

Madison, DIS 2005

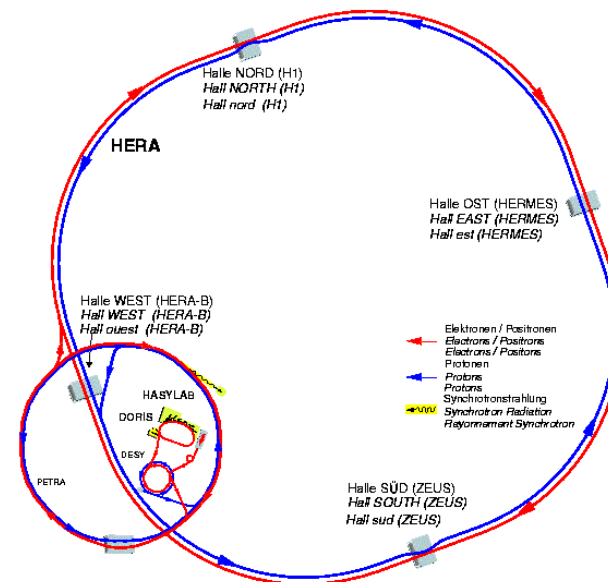
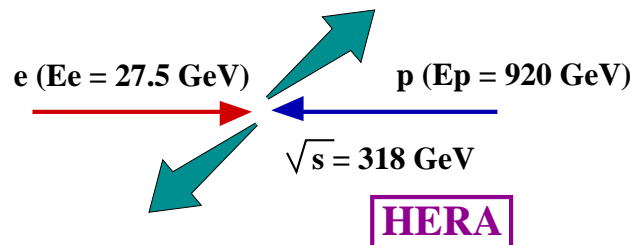
April 27th, 2005

PROTON PDFS USING STRUCTURE FUNCTION AND JET DATA FROM ZEUS

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ZEUS Collaboration



Determination of the proton PDFs

- **Observables used in the fits to determine the proton PDFs:**

→ **Inclusive measurements of deep inelastic lN scattering**

- **Advantages:**

→ **inclusive (only final-state lepton is tagged)**

⇒ **no QCD corrections associated to the final-state signature (lepton!)**

⇒ **no hadronisation corrections**

associated to the final-state signature (lepton!)

- **Disadvantages:**

→ **the gluon distribution contributes indirectly through higher-order terms**

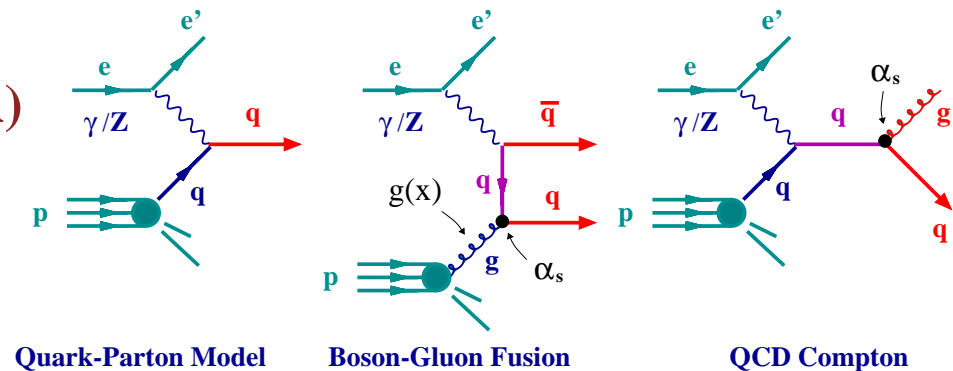
- **Observables based on jets have hardly been used (except jet production in $p\bar{p}$ at TeVatron)**

→ **large QCD corrections and hadronisation corrections**

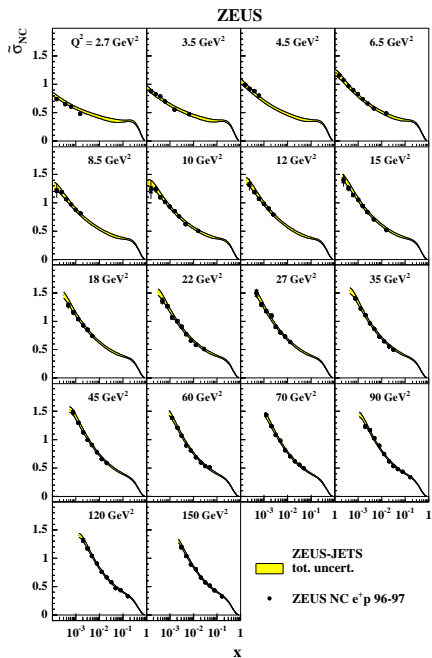
- **That's the past! NOW there are measurements of jet cross sections at HERA**

⇒ **directly sensitive to the gluon density in the proton**

⇒ **with small experimental and theoretical (terms beyond NLO, hadronisation) uncertainties!**

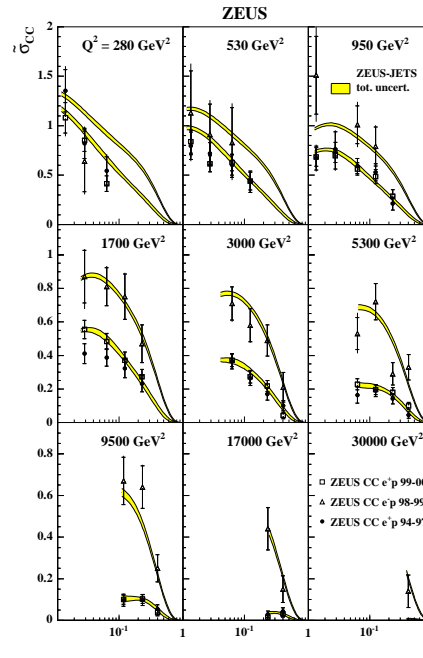
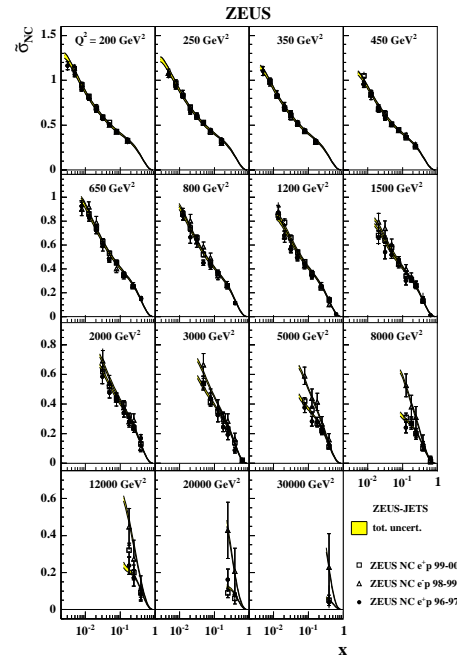


Determination of PDFs using structure function and jet data from ZEUS

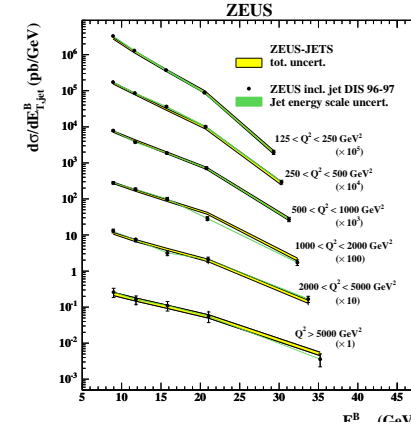


low- Q^2 NC DIS

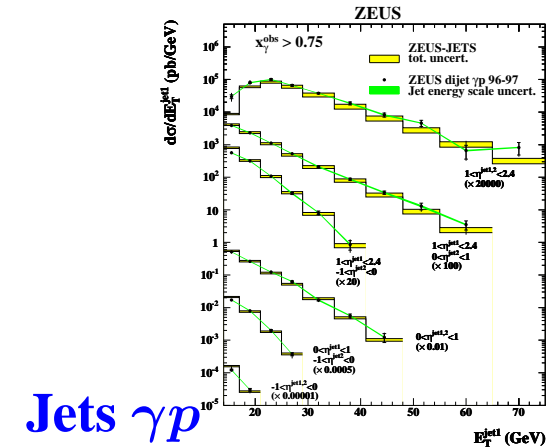
high- Q^2 NC DIS



high- Q^2 CC DIS



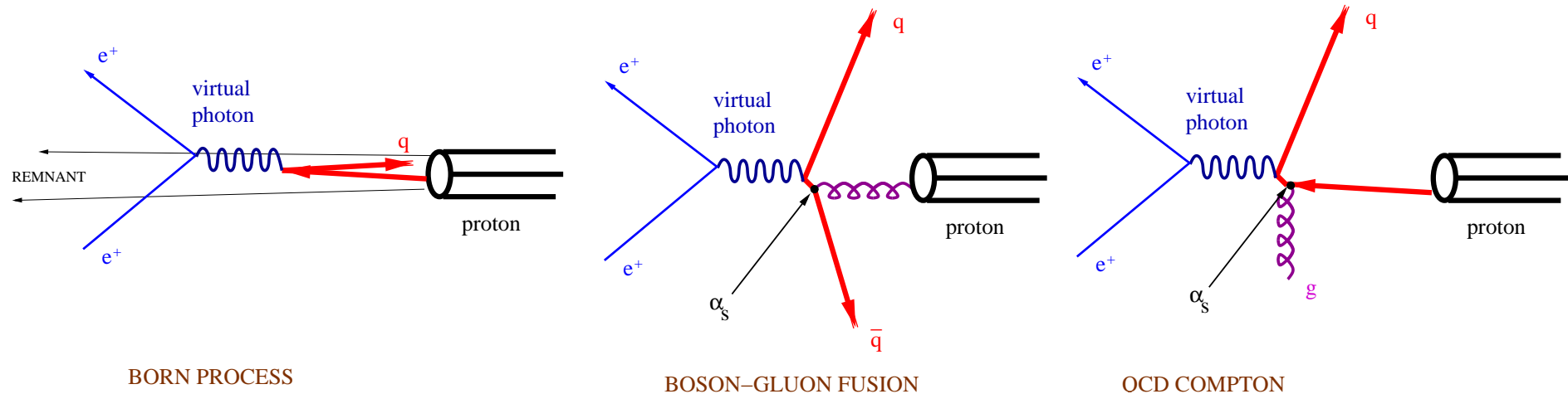
Jets NC DIS



Jets γ

- Determination of the proton PDFs using inclusive measurements of NC and CC DIS $e^\pm p$ in a large level arm of x and $Q^2 \Rightarrow 6.3 \cdot 10^{-5} < x < 0.65, 2.7 < Q^2 < 30000 \text{ GeV}^2$ and measurements of jet cross sections in NC DIS and γp collisions from ZEUS only (!)
- Sufficient sensitivity to determine the proton PDFs within a single experiment

Inclusive Production of High- E_T Jets in NC DIS in the Breit Frame



- In the Breit frame the virtual boson collides head-on with the proton

- High- E_T jet production in the Breit frame

- suppression of the Born contribution (struck quark has zero E_T)

- suppression of the beam-remnant jet (zero E_T)

- lowest-order contributions from

- Boson-gluon fusion process $\gamma^* g \rightarrow q\bar{q}$

- QCD-Compton process $\gamma^* q \rightarrow qg$

- ⇒ directly sensitive to α_s and the gluon density in the proton

Inclusive Jet Cross Sections in NC DIS at $Q^2 > 125 \text{ GeV}^2$

- Measurement of inclusive jet cross sections in the kinematic region defined by $Q^2 > 125 \text{ GeV}^2$ and $-0.7 < \cos \gamma < 0.5$ for jets with

$$E_{T,jet}^B > 8 \text{ GeV and } -2 < \eta_{jet}^B < 1.8$$

$$\text{where } \cos \gamma = \frac{(1-y)x E_p - y E_e}{(1-y)x E_p + y E_e}$$

- Jets identified using the longitudinally invariant k_T cluster algorithm in the Breit frame
- Inclusive jet production: smaller uncertainties than for dijet production

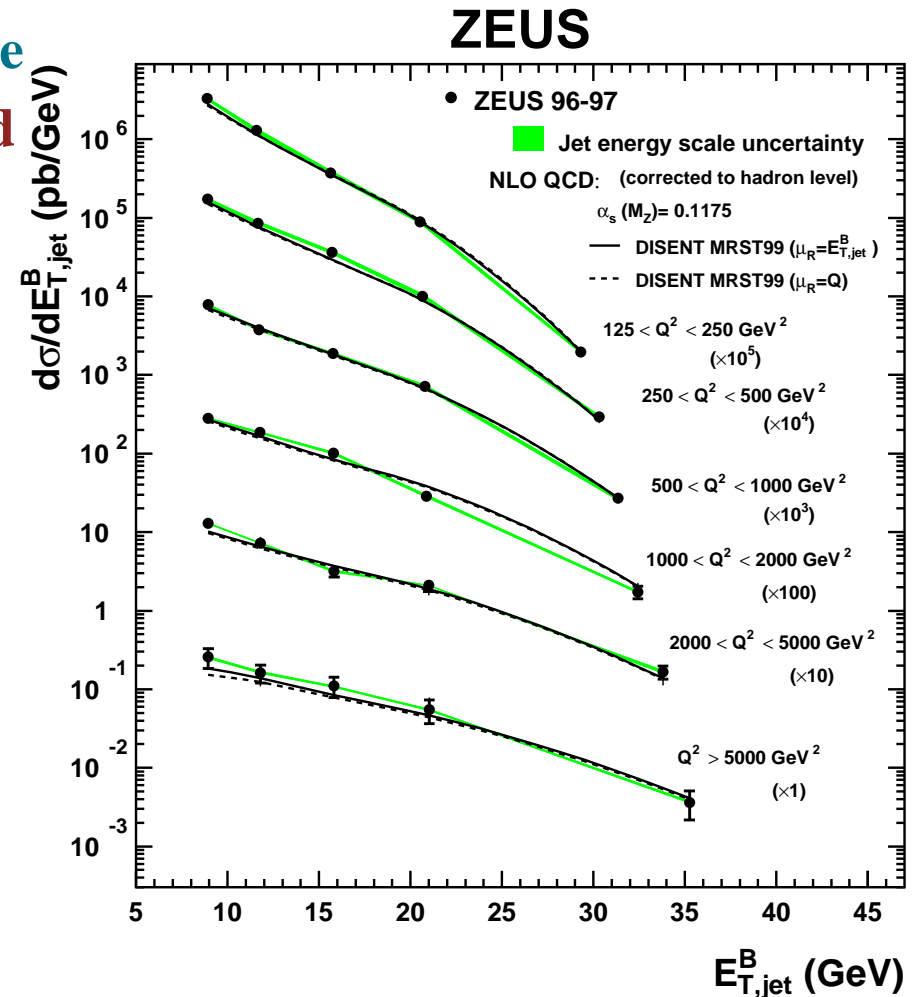
- Small experimental uncertainties:

→ jet energy scale (1% for $E_{T,jet} > 10 \text{ GeV}$)

⇒ $\sim \pm 5\%$ on the cross sections

- Small parton-to-hadron corrections (C_{had}): $< 10\%$

→ $< 5\%$ for $Q^2 > 250 \text{ GeV}^2$ except for lowest- $E_{T,jet}^B$ point



Inclusive Jet Cross Sections in NC DIS at $Q^2 > 125 \text{ GeV}^2$

- Several NLO QCD ($\mathcal{O}(\alpha_s^2)$) programs are available for performing jet-cross-section calculations: MEPJET, DISENT, DISASTER++, NLOJET

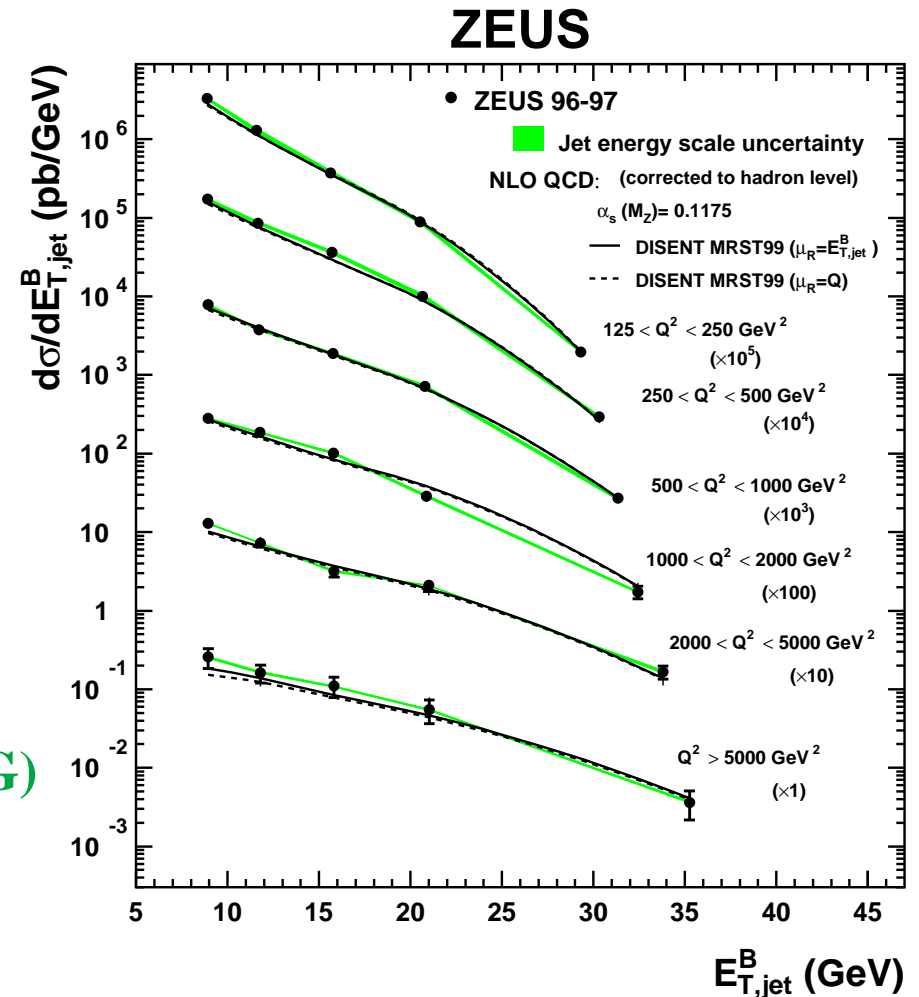
- **Small theoretical uncertainties** :

→ higher-order terms ($> \text{NLO}$); varying μ_R

between $\frac{1}{2} \cdot E_{T,jet}^B$ and $2 \cdot E_{T,jet}^B \Rightarrow \pm 5\%$

→ hadronisation corrections; variance of C_{had} values (ARIADNE, LEPTO-MEPS, HERWIG)

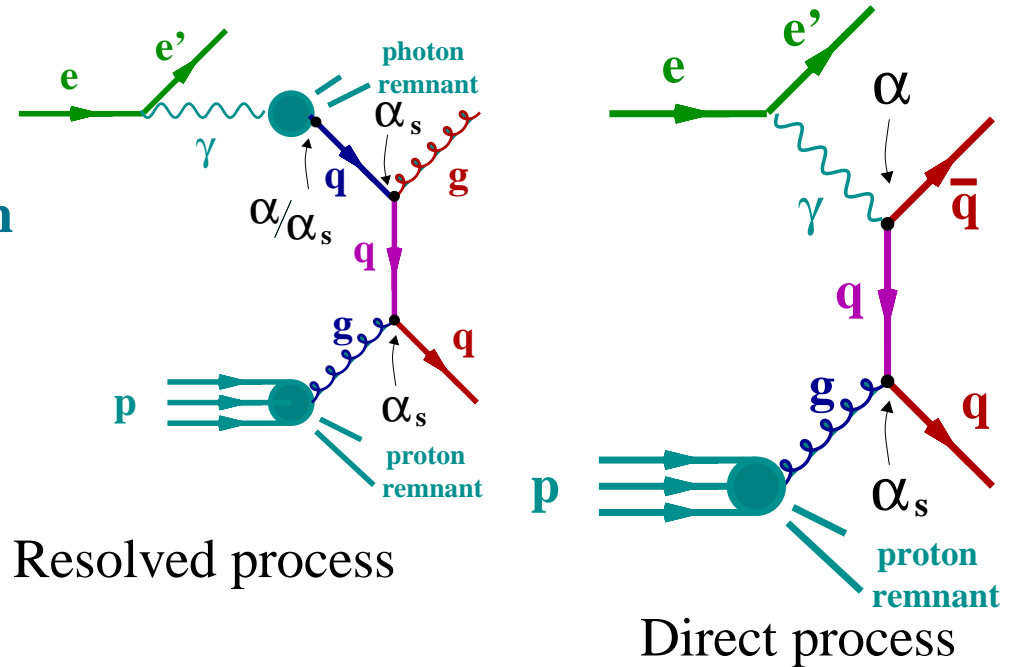
$\Rightarrow < 1\%$



- Inclusive jet cross sections in NC DIS in the Breit frame provide direct sensitivity to the gluon density in the proton with small experimental and theoretical uncertainties

Dijet Photoproduction

- Production of jets in γp collisions has been measured via ep scattering at $Q^2 \approx 0$
- At lowest order QCD, two hard scattering processes contribute to jet production \Rightarrow
- pQCD calculations of jet cross sections



$$\sigma_{2jet} = \sum_{a,b} \int_0^1 dy f_{\gamma/e}(y) \int_0^1 dx_\gamma f_{a/\gamma}(x_\gamma, \mu_{F\gamma}^2) \int_0^1 dx_p f_{b/p}(x_p, \mu_{Fp}^2) \hat{\sigma}_{ab \rightarrow jj}$$

longitudinal momentum fraction of γ/e^+ (y), **parton a/γ** (x_γ), **parton b /proton** (x_p)

$\rightarrow f_{\gamma/e}(y) =$ flux of photons in the positron (WW approximation)

$\rightarrow f_{a/\gamma}(x_\gamma, \mu_{F\gamma}^2) =$ parton densities in the photon (for direct processes $\delta(1 - x_\gamma)$)

$\rightarrow f_{b/p}(x_p, \mu_{Fp}^2) =$ parton densities in the proton

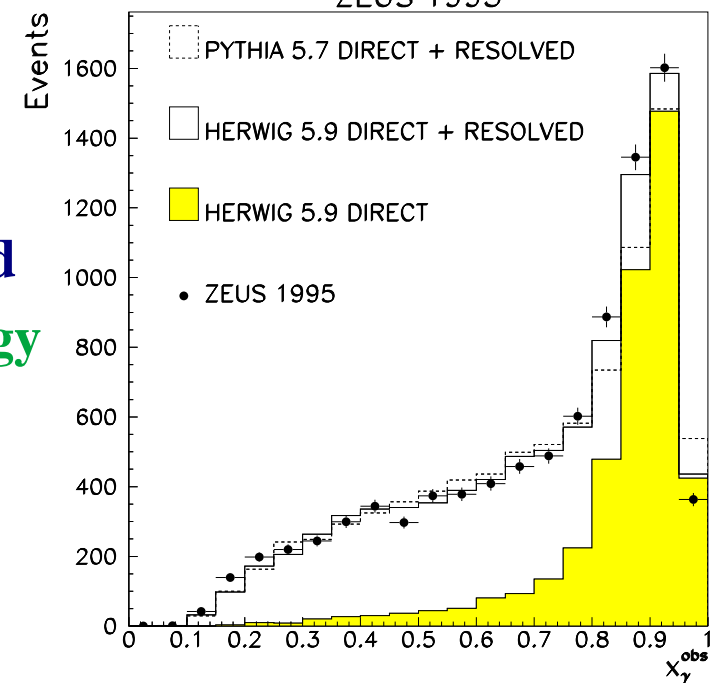
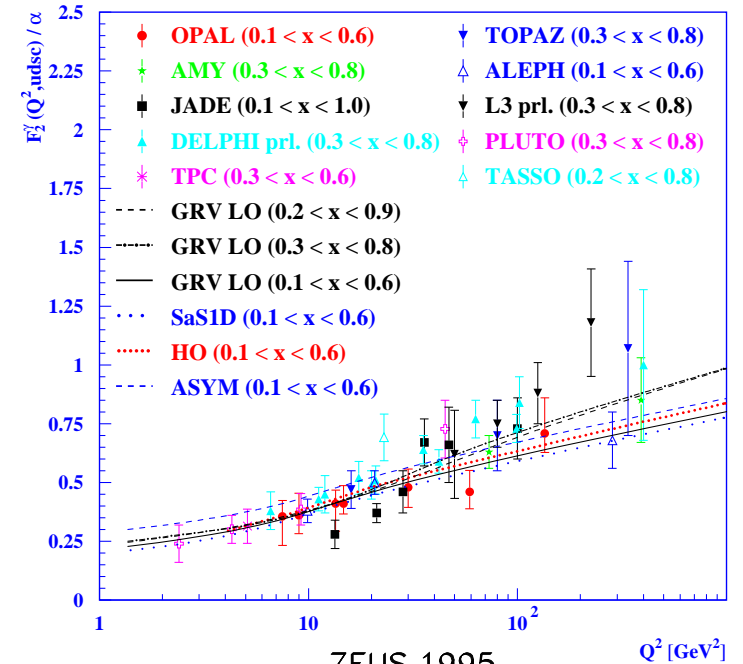
$\rightarrow \sigma_{ab \rightarrow jj}$ subprocess cross section; short-distance structure of the interaction

Dijet Photoproduction

- Measurements of dijet photoproduction provide
 - **direct sensitivity to α_s and gluon density in the proton**
- However, resolved processes are also sensitive
 - to both the quark and gluon densities in the photon
 - **Photon structure: information on quark densities from F_2^γ in e^+e^- ; gluon density poorly constrained!**
- To suppress the dependence on the photon PDFs
 - **measurements restricted to the region where direct processes dominate**
 - ⇒ **Observable to separate the contributions from resolved and direct processes: the fraction of the photon's energy participating in the production of the dijet system**

$$x_\gamma^{obs} = \frac{1}{2E_\gamma} \sum_{i=1}^2 E_T^{jet_i} e^{-\eta^{jet_i}}$$

→ **Direct-process-enriched region: $x_\gamma^{OBS} > 0.75$**



Dijet Photoproduction Cross Sections for $x_\gamma^{obs} > 0.75$

- Measurement of the differential cross section

section $d\sigma/dE_T^{jet,1}$ for dijet events with

$$E_T^{jet,1} > 14 \text{ GeV}, E_T^{jet,2} > 11 \text{ GeV},$$

$$-1 < \eta^{jet} < 2.4 \text{ (both jets) and } x_\gamma^{obs} > 0.75$$

in the kinematic region

$$Q^2 < 1 \text{ GeV}^2 \text{ and } 134 < W_{\gamma p} < 277 \text{ GeV}$$

- Jets identified using the longitudinally invariant k_T cluster algorithm in the laboratory frame

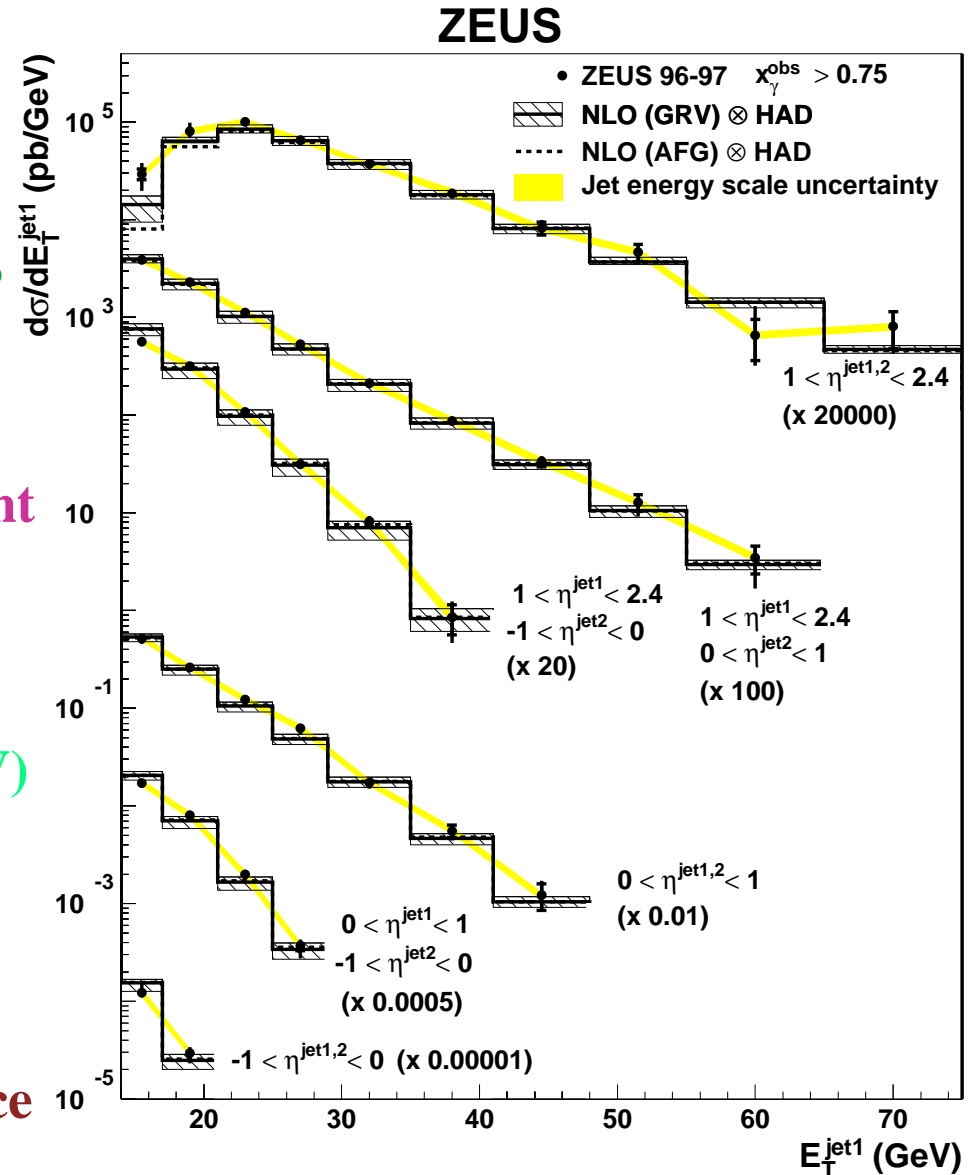
- Small experimental uncertainties:

→ jet energy scale (1% for $E_{T,jet} > 10 \text{ GeV}$)

⇒ $\sim \pm 5\%$ on the cross sections

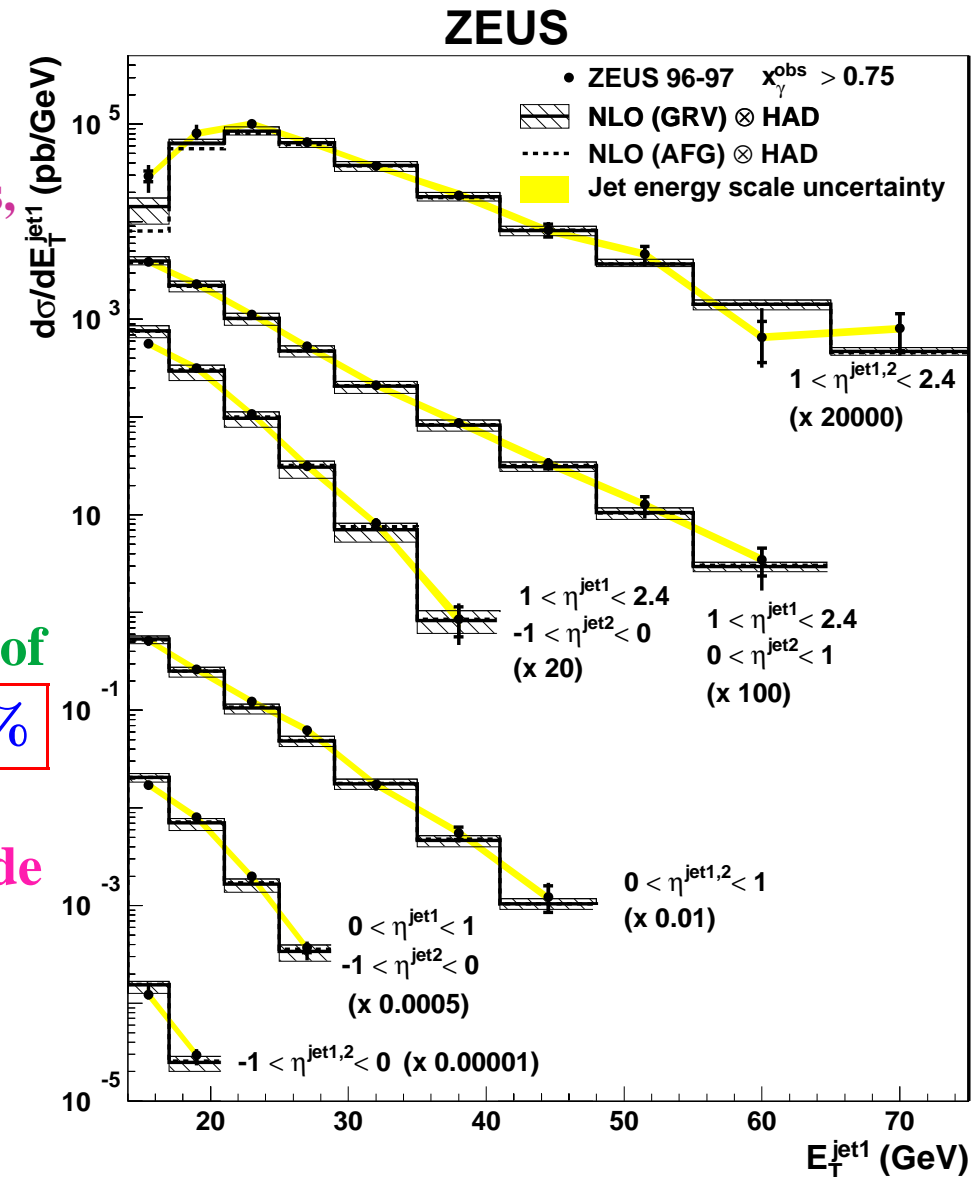
- Small parton-to-hadron corrections (C_{had}):

→ $< 10\%$ except at the edges of phase space



Dijet Photoproduction Cross Sections for $x_\gamma^{obs} > 0.75$

- Several NLO QCD ($\mathcal{O}(\alpha_s^2)$) programs are available for performing jet-cross-section calculations: Klasen & Kramer, Harris & Owens, Aurenche et al, Frixione & Ridolfi
- **Small theoretical uncertainties**:
 - higher-order terms ($> \text{NLO}$); varying μ_R between $\frac{1}{4} \cdot E_T^{\text{sum}}$ and $E_T^{\text{sum}} \Rightarrow \pm 10\%$
 - hadronisation corrections; half the spread of C_{had} values (PYTHIA, HERWIG) $\Rightarrow 2 - 3\%$
- Dijet cross sections in photoproduction provide **direct sensitivity to the gluon density in the proton with small experimental and theoretical uncertainties**



Determination of the proton PDFs

- Parametrization of the proton PDFs

→ u valence (xu_V), d valence (xd_V), gluon (xg), total sea (xS) and $x\Delta = x(\bar{d} - \bar{u})$ at $Q_0^2 = 7 \text{ GeV}^2$ by the functional form

$$xf(x) = p_1 x^{p_2} (1 - x)^{p_3} (1 + p_4 x)$$

- Constraints on the parameters $\{p_i\}$:

→ momentum and number sum rules $\Rightarrow p_{1,g}, p_{1,u_V}, p_{1,d_V}$

→ no sensitivity to difference on low- x behaviour of u and d valence: $p_{2,u_V} = p_{2,d_V}$

→ no sensitivity to flavour structure of light-quark sea: fix $p_{i,\Delta}$ consistent with Gottfried sum rule and Drell-Yan data

→ suppression of strange sea in accordance with dimuon data from CCFR-NuTeV

- The results of the fit are not sensitive to reasonable variations of these assumptions

→ it is possible to extract a flavour-averaged sea distribution from ZEUS data alone

- Heavy quarks: variable flavour-number scheme of Roberts and Thorne

\Rightarrow 11 free parameters ($+\alpha_s$ when free; otherwise $\alpha_s(M_Z) = 0.118$)

- Evolution of the PDFs with the energy scale: DGLAP equations at NLO (\overline{MS} scheme)

Inclusion of jet cross sections in a fit of the proton PDFs

- In contrast with the evaluation of the structure functions from the evolved PDFs,
 - the calculation of jet cross sections at NLO requires much more CPU time:
 - $\mathcal{O}(10)$ hours per PDF set \Rightarrow Unaffordable in a fit of the proton PDFs
- Deconvolution of the PDFs and α_s from the matrix elements in the calculations
 - construction of grids such that the calculations for the jet cross sections can be performed sufficiently fast (and accurately) $\Rightarrow \mathcal{O}(1 \text{ second})$ per PDF set!

$$\sigma_{jet}(\text{bin } i) = \sum_{n=1,2} \sum_{a=q,\bar{q},g} \sum_j \sum_k f_a(\xi_k, Q_j^2) \cdot \alpha_s^n(Q_j^2) \cdot G_{i,a,n,\xi_k,Q_j^2}$$

- G_{i,a,n,ξ_k,Q_j^2} grid of weights in the “($x = \xi, Q^2$)” plane for a given bin of the cross section (i), parton species (a) and order (n); typical size in $(\xi, Q^2) \rightarrow 100 \times 100$
- \Rightarrow jet-cross-section calculations can be performed for ANY PDF set and ANY value of α_s in a fast way and with an accuracy better than 0.5%

Determination of the proton PDFs

- Data sets used in the fit (577 data points):

→ Structure function measurements: reduced double differential cross sections in x and Q^2 for

- neutral current DIS e^+p and e^-p
- charged current DIS e^+p and e^-p

→ Jet cross section measurements:

- inclusive jet production in NC DIS
- dijet production in γp collisions

⇒ $6.3 \cdot 10^{-5} < x < 0.65, 2.7 < Q^2 < 30000 \text{ GeV}^2$

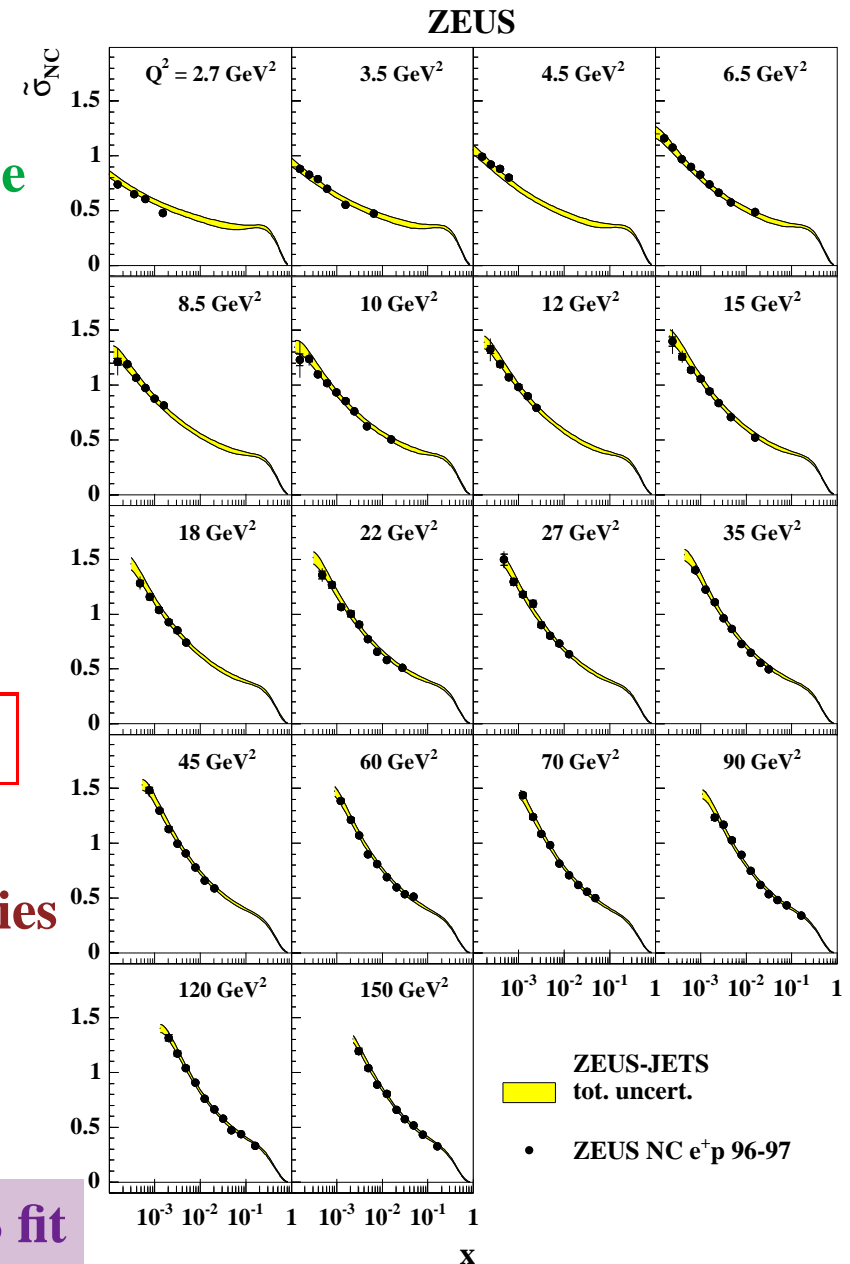
→ $W^2 > 20 \text{ GeV}^2$

- Full account of correlated experimental uncertainties using the offset method

- A good description of the data is obtained:

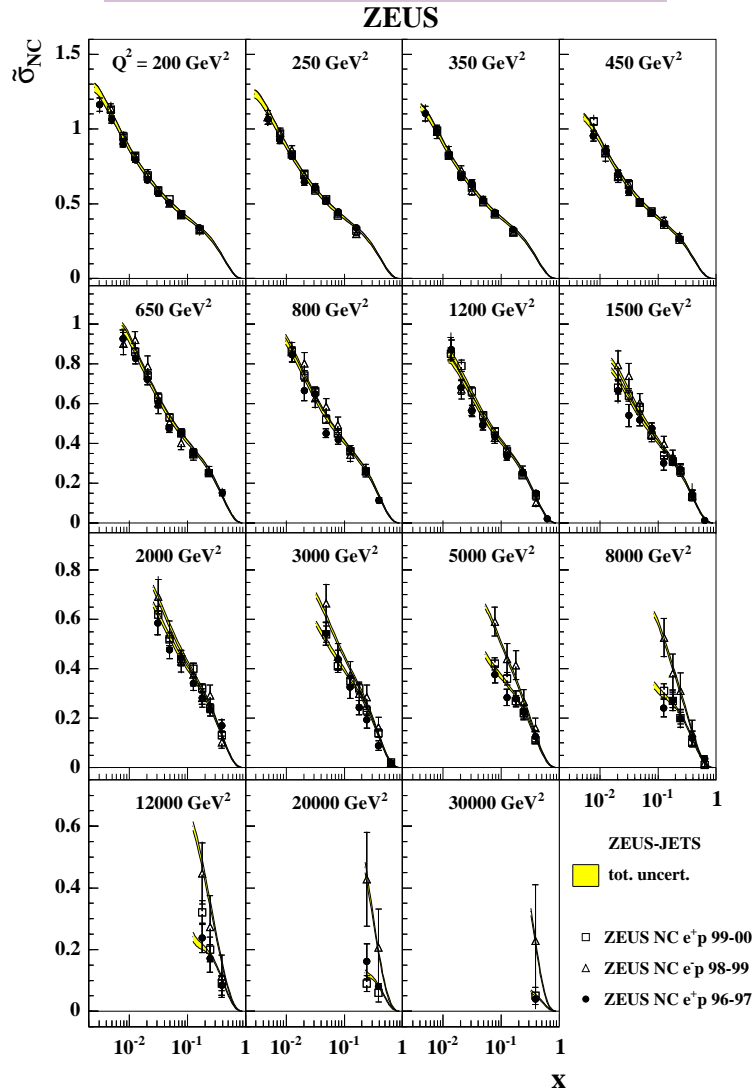
$\chi^2 = 470$ for 577 data points

Low- Q^2 NC DIS e^+p data vs ZEUS-JETS fit

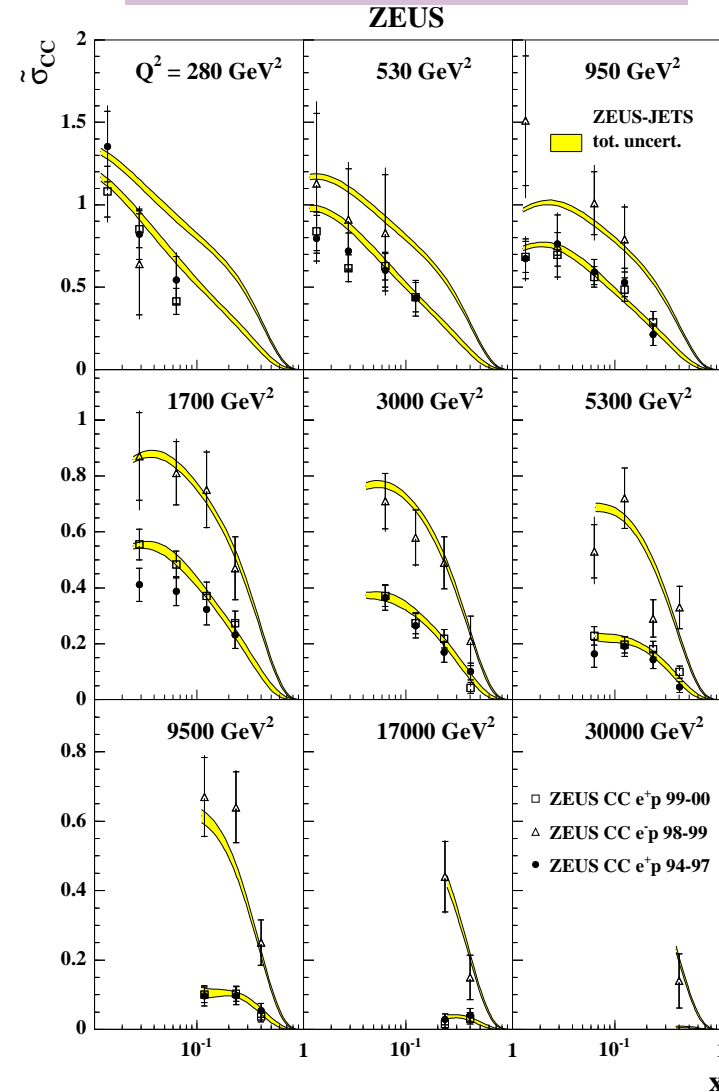


High- Q^2 NC and CC DIS $e^\pm p$ data vs ZEUS-JETS fit

High- Q^2 NC DIS $e^\pm p$



High- Q^2 CC DIS $e^\pm p$



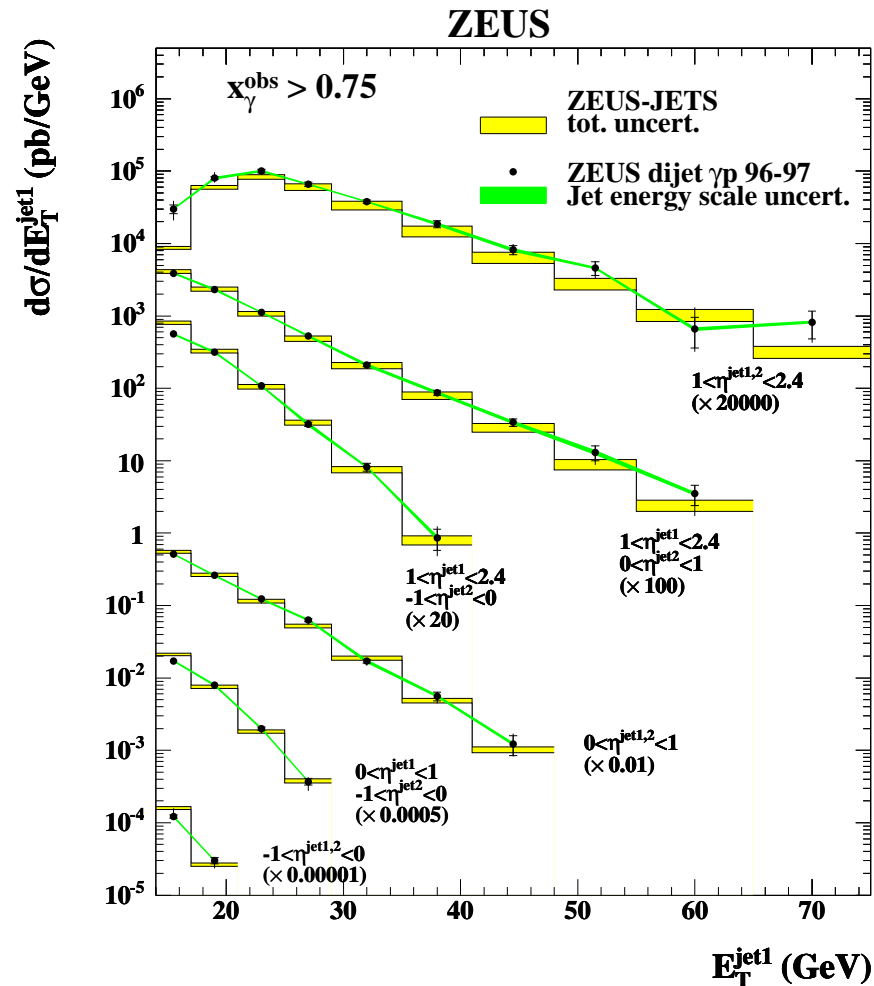
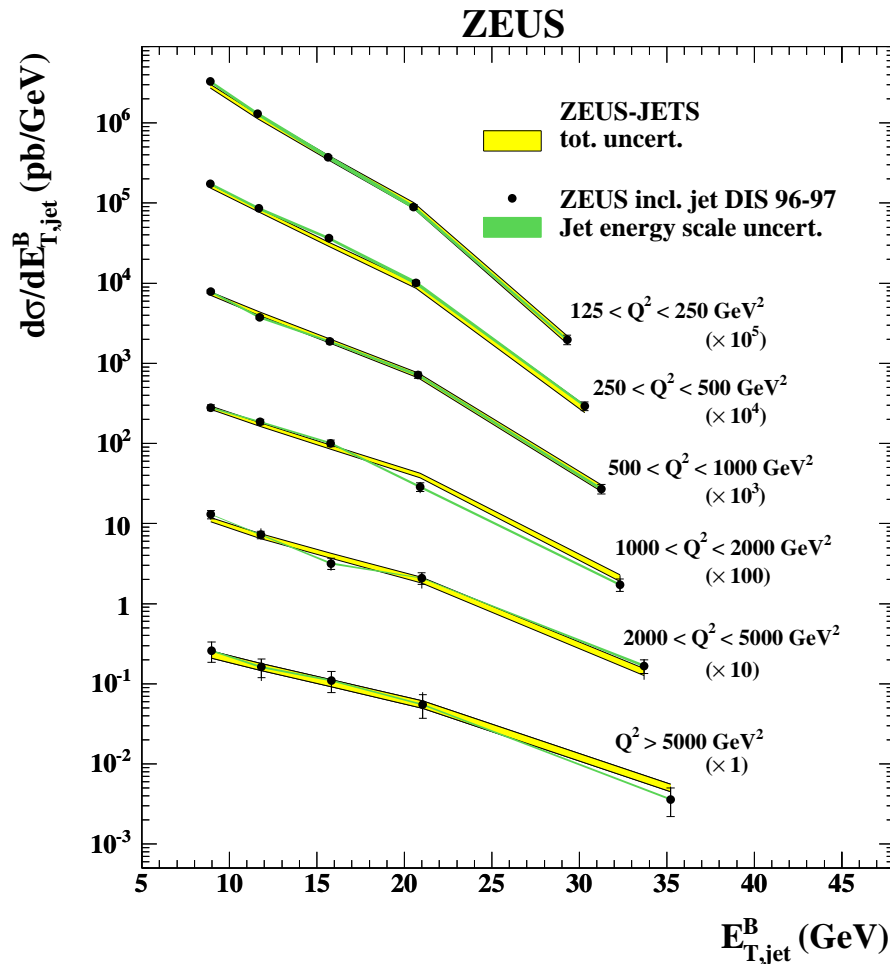
Jet cross sections vs ZEUS-JETS fit

Inclusive Jet Production in NC DIS e^+p

$$Q^2 > 125 \text{ GeV}^2, E_{T,jet}^B > 8 \text{ GeV}$$

Dijet Production in γp collisions

$$E_T^{jet,1(2)} > 14 (11) \text{ GeV}, x_\gamma^{obs} > 0.75$$

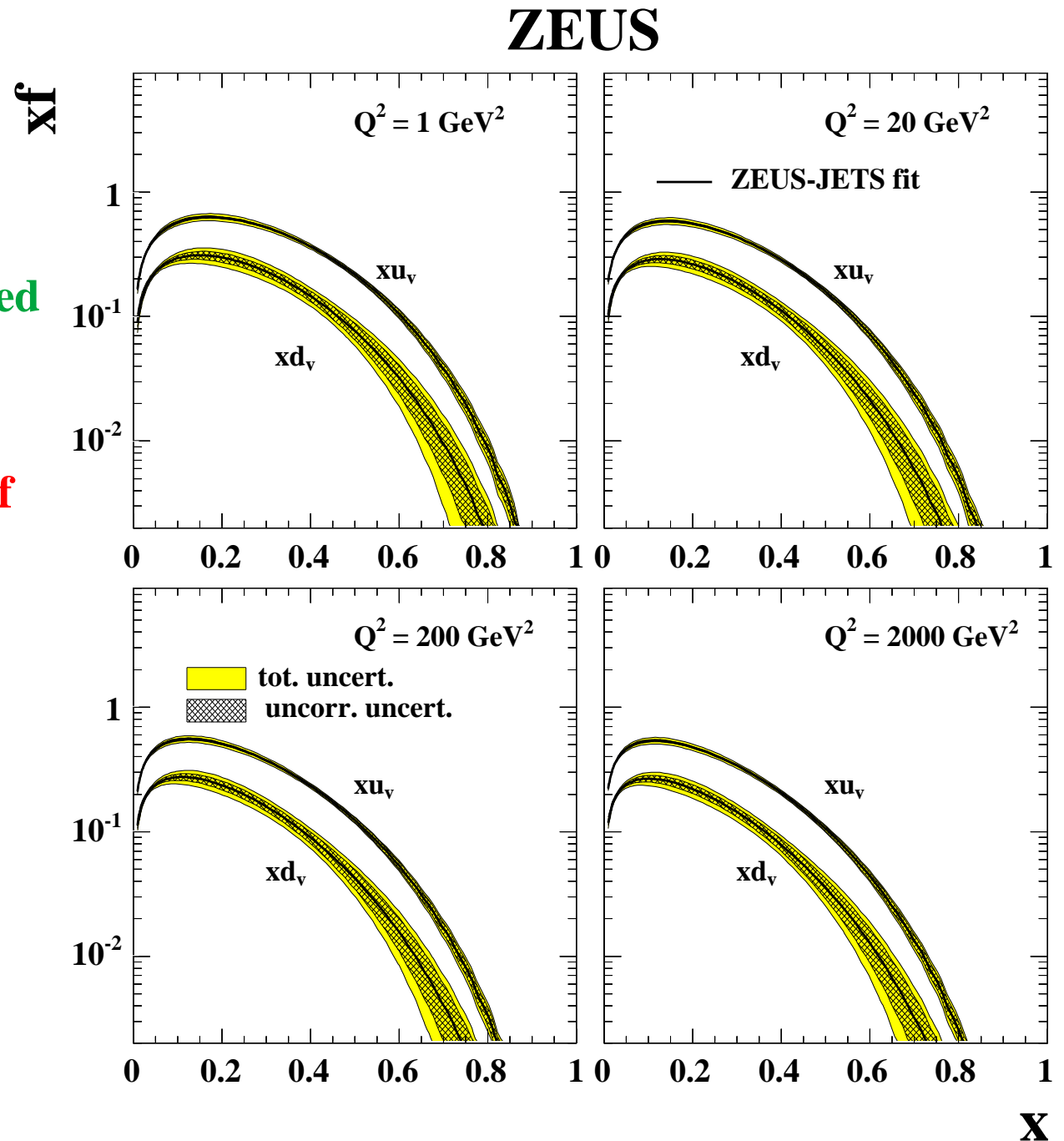


Valence distributions

- Precision of the valence quark distributions

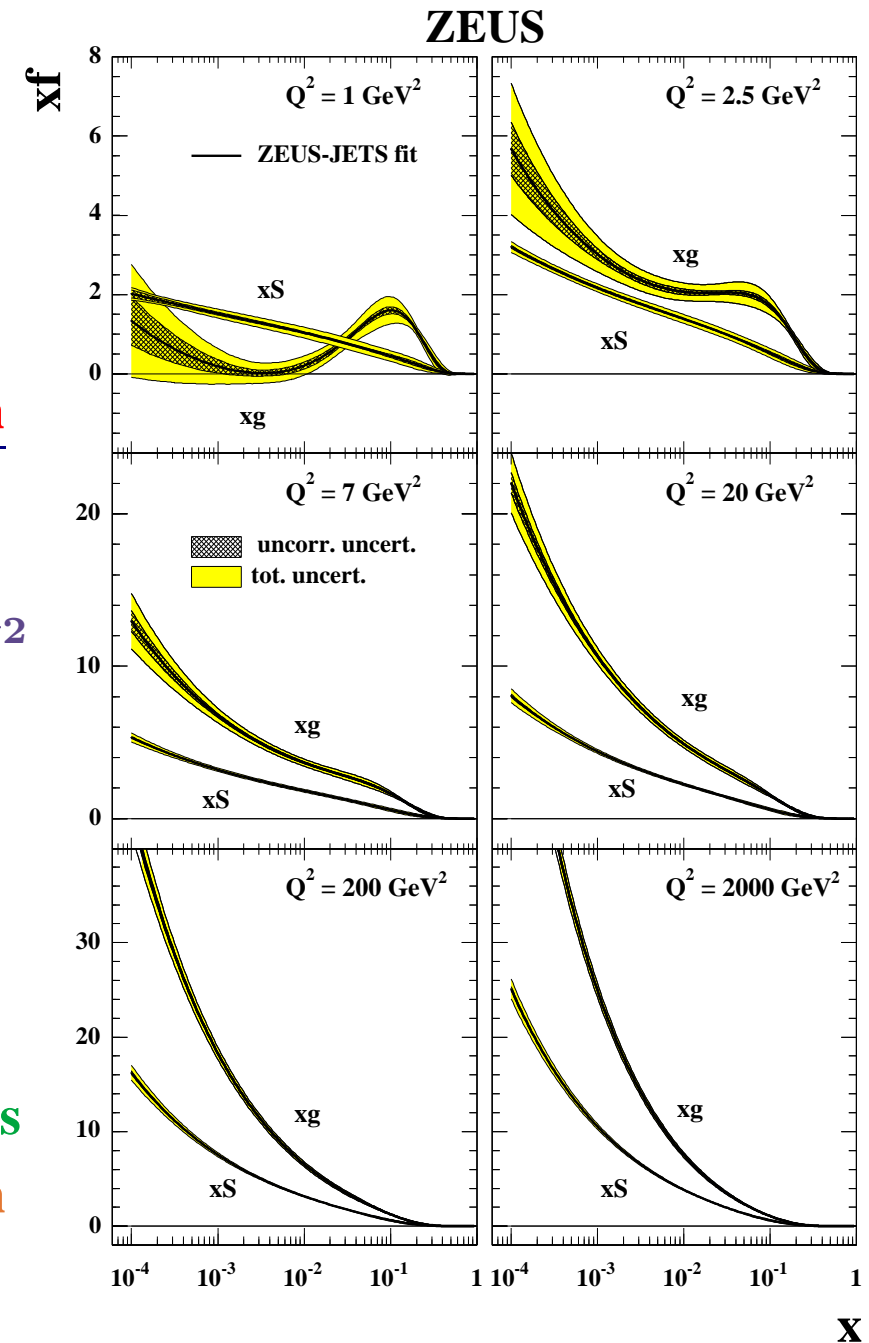
→ at high- x not as well constrained as in global PDF fits which include fixed-target DIS data

→ competitive + the advantage of being free from heavy-target corrections, higher twists and isospin-symmetry assumptions



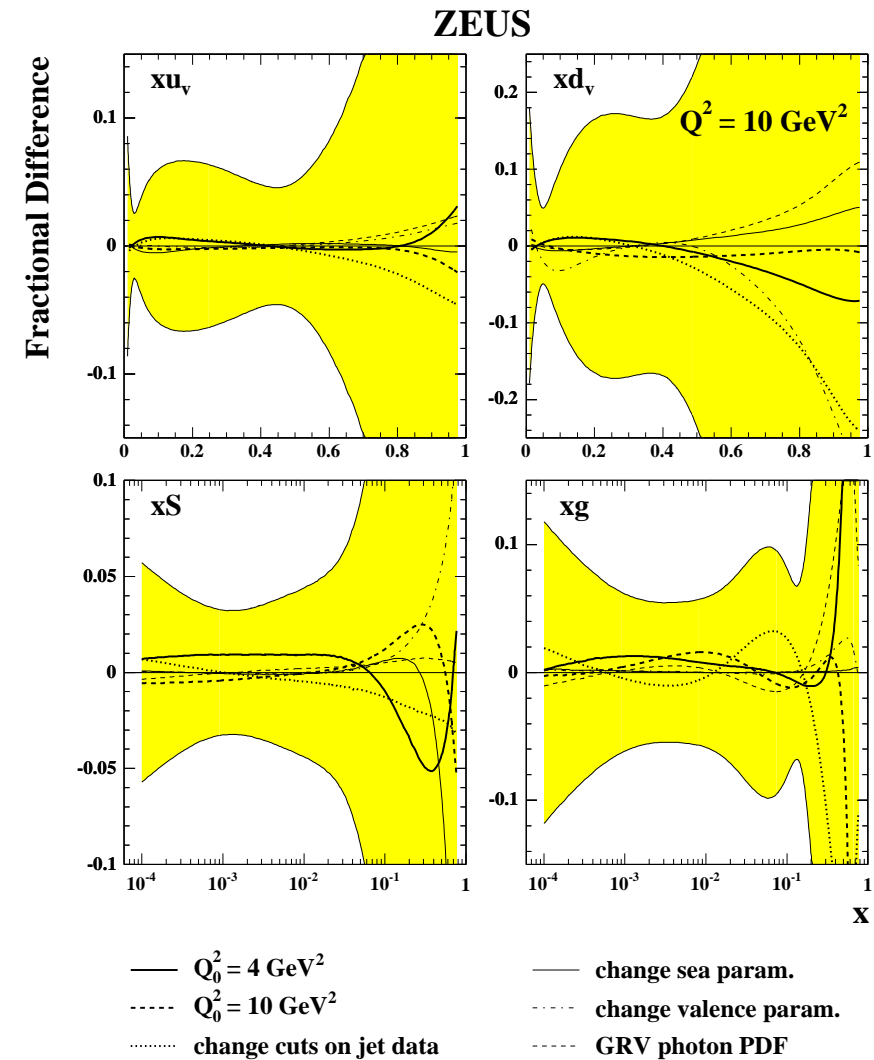
Gluon and sea distributions

- Precision of the gluon and sea distributions
 - at low- x : as well determined as in global PDF fits which make use of HERA data
 - at higher- x : the uncertainty of the gluon has been reduced by the addition of jet data
- Model uncertainties in PDFs:
 - variation of Q_0^2 in the range $4 < Q_0^2 < 10 \text{ GeV}^2$
 - modification $(1 + p_4 x) \rightarrow (1 + p_5 \sqrt{x})$ in parametrizations of valence quarks
 - choice of constraints in parametrization of sea
 - tighter cuts on jet data; $E_{T,jet}^B > 10 \text{ GeV}$ in NC DIS and $E_T^{jet,1} > 17 \text{ GeV}$ in γp
 - parton-to-hadron corrections to jet cross sections
 - residual uncertainty due to photon PDFs used in the dijet calculations (AFG \rightarrow GRV)



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⇒ Effects much smaller than the experimental uncertainties

Improving the gluon distribution: jet data

- Comparison of gluon distributions from fits with and without jet data

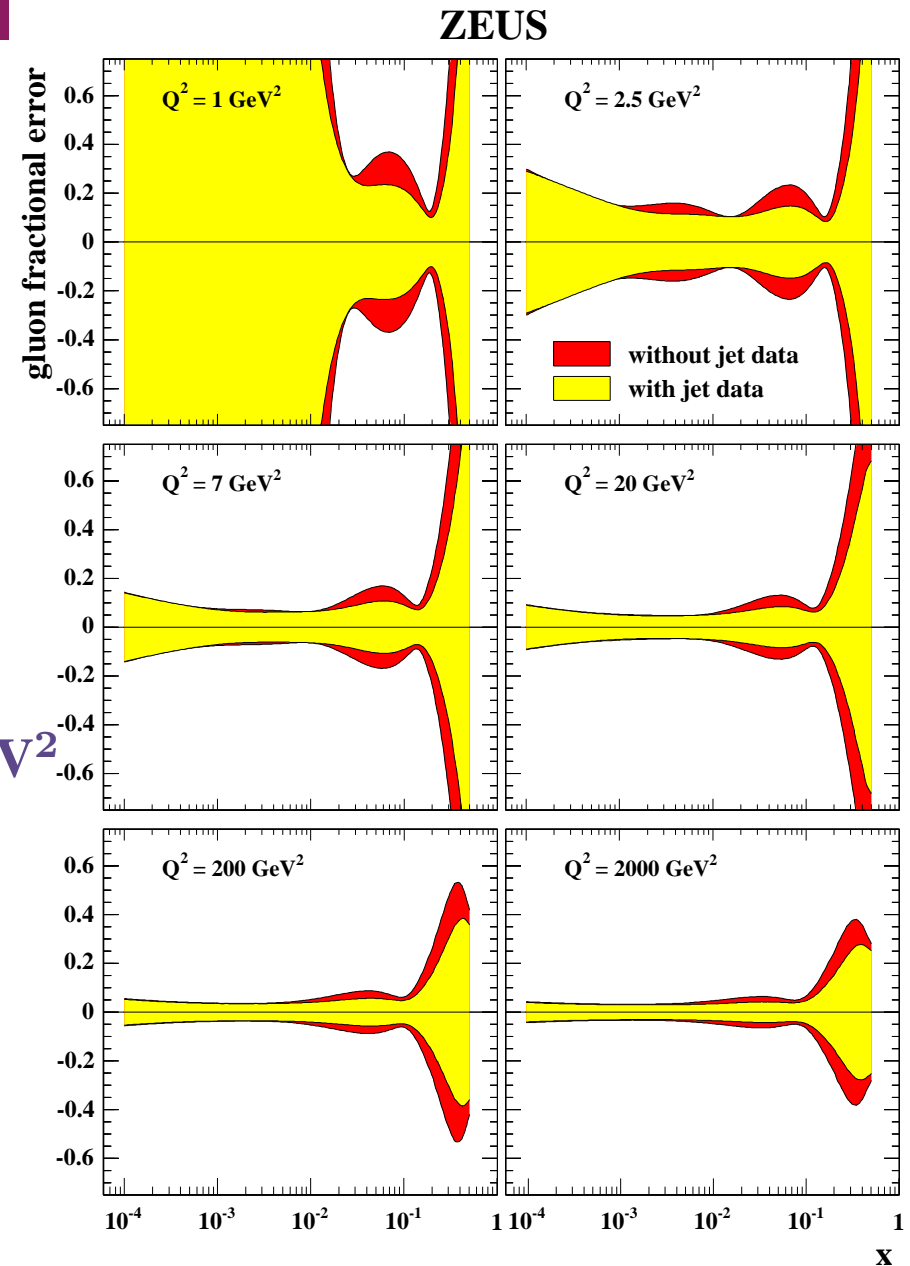
→ no significant change of the shape: no tension between jet and inclusive data

→ the jet cross sections constrain the gluon density in the range 0.01 – 0.4

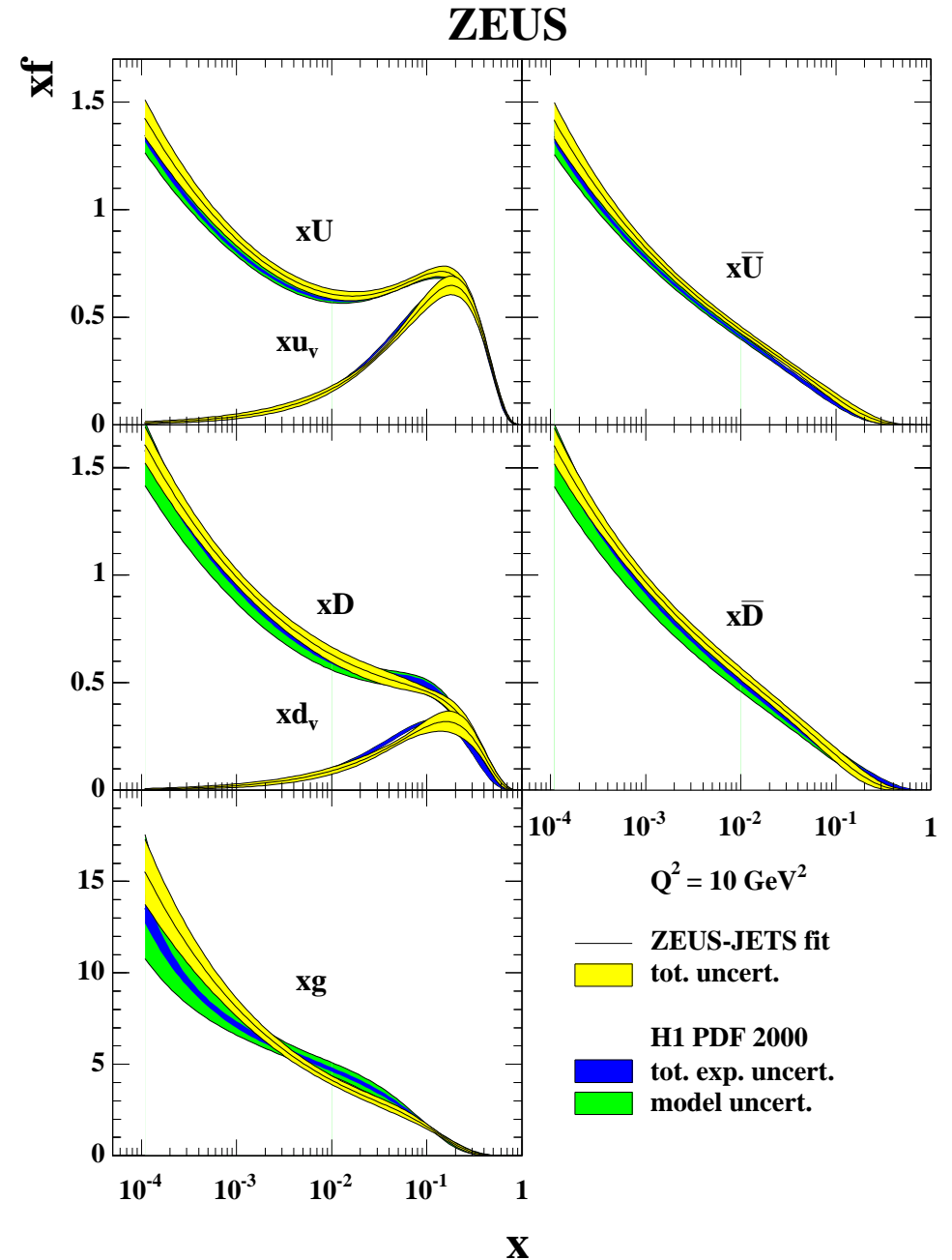
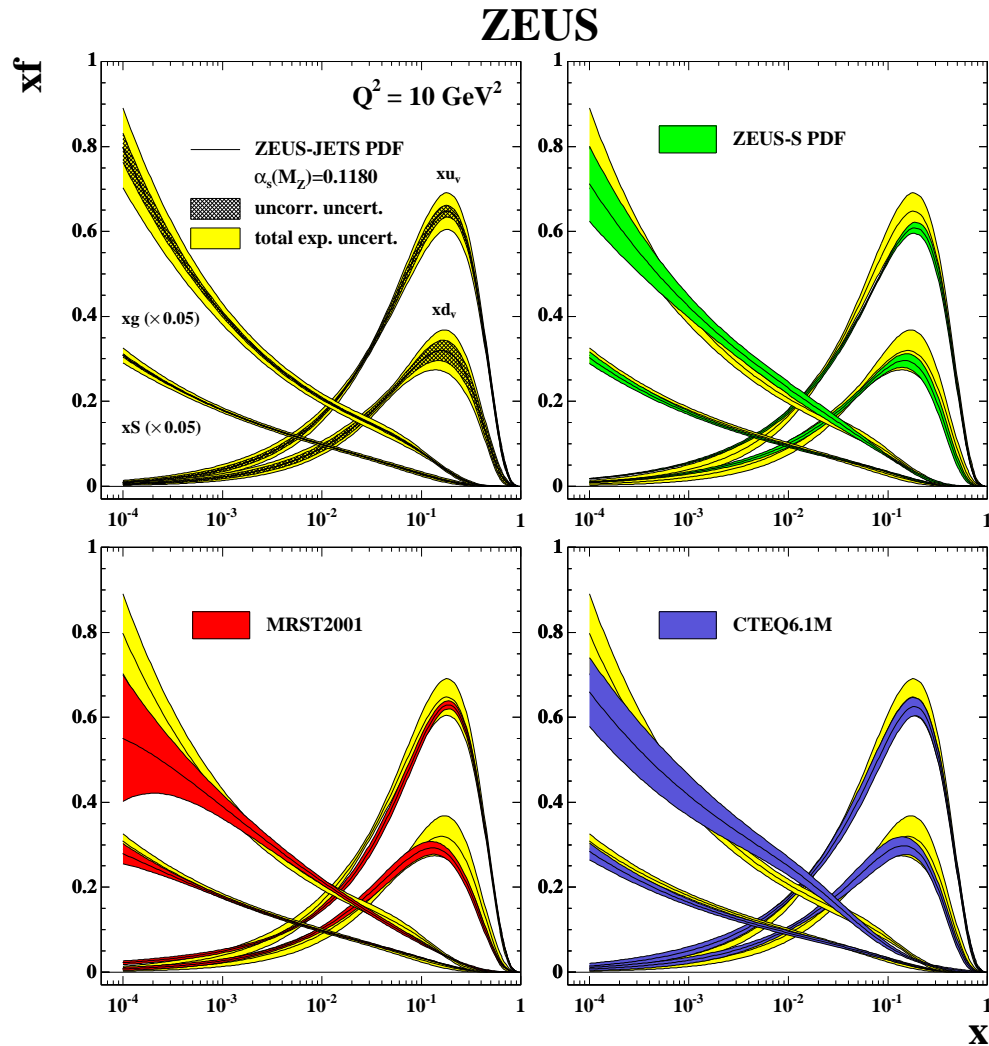
→ **Sizeable reduction of the gluon uncertainty**

e.g. from 17% to 10% at $x = 0.06$ and $Q^2 = 7 \text{ GeV}^2$

→ similar reduction by a factor of two in the mid- x region over the full Q^2 range



Comparison of proton PDFs



- Compatible with MRST2001 and CTEQ6.1M
- Compatible with H1 analysis

Determination of $\alpha_s(M_Z)$

- Simultaneous determination of the proton PDFs and $\alpha_s(M_Z)$
- Jet cross sections are directly sensitive to $\alpha_s(M_Z)$ via $\gamma^{(*)}g \rightarrow q\bar{q}$ (coupled to gluon density) and via $\gamma^{(*)}q \rightarrow qg$ (NOT coupled to gluon density)
- ⇒ The inclusion of the jet cross sections allows an extraction of $\alpha_s(M_Z)$ that is NOT strongly correlated to the gluon density

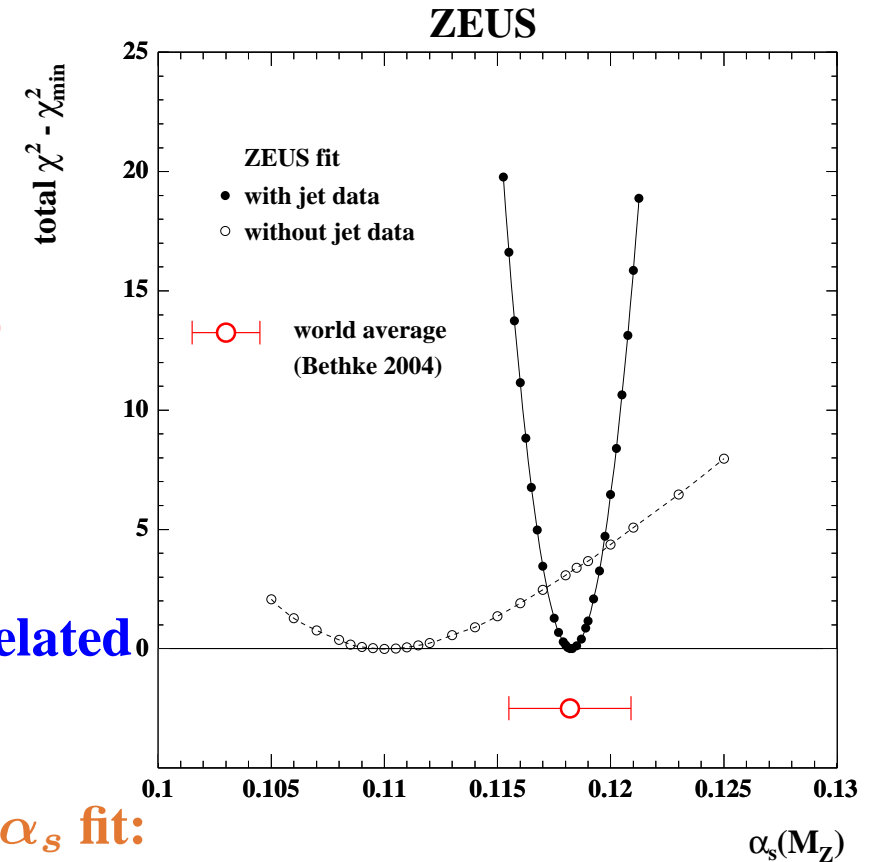
- Determination of $\alpha_s(M_Z)$ from the ZEUS-JETS- α_s fit:

$$\alpha_s(M_Z) = 0.1183 \pm 0.0007 \text{ (uncorr.)} \pm 0.0022 \text{ (corr.)} \\ \pm 0.0016 \text{ (norm.)} \pm 0.0008 \text{ (model)}$$

+ estimation of the uncertainty due to terms beyond NLO $\rightarrow \Delta\alpha_s(M_Z) = \pm 0.0050$

⇒ Precise determination $\alpha_s(M_Z) = 0.1183 \pm 0.0058$ from ZEUS data alone

→ in agreement with world average $\alpha_s(M_Z) = 0.1182 \pm 0.0027$ (Bethke, 2004)



Summary

- Due to the **precision and kinematic coverage** of the ZEUS measurements on structure functions and **jet cross sections**
⇒ **determination of α_s and the proton PDFs within a single experiment (ZEUS!) with minimal external input**
- Feasibility of including rigorously jet cross-section measurements in a global fit
- Advantages of including jet cross-section measurements which have been selected according to → **direct sensitivity to α_s and the gluon density in the proton and small experimental and theoretical uncertainties**
⇒ **sizeable reduction of the uncertainty of the gluon density in the mid- to high- x region**
⇒ **precise determination of α_s from HERA data alone**

$$\alpha_s(M_Z) = 0.1183 \pm 0.0028 \text{ (exp.)} \pm 0.0008 \text{ (model)} \pm 0.0050 \text{ (th.)}$$