

# Direct measurement of the longitudinal structure functions at HERA

New Trends in HERA Physics 2005

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# Outlook

- Motivation
- Feasibility of a direct measurement
- Undirect determination
- Additional benefits from running at low  $E_p$
- A possible strategy and conclusion

$$\frac{d^2\sigma}{dx dQ^2} = \frac{2\pi\alpha^2 Y_+}{Q^4 x} \left[ F_2(x, Q^2) - f(y) \cdot F_L(x, Q^2) \right]$$

with  $f(y) = \frac{y^2}{Y_+}$  and  $Y_+ = \left[ 1 + (1 - y)^2 \right]$

The structure function  $F_L$  is a basic Structure Function which has to be measured.

The measurement is difficult and HERA experiments are probably the best experiments which have ever was to perform it.

It would be a text book measurement.

# Physics motivations

(2/3)

$F_L$  simply related to  $\sigma_L$ , the inclusive cross section of longitudinally polarised photons :

$$F_L = \left( \frac{Q^2}{4\pi^2 \alpha} \right) \sigma_L$$

In QPM:  $\sigma_L = 0$

$F_L$  gets its value from perturbative QCD,

In LO :

$$F_L = \frac{\alpha_s}{4\pi} x^2 \int_x^1 \frac{dz}{z^3} \left[ \frac{16}{3} F_2 + 8 \sum e_q^2 \left(1 - \frac{x}{z}\right) z g \right]$$

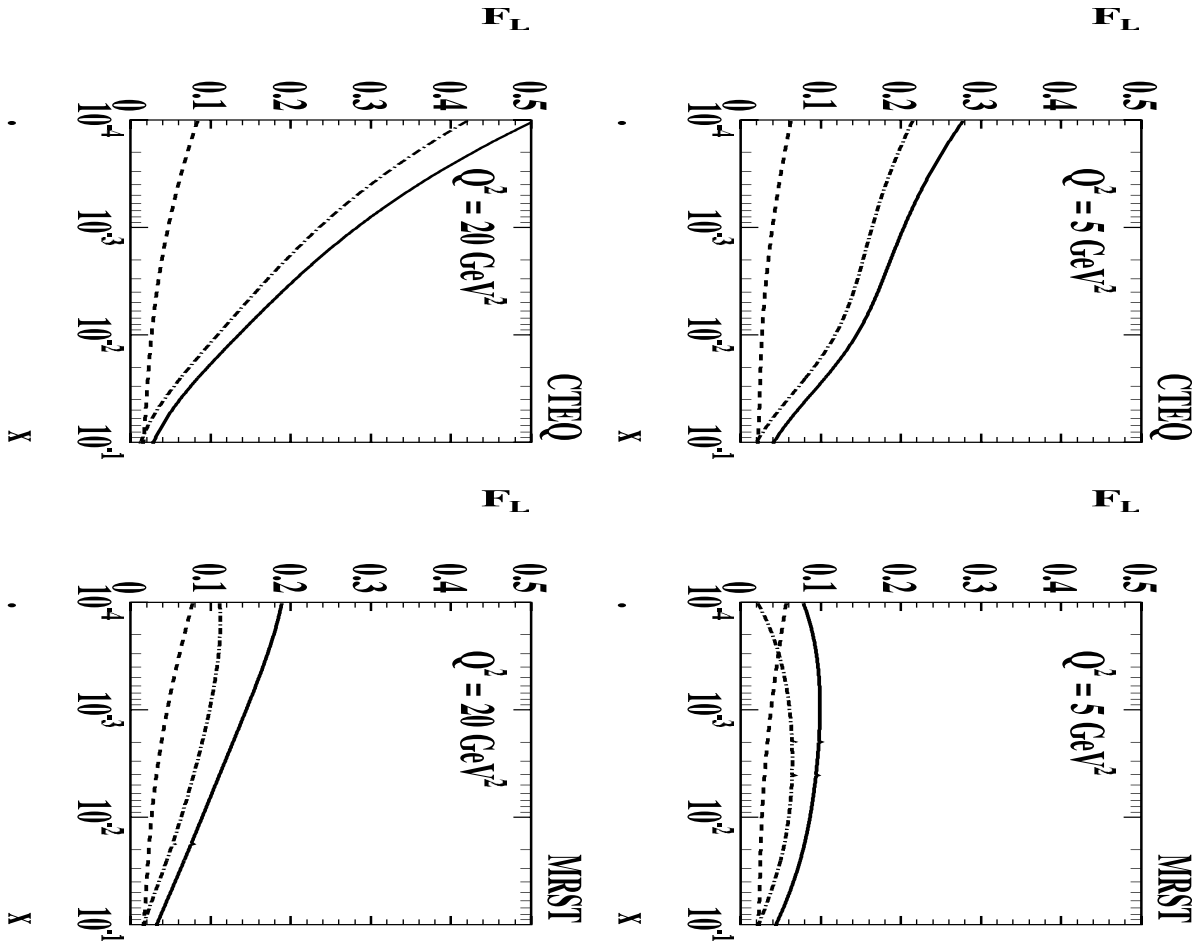
At low  $x$ , the gluon density dominates.  $F_L$  is a clean probe of the gluon distribution .

Measurement of  $F_L$  would be an important input to QCD fits of parton distributions and  $\alpha_s$  (cf R. Thorne).

# Motivations

(3/3)

CTEQ and MRST do fit  $F_2(x, Q^2)$  data from H1 at low  $x$  but  $F_L$  is poorly constrained by present data.



Measuring  $F_L$  at  $x$  from  $10^{-4}$  to a few  $10^{-3}$  would provide an additional constraint on gluon density for Higgs and W production at LHC

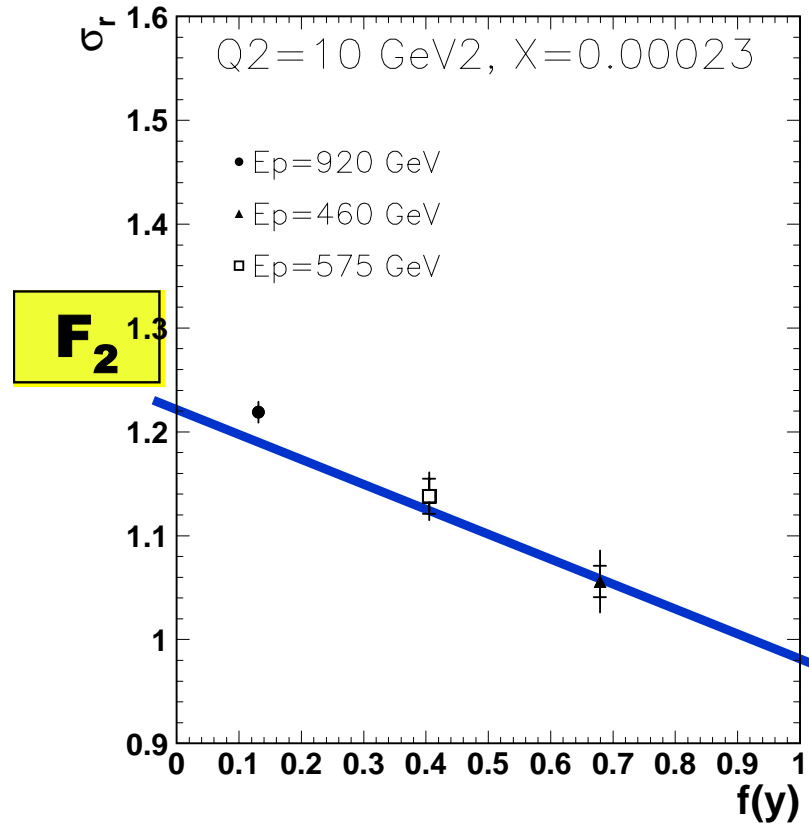
# Direct measurement of $F_L(Q^2, x)$

$$\frac{d^2\sigma}{dx dQ^2} = \frac{2\pi\alpha^2 Y_+}{Q^4 x} [F_2(x, Q^2) - f(y) \cdot F_L(x, Q^2)] = \frac{2\pi\alpha^2 Y_+}{Q^4 x} \sigma_r$$

Measure at the same  $(Q^2, x)$  reduced cross sections from different beam energies, i.e. different  $y$ .

Perform straight line fit of  $\sigma_r$  to extract  $F_2$  and  $F_L$

$F_L$  is very sensitive to small relative shifts on cross sections



# Direct measurement of $F_L(Q^2, x)$

Requirements :

- At least two beam settings which overlap in the  $(Q^2, x)$  plane
- A large  $y$  difference
- The highest possible  $y$  at low beam energies (error on  $F_L \sim 1/y^2$ )
- Enough luminosity

To have both measurements in the same part of the apparatus,

To access the largest  $y$ ,

Better to only reduce the proton beam energy

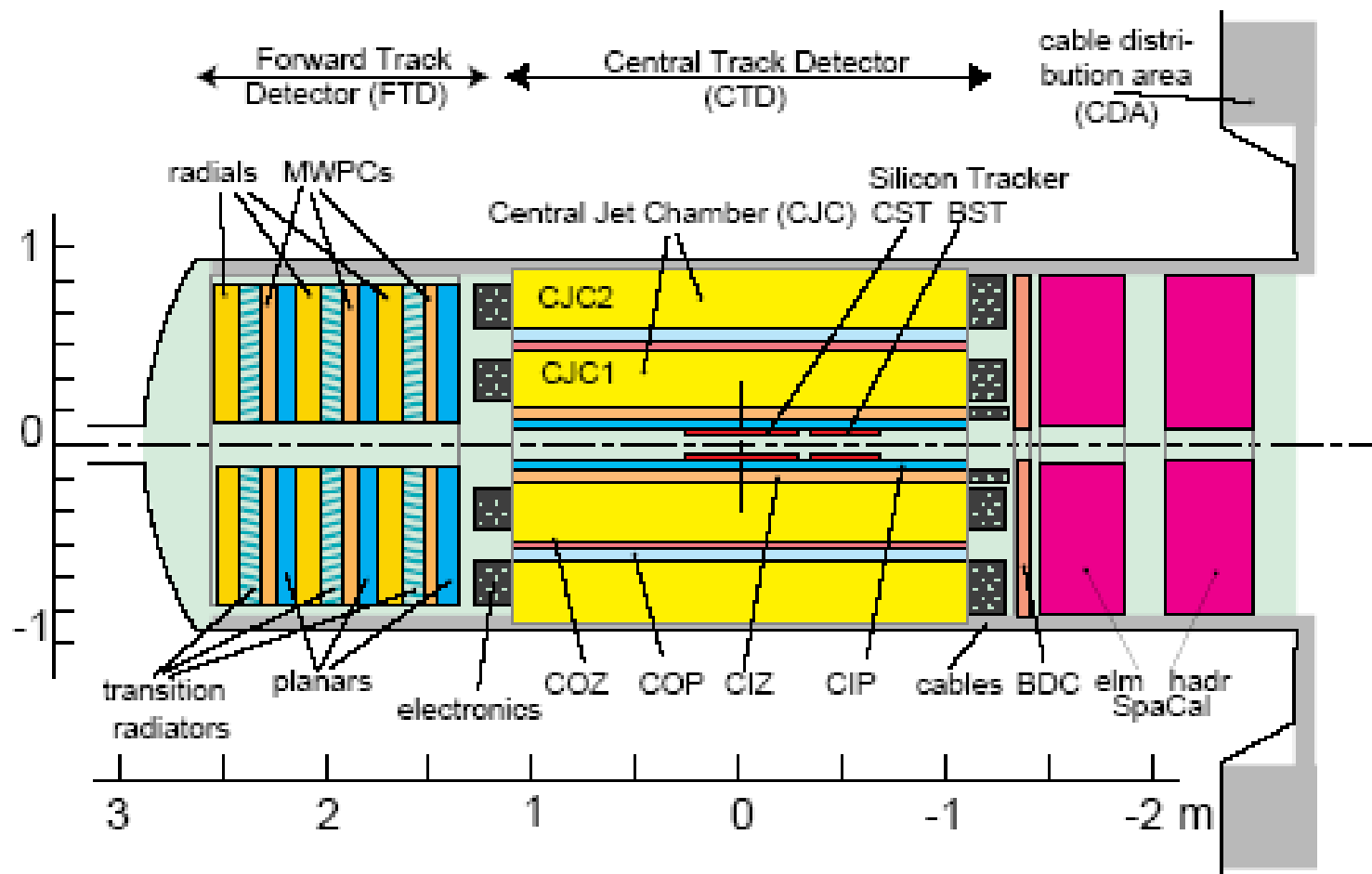
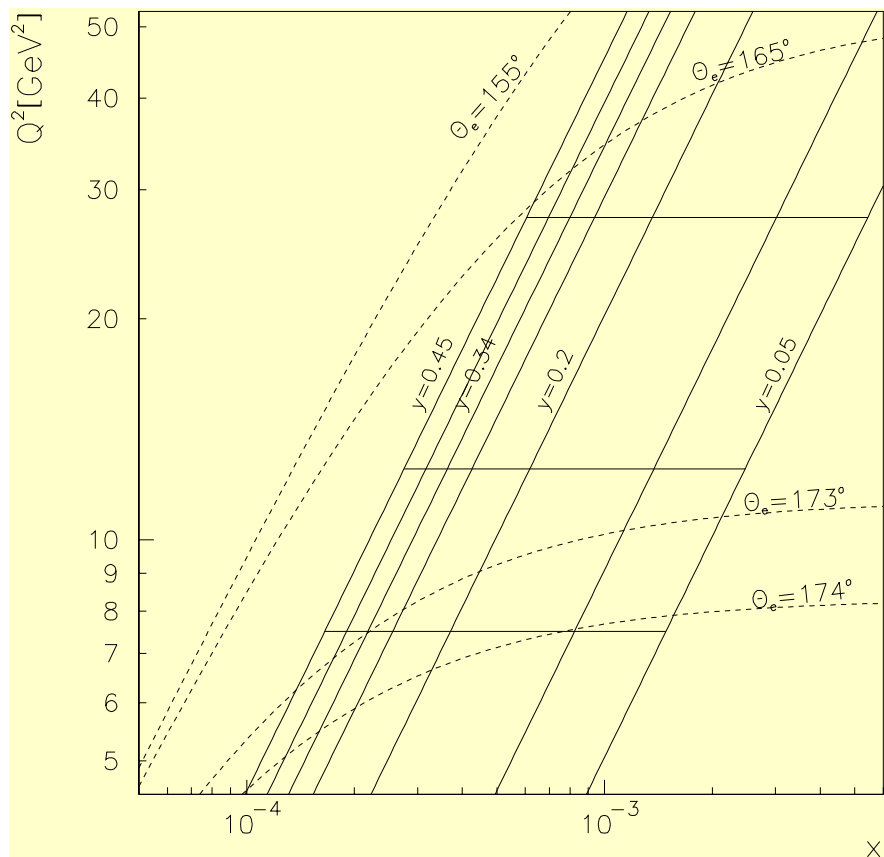


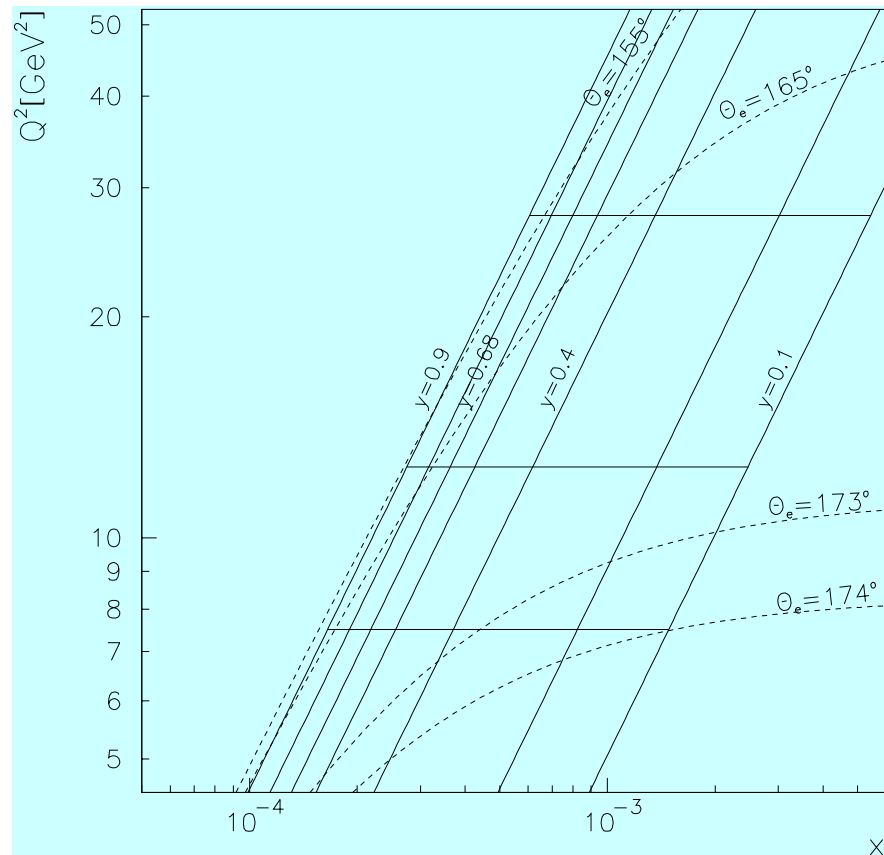
Figure 3.3: The H1 tracking detectors.



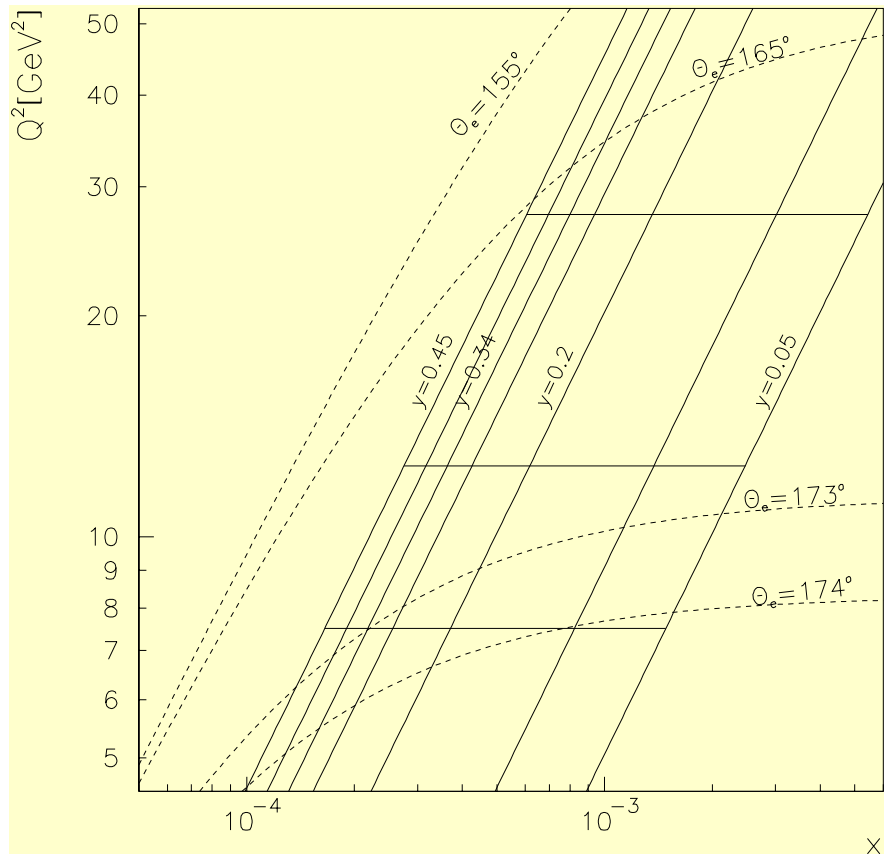
$E_p = 920 \text{ GeV}$



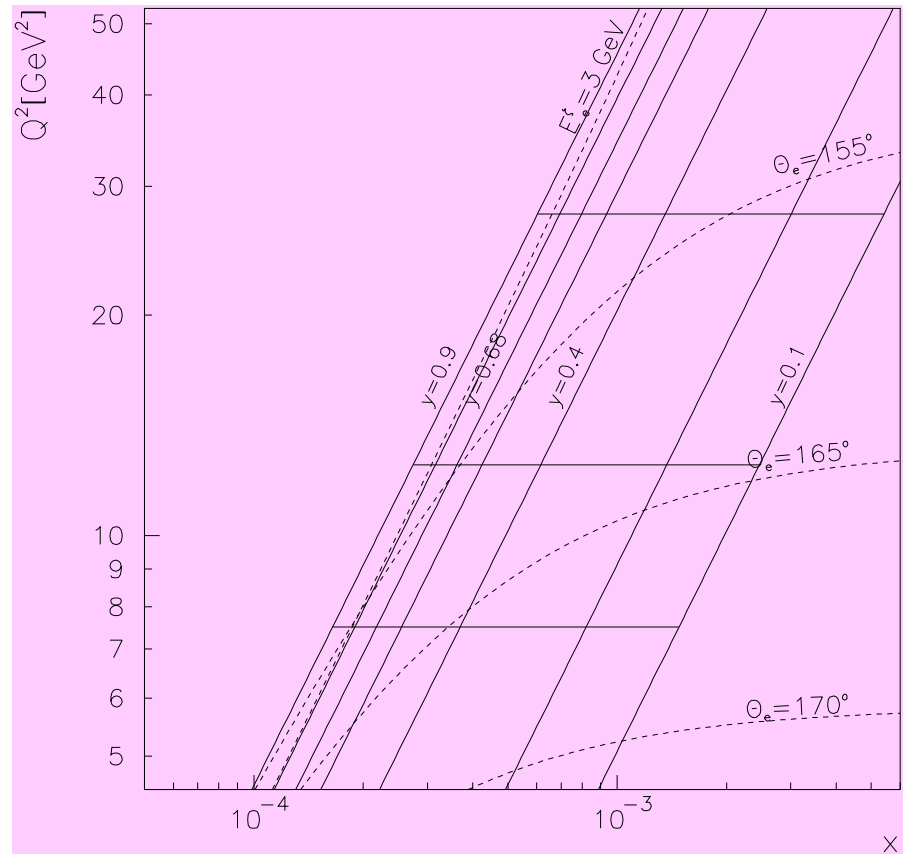
$E_p = 460 \text{ GeV}$



$E_e = 27.6 \text{ GeV}$



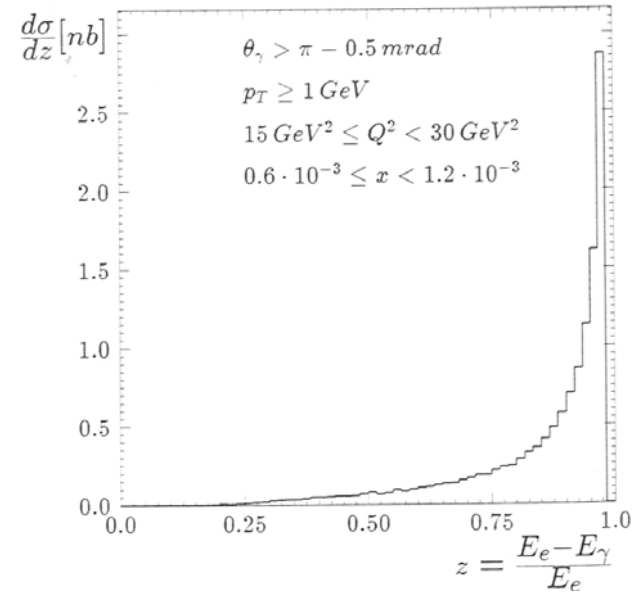
$E_e = 13.8 \text{ GeV}$



# Why not using the radiative events?

$$e + p \rightarrow e + \gamma + X$$

- We get radiative events for free !
  - Need a huge statistics ( $\sim 200 \text{ pb}^{-1}$ )  
[Krasny, Placzek, Spiesberger, 1991]
  - However :
    - For a fixed ( $Q^2, x$ ) bin, at different  $y$ , the overlap in the same part of the detector is quite small
    - No access to very high  $y$  ( $E'_e > 3 \text{ GeV}$ )
    - Severe pile up of Bethe-Heitler events ( $e p \rightarrow e p \gamma$ ) in the gamma detector [Favart, Maracek, 1996]
- ➔ huge errors on  $F_L \sim 50\text{-}100 \%$ !



# Direct measurement of $F_L(Q^2, x)$

Present assumptions (to be tuned when Luminosity is better known) :

$E_p(\text{GeV})$	920	460
$L(\text{pb}^{-1})$	30	10

A possible option :

$E_p(\text{GeV})$	920	460	575
$L(\text{pb}^{-1})$	30	5	3.5

# Systematic errors

$F_L$  is only sensitive to relative shifts of cross section. Errors based on cumulated expertise in analysing  $F_2$  data and anticipation of BST performances.

- Correlated errors

- Energy of scattered electron : from 2% at 3 GeV to 0.2% at 30 GeV.
- Angle of scattered electron : 0.2 mrad in BST and 1 mrad at  $\Theta_e < 165^\circ$ .
- Residual photoproduction background (from a fit on negative tracks in positron run)  
 $0.267 - 0.8 y + 0.6 y^2$  at  $y > 0.65$

- Uncorrelated efficiencies:

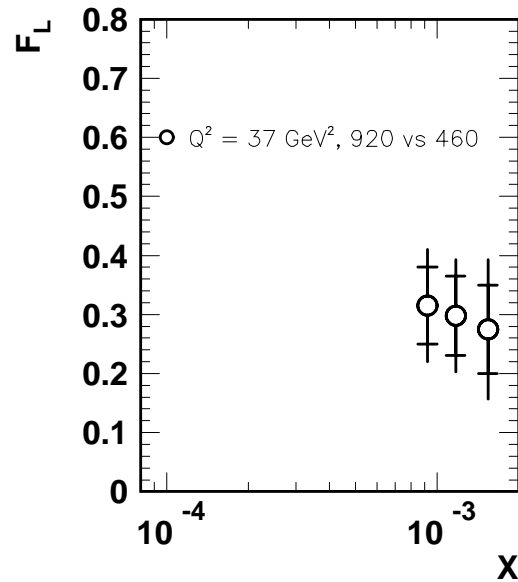
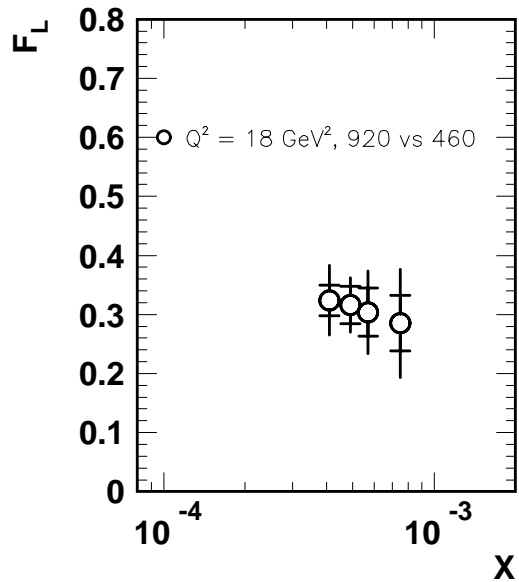
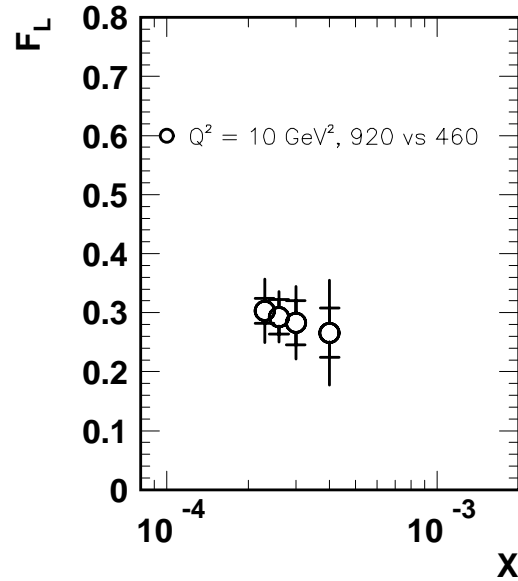
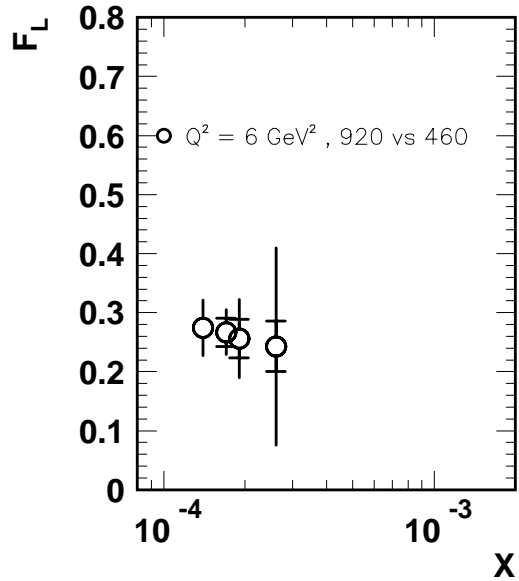
electron identification, trigger, vertex, radiative corrections : 1%

# Estimates of errors on $F_L(x, Q^2)$

- Fast simulations based on  $F_2$  and  $F_L$  parametrizations from H1 QCD fits (2000)
  - Two methods have been used :
    - Analytic calculations (MK)
    - Fast montecarlo (JF)
- Excellent agreement

ERRORS IN PERCENT,  $Q^2 \sim 9.4 \text{ GeV}^2$ , PROTON BEAM ENERGIES 920 vs 460 GeV

<X>	0.00023	0.00026	0.00030	0.00040
<Y> at 460 GeV	0.835	0.728	0.628	0.483
<FL>	0.303	0.293	0.283	0.266
STATISTICAL ERROR	7.0	9.9	13.1	15.8
SYST: EFFICIENCIES	6.4	9.3	13.8	28.0
SYST: GAMMA-P	13.0	4.6	1.0	0.0
SYST: ELEC. ENERGY	4.7	2.9	9.3	4.7
SYST: ELEC. ANGLE	6.2	0.3	3.2	7.5
SYST: TOTAL	16.4	10.8	17.0	29.4
SYST+STAT	17.9	14.6	21.4	33.4

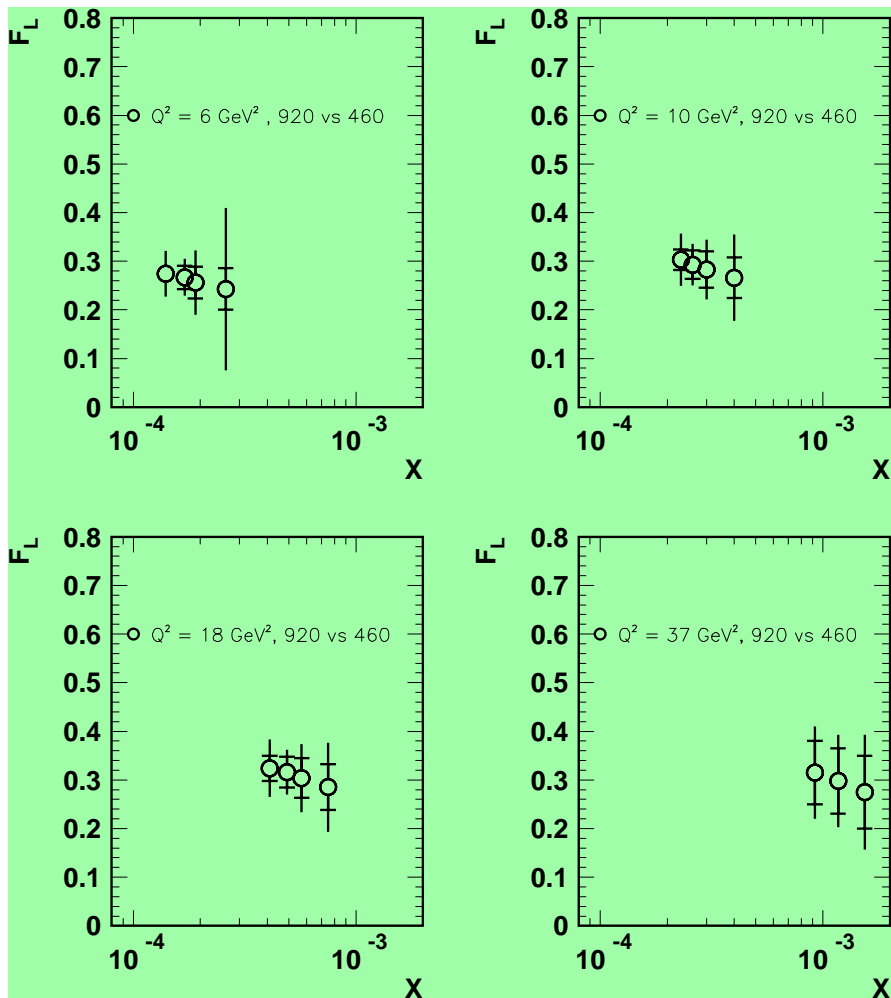


30 pb<sup>-1</sup>, E<sub>p</sub>=920 GeV

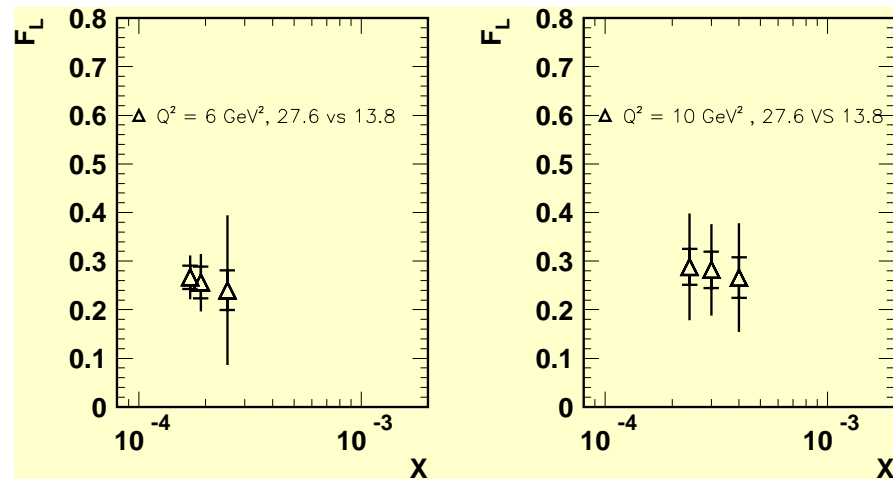
10 pb<sup>-1</sup>, E<sub>p</sub>=460 GeV

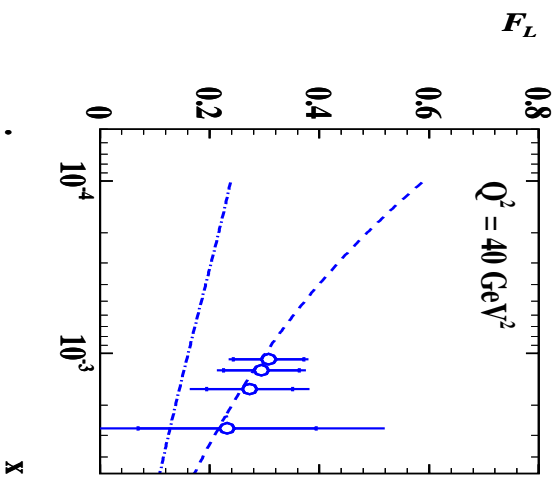
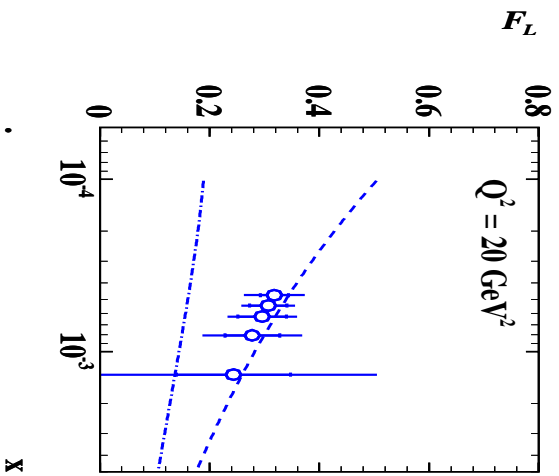
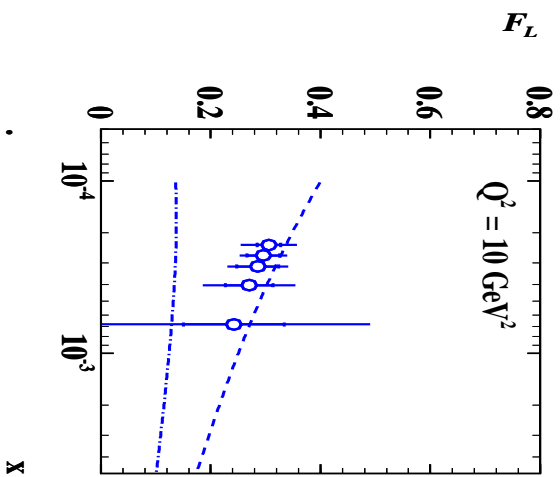
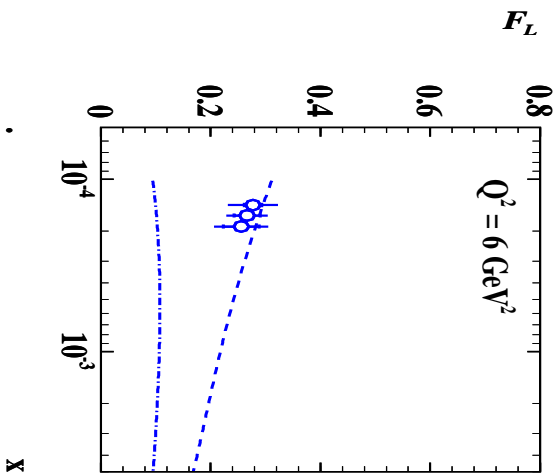


# 920 x 27.5 vs 460 x 27.5



# 920 x 27.5 vs 920 x 13.8

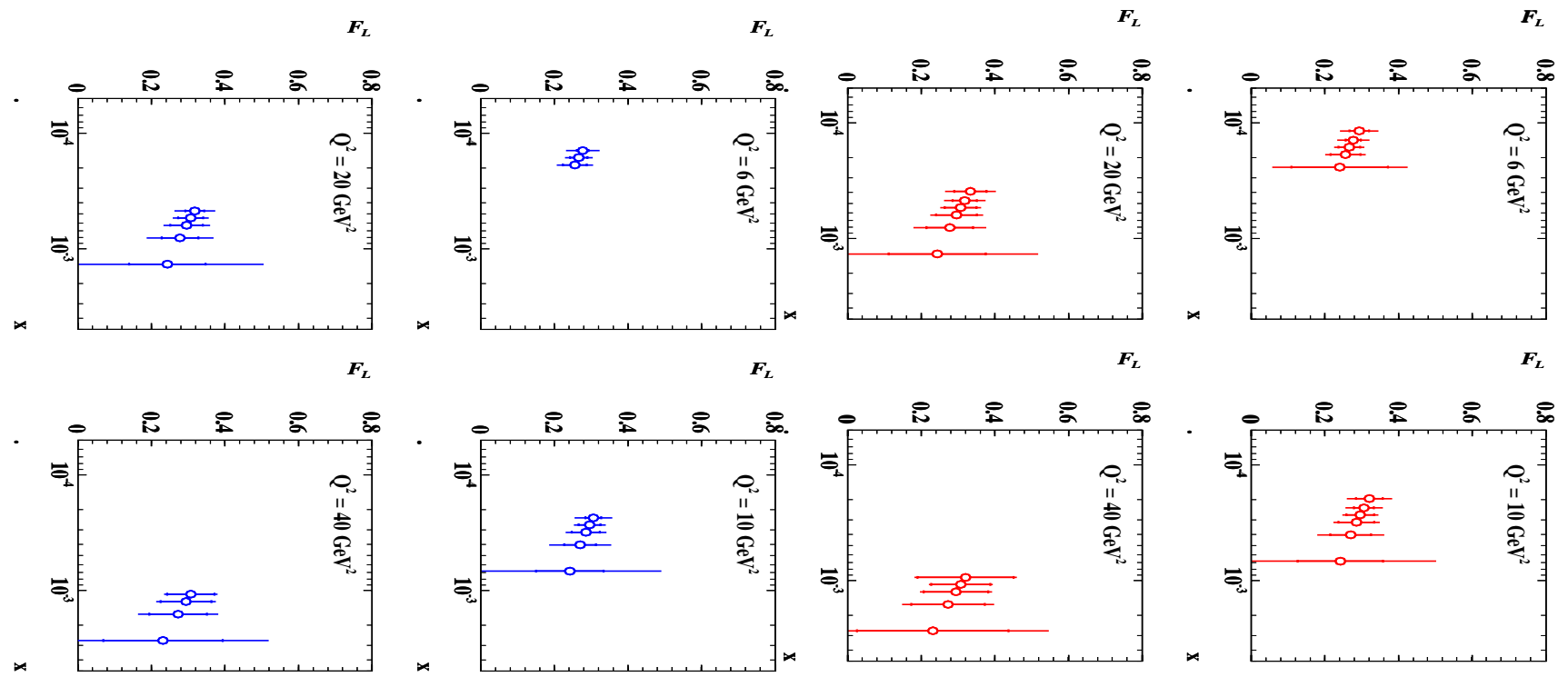




30 pb<sup>-1</sup>, E<sub>p</sub>=920 GeV  
 10 pb<sup>-1</sup>, E<sub>p</sub>=460 GeV

10 pb<sup>-1</sup> at E<sub>p</sub>=460GeV

5 pb<sup>-1</sup> at E<sub>p</sub> = 460 GeV and  
3.5 pb<sup>-1</sup> at E<sub>p</sub> = 575 GeV



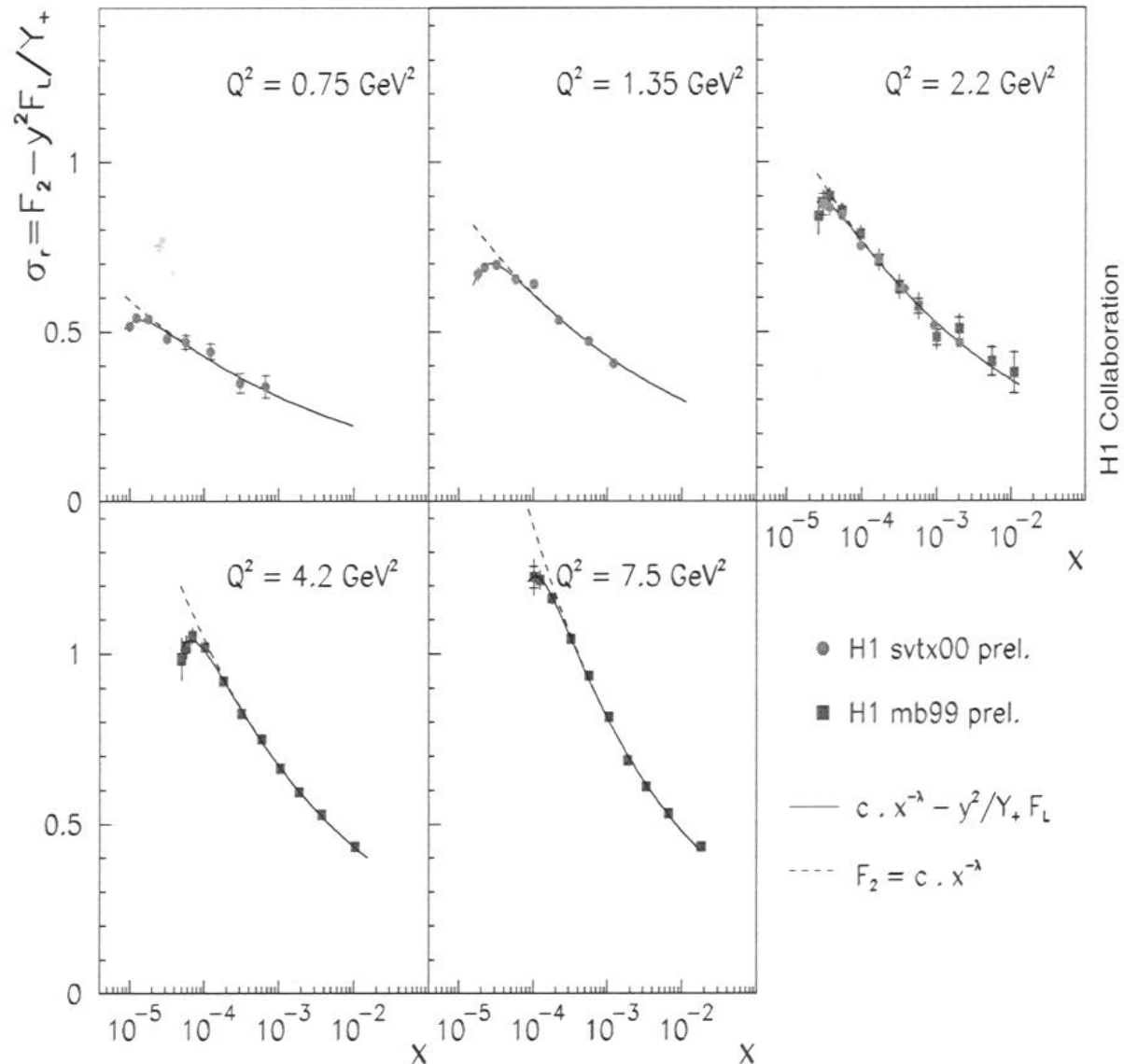
# Undirect determination of $F_L$ ?

H1 has invented three methods to determine  $F_L$  (not a measurement) at fixed beam energies :

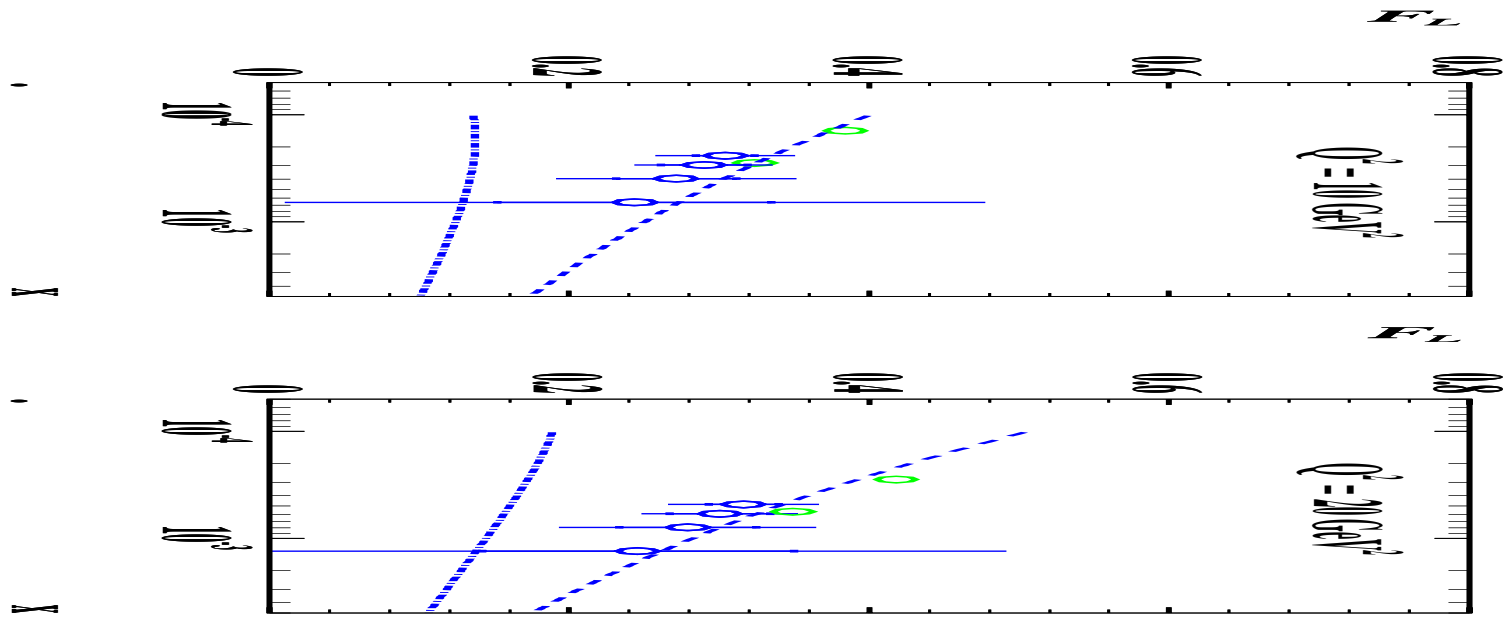
‘extrapolation’, ‘derivative’ and ‘shape’

The methods provide an indirect determination of  $F_L$  somewhat model dependent (cf Robert Thorne) and with a modest precision (two to three sigmas).

The shape method is based on a simple parametrization of  $\sigma_r$  at fixed  $Q^2$



# Undirect determination of $F_L$



A new undirect determination of  $F_L$  at lower proton beam energy would provide an interesting cross check based on quite different systematics.

# The Diffractive Longitudinal Structure Function $F_L^D$

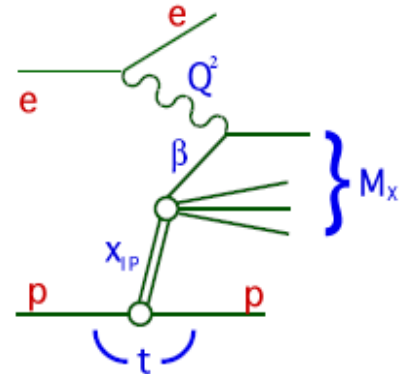
By analogy with inclusive case,

$$\frac{d^3 \sigma^{ep \rightarrow eXY}}{dx_{\mathbb{P}} d\beta dQ^2} = \frac{2\pi\alpha^2}{\beta Q^4} \cdot Y_+ \cdot \sigma_r^D(x_{\mathbb{P}}, x, Q^2)$$

where  $\sigma_r^D = F_2^D - \frac{y^2}{Y_+} F_L^D$  and  $Y_+ = 1 + (1 - y)^2$

Several different measurement possibilities:

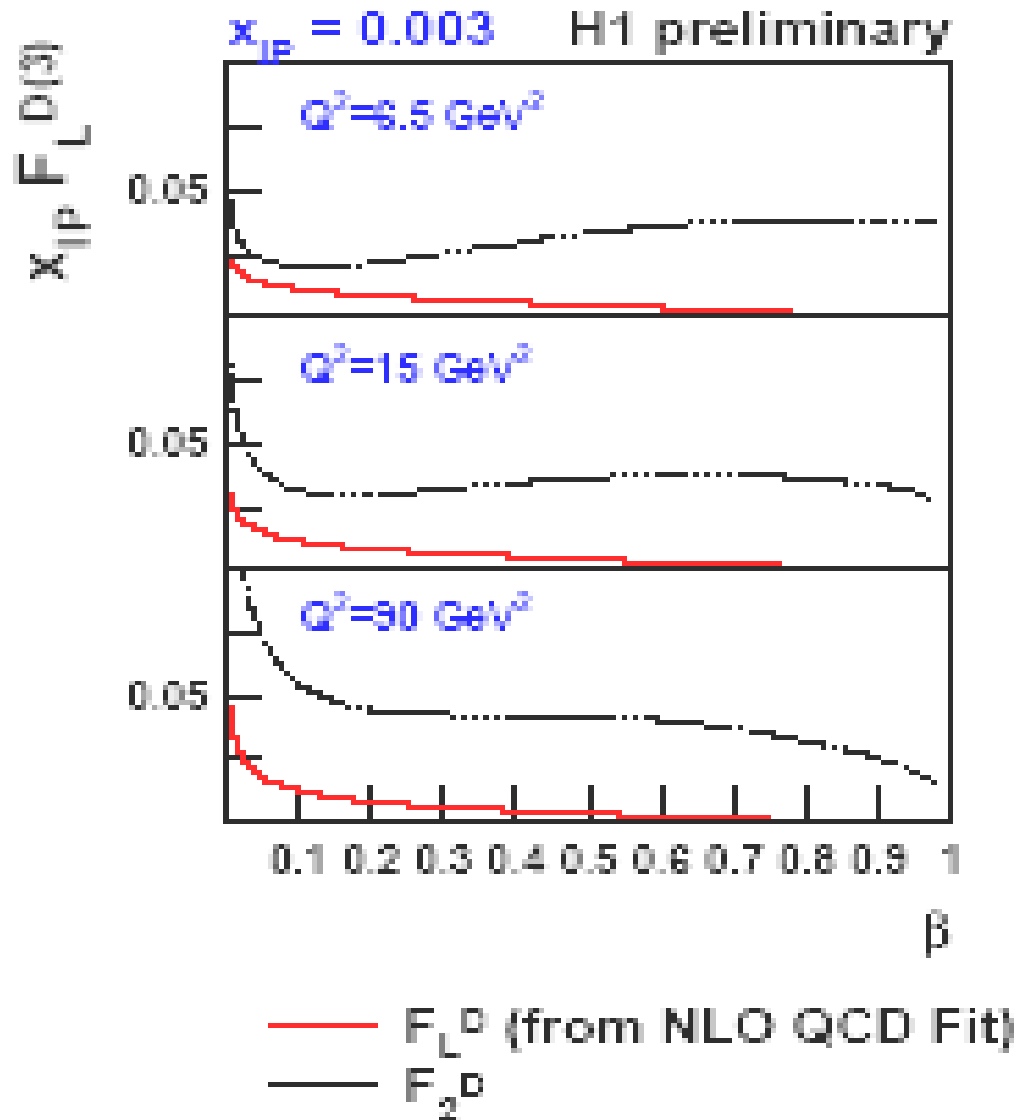
- “Shape” method as in inclusive case:
  - ... requires knowledge of  $x_{\mathbb{P}}$  dependence of  $F_2^D$ , which is poorly constrained by theory and complicated by meson effects, interference etc.
- Exploit interference between transverse and longitudinal contributions leading to modulation in  $\Delta\phi$  between lepton and proton scattering planes.
  - ... Current results (ZEUS) consistent with zero due to poor statistics and large systematics.
  - ... Maybe interesting with VFPS?
- Varying  $y$  at fixed  $Q^2, \beta, x_{\mathbb{P}}$  by changing  $\sqrt{s}$  promises model independent result.



$F_L^D$  has never been measured

$F_L^D$  predicted from QCD fits to be large at low  $\beta$ !

Inclusive diffraction cannot be fully understood without separating out  $F_L^D$  contribution.





# Results and Uncertainties for $Q^2 = 12 \text{ GeV}^2, \beta = 0.23$

$y_{400}$	$x_P$	$\delta_{\text{unc}}$	$\delta_{\text{pdiss}}$	$\delta E'_e$	$\delta \theta'_e$	$\delta M_X$	$\delta \gamma p$	$\delta_{\text{syst}}$	$\delta_{\text{stat}}$	$\delta_{\text{tot}}$
0.5 – 0.7	0.0020	34	6	8	2	7	0	36	20	41
0.7 – 0.8	0.0016	19	6	3	2	5	6	22	17	28
0.8 – 0.9	0.0014	14	6	6	1	2	13	21	13	25

Uncertainties correlated between beam energies:

- $\delta E'_e = 0.2\%$  (kinematic peak) ...  $2\%$  ( $E'_e = 3 \text{ GeV}$ )
- $\delta \theta'_e = 0.2 \text{ mrad}$
- Hadronic energy scale  $\delta M_X = 4\%$  (as now)
- Photoproduction background  $\delta \gamma p = 25\%$  (as now)
- Proton dissociation corrections  $\delta_{\text{pdiss}} = 6\%$  (as now, assumed 100% correlated)

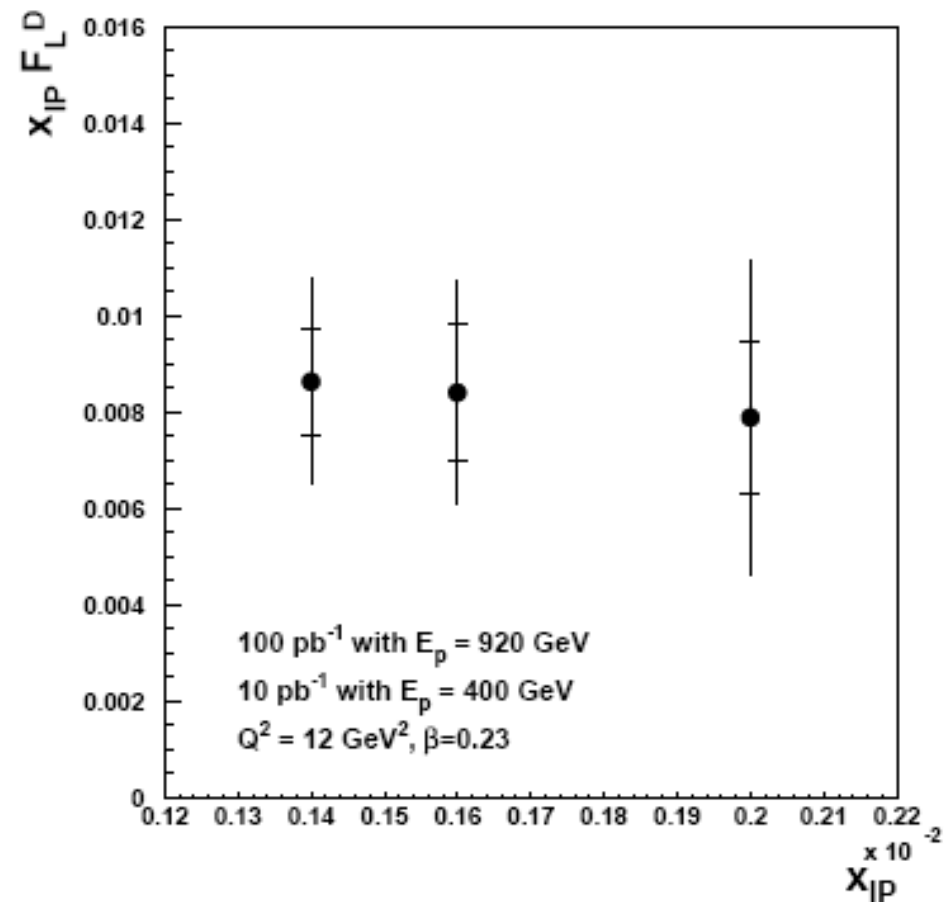
Uncorrelated uncertainty =  $2.4\%$ , mainly from acceptance corrections with RAPGAP

Encouraging result !

$F_L^D$  extracted to 3 – 4  $\sigma$  in 3 bins

...

What about two bins of  $t$  with VFPS?



# Possible scenarii for a direct measurement of $F_L$

Combine 30 pb<sup>-1</sup> at  $E_p = 920$  GeV with data taken in the last couple of months of HERA running at low proton beam energy(s):

I) 15 pb<sup>-1</sup> at  $E_p = 460$  GeV :

Estimated time (F.Willeke):

3 weeks to tune the machine plus 10 weeks of data taking.

→ 10 pb<sup>-1</sup> of good data.

More precision at  $x \sim 10^{-3}$ . Good safety factor.

Would provide a first measurement of  $F_L^D$ .

II) 7pb<sup>-1</sup>(→ 5pb<sup>-1</sup> ) at  $E_p = 460$  GeV

plus 5 pb<sup>-1</sup> (→ 3.5 pb<sup>-1</sup>) at  $E_p = 575$  GeV

Estimated time :

(3 + 3) weeks to tune HERA plus (5 + 2.5) weeks of data taking.

An excellent check of the systematics and extension of the x range.

# Conclusion

Based on the excellent performance of the HERA detectors and on the cumulated expertise on systematic effects since 1992, a direct measurement of  $F_L$  at  $x$  from  $10^{-4}$  to  $10^{-3}$  at the 5 sigmas level of precision can be reached by running HERA for a few weeks at low proton beam energy(s).

It could also provide the first measurement of  $F_L^D$  at the 3-4 sigmas level.

An attractive added value to the legacy of HERA at small  $x$ !

Be open minded and prepared for a final decision by end 2006.