{m_b^2} \epsilon(c) \sim 0.1 \epsilon(c)$$

pQCD is applicable to “initial” Q-fragmentation: LO, NLO, LL, NLL, ...

anyhow, some parameterisation is needed for the non-perturbative (NP) rest

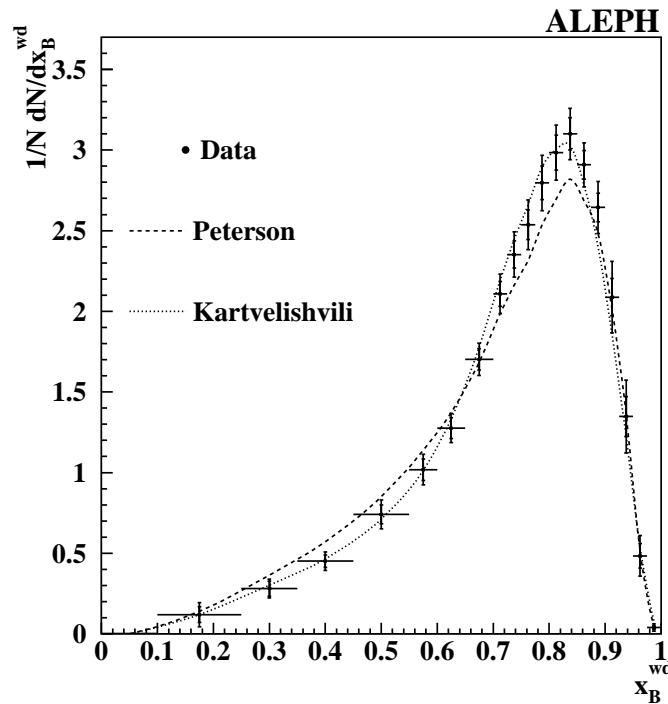
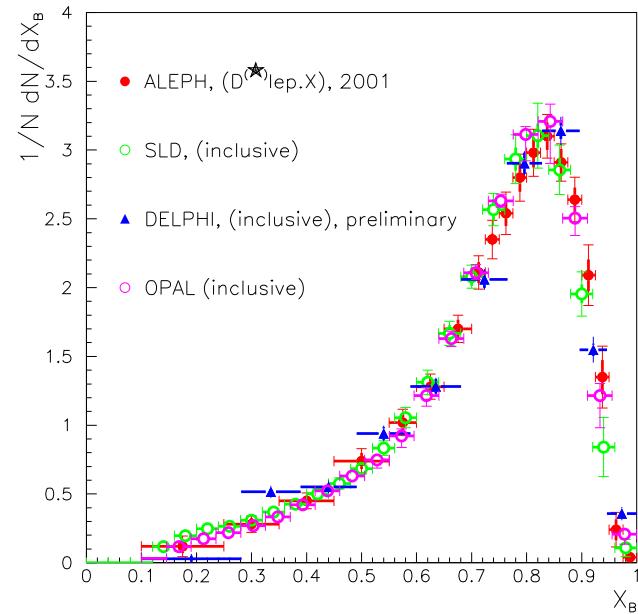
the NP parameterisation is strongly dependent from the perturbative core

(it is wrong to use MC fragmentation for NLO w/o full retuning the fragm. parameters)

the NP parameterisation can include some decays

FB's are expected to be independent from the perturbative core

b fragmentation function, FF with LL MC



measured at LEP and SLD
from sec. vertices or s/l decays (ALEPH)

in terms of the scaled energy: $x \equiv E_{\text{hadron}}/E_{\text{beam}}$
for weakly decaying B hadrons

$$\langle x \rangle = 0.7193 \pm 0.0016^{+0.0038}_{-0.0033} \text{ (OPAL)}$$

FF with LL Monte Carlo

Bowler, $\frac{1}{z^{1+bm_\perp^2}}(1-z)^a \exp\left(\frac{-bm_\perp^2}{z}\right)$, and

Lund symmetric, $\frac{1}{z}(1-z)^a \exp\left(\frac{-bm_\perp^2}{z}\right)$, are
in good agreement with data (2 parameters)

Kartvelishvili et al., $z^\alpha(1-z)$, is o.k. with 1 parameter

Collins-Spiller, $(\frac{1-z}{z} + \frac{(2-z)\epsilon}{1-z})(1+z^2)(1 - \frac{1}{z} - \frac{\epsilon}{1-z})^{-2}$, and
Peterson are too broad (1 par.)

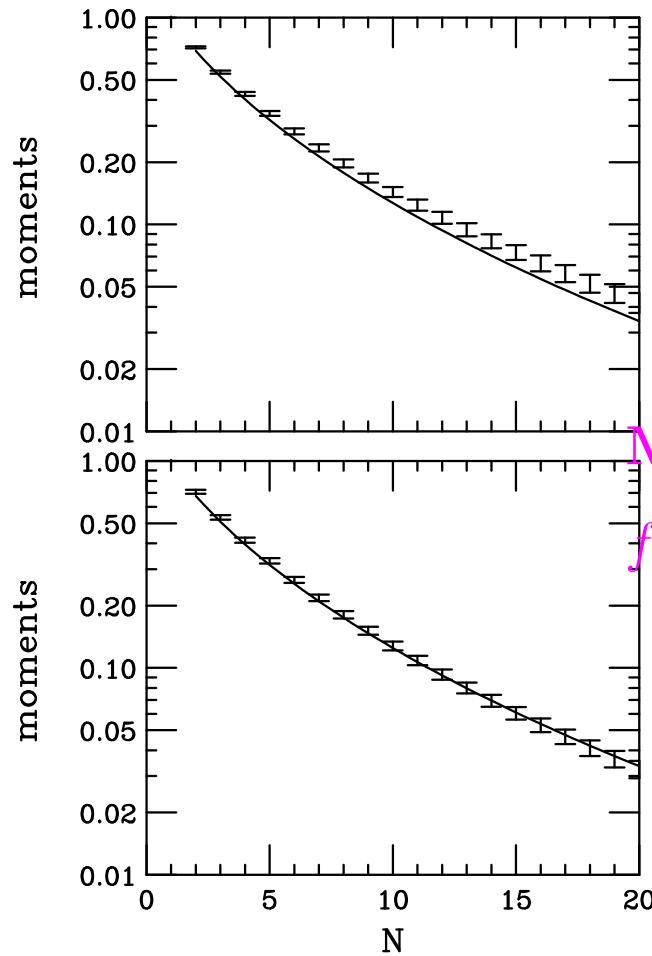
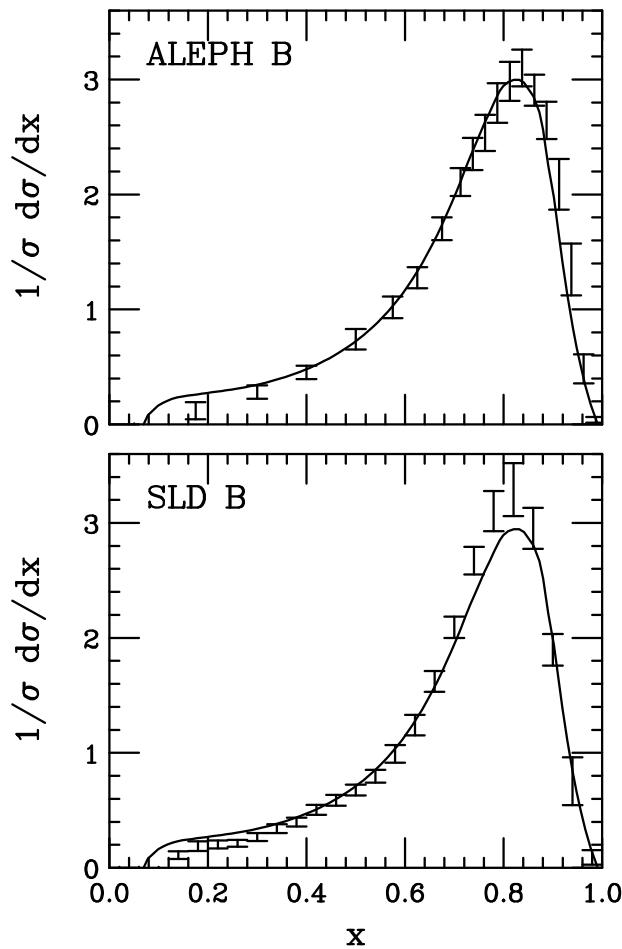
$$z \equiv (E + p_{||})_{\text{hadron}} / (E + p)_{\text{quark}}$$

HERWIG cluster model disfavoured

(what about MC@NLO ?)

b FF with NLO+NLL+Sudakov pQCD

Cacciari, Nason, Oleari



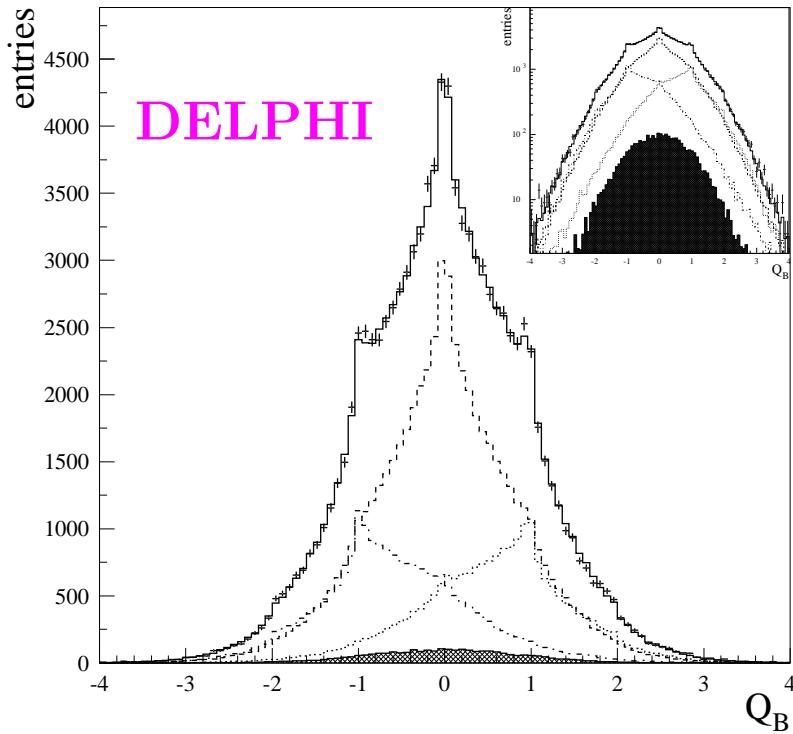
Mellin transform:

$$f(N) \equiv \int_0^1 x^{N-1} f(x) dx$$

$$D_{NP}(x) \propto (1-x)^a x^b \text{ (Colangelo-Nason)}$$

provides reasonable description with $a = 24 \pm 2$, $b = 1.5 \pm 0.2$

b fragmentation branchings



HFAG from $B_s^0 \rightarrow D_s^- l^+ \nu_l X$, $\Lambda_b^0 \rightarrow \Lambda_c^+ l^- \bar{\nu}_l X$, $\Xi_b^- \rightarrow \Xi^- l^- \bar{\nu}_l X$

using $f_u = f_d$ and $f_u + f_d + f_s + f_{\text{baryon}} = 1$:

$$f_u = f_d = (40.3 \pm 1.1)\%$$

$$f_s = (8.8 \pm 2.1)\%$$

$$f_{\text{baryon}} = (10.7 \pm 1.8)\%$$

using time-integrated mixing probabilities (f_d and f_s) \Rightarrow

$$f_u = f_d = (39.7 \pm 1.1)\%$$

$$f_s = (10.7 \pm 1.1)\%$$

$$f_{\text{baryon}} = (9.9 \pm 1.7)\%$$

fit of the secondary vertex charge:

$$f^+ = (42.09 \pm 0.82 \pm 0.89)\%$$

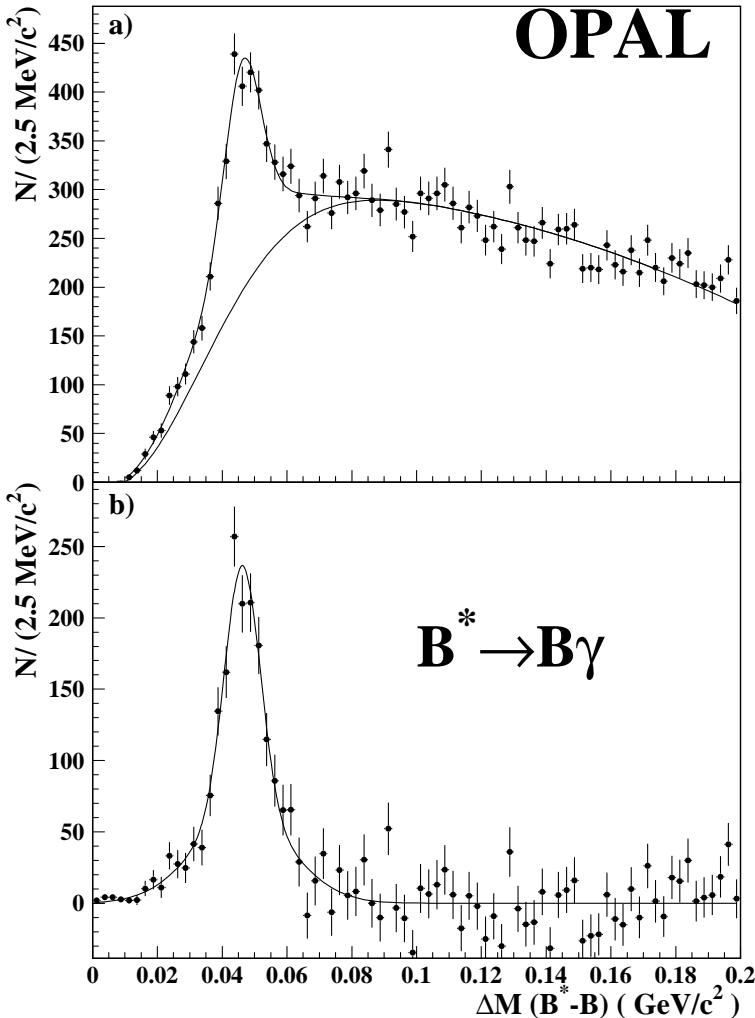
$$\text{using } f(b \rightarrow \Xi_b^-) = (1.1 \pm 0.5)\%$$

from LEP measurements of $\Xi_b^- \rightarrow \Xi^- l^- \bar{\nu}_l X$ and neglecting Ω_b^-

$$f_u = (40.99 \pm 0.82 \pm 1.11)\%$$

b fragmentation ratios

- a) $R_{u/d} = f_u/f_d \equiv 1$ by construction (agrees with DELPHI's f_u measurement)
- b) $\gamma_s = \frac{2f_s}{f_u+f_d} = 0.27 \pm 0.03$ B_s production suppressed by factor ≈ 3.7



c) P_v



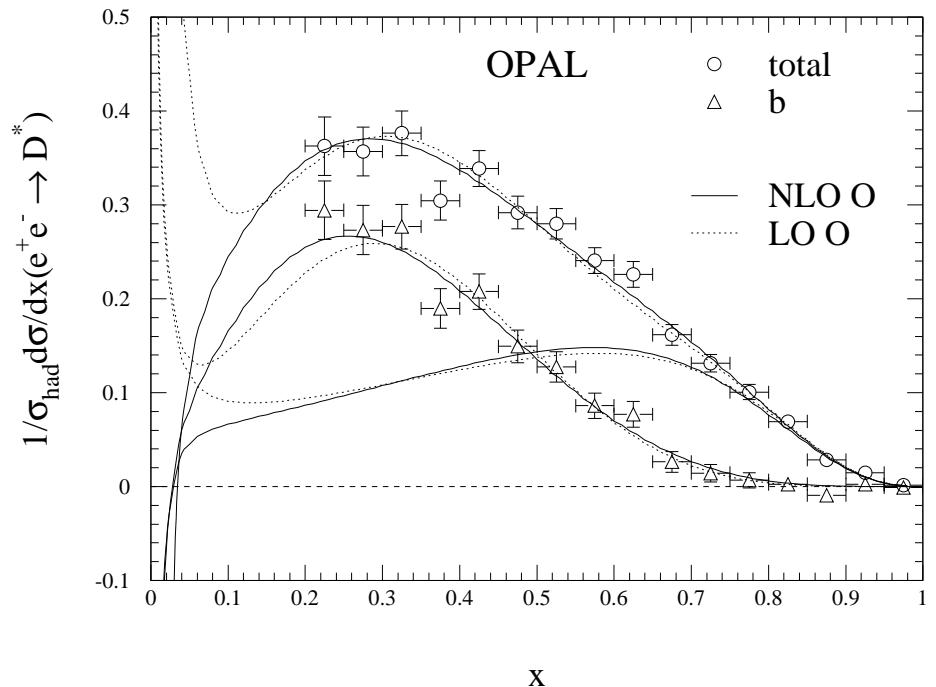
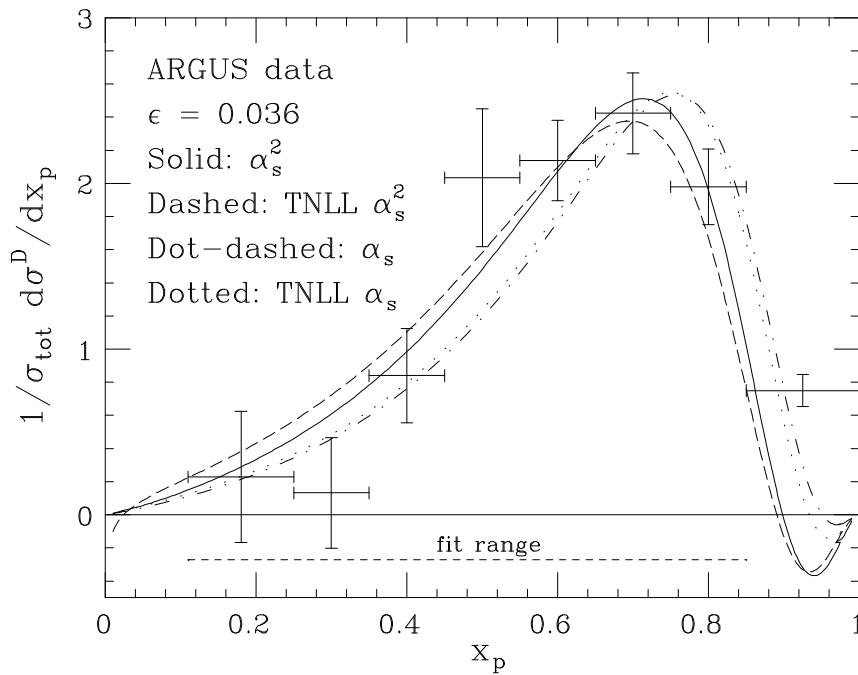
$$P_v = \sigma(B^*)/\sigma(B)$$

OPAL, ALEPH, DELPHI, L3 :

$$P_v = 0.75 \pm 0.04$$

agrees with spin counting

c fragmentation function, NLO with Peterson FF



Fixed-order approach (NLO fits of P. Nason and C. Oleari) :

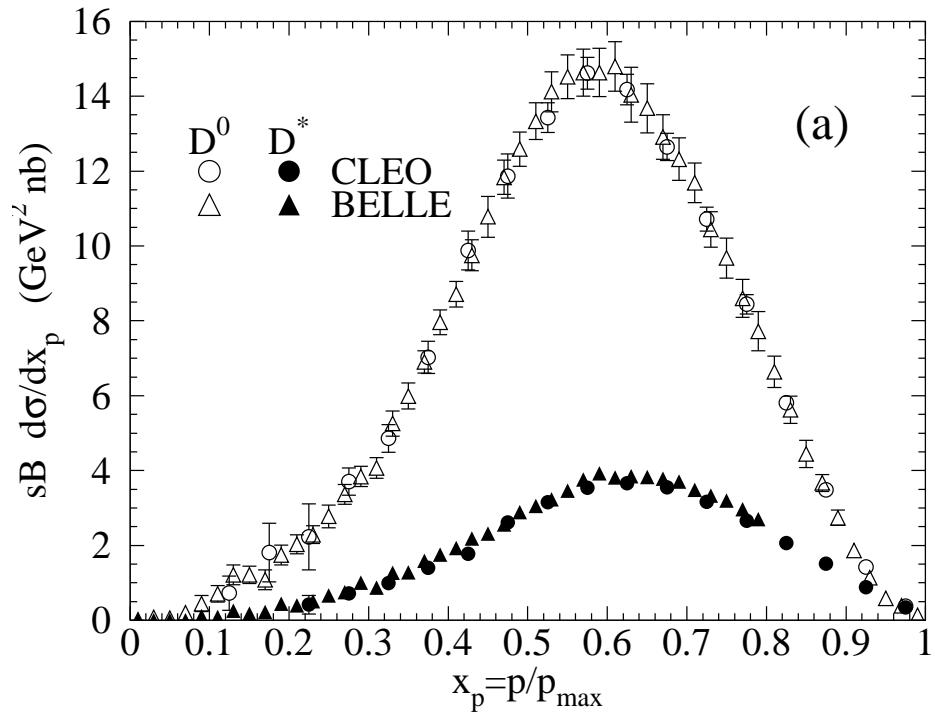
$$\epsilon(D^*, D_s) = 0.035 \quad \text{ARGUS data} \Leftarrow \text{Recommended !}$$

Resummed approach (LEP I data fit) :

$$\text{Kniehl et al.} \quad \epsilon(D^*) = 0.116$$

(fit results depend from the perturbative core)

c fragmentation function with LL MC



FF (Belle, D^{*+}) with LL Monte Carlo

Bowler : $\chi^2/NDF = 541.8/55$

Lund : $\chi^2/NDF = 965.6/55$

Kartvelishvili : $\chi^2/NDF = 1271.1/54$

Collins-Spiller : $\chi^2/NDF = 1540.7/54$

Peterson : $\chi^2/NDF = 3003.0/54$

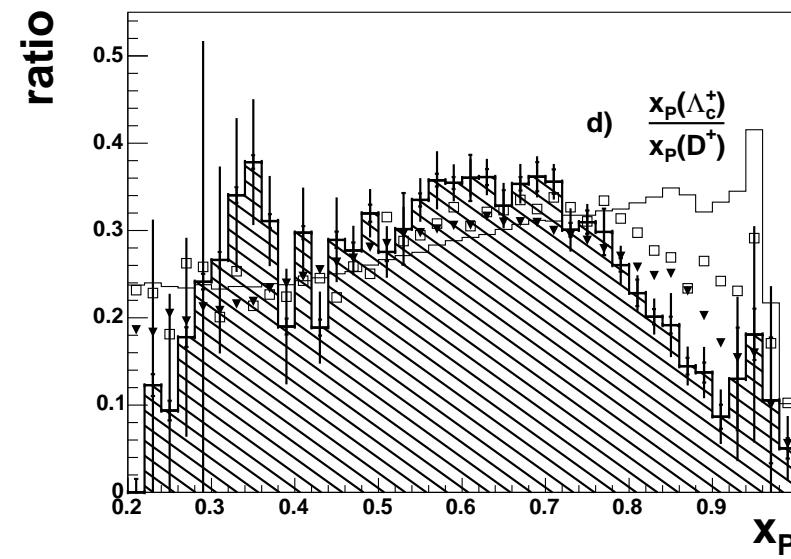
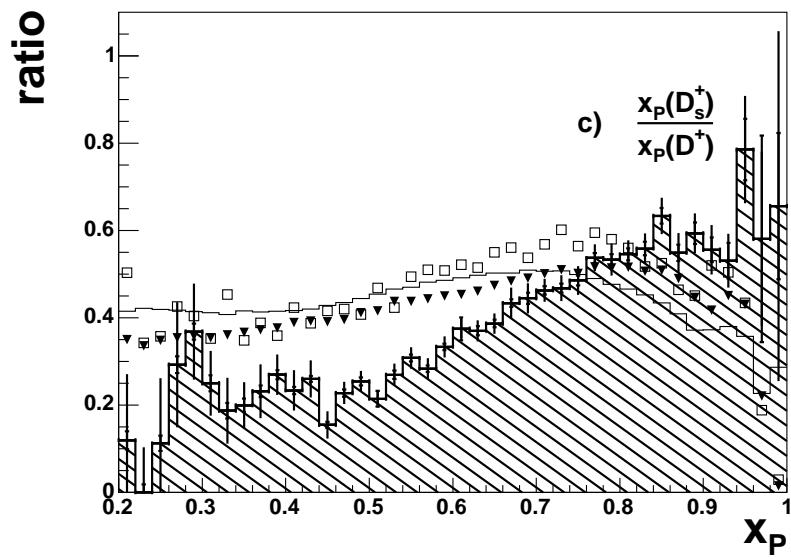
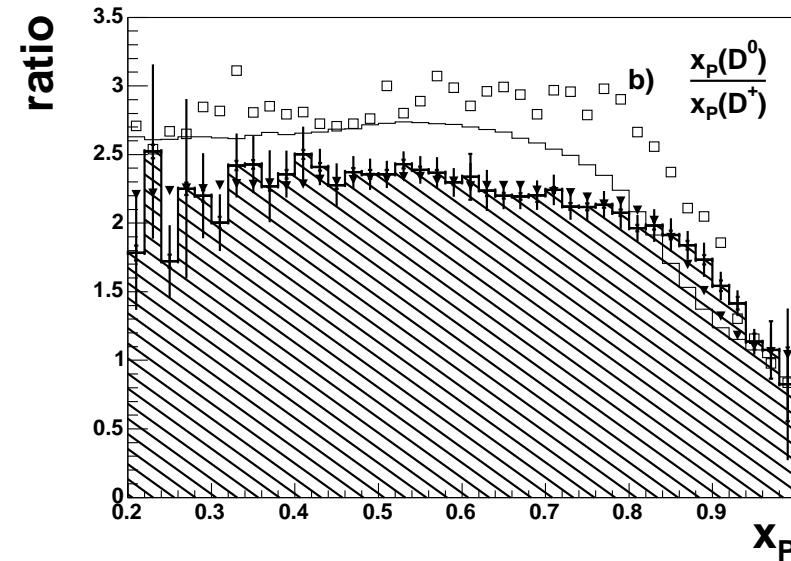
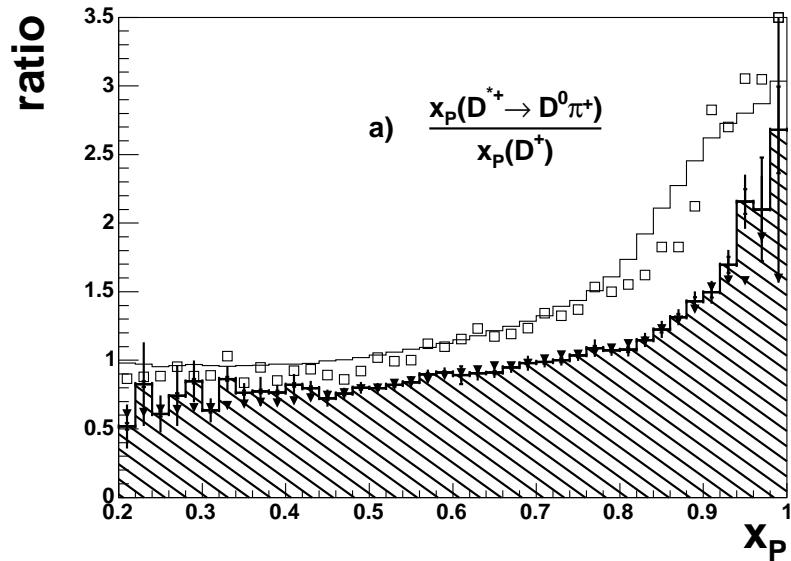
Recent precise measurements

from CLEO and Belle

$\langle x_p \rangle = 0.611 \pm 0.007 \pm 0.004$ (CLEO, D^{*+})

Qualitatively, the same picture
as for b FF with LL MC

differential fragmentation ratios (Belle)



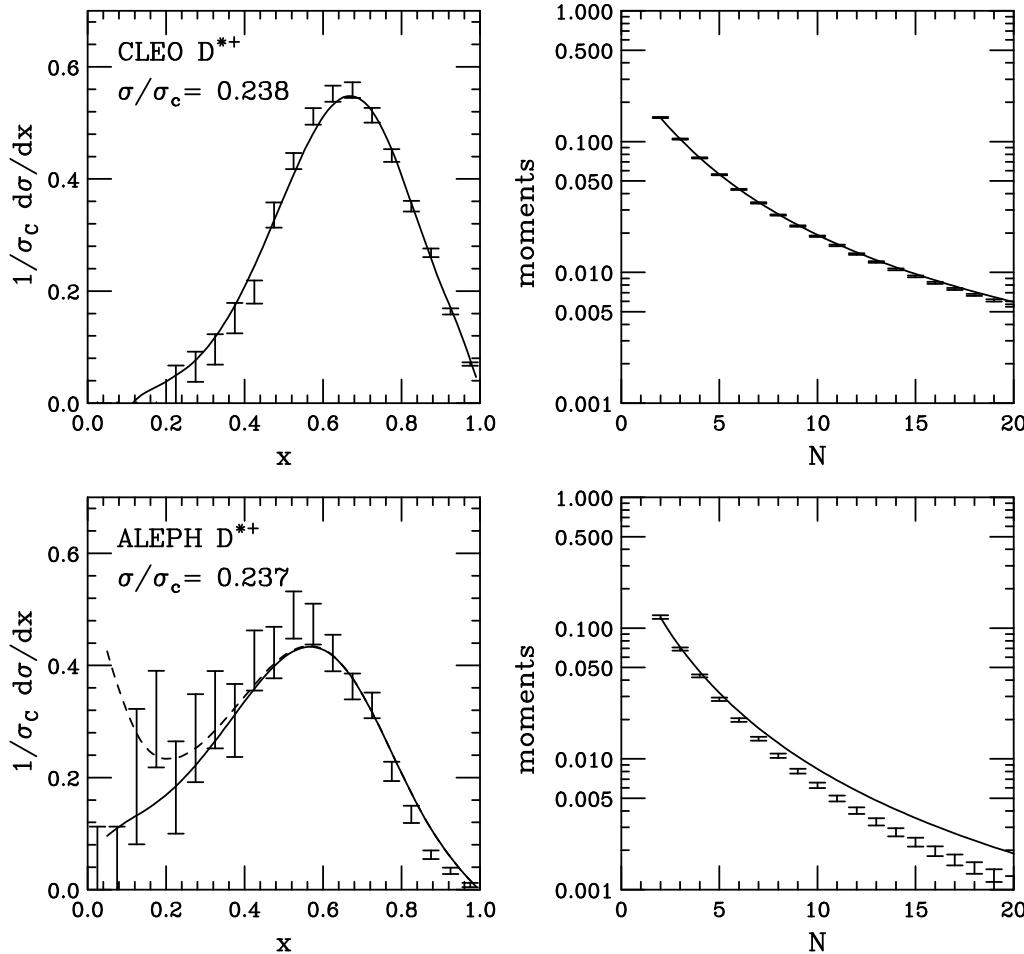
hatched histograms: data

full triangles: tuned Bowler with PARP(13)=0.59 (P_v)

c FF with NLO+NLL+Sudakov pQCD

Cacciari, Nason, Oleari

$$D_{NP}(x) \propto \frac{1}{1+c} [\delta(1-x) + c N_{a,b}^{-1} (1-x)^a x^b]$$



CLEO, Belle described with
 $a = 1.8 \pm 0.2$, $b = 11.3 \pm 0.6$, $c = 2.46 \pm 0.07$

above fit does not describe ALEPH

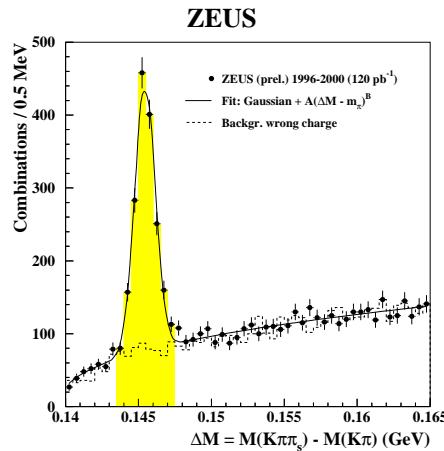
fit to ALEPH yields

$a = 2.4 \pm 1.2$, $b = 13.9 \pm 5.7$, $c = 5.9 \pm 1.7$

the difference for D^{*+} hadroproduction up to 20%

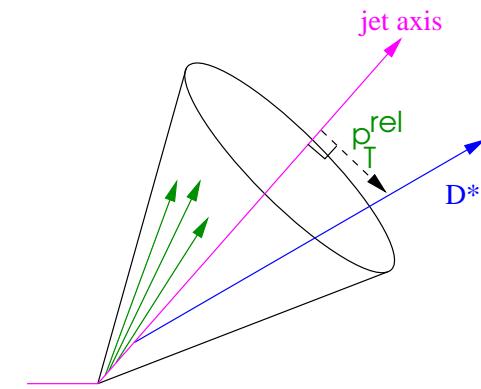
experimental info on fragmentation in hadroproduction ?

Measurement of $c \rightarrow D^{*+}$ fragmentation function



$$\mathcal{L}_{int} = 120 \text{ pb}^{-1}$$

$$N(D^{*\pm}) = 1268 \pm 52$$



In e^+e^- annihilations, $D^{*\pm}$ energy is related to $\sqrt{s}/2$. In ep ?

- 1) ZEUS: find jet containing $D^{*\pm}$ and relate the $D^{*\pm}$ energy to the energy of this jet: $Q^2 < 1 \text{ GeV}^2$, $P_T(D^{*\pm}) > 2 \text{ GeV}$, $E_T^{\text{jet}} > 9 \text{ GeV}$

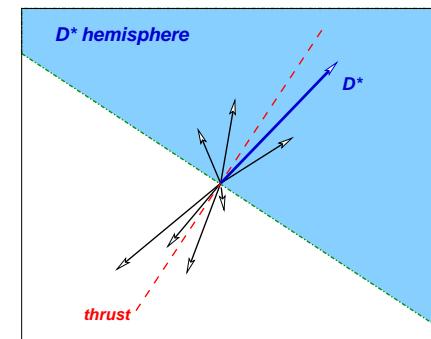
$$z = (E + p_{||})^{D^*} / (E + p_{||})^{\text{jet}} \equiv (E + p_{||})^{D^*} / 2 E^{\text{jet}}$$

- 2) H1, jet method: $Q^2 > 2 \text{ GeV}^2$, $P_T(D^{*\pm}) > 1.5 \text{ GeV}$, $E_T^{\text{jet}} > 3 \text{ GeV}$

$$z_{\text{jet}} = (E + p_{||})^{D^*} / (E + p)^{\text{jet}} \text{ in } \gamma^* p$$

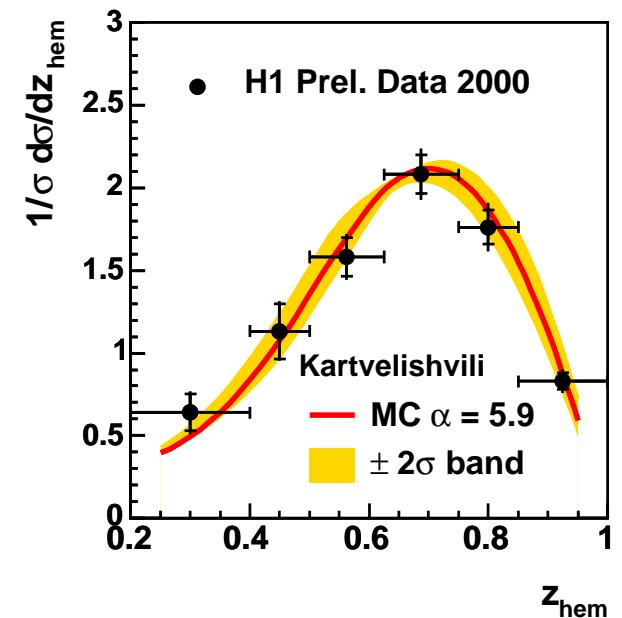
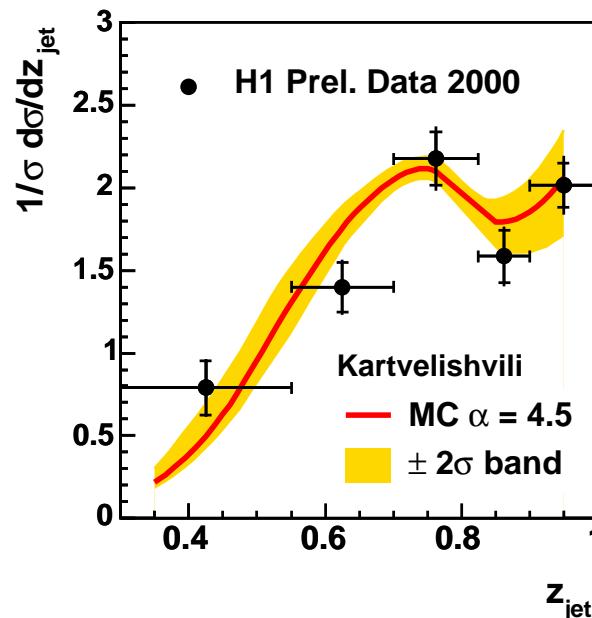
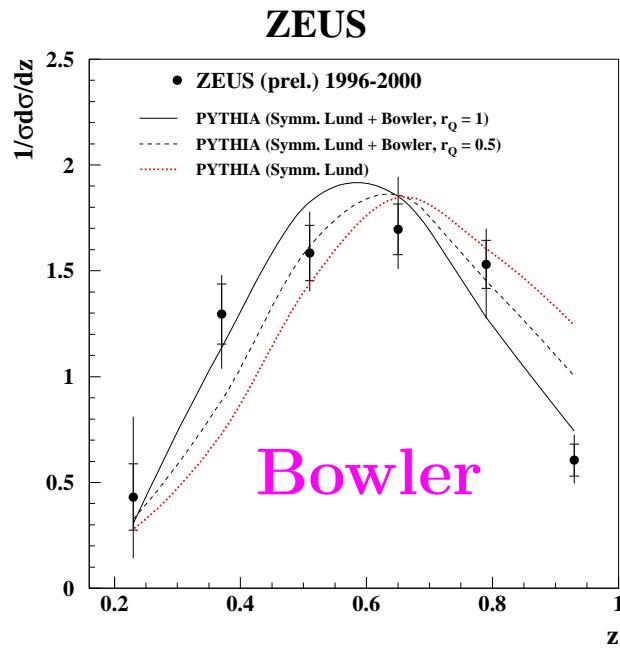
- 3) H1, hemisphere method:

$$z_{\text{hem}} = (E + p_{||})^{D^*} / \sum_{\text{hem}} (E + p) \text{ in } \gamma^* p$$



Bowler and Kartvelishvili parameterizations

Parameters are extracted using MC (PYTHIA or RAPGAP+PYTHIA),
i.e. they are optimized input parameters of the MC simulations



$$\frac{1}{z^{1+r_Q b m_Q^2}} (1-z)^a \exp\left(\frac{-bm_\perp^2}{z}\right)$$

$r_Q = 1$ (default) is preferable

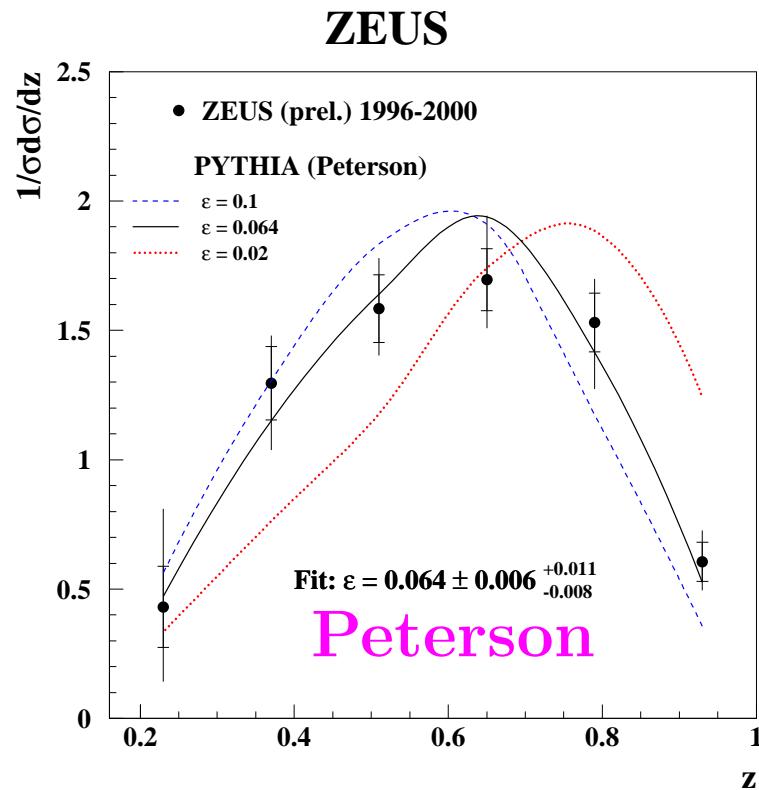
$$f(z) \propto z^\alpha (1-z)$$

$\alpha = 4.5 \pm 0.5$ (H1 jet method)

$\alpha = 5.9^{+0.9}_{-0.6}$ (H1 hem. method)

$4.0 < \alpha < 6.8$ (H1 prel.)

Peterson parameterization: $f(z) \propto \frac{1}{z(1-1/z-\epsilon/(1-z))^2}$



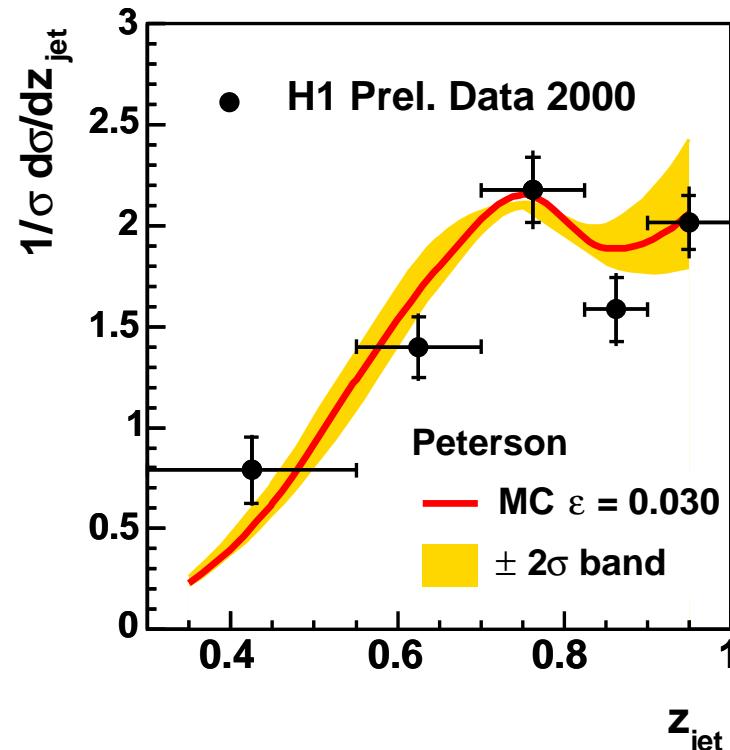
$$\epsilon = 0.064 \pm 0.006^{+0.011}_{-0.008} \text{ (ZEUS prel.)}$$

$\epsilon = 0.05$ (PYTHIA default)

$\epsilon = 0.053$ (LL fit to ARGUS data

by Nason and Oleari)

uncorrected for D^{**} decays



$$\epsilon = 0.030^{+0.006}_{-0.005} \text{ (H1 jet method)}$$

$\epsilon = 0.018 \pm 0.004$ (hem. method)

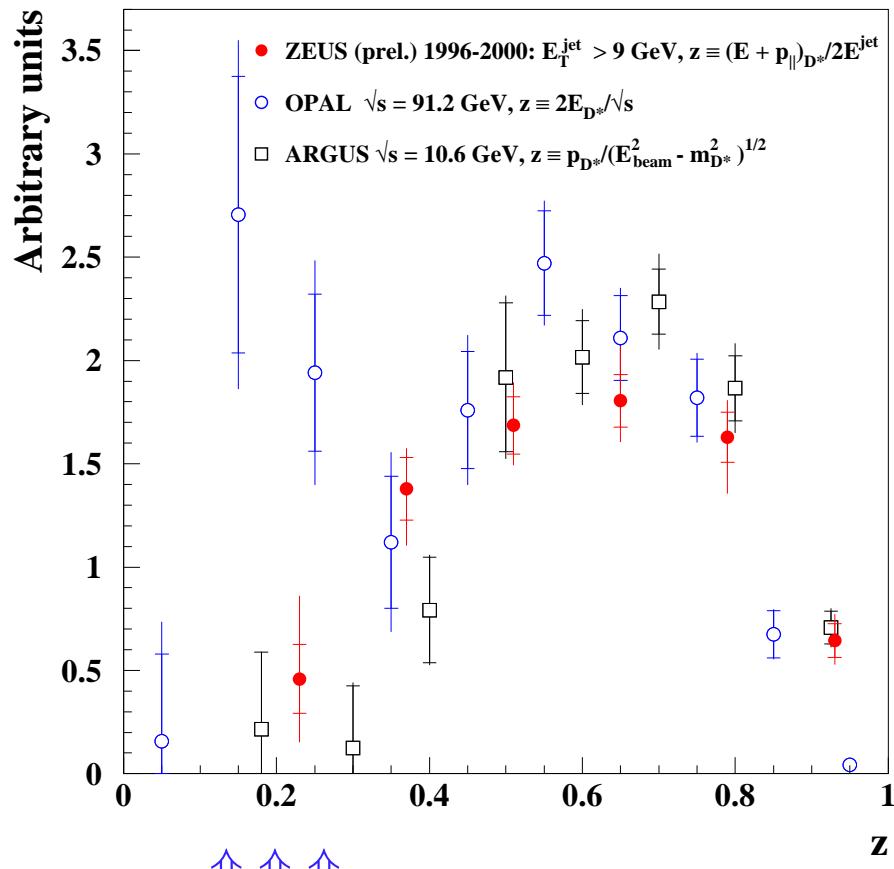
$0.014 < \epsilon < 0.036$ (H1 prel.)

corrected for D^{**} decays

NLO fits are expected

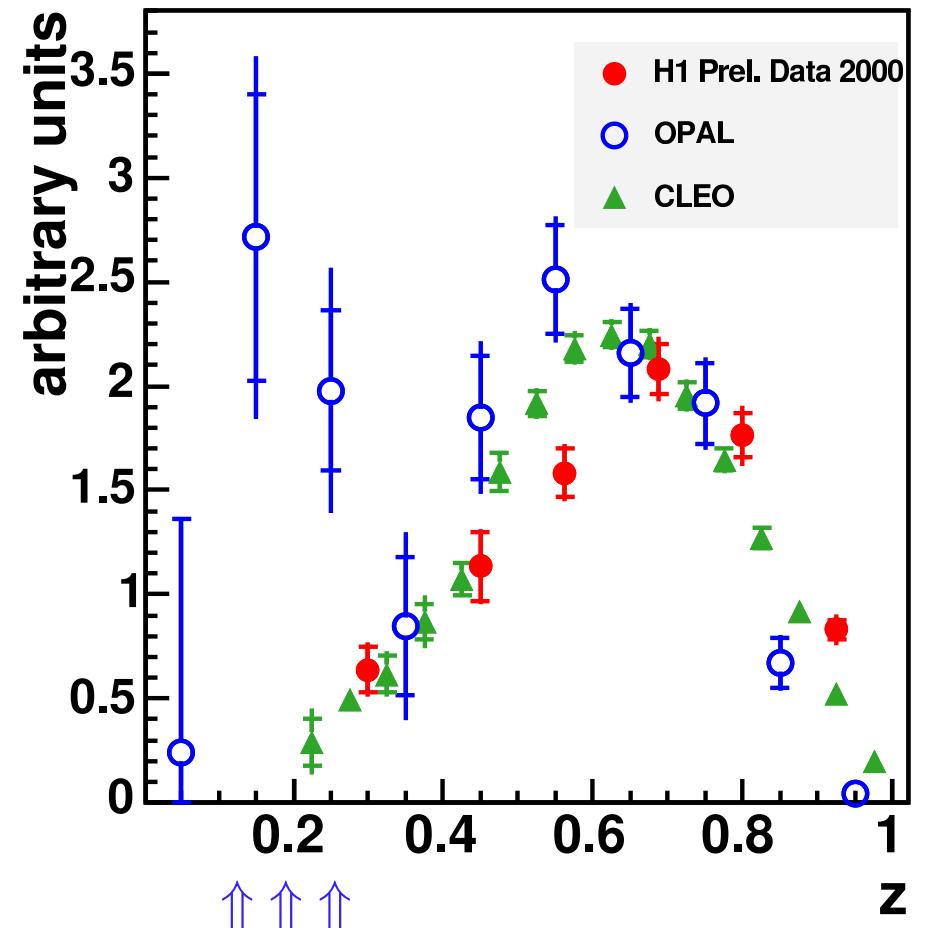
Charm fragmentation function in ep and e^+e^- collisions

ZEUS



↑↑↑

no gluon-splitting component in low-energy data



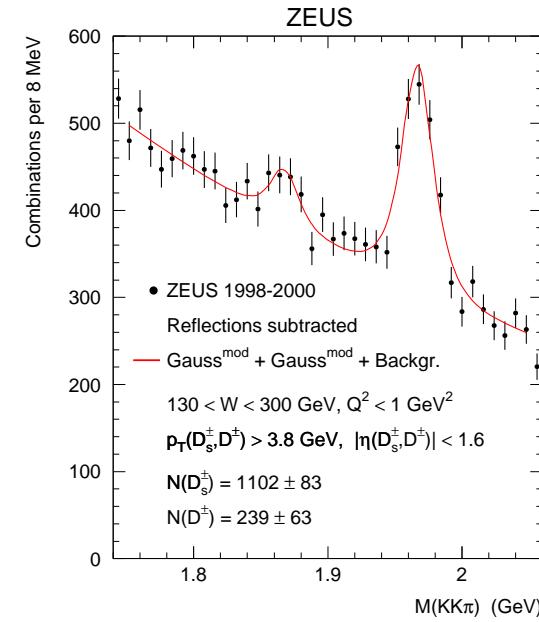
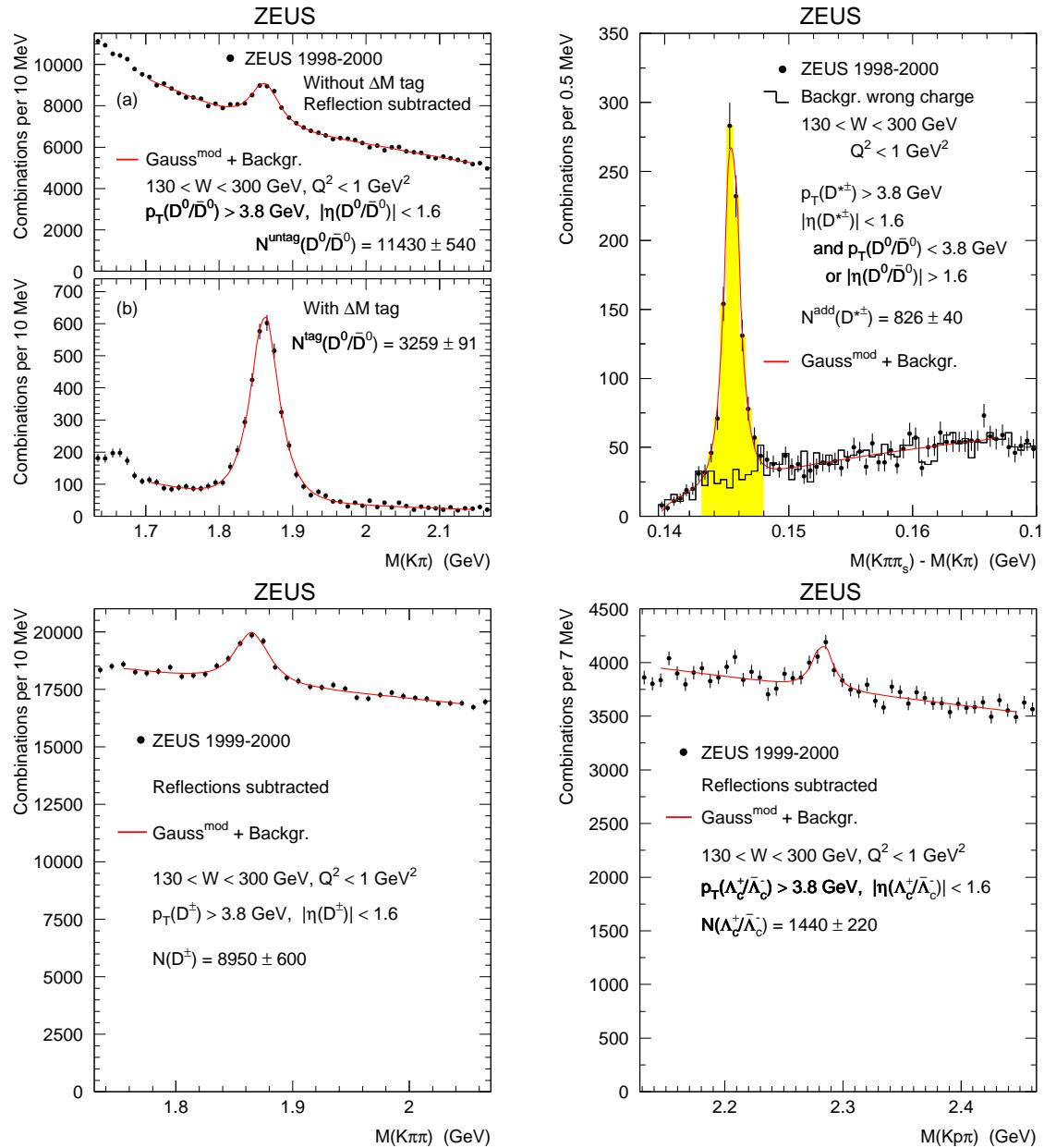
↑↑↑

different z definitions

qualitative agreement

Measurement of c -fragmentation ratios and branchings

$D^{*\pm}$ and c ground states: D^0 , D_s^\pm , D^\pm and Λ_c^\pm



Kinematic range of ZEUS γp analysis:

$P_T(D, \Lambda) > 3.8 \text{ GeV}, |\eta(D, \Lambda)| < 1.6$
 $130 < W < 280 \text{ GeV}, Q^2 < 1 \text{ GeV}^2$

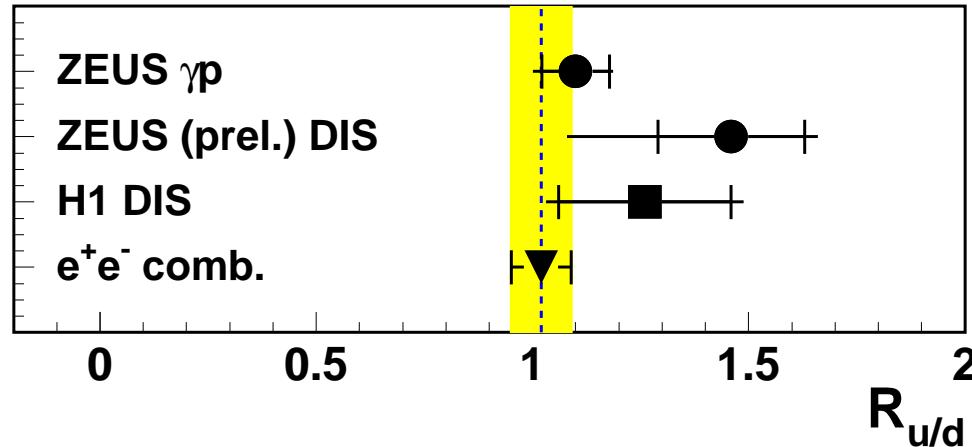
“Measured” x-sections:

$\sigma^{\text{untag}}(D^0), \sigma^{\text{tag}}(D^0), \sigma^{\text{add}}(D^{*\pm})$
 $\sigma(D_s^\pm), \sigma(D^\pm), \sigma(\Lambda_c^\pm)$

$R_{u/d}$ measurement

$$\begin{aligned}
 R_{u/d} &= \frac{c\bar{u}}{c\bar{d}} = \frac{\sigma^{dir}(D^{0,*0})}{\sigma^{dir}(D^{\pm,*\pm})} = \frac{\sigma(D^0) - \sigma(D^{*\pm}) \times BR}{\sigma(D^\pm) - \sigma(D^{*\pm}) \times (1-BR) + \sigma(D^{*\pm})} \\
 &= \frac{\sigma(D^0) - \sigma(D^{*\pm}) \times BR}{\sigma(D^\pm) + \sigma(D^{*\pm}) \times BR} = \frac{\sigma^{\text{untag}}(D^0)}{\sigma(D^\pm) + \sigma^{\text{tag}}(D^0)} , \quad BR = B_{D^{*+} \rightarrow D^0 \pi^+} = (67.7 \pm 0.5)\%
 \end{aligned}$$

$$R_{u/d} = 1.100 \pm 0.078 \text{ (stat)} {}^{+0.038}_{-0.061} \text{ (syst)} {}^{+0.047}_{-0.049} \text{ (br)} \quad (\text{ZEUS } \gamma p)$$



$$\begin{aligned}
 &= \frac{f(c \rightarrow D^0) - f(c \rightarrow D^{*+}) \times BR}{f(c \rightarrow D^+) + f(c \rightarrow D^{*+}) \times BR} \\
 &\text{for H1 and } e^+e^-
 \end{aligned}$$

consistent with isospin invariance

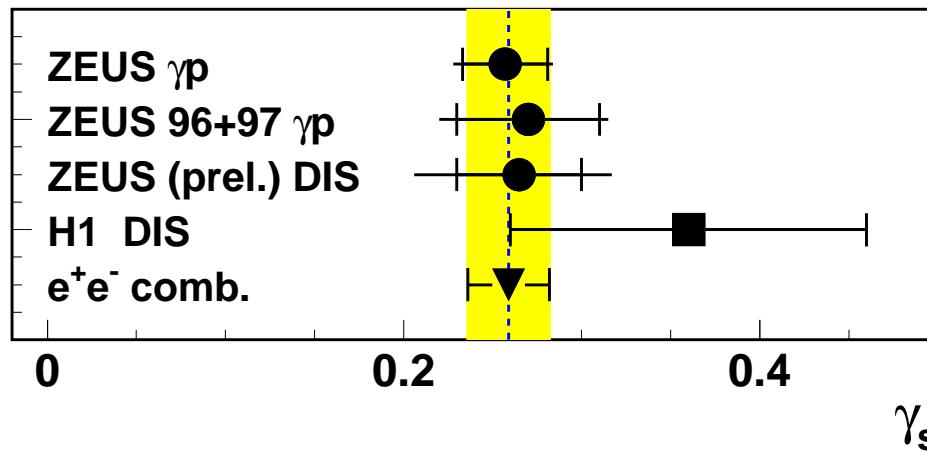
u and d quarks are produced equally in charm fragmentation

more precise measurement in DIS ?

γ_s measurement

$$\gamma_s = \frac{2 c \bar{s}}{c d + c \bar{u}} = \frac{2 \sigma(D_s^\pm)}{\sigma(D^\pm) + \sigma^{\text{untag}}(D^0) + \sigma^{\text{tag}}(D^0) + \sigma^{\text{add}}(D^{*\pm}) \cdot (1 + R_{u/d})}$$

$$\gamma_s = 0.257 \pm 0.024 \text{ (stat)}^{+0.013}_{-0.016} \text{ (syst)}^{+0.078}_{-0.049} \text{ (br)} \quad (\text{ZEUS } \gamma p)$$



$$= \frac{2 f(c \rightarrow D_s^+)}{f(c \rightarrow D^+) + f(c \rightarrow D^0)}$$

for H1 and e^+e^-

perfect agreement between measurements

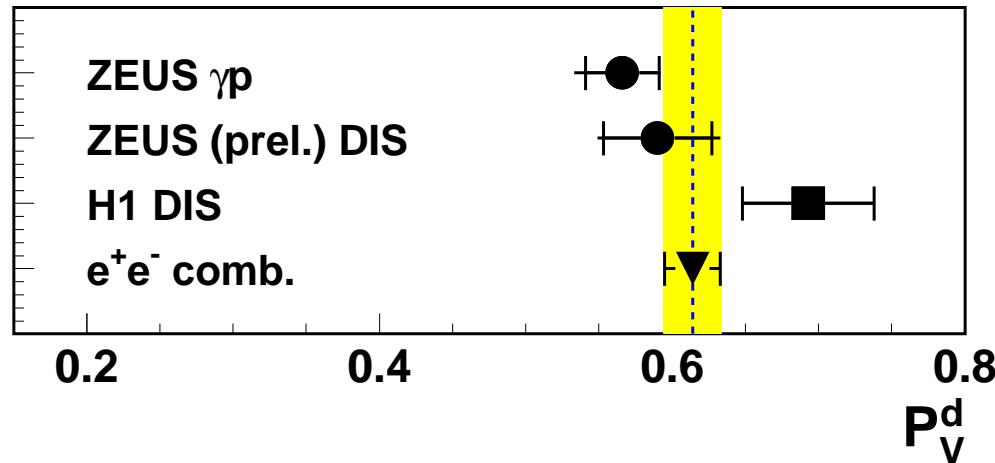
D_s production suppressed by factor ≈ 3.9 in c -fragmentation

note: excited charm-strange mesons like to decay to non-strange D mesons
 \Rightarrow Lund strangeness-suppression parameter is 10 – 30% larger
than the observable γ_s

P_v^d measurement ($P_v^d \equiv P_v$ for $c\bar{d}/\bar{c}d$ mesons)

$$P_v^d = \frac{V}{V+PS} = \frac{\sigma(D^{*\pm})}{\sigma(D^{*\pm}) + \sigma^{dir}(D^\pm)} = \frac{\sigma^{\text{tag}}(D^0)/BR + \sigma^{\text{add}}(D^{*\pm})}{\sigma(D^\pm) + \sigma^{\text{tag}}(D^0) + \sigma^{\text{add}}(D^{*\pm})}$$

$$P_v^d = 0.566 \pm 0.025 \text{ (stat)}^{+0.007}_{-0.022} \text{ (syst)}^{+0.022}_{-0.023} \text{ (br)} \quad (\text{ZEUS } \gamma p)$$



(recent precise Belle results: $P_v^d = 0.564$)

$$= \frac{f(c \rightarrow D^{*+})}{f(c \rightarrow D^+) + f(c \rightarrow D^{*+}) \cdot BR}$$

for H1 and e^+e^-

$P_v \neq 0.75 \implies$ naive spin counting does not work for charm

challenge for fragmentation models:

thermodynamics and string fragmentation predict 2/3

BKL predicts ≈ 0.6 for e^+e^- where only fragmentation diagrams contribute for ZEUS γp kinematic range, BKL prediction is ≈ 0.66

Charm fragmentation branchings, $f(c \rightarrow D, \Lambda_c) = \sigma(D, \Lambda_c)/\sigma_{\text{gs}}$

ZEUS γp

$$f(c \rightarrow D^+) = 0.217 \pm 0.014^{+0.013+0.014}_{-0.005-0.016}$$

$$f(c \rightarrow D^0) = 0.523 \pm 0.021^{+0.018+0.022}_{-0.017-0.032}$$

$$f(c \rightarrow D_s^+) = 0.095 \pm 0.008^{+0.005+0.026}_{-0.005-0.017}$$

$$f(c \rightarrow \Lambda_c^+) = 0.144 \pm 0.022^{+0.013+0.037}_{-0.022-0.025}$$

$$f(c \rightarrow D^{*+}) = 0.200 \pm 0.009^{+0.008+0.008}_{-0.006-0.012}$$

Combined e^+e^- data

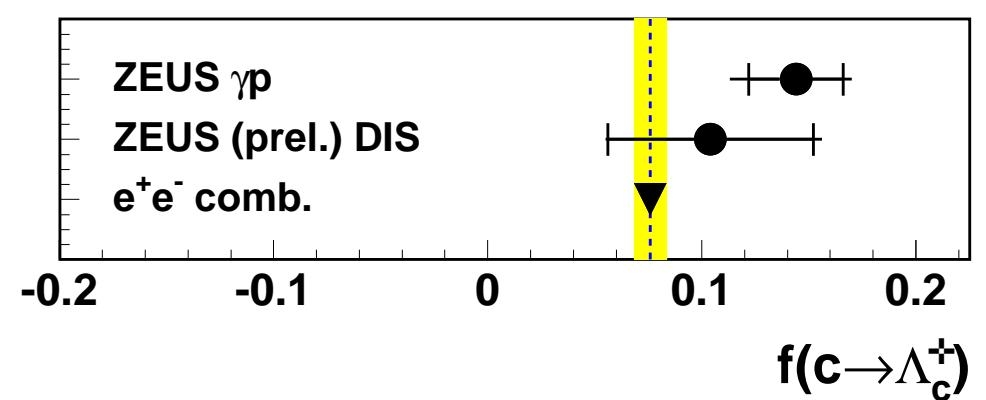
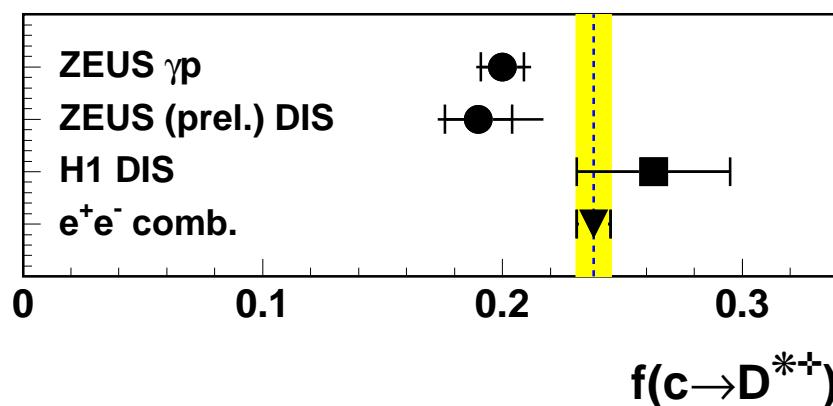
$$0.226 \pm 0.010^{+0.016}_{-0.014}$$

$$0.557 \pm 0.023^{+0.014}_{-0.013}$$

$$0.101 \pm 0.009^{+0.034}_{-0.020}$$

$$0.076 \pm 0.007^{+0.027}_{-0.016}$$

$$0.238 \pm 0.007 \pm 0.003$$



consistent with universality of charm fragmentation branchings

a half of the difference in $f(c \rightarrow D^{*+})$ is due to the difference in $f(c \rightarrow \Lambda_c^+)$

Summary

Measurements of HQ fragmentation

- test pQCD calculations
- provide non-perturbative input

LL Monte Carlo with Bowler FF is generally able to describe b/c fragmentation

fails for $x_p(D_s^+)/x_p(D^+)$ and $x_p(\Lambda_c^+)/x_p(D^+)$

NLO/NLL pQCD calculations are compatible
with high-precision FF measurements

sizeable difference between fits to CLEO+Belle and ALEPH data observed
measurements of b/c fragmentation ratios suggest

u and d quarks are produced equally in HQ fragmentation

s -quark production suppressed by factor $3 - 4$ in c -fragmentation

$P_v(b) = 0.75 \implies$ naive spin counting works for beauty

$P_v(c) \neq 0.75 \implies$ naive spin counting does not work for charm. Why ?

Measurements of charm fragmentation at HERA generally support
the hypothesis that fragmentation proceeds independently
of the hard sub-process

Estimates of extrapolation factors

factors which correct the ZEUS ratios and branchings measured in the accepted $P_T(D, \Lambda_c)$ and $\eta(D, \Lambda_c)$ region to the full phase space

	Peterson (PYTHIA)	Bowler (PYTHIA)	Cluster model (HERWIG)
$R_{u/d}$	$0.99^{+0.02}_{-0.00}$	$0.99^{+0.02}_{-0.00}$	$1.00^{+0.01}_{-0.00}$
γ_s	$1.04^{+0.04}_{-0.07}$	$1.00^{+0.05}_{-0.04}$	$1.18^{+0.07}_{-0.05}$
P_v^d	1.00 ± 0.02	$0.97^{+0.01}_{-0.00}$	$0.96^{+0.02}_{-0.01}$
$f(c \rightarrow D^+)$	$1.00^{+0.02}_{-0.01}$	$1.02^{+0.01}_{-0.02}$	$0.99^{+0.01}_{-0.03}$
$f(c \rightarrow D^0)$	0.99 ± 0.01	0.98 ± 0.01	$0.96^{+0.00}_{-0.02}$
$f(c \rightarrow D_s^+)$	$1.03^{+0.03}_{-0.06}$	$1.00^{+0.04}_{-0.03}$	$1.15^{+0.06}_{-0.05}$
$f(c \rightarrow \Lambda_c^+)$	$1.01^{+0.02}_{-0.05}$	$1.08^{+0.03}_{-0.02}$	$1.46^{+0.03}_{-0.09}$
$f(c \rightarrow D^{*+})$	$1.00^{+0.02}_{-0.03}$	$0.96^{+0.00}_{-0.02}$	$0.93^{+0.01}_{-0.02}$

large extrapolation factors are not expected

Fragmentation branchings for excited D mesons

Using world average for $f(c \rightarrow D^{*+})$:

	$f(c \rightarrow D_1^0)$ [%]	$f(c \rightarrow D_2^{*0})$ [%]	$f(c \rightarrow D_{s1}^+)$ [%]
ZEUS (prel.)	$1.46 \pm 0.18^{+0.33}_{-0.27} \pm 0.06$	$2.00 \pm 0.58^{+1.40}_{-0.48} \pm 0.41$	$1.24 \pm 0.18^{+0.08}_{-0.06} \pm 0.14$
CLEO	1.8 ± 0.3	1.9 ± 0.3	
OPAL	2.1 ± 0.8	5.2 ± 2.6	$1.6 \pm 0.4 \pm 0.3$
ALEPH	1.6 ± 0.5	4.7 ± 1.0	$0.94 \pm 0.22 \pm 0.07$
DELPHI	1.9 ± 0.4	4.7 ± 1.3	

1) the same amounts of excited D mesons in e^+e^- and ep data

2) situation with $f(c \rightarrow D_2^{*0})$ is not clear

3) $f(c \rightarrow D_{s1}^+)$ is twice as large as the expectation :

$$\gamma_s \times f(c \rightarrow D_1^0) \approx 0.3 \times 2\% = 0.6\%$$

Why $f(c \rightarrow D_{s1}^+)$ is so large ?

Is it connected with its strange helicity ?