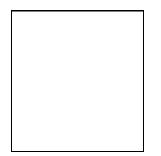
### VECTOR MESON PRODUCTION AT HERA

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At the HERA collider the experiments H1 and ZEUS have studied the exclusive production of vector mesons over a wide kinematical range. The recent measurements and their discussion within the framework of color dipole models and pQCD are reported.

## 1 Introduction

Exclusive photo- and electroproduction of light  $(\rho, \omega, \phi)$  and heavy  $(J/\psi, \psi', \Upsilon)$  vector mesons (VM) have been subject of intensive studies at HERA. The accelerator and its general purpose detectors H1 and ZEUS provide a unique opportunity to measure the exclusive diffractive production of vector mesons with different masses  $M_{VM}$  in photo- and electroproduction.

#### 1.1 Diffractive vector meson production

The process  $ep \to e(VM)p$ , drawn in fig. 1, can be described as a two step process. The incoming electron emits a photon. This photon fluctuates into a  $q\bar{q}$  state which scatters with the proton by exchanging nothing but momentum.

The kinematical variables which are used to characterize the process are: the 4-momentum transfer squared  $Q^2$  at the electron vertex, the center of mass energy of the  $\gamma$ -proton system W and the 4-momentum transfer squared t at the proton vertex. The dependences of the cross sections on these variables are presented in the ranges:  $2 < Q^2 < 100 \ GeV^2$ ,  $30 < W < 260 \ GeV$  and  $|t| < 1 \ GeV^2$ .

In the absence of a hard scale the colorless exchange between the vector meson and the proton can be modeled by a soft pomeron trajectory using the Regge approach. In this approach the cross section is predicted to rise slowly with W. In the presence of a hard scale, the vector meson production can be calculated using pertubative QCD (pQCD). In this case the colorless

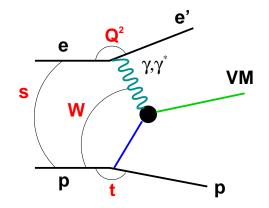


Figure 1: Diagram of exclusive vector meson (VM) production at HERA.

exchange is modeled in leading order by a pair of gluons, the cross section is proportional to the square of the gluon density. This predicts a steep rise with increasing values of W.

## 2 W dependence

For exclusive photoproduction  $(Q^2 \approx 0 \ GeV^2)$  of vector mesons the W dependence is shown in fig.2a). The lines indicate the rising of the cross sections assuming the form  $W^{\delta}$ . For the light vector mesons  $\rho$ ,  $\omega$  and  $\phi$  the slope is  $\delta \approx 0.22$ . This value is very similar to the total photoproduction cross section and is predicted by the Regge approach. For the heavier vector mesons  $J/\psi$ ,  $\psi(2s)$  and  $\Upsilon$  the observed slope is higher ( $\delta \gtrsim 0.8$ ).

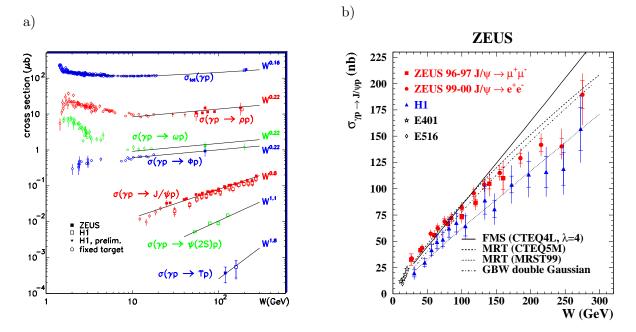


Figure 2: a) Compilation of W dependences for exclusive photoproduction of light and heavy vector mesons. The lines show the behaviour of the cross section assuming the form  $W^{\delta}$ . The values for  $\delta$  are given at the right edge. b) W dependence of the  $J/\psi$  photoproduction cross section in comparison with theoretical models based on pQCD.

For the heavy vector mesons the masses of the charm and the bottom quarks provide a hard scale which allows the use of pQCD to calculate the cross sections. Such models are able to describe the  $J/\psi$  <sup>1,2</sup> cross section as it is shown in fig.2b). In particular the steep rise as a

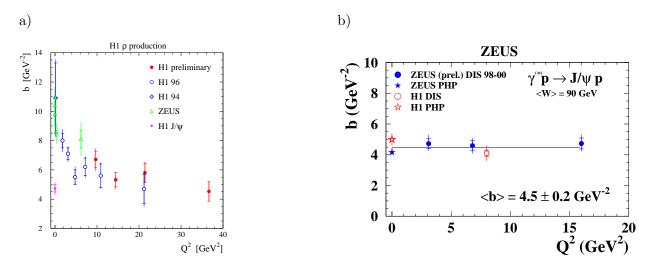


Figure 3: a) Fitted values of the *b*-slope as a function of  $Q^2$  in  $\rho$  photo- and electroproduction. b) Fitted values of the *b*-slope as a function of  $Q^2$  in  $J/\psi$  photo- and electroproduction.

function of the energy W is well described by  $pQCD^3$ .

# **3** $Q^2$ and |t| dependence

The |t| dependence of the cross section for exclusive vector meson production is well described by the form  $\frac{d\sigma}{dt} \propto e^{-b|t|}$  for small values of t ( $|t| < 1 \ GeV^2$ ). Fig.3a+b) show the dependence of the slope b as a function of  $Q^2$  for  $\rho^4$  and  $J/\Psi^5$  in photo- and electroproduction. The slope of the  $\rho$  cross section decreases with increasing  $Q^2$ . This indicates that the size of the interacting region is changing with  $Q^2$ . On the right hand side one can see that, in contrast to the  $\rho$ , the  $J/\psi$  has is no change in the slope with  $Q^2$ . The production mechanism for  $J/\psi$  at the photoproduction limit is already the same as in the higher  $Q^2 > 0$  range. This is interpreted as due to the fact that the  $J/\psi$  mass already provides a hard scale at  $Q^2 = 0$ , in contrast to exclusive  $\rho$  electroproduction.

In the diffractive picture b is related to the radii of the colliding objects i.e. of the proton and the VM:  $b \propto r_{VM}^2 + R_{\text{proton}}^2$ . The values of  $b \approx 4.5 \ GeV^{-2}$  measured in the hard regime implies a combined radius of the order of the size of the proton. This observation suggests that the transverse size of the  $q\bar{q}$  fluctuation producing the  $J/\psi$  in photoproduction and the  $\rho$  at high  $Q^2$  is smaller than that of the proton.

### 4 Decay angular distributions

The production and decay of a VM into a pair of oppositely charged particles can be described in terms of three angles:  $\Phi_h$ , the angle between the VM production plane and the lepton scattering plane;  $\theta_h$  and  $\phi_h$ , the polar and azimuthal angles of the positively charged decay lepton in the *s*-channel helicity frame.

Under the assumption of s-channel helicity conservation (SCHC), the angular distribution for the decay of the VM depends only on two angles,  $\theta_h$  and  $\psi_h = \phi_h - \Phi_h$ . From the  $\theta_h$  distribution the spin density matrix element  $r_{00}^{04}$  can be extracted which is proportional to the the helicity amplitude  $T_{00}$ .  $T_{00}$  corresponds to an amplitude, where a longitudinally polarized photon yields a longitudinally polarized VM. Also under the assumption of SCHC the ratio R of cross sections for longitudinally and transverse polarized photons can be calculated using  $r_{00}^{04}$ ;  $(R = \frac{1}{\epsilon} \frac{r_{00}^{04}}{1-r_{00}^{04}})$ . This measurement for  $\rho$  electroproduction is shown in fig.4). The left plot shows the rise of R;

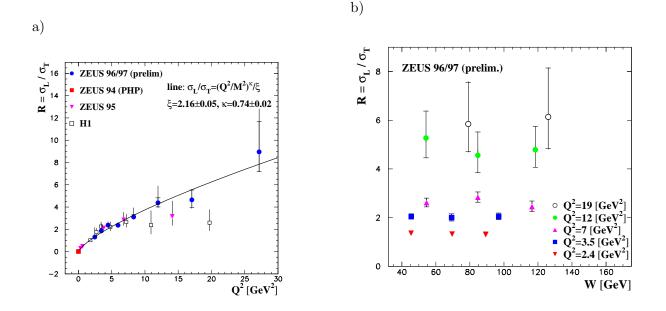


Figure 4: a)  $Q^2$  dependence of R in  $\rho$  electroproduction. b) W dependence of R in  $\rho$  electroproduction in different bins of  $Q^2$ .

it indicates that the longitudinal dipole configuration becomes more dominant when increasing the value of  $Q^2$ . Within pQCD models a small dipole is most likely to be produced if the virtual photon is longitudinally polarized, which predicts that  $|T_{00}|$  is the dominant amplitude.

Fig.4b) shows the ratio R in  $\rho$  electroproduction for different values of  $Q^2$  regions as a function of W. The data is consistent with no W dependence.

### Acknowledgments

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