# EXPERIMENTAL RESULTS ON HEAVY QUARK FRAGMENTATION \*

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Experimental results on c- and b-quark fragmentation are reviewed. The discussion is concentrated on measurements of heavy-quark fragmentation functions and fragmentation fractions. Measurements of various heavy-quark fragmentation ratios are also discussed. The experimental results are compared with theoretical expectations and model predictions.

## 1. Introduction

The initial stage of charm/bottom quark fragmentation can be described by perturbative QCD (pQCD) calculations <sup>1</sup>. A non-perturbative (NP) parameterisation is needed to describe the final heavy-quark transformation to a particular charmed or bottom hadron. Such parameterisation can include effects producing by the excited states decaying to a given hadron. The NP fragmentation parameterisation can be splited in two parts: fragmentation function and fragmentation fraction. Fragmentation functions are used to parameterise the transfer of the quark's energy to a given meson; they can be different for different pQCD calculations used to describe the initial fragmentation. Fragmentation fractions are the fractions of c/bquarks hadronising as a particular charmed/bottom hadron; they are expected to be universal for all pQCD calculations.

Measurements of the heavy quark fragmentation allow testing pQCD calculations and extracting fragmentation functions and fractions. A deeper phenomenological understanding of the heavy quark fragmentation can be obtained by measuring various heavy-quark fragmentation ratios. In partic-

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ular, we will discuss the ratio of neutral and charged D/B meson production rates,  $R_{u/d}$ , the strangeness-suppression factor,  $\gamma_s$ , and the fraction of D/Bmesons produced in a vector state,  $P_v$ .

## 2. Bottom quark fragmentation

The *b*-quark fragmentation function was measured at LEP  $^{2,3,4}$  and SLD <sup>5</sup>. The measured spectra were compared with predictions of the leading-logarithmic (LL) JETSET 7.4 <sup>6</sup> Monte Carlo (MC) using different parametrisations for the fragmentation function. The best description of the data with a parametrisation with one free parameter was obtained using the parametrisation of Kartvelishvili et al. <sup>7</sup>. The Bowler <sup>8</sup> and symmetric LUND <sup>9</sup> parametrisations with two free parameters provided a better data description. The Peterson <sup>10</sup> and Collins-Spiller <sup>11</sup> parameterisations, and the HERWIG cluster model <sup>12</sup> predictions were found to be too broad to describe the data. The *b*-quark fragmentation function measurements were also used for fitting the NP parametrisation with the next-to-leading-order (NLO) calculations <sup>13,14</sup>.

The *b*-quark fragmentation fractions were obtained by combining of all published LEP and CDF results on production of the weakly decaying *B* hadrons with measurements of the time-integrated mixing probabilities <sup>15,16</sup>. The isospin invariance, i.e.  $R_{u/d} = 1$ , was assumed in this procedure. Using the measured fragmentation fractions, the strangenesssuppression factor for bottom mesons is

$$\gamma_s = 2f(\bar{b} \to B_s^0) / [f(\bar{b} \to B^0) + f(\bar{b} \to B^+)] = 0.27 \pm 0.03.$$

Thus, bottom-strange meson production is suppressed by a factor  $\approx 3.7$ . The combined LEP value for the fraction of *B* mesons, produced in a vector state, is  $P_{\rm v} = 0.75 \pm 0.04^{-16}$ , that is in perfect agreement with the naive spin counting expectation (0.75).

#### 3. Charm quark fragmentation

The *c*-quark fragmentation function has been recently measured with high precision by the CLEO  $^{17}$  and BELLE  $^{18}$  collaborations. The data comparison with the JETSET MC predictions revealed the same picture as for the *b*-quark fragmentation. The best description of the data was obtained using the Bowler parametrisation with two free parameters, and the parametrisation of Kartvelishvili et al. with one free parameter.

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A discrepancy between the NP parametrisations obtained with the CLEO/BELLE data and earlier ALEPH measurement <sup>19</sup> has been observed using the NLO initial conditions, next-to-leading logarithmic (NLL) evolution, NLO coefficient functions and NLL Sudakov resummation <sup>1,13</sup>. The difference, which was attributed to the evolution between the  $\Upsilon(4S)$  and  $Z^0$  energies, results in an additional uncertainty in predictions for  $D^{*\pm}$  hadroproduction of the order 20%. To reduce the uncertainty direct measurements of the charm fragmentation function at hadronic machines would be useful. Such measurements were already performed in *ep* interactions at HERA by the ZEUS <sup>20</sup> and H1 <sup>21</sup> collaborations; their results were found to be in qualitative agreement with those obtained in  $e^+e^-$  annihilations.

Table 1. The fractions of c quarks hadronising as a particular charm hadron,  $f(c \to D, \Lambda_c)$ . The fractions are shown for the  $D^+$ ,  $D^0$ ,  $D_s^+$  and  $\Lambda_c^+$  charm ground states and for the  $D^{*+}$  state.

	ZEUS $(\gamma p)^{22}$	Combined $e^+e^-$ data <sup>24</sup>	H1 (DIS) $^{23}$
	stat. syst. br.	stat. $\oplus$ syst. br.	total
$f(c \to D^+)$	$0.217 \pm 0.014 \begin{array}{c} +0.013 + 0.014 \\ -0.005 - 0.016 \end{array}$	$0.226 \pm 0.010 \ ^{+0.016}_{-0.014}$	$0.203 \pm 0.026$
$f(c \to D^0)$	$0.523 \pm 0.021 \begin{array}{c} +0.018 + 0.022 \\ -0.017 - 0.032 \end{array}$	$0.557 \pm 0.023 \ \substack{+0.014 \\ -0.013}$	$0.560 \pm 0.046$
$f(c \to D_s^+)$	$0.095 \pm 0.008 \ {}^{+0.005 \ +0.026}_{-0.005 \ -0.017}$	$0.101 \ \pm 0.009 \ \ {}^{+0.034}_{-0.020}$	$0.151 \pm 0.055$
$f(c \to \Lambda_c^+)$	$0.144 \pm 0.022 \begin{array}{c} +0.013 + 0.037 \\ -0.022 - 0.025 \end{array}$	$0.076 \pm 0.007 \stackrel{+0.027}{_{-0.016}}$	
$f(c \to D^{*+})$	$0.200 \pm 0.009 \begin{array}{c} +0.008 + 0.008 \\ -0.006 - 0.012 \end{array}$	$0.238 \pm 0.007 \ ^{+0.003}_{-0.003}$	$0.263 \pm 0.032$

Table 1 compares the *c*-quark fragmentation fractions measured in *ep* interactions at HERA by the ZEUS <sup>22</sup> and H1 <sup>23</sup> collaborations with those obtained in  $e^+e^-$  annihilations. The latter values were compiled previously <sup>24</sup> and updated with the recent branching ratio values <sup>16</sup>. The measurements performed in  $e^+e^-$  and *ep* interactions are consistent. Measurements of the  $R_{u/d}$  value in charm fragmentaion confirmed isospin invariance <sup>22,23,24,25</sup>. Measurements of the strangeness-suppression factor in charm fragmentation showed that charmed-strange meson production is suppressed by a factor  $\approx 3.9$  (similar to the suppression in bottom fragmentation). The fraction of charged *D* mesons produced in a vector state,  $P_v^d$ , was found to be  $\approx 0.6^{22,23,24,25}$  in both  $e^+e^-$  and *ep* interactions. The value is significantly smaller than that obtained in bottom fragmentation and does not agree with the naive spin counting expectation (0.75).

#### 4. Summary

The *b*-qaurk fragmentation function and fractions were measured in  $e^+e^$ annihilations, while the *c*-quark fragmentation was studied in both  $e^+e^-$  and 4

ep interactions. Comparison of the charm fragmentation characteristics, obtained in  $e^+e^-$  and ep interactions, generally supports the hypothesis that fragmentation proceeds independently of the hard sub-process.

The fraction of charged D mesons produced in a vector state,  $P_v^d$ , in charm fragmentation was found to be  $\approx 0.6$  in both  $e^+e^-$  and ep interactions. The value is significantly smaller than that obtained in bottom fragmentation and does not agree with the naive spin counting expectation (0.75).

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