

ZEUS PDF fits

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Published **GLOBAL ZEUS-S** fits to 30 pb⁻¹ of ZEUS 96/97 NC e⁺ differential cross-section data **and fixed target DIS structure function data**

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Central PDFs **and error analysis** available on
<http://durpdg.dur.ac.uk/hepdata/zeus2002.html>

as **eigenvector PDF sets** in **LHAPDF** compatible format

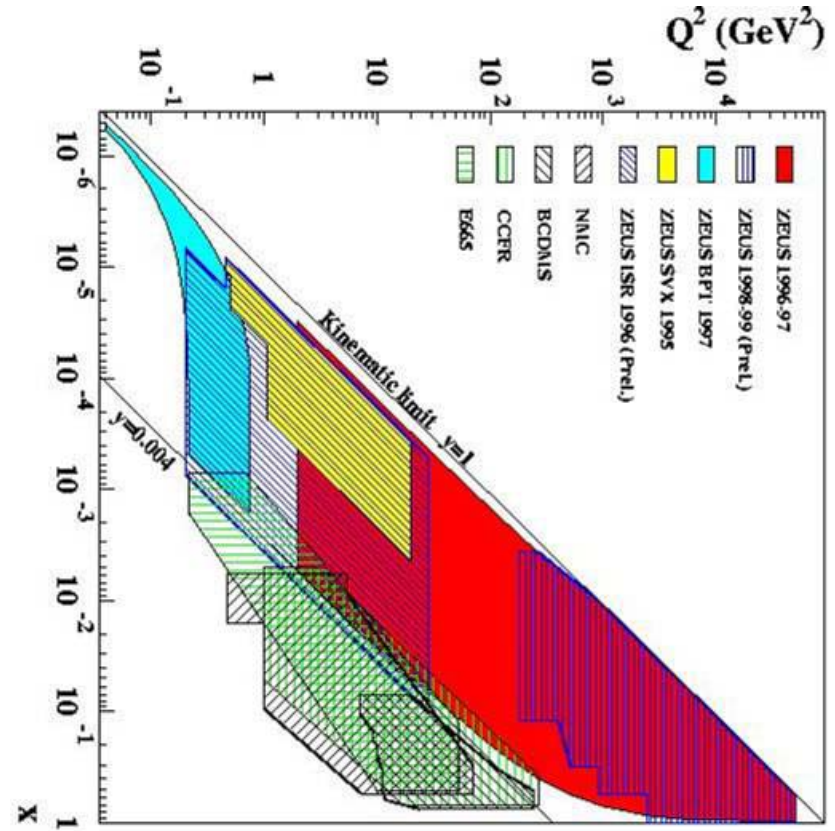
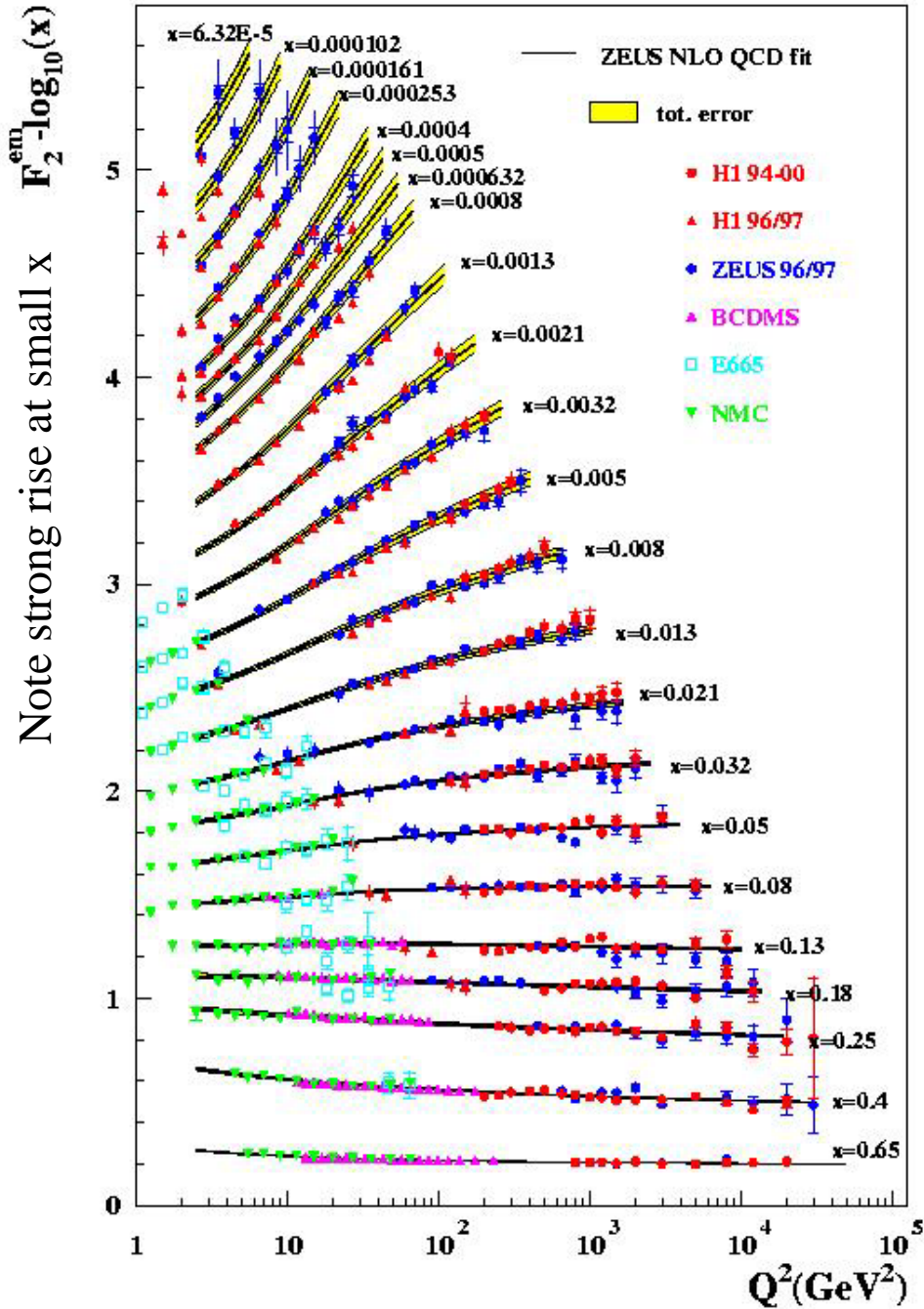
Preliminary **ZEUS-Only** fits to 109 pb⁻¹ of HERA-I data: **94-97 NC/CC e⁺/e⁻ inclusive differential cross-section data**

Proton target data from a single experiment

Discussion of ways of treating correlated systematic errors

Use of jet data as well as inclusive xsecs

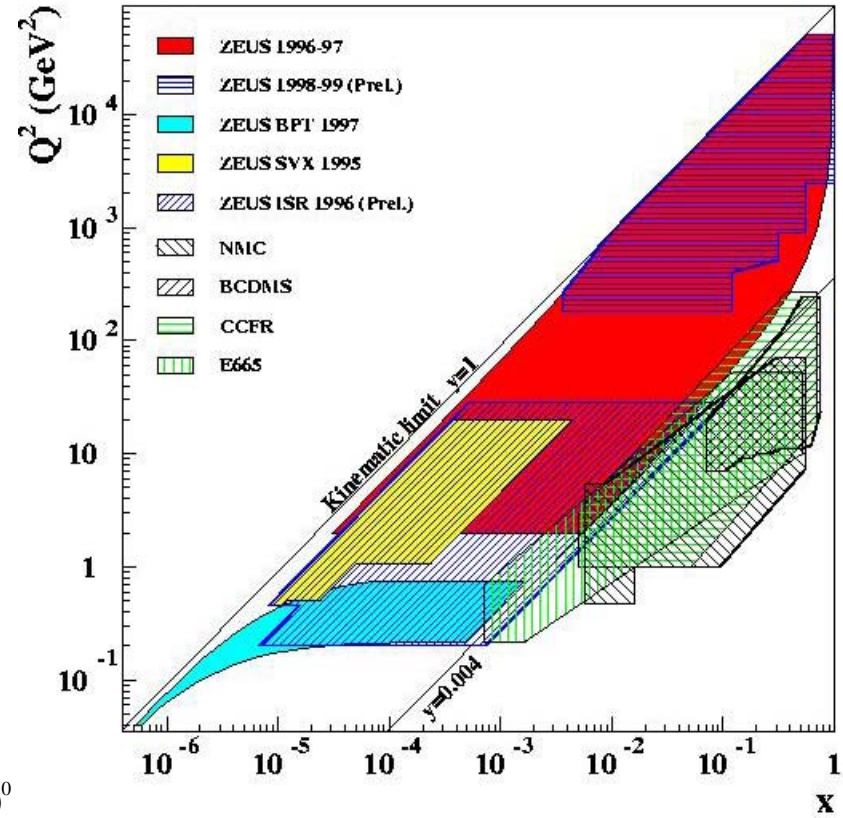
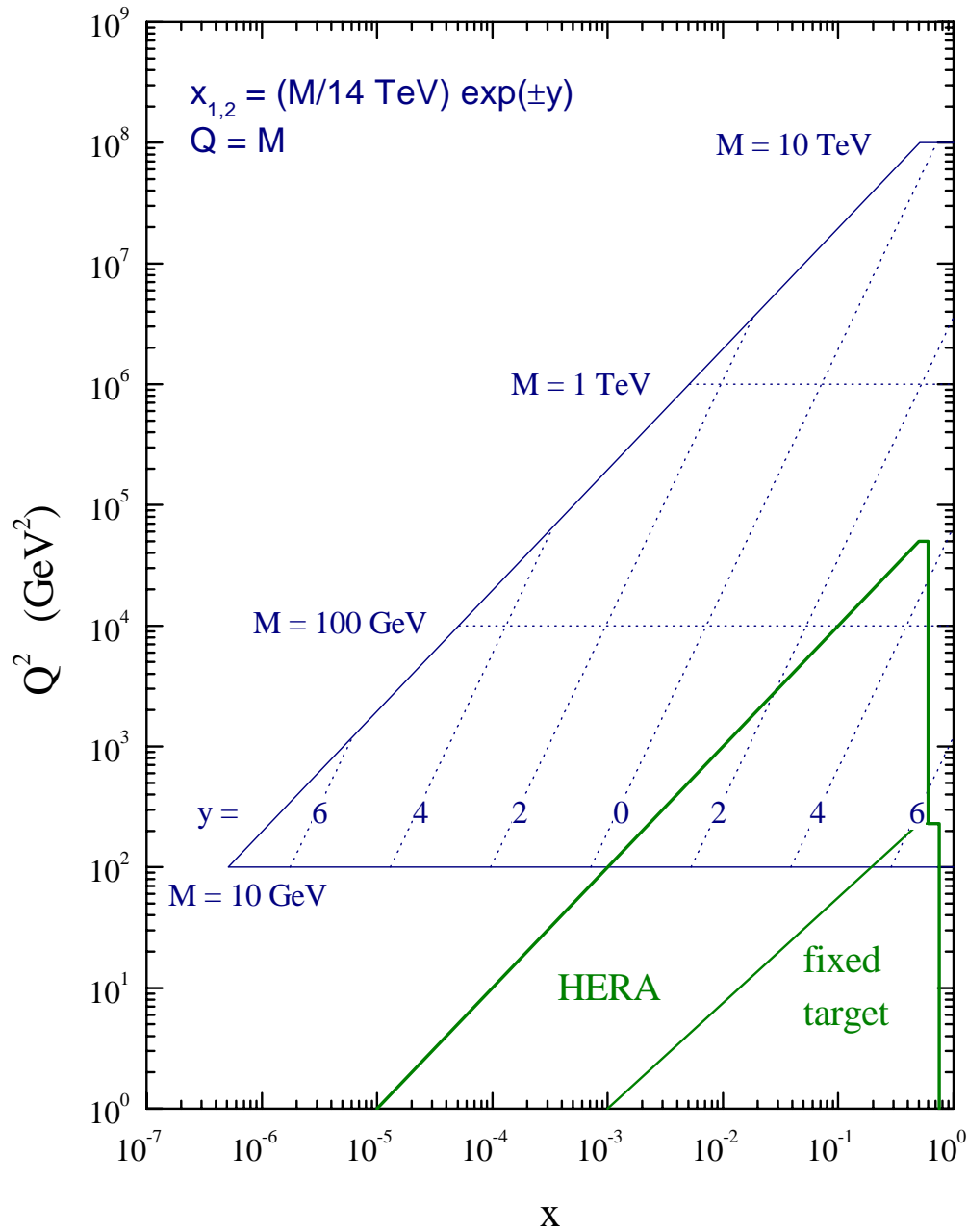
Note strong rise at small x



Terrific expansion in measured range across the x, Q^2 plane due to HERA data

Pre HERA fixed target $\bar{\mu}p, \mu D$
 NMC, BDCMS, E665 and $\nu, \bar{\nu}$ Fe CCFR

LHC parton kinematics



Parametrize parton distributions at Q^2_0

Evolve in Q^2 using NLO DGLAP (QCDNUM)

Convolute with coefficient functions \implies structure functions \implies cross-sections

Treatment of Heavy Quarks by Thorne-Roberts Variable Flavour Number

Cuts, $W^2 > 20$ (to remove higher twist), $30,000 > Q^2 > 2.7$, $x > 6.3 \cdot 10^{-5}$

χ^2 fit to 1263 data points \implies errors on params \implies errors on extracted PDF shapes, predicted structure functions and cross-sections

Accounting for correlated systematic errors by Offset method

- $x_{uv}(x) = A_u x^{a_v} (1-x)^{b_u} (1 + C_u x)$
- $x_{dv}(x) = A_d x^{a_v} (1-x)^{b_d} (1 + C_d x)$
- $x_S(x) = A_s x^{a_s} (1-x)^{b_s} (1 + C_s x)$
- $x_g(x) = A_g x^{a_g} (1-x)^{b_g} (1 + C_g x)$
- $x_{\Delta}(x) = A_{\Delta} x^{a_v} (1-x)^{b_s+2}$

\longleftarrow Model choices \implies Form of parametrization at Q^2_0 , value of Q^2_0 , flavour structure of sea, cuts applied, heavy flavour scheme

These parameters control the low-x shape

These parameters control the high-x shape

These parameters control the middling-x shape

A_u, A_d, A_g are fixed by the number and momentum sum-rules

Treatment of correlated systematic errors

$$\chi^2 = \sum_i \frac{[F_i^{\text{QCD}}(\mathbf{p}) - F_i^{\text{MEAS}}]^2}{(\sigma_i^{\text{STAT}})^2 + (\Delta_i^{\text{SYS}})^2}$$

Errors on the fit parameters, \mathbf{p} , evaluated from $\Delta\chi^2 = 1$,

THIS IS NOT GOOD ENOUGH if experimental systematic errors are correlated between data points- e.g. **Normalisations**

BUT there are more subtle cases- e.g. **Calorimeter energy scale/angular resolutions** can move events between x, Q^2 bins and thus **change the shape** of experimental distributions

$$\chi^2 = \sum_i \sum_j [F_i^{\text{QCD}}(\mathbf{p}) - F_i^{\text{MEAS}}] V_{ij}^{-1} [F_j^{\text{QCD}}(\mathbf{p}) - F_j^{\text{MEAS}}]$$

$$V_{ij} = \delta_{ij}(\sigma_i^{\text{STAT}})^2 + \sum_{\lambda} \Delta_{i\lambda}^{\text{SYS}} \Delta_{j\lambda}^{\text{SYS}}$$

Where $\Delta_{i\lambda}^{\text{SYS}}$ is the correlated error on point i due to systematic error source λ

It can be established that this is equivalent to

$$\chi^2 = \sum_i \frac{[F_i^{\text{QCD}}(\mathbf{p}) - \sum_{\lambda} s_{\lambda} \Delta_{i\lambda}^{\text{SYS}} - F_i^{\text{MEAS}}]^2}{(\sigma_i^{\text{STAT}})^2} + \sum_{\lambda} s_{\lambda}^2$$

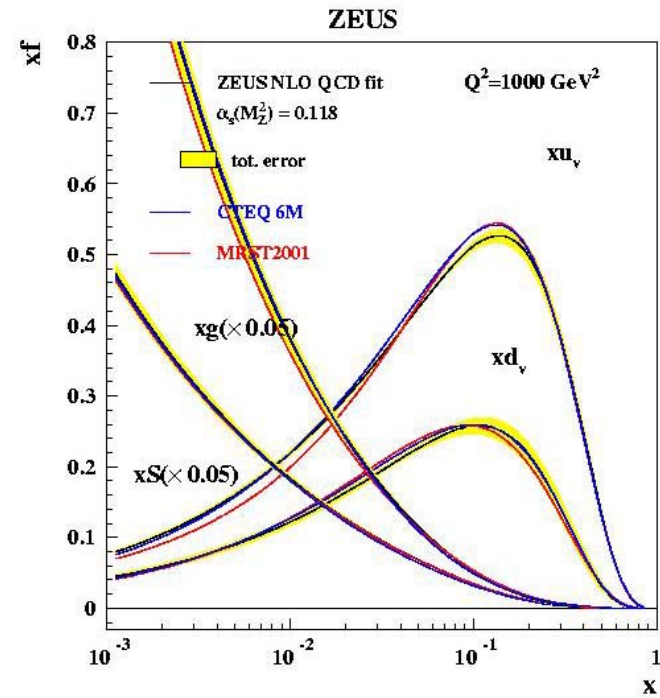
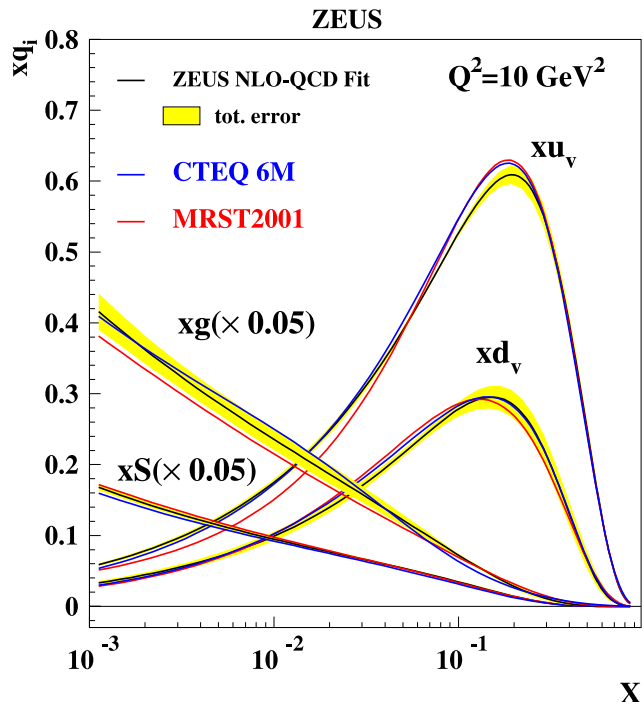
Where s_{λ} are systematic uncertainty fit parameters of zero mean and unit variance

This has modified the fit prediction by each source of systematic uncertainty

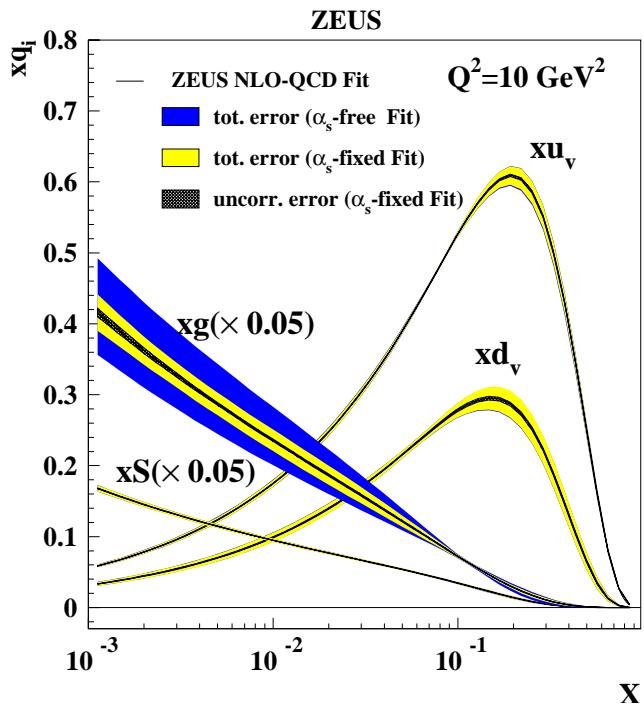
How do experimentalists usually proceed: OFFSET method

1. Perform fit without correlated errors ($s_\lambda = 0$) for central fit
2. Shift measurement to upper limit of one of its systematic uncertainties ($s_\lambda = +1$)
3. Redo fit, record differences of parameters from those of step 1
4. Go back to 2, shift measurement to lower limit ($s_\lambda = -1$)
5. Go back to 2, repeat 2-4 for next source of systematic uncertainty
6. Add all deviations from central fit in quadrature (positive and negative deviations added in quadrature separately)
7. This method does not assume that correlated systematic uncertainties are Gaussian distributed

Fortunately, there are smart ways to do this (Pascaud and Zomer LAL-95-05, Botje hep-ph-0110123)



Evolve in $Q^2 \Rightarrow$
low-x uncertainties
of sea/gluon
decrease



Value of α_s and shape of gluon are correlated in
DGLAP evolution

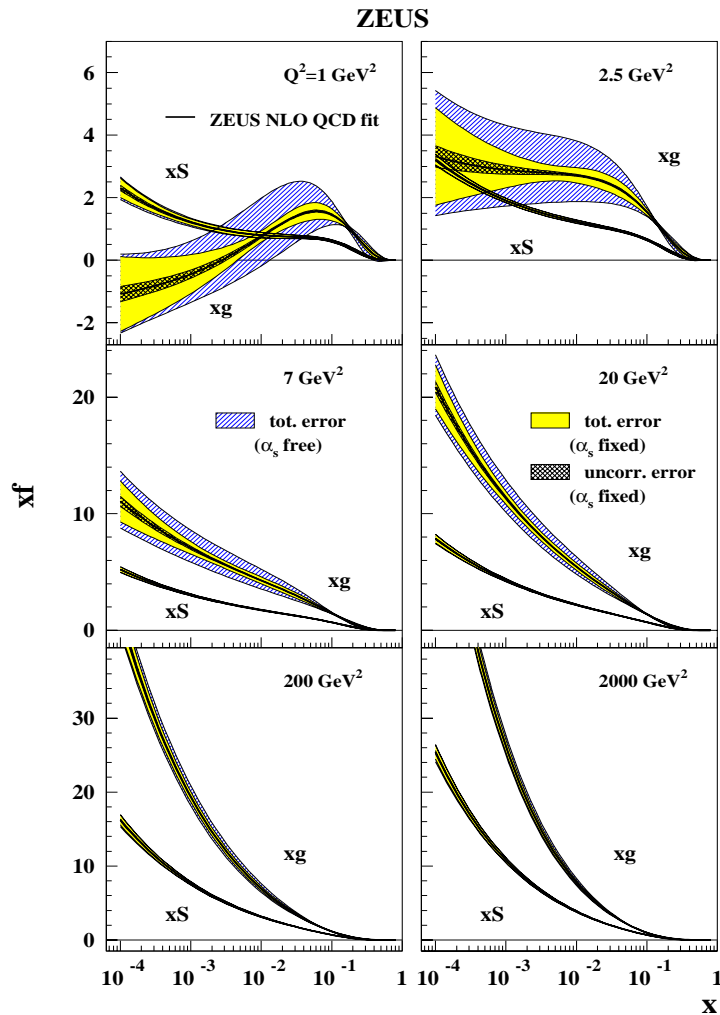
α_s increases \Rightarrow harder gluon

So fit α_s and PDF parameters
simultaneously
Uncertainty on gluon increases

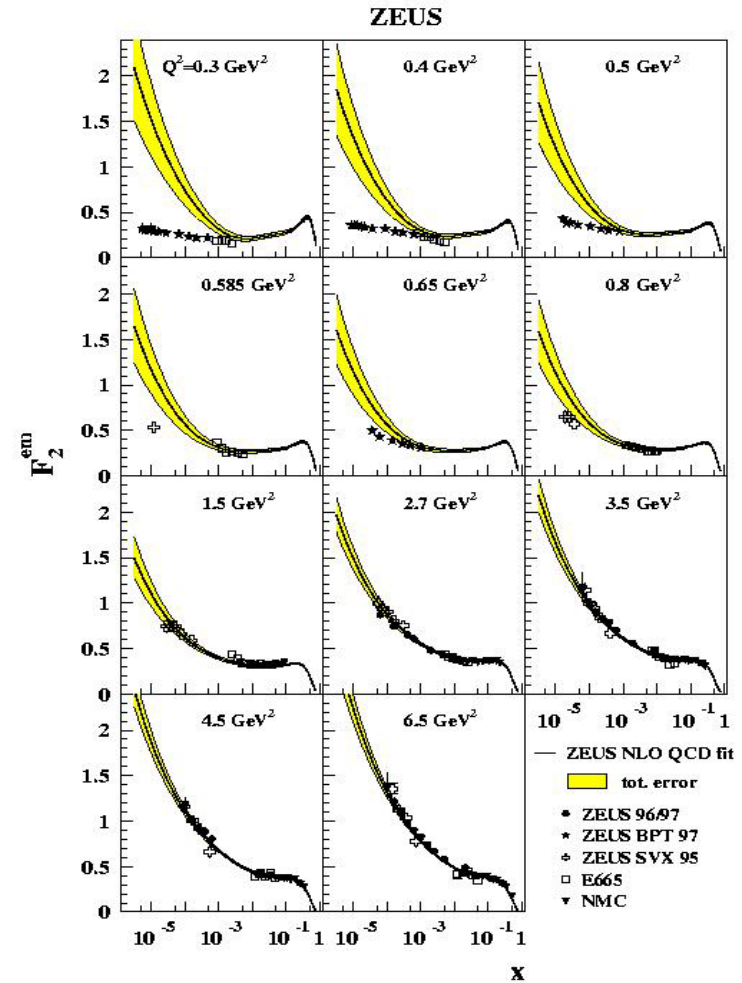
$$\alpha_s = 0.1166 \pm 0.0008 \pm 0.0032 \pm 0.0036 \pm 0.0018$$

stat. sys. norms. model

Look more closely at small-x



BUT below $Q^2 \sim 5 \text{ GeV}^2$ the gluon is no longer steep at small x – in fact its becoming valence-like, and then negative!



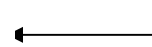
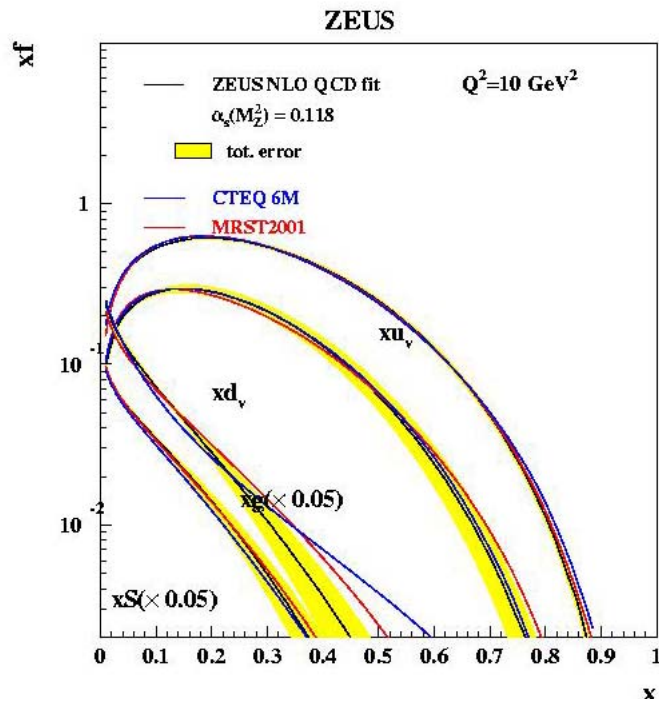
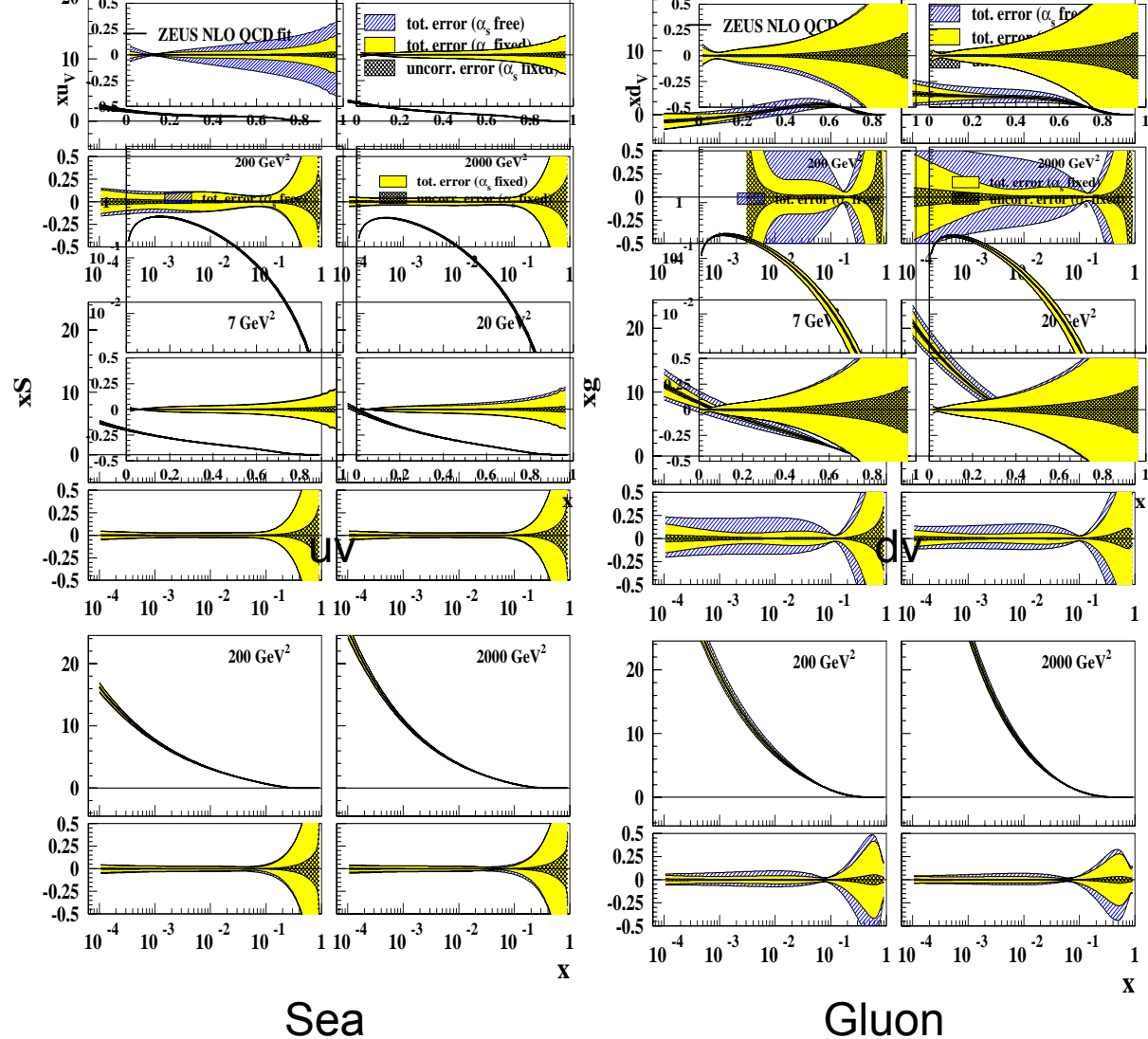
It was a surprise to see F_2 still steep at small x - even for $Q^2 \sim 1 \text{ GeV}^2$ should perturbative QCD work? α_s becoming large

Look more closely at high-x

u_v much better measured than d_v

Valence much better measured than sea/gluon

Uncertainties at high-x do not decrease so much with Q^2 evolution



Compare PDFs at high-x

High-x gluon \Rightarrow High ET jet production at Tevatron/LHC

There are other ways to treat correlated systematic errors- HESSIAN method (covariance method)

Allow $s\lambda$ parameters to vary for the central fit –there are smart ways to do this CTEQ
hep-ph/0101032

If we believe the theory why not let it calibrate the detector(s)? The fit determines the optimal settings for correlated systematic shifts.

The resulting estimate of PDF errors is much smaller than for the Offset method for $\Delta\chi^2 = 1$

We must be very confident of the theory – but more dubiously we must be very confident of the model choices we made in setting boundary conditions

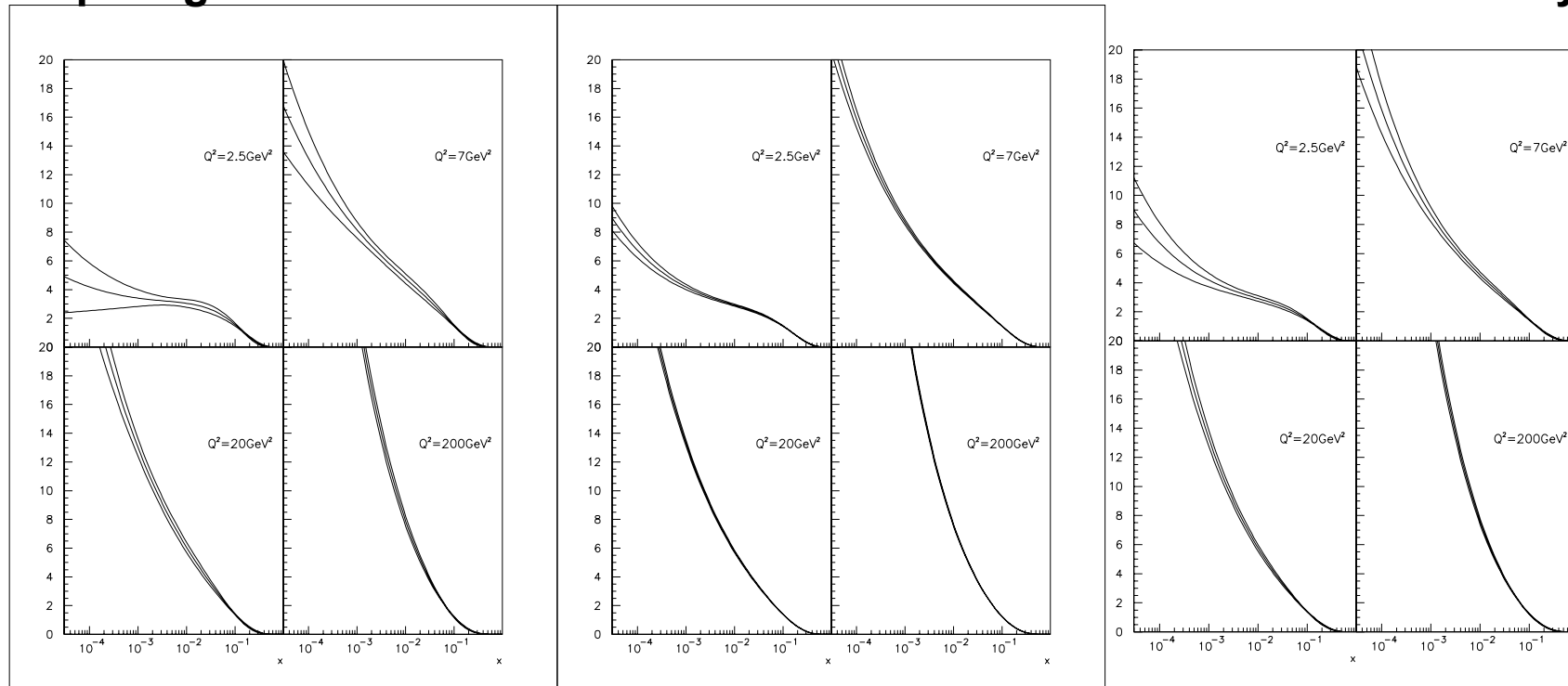
In a global fit the best fit parameters can be far from those which would be acceptable for some of the individual experiments- **data inconsistencies?**

One could **restrict the data sets** to those which are sufficiently consistent that these problems do not arise – (e.g. Giele, Keller, Kosover, FNAL)

But one **loses information** since partons need constraints from many different data sets – no single experiment has sufficient kinematic range / flavour info.

CTEQ use an increased χ^2 tolerance, $\Delta\chi^2 = T^2$, $T = 10$ to make an estimate of the PDF error which allows for this level of inconsistency in the data
MRST have also used increased tolerances in recent fits

Compare gluon PDFs for Hessian and Offset methods for the ZEUS fit analysis



Offset method

Hessian method $T=1$

Hessian method $T=7$

The Hessian method gives comparable size of error band as the Offset method, when the tolerance is raised to $T \sim 7$ – (similar ball park to CTEQ, $T=10$)

Note this makes the error band large enough to encompass reasonable variations of model choice since the criterion for acceptability of an alternative hypothesis, or model, is that χ^2 lie in the range $N \pm \sqrt{2N}$, where N is the number of degrees of freedom. For the ZEUS global fit $\sqrt{2N}=50$.

To do better investigate the possibility of using ZEUS data alone

Where does the information come from in a global PDF fit to DIS data?

Valence: from fixed target data – CCFR ν Fe xF_3 , NMC D/p ratio at high-x –
HEAVY target corrections

Sea: Low-x from HERA F2 data

High-x from fixed target F2 data

Gluon: Low-x from HERA $dF_2/d\ln Q^2$ data

High-x from mom-sum rule only- (UNLESS we put in JET DATA!)

Where does the information come from in a ZEUS-Only fit ?

Valence: HERA High- Q^2 cross-sections CC/NC $e^+/-$

Sea: Low-x from HERA F2 data

Gluon: Low-x from HERA $dF_2/d\ln Q^2$ data

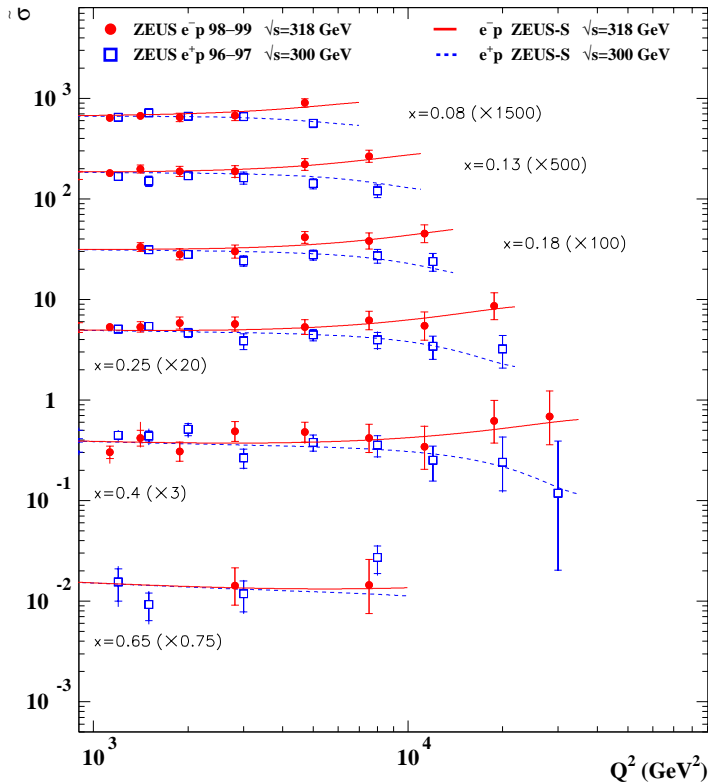
High-x from mom-sum rule only- (UNLESS we put in JET DATA!)

Advantages:

Pure proton target- no heavy target correction or deuterium corrections

Single experiment - correlated systematic errors well understood

ZEUS



Use ALL HERA-I data on NC/CC e+/e- high- Q^2 differential cross-sections $\sim 109\text{pb}^{-1}$, 509 data points

HERA at high $Q^2 \Rightarrow Z^0$ and $W^{+/-}$ exchanges become important

for NC processes

$$F_2 = \sum_i A_i(Q^2) [xq_i(x, Q^2) + x\bar{q}_i(x, Q^2)]$$

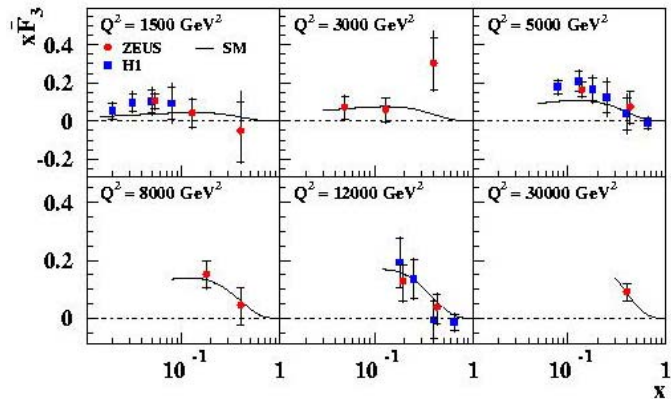
$$xF_3 = \sum_i B_i(Q^2) [xq_i(x, Q^2) - x\bar{q}_i(x, Q^2)]$$

$$A_i(Q^2) = e_i^2 - 2 e_i v_i v_e P_Z + (v_e^2 + a_e^2)(v_i^2 + a_i^2) P_Z^2$$

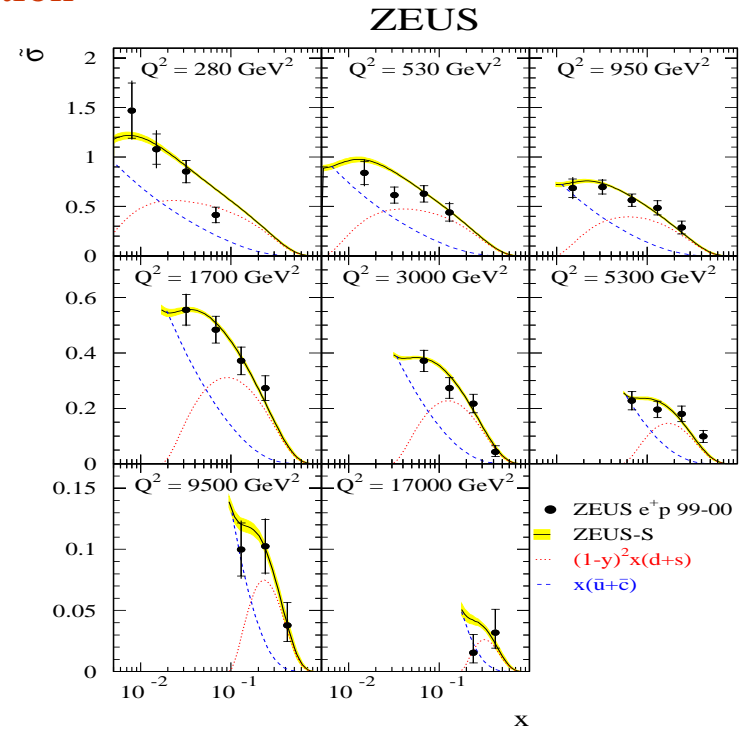
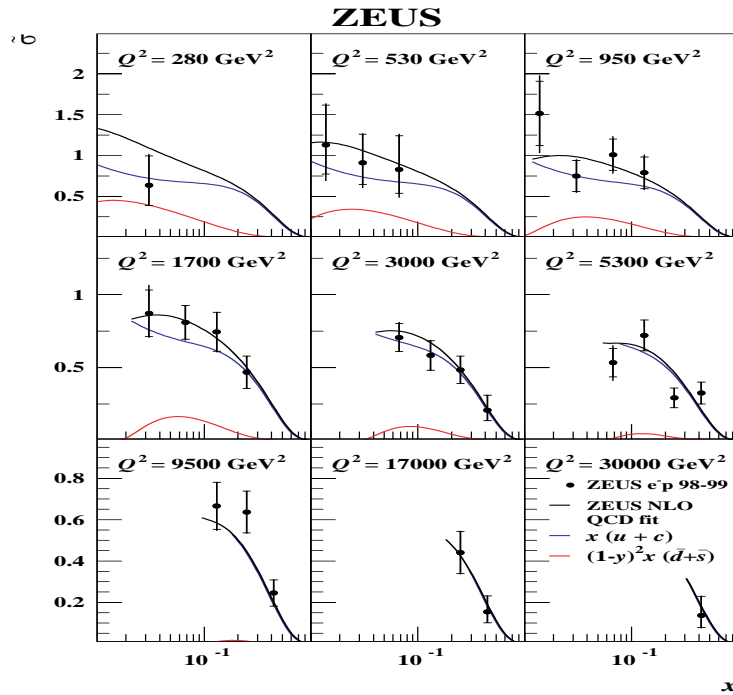
$$B_i(Q^2) = -2 e_i a_i a_e P_Z + 4a_i a_e v_i v_e P_Z^2$$

$$P_Z^2 = Q^2 / (Q^2 + M_Z^2) 1 / \sin^2 \theta_W$$

\Rightarrow Z exchange gives a new valence structure function xF_3 measurable from low to high x- on a pure proton target



CC processes give flavour information



$$\frac{d^2\sigma(e^-p)}{dx dy} = \frac{G_F^2 M_W^4}{2\pi x(Q^2 + M_W^2)^2} [x(u+c) + (1-y)^2 x(\bar{d} + \bar{s})]$$

M_W information

u_v at high x

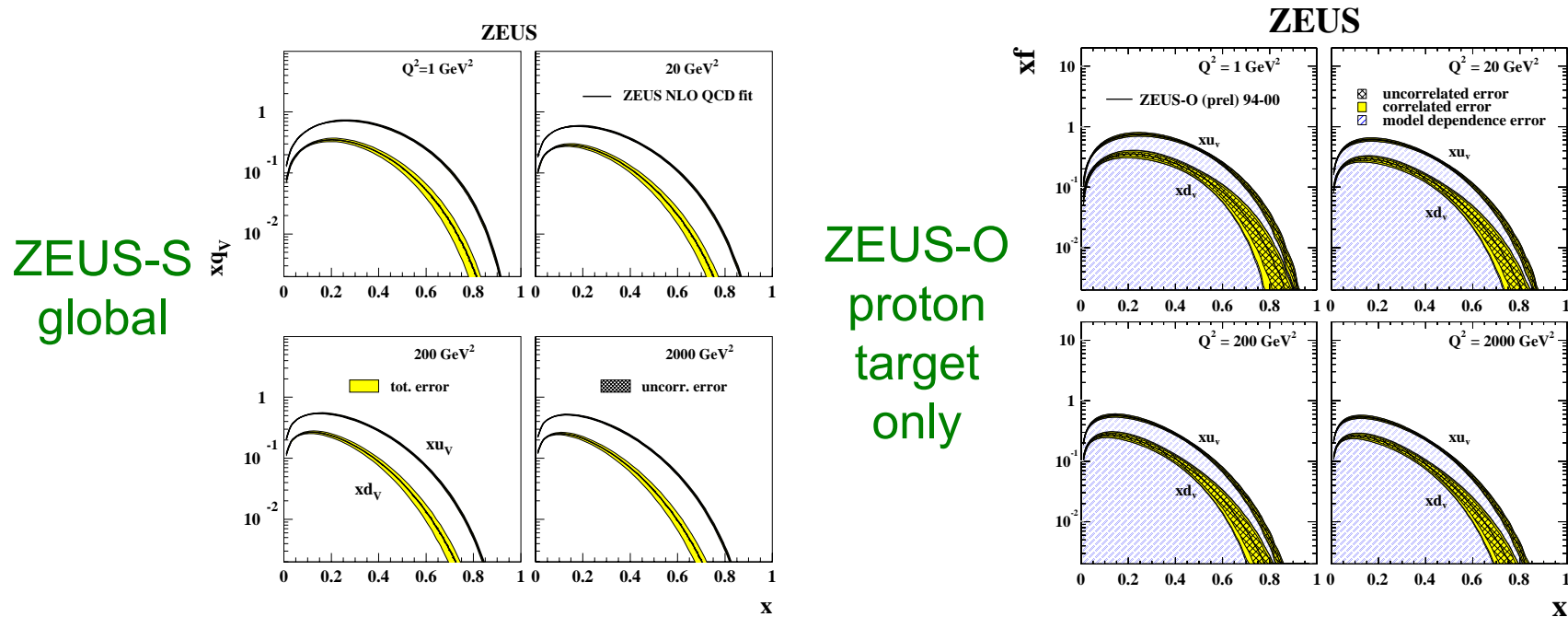
$$\frac{d^2\sigma(e^+p)}{dx dy} = \frac{G_F^2 M_W^4}{2\pi x(Q^2 + M_W^2)^2} [x(\bar{u} + \bar{c}) + (1-y)^2 x(d+s)]$$

d_v at high x

Measurement of high x , d -valence on a pure proton target. Most processes dominantly measure u -valence, only $\nu\text{Fe } xF_3$ and $\mu\text{D/p}$ give d -valence info.

And these are heavy target - even Deuterium needs corrections, does $d_v/u_v \rightarrow 0$, as $x \rightarrow 1$?

Compare valence PDFs for ZEUS-Only and ZEUS-S global fits

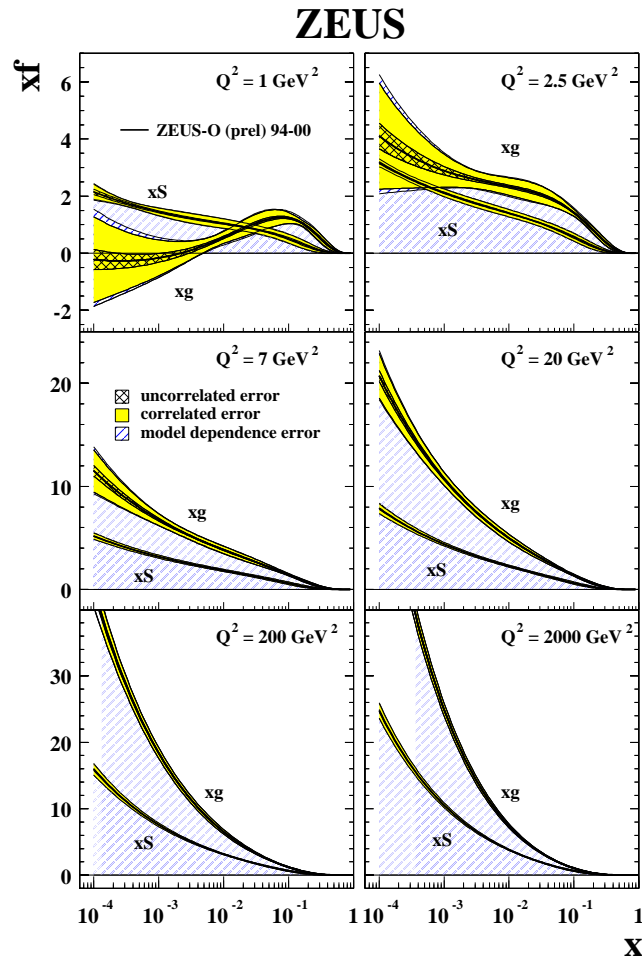


ZEUS-O fit precision is becoming competitive – and is on a proton target- statistical precision will improve with HERA-II data. ZEUS-S global fit precision is already systematics dominated

The precision on d_v is much worse than for u_v because most cross-sections measure u_v , but HERA high Q^2 CC e^+ measures d_v on a proton target

In HERA-II things can only get better! – more high Q^2 CC plus NC xF_3 data

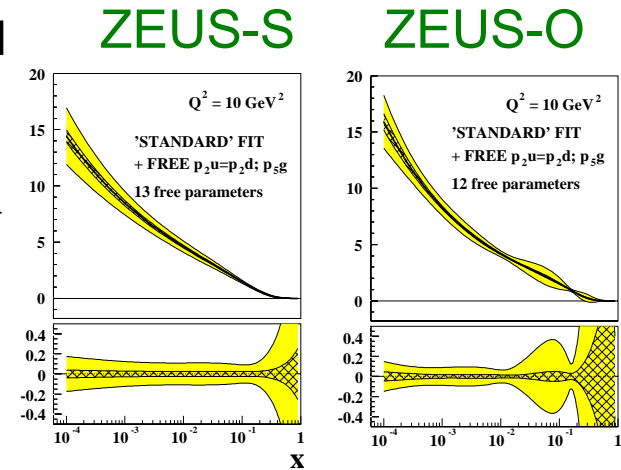
Compare sea and gluon PDFs for ZEUS-Only and ZEUS-S global fits



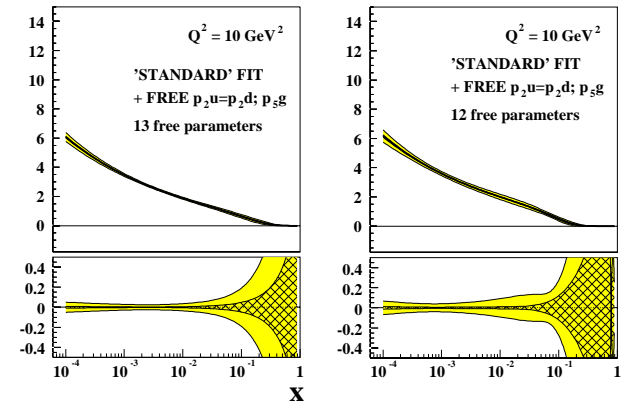
Low- x precision is comparable – info. In global fit came from ZEUS data
 High- x precision is worse.

Interim solution: simplify sea/glue high- x param.
 Eigenvector procedure gives information on fit stability and parameter correlations- tells you which params are constrained best/which you need

Long term solution: HERA-II data
 Medium term solution: use ZEUS jet data from HERA-I \Rightarrow impacts on gluon $0.01 < x < 0.1 \Rightarrow$ via momentum sum-rule on higher- x gluon



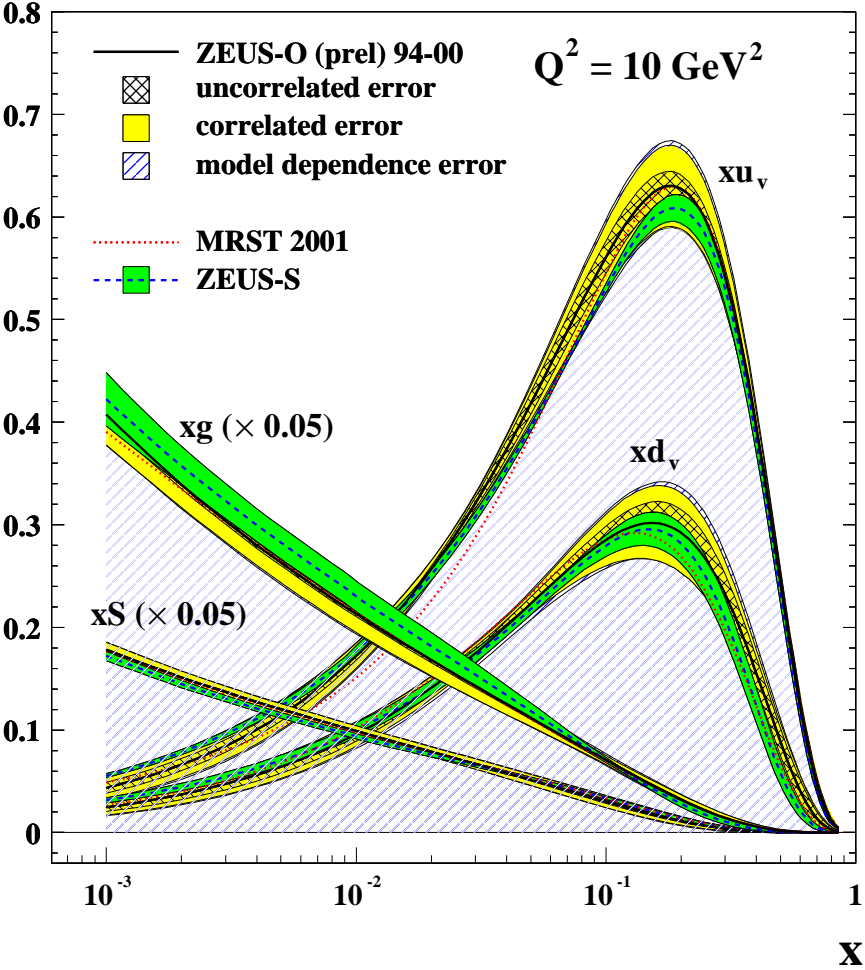
Glue



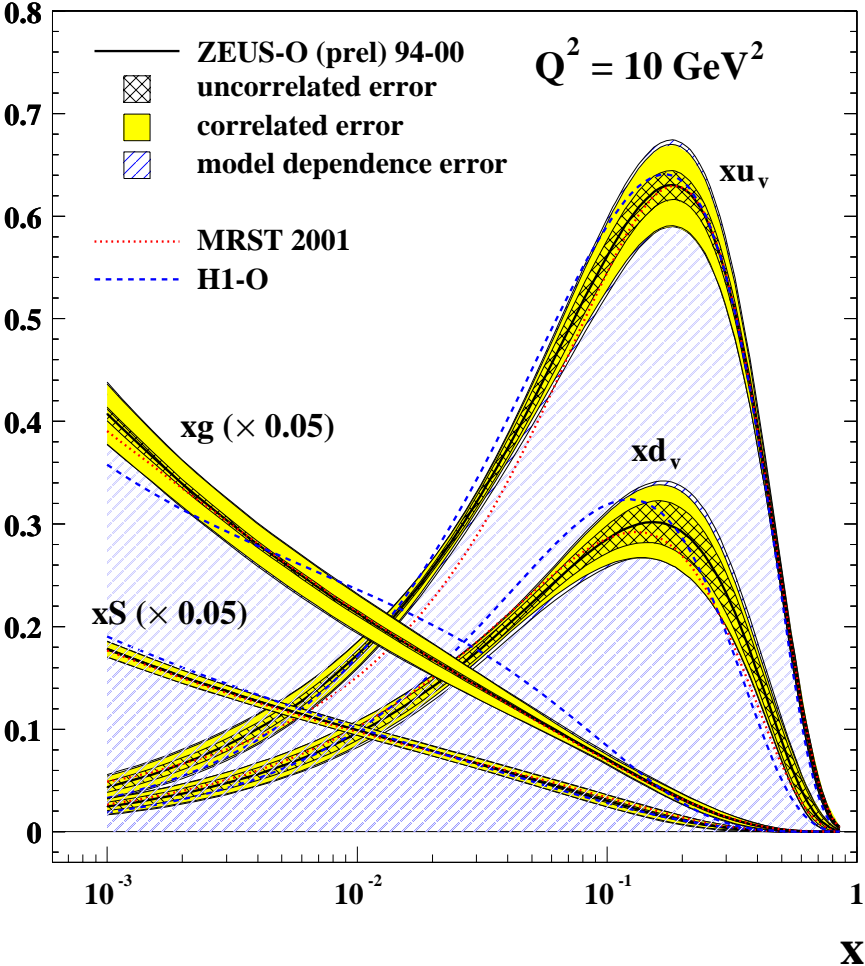
Sea

Interim solution: Compare HERA-I ZEUS-Only PDFs extracted from inclusive cross-section data to published ZEUS-S Global PDFs, and to MRST and H1 PDFS

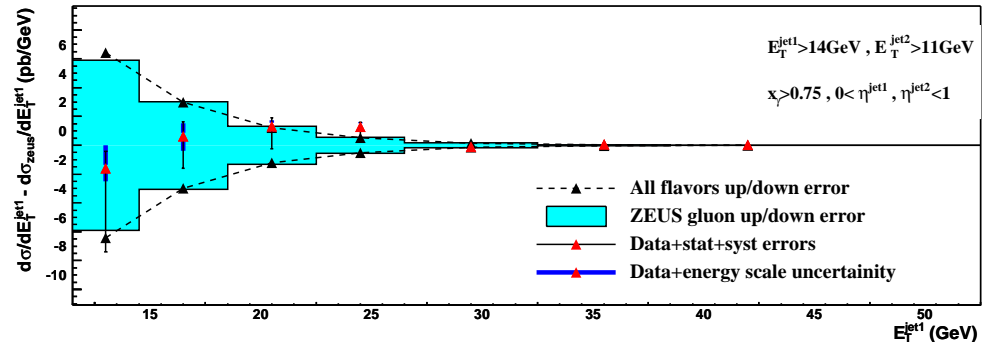
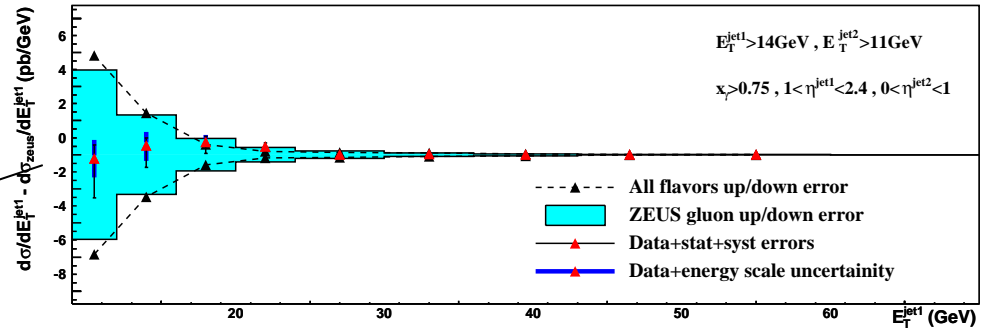
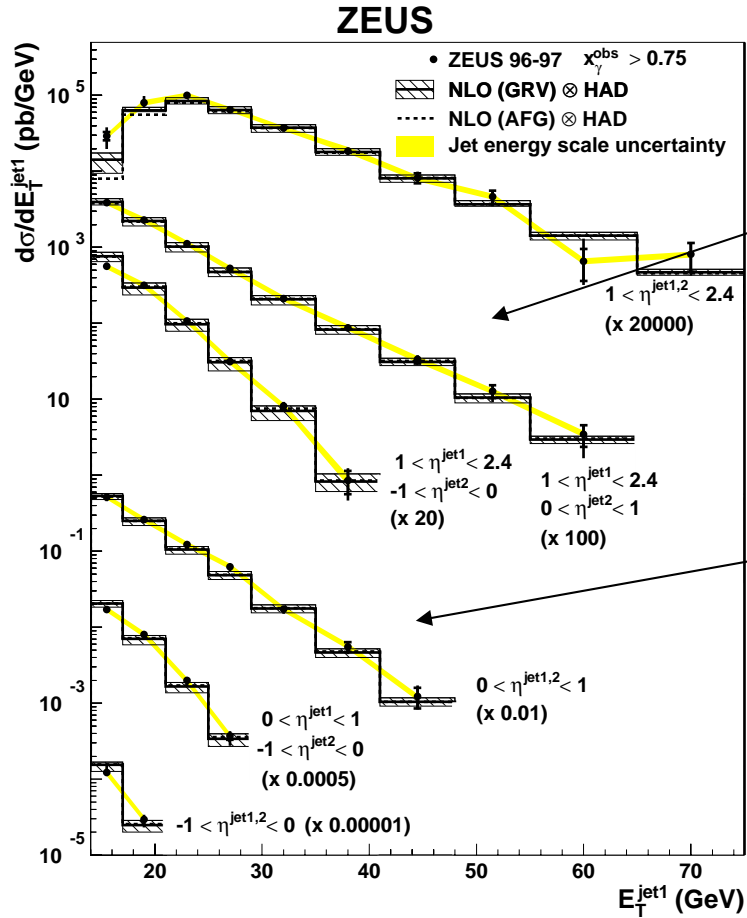
ZEUS



ZEUS



Mid term solution



Photoproduction dijet cross-sections vs ET

Have significant impact on the uncertainties of the ZEUS-Only fit

Many jet cross-sections can be exploited to improve gluon measurement pre-LHC