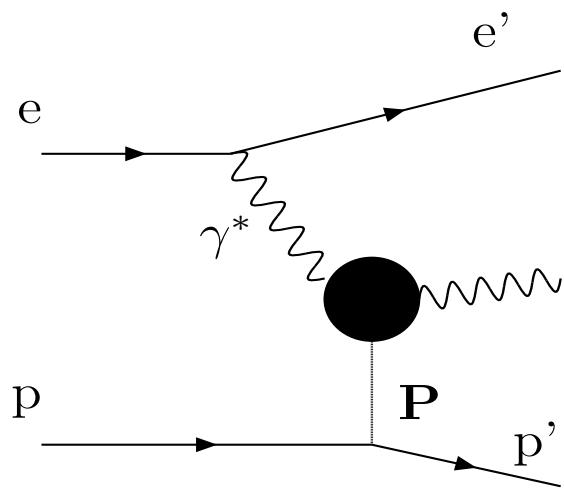


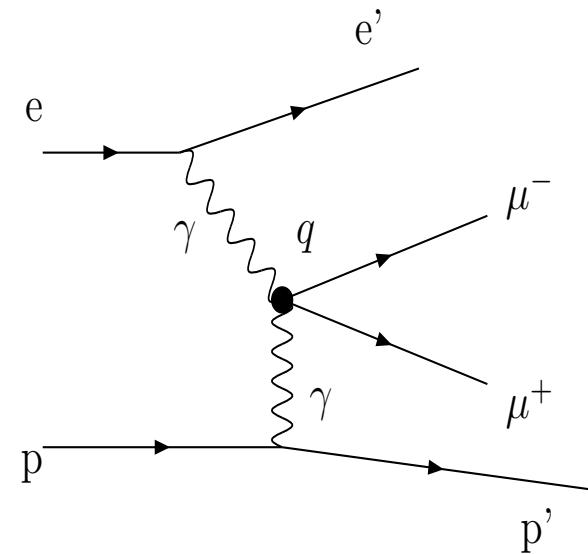
*Deeply Virtual Compton Scattering
and
Light Cone Wave Function of the Photon*

Justyna Ukleja, Warsaw University

for the ZEUS Collaboration



DVCS



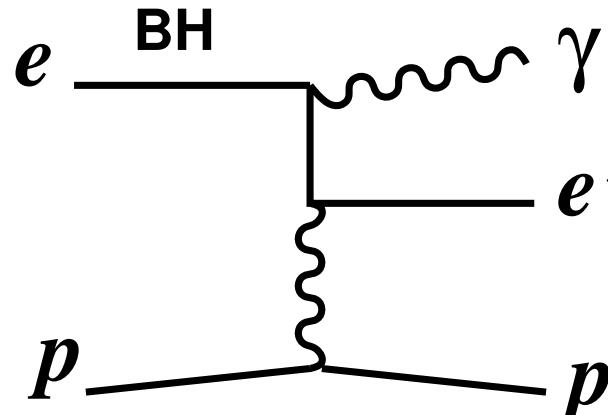
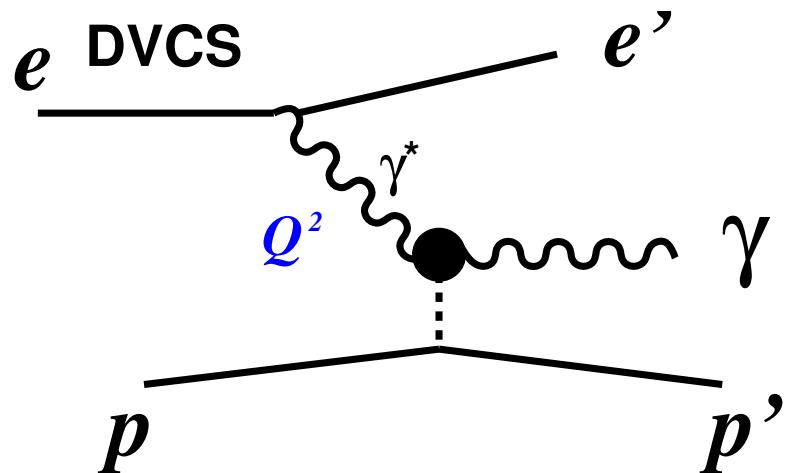
LCWF

Deeply Virtual Compton Scattering

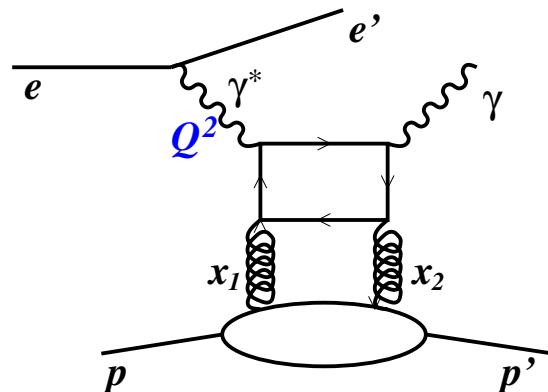
Deeply Virtual Compton Scattering (DVCS)

→ diffractive scattering of the virtual photon

- Interference between QCD with QED (BH) amplitudes → rich structure in ϕ , angle between the hadronic and leptonic planes → asymmetries (angular, charge)
→ gives access to the Generalized Parton Distributions GPDs



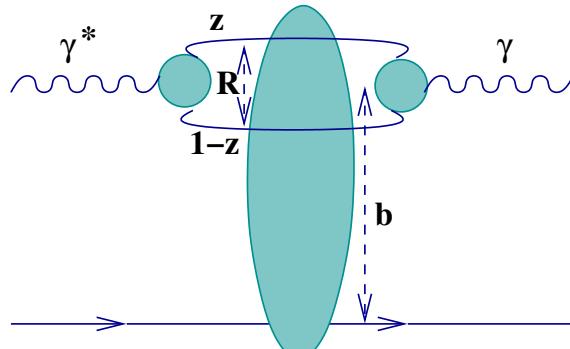
QCD Models for DVCS



QCD-based Model - Frankfurt, Freund and Strikman (FFS):

$$\frac{d^3\sigma_{\text{DVCS}}}{dx dQ^2 dt} = \frac{\pi^2 \alpha^3}{2x R^2 Q^6} [1 + (1-y)^2] e^{bt} F_2^2(x, Q^2) (1 + \rho^2)$$

NLO calculation Freund, McDermott



Colour Dipole Model - Donnachie and Dosch (DD), Frankfurt, Guzey and Kerley and Shaw (FKS), McDermott, Frankfurt, Guzey and Strikman (MFGS)

$$\mathcal{A} \sim \int_{R,z} \psi_{in}^{\gamma^*} \sigma_D \psi_{out}^{\gamma}$$

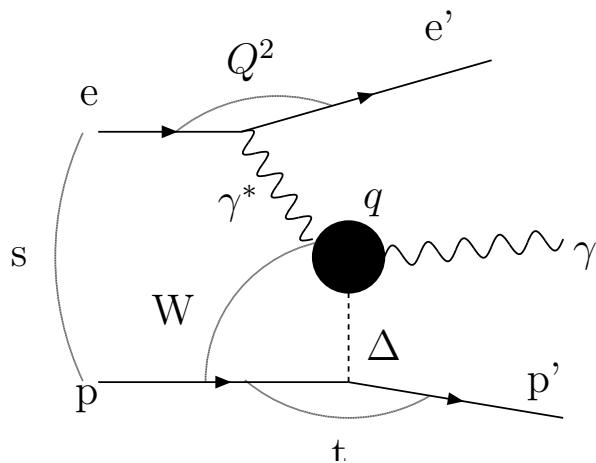
σ_D - dipole- p cross section

$\psi_{in}^{\gamma^*}$ - light-cone wave function of the incoming photon

ψ_{out}^{γ} - light-cone wave function of the outgoing photon

DVCS - event selection

ZEUS: $e^+ p$ sample, 95 pb^{-1} , $e^- p$ sample 17 pb^{-1}



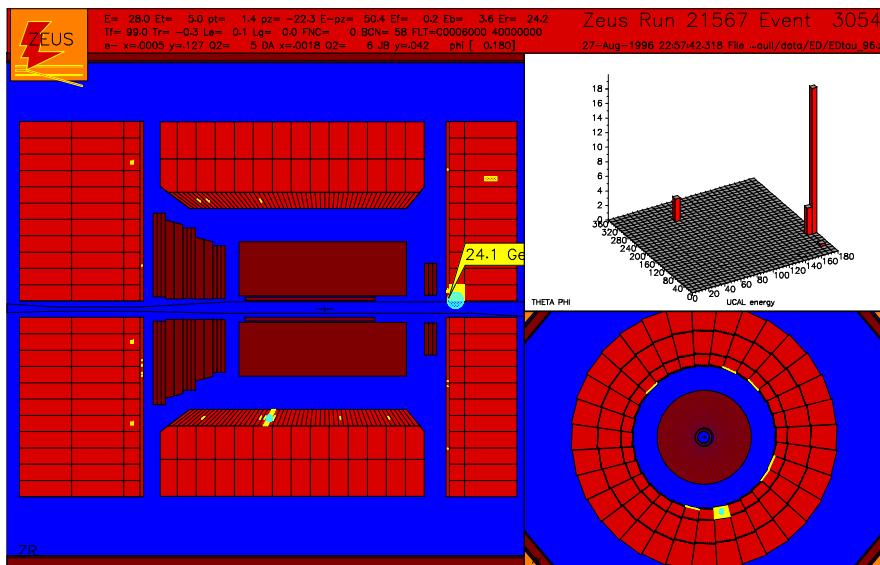
HERA

$$\sqrt{s} = 300, 320 \text{ GeV}$$

$$5 < Q^2 < 100 \text{ GeV}^2$$

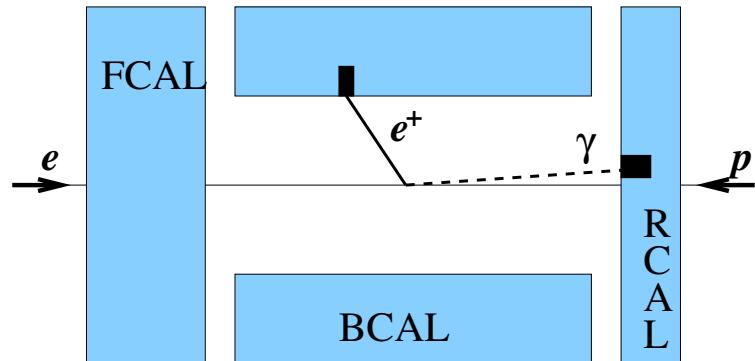
$$40 < W < 140 \text{ GeV}$$

$$10^{-4} < x < 10^{-2}$$



- 2 EM clusters
- 0 or 1 Track; if 1, must be matched to EM cluster
- No other activity above noise level to reject events where the proton dissociates

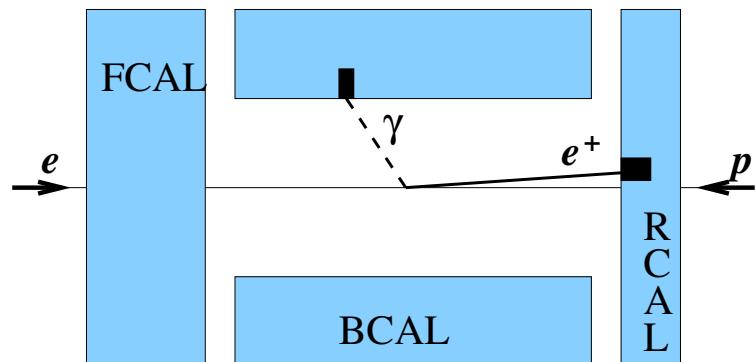
DVCS - event topology



γ from QED bremsstrahlung emitted
in e -beam direction

e^+ sample CONTROL SAMPLE

only BH contributes, 7000 events



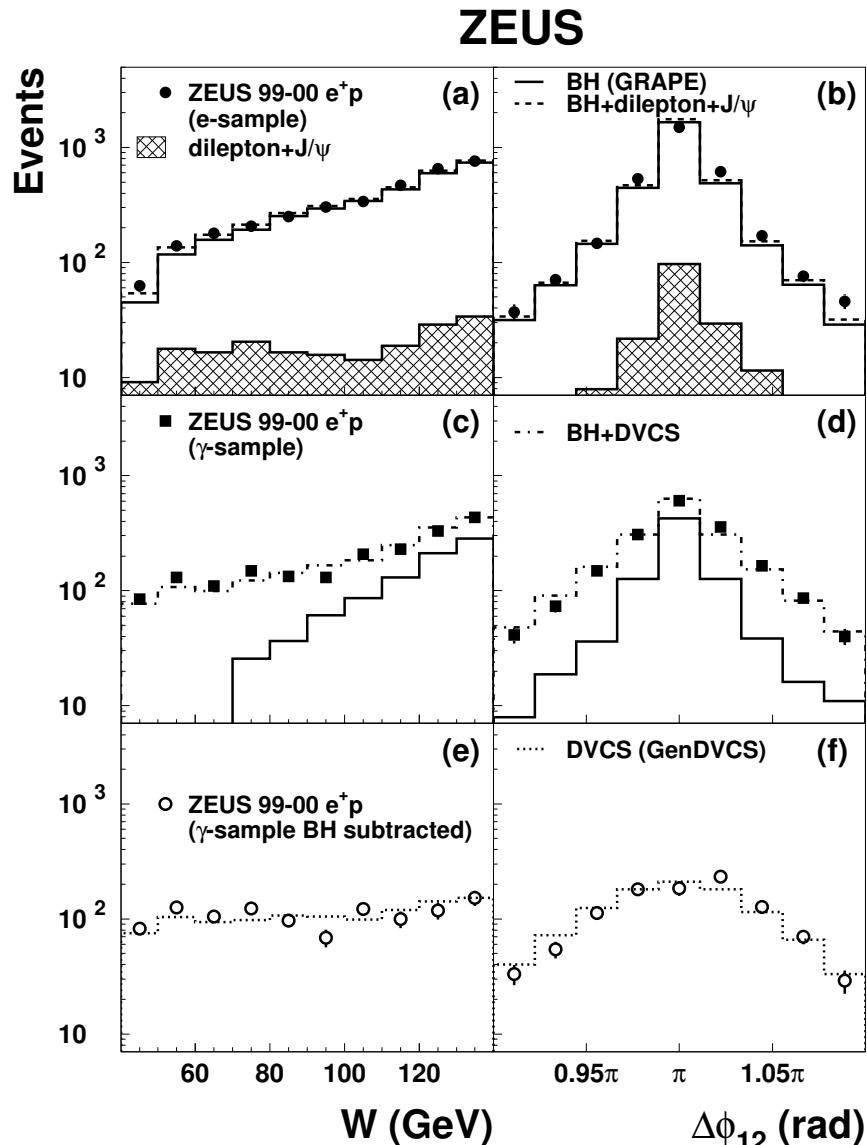
γ sample ENRICHED WITH DVCS

both BH and DVCS contribute, 4000 events

BH MC describes data in control sample

→ use MC to subtract BH in enriched DVCS sample

DVCS - control plots

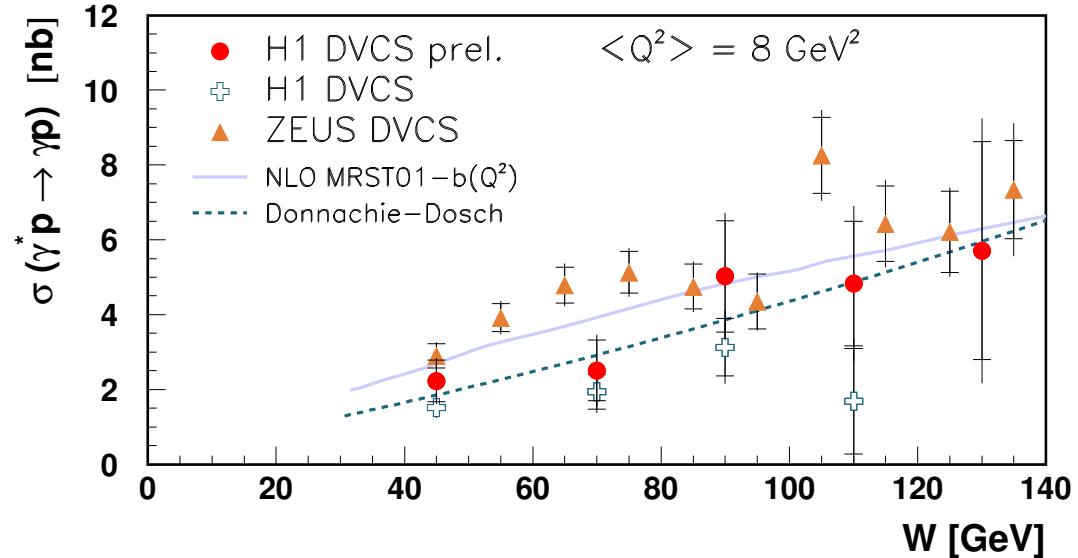
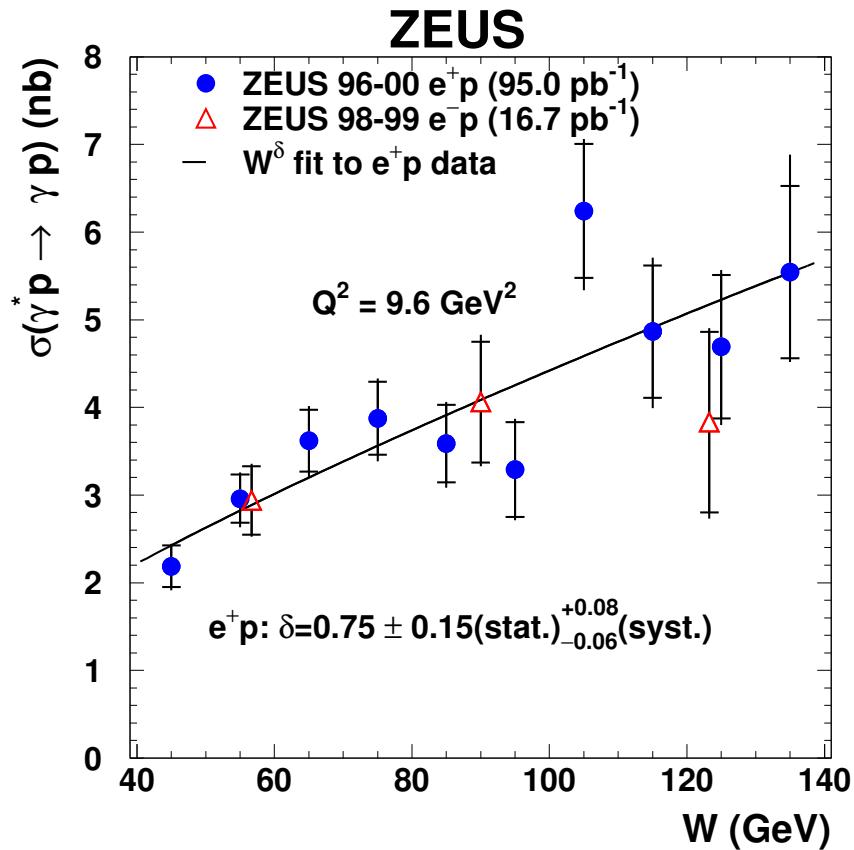


e^+ sample BH control sample

γ sample DVCS + BH

γ sample after BH subtraction DVCS

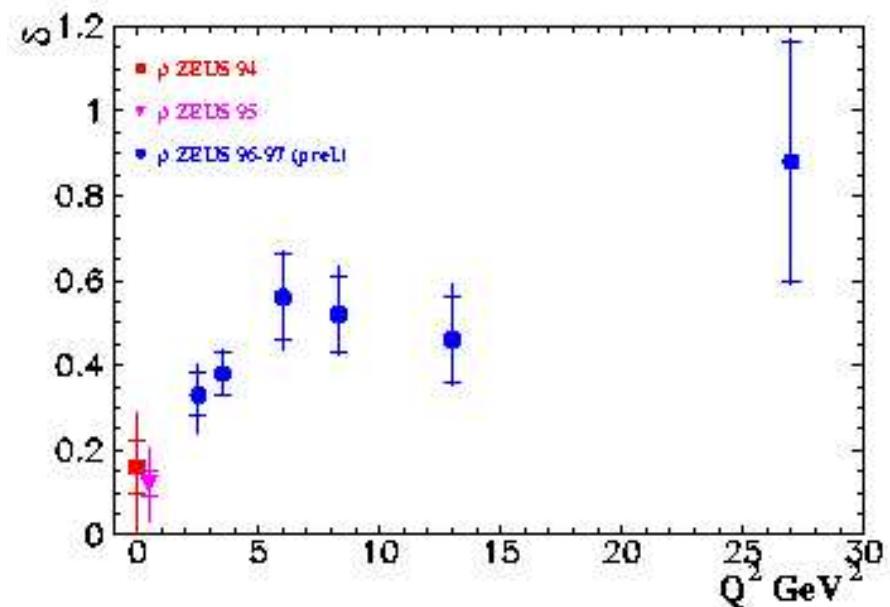
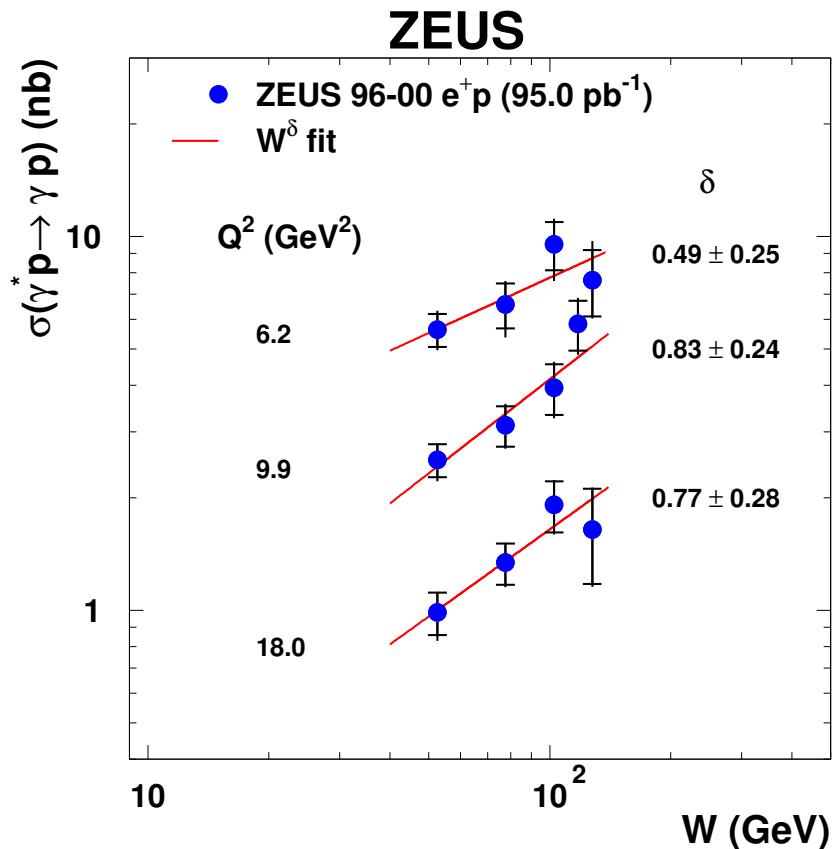
DVCS - W dependence



steep rise of σ with W
 → DVCS is a hard process

No differences between $e^+ p$
 and $e^- p$ samples

DVCS - W dependence

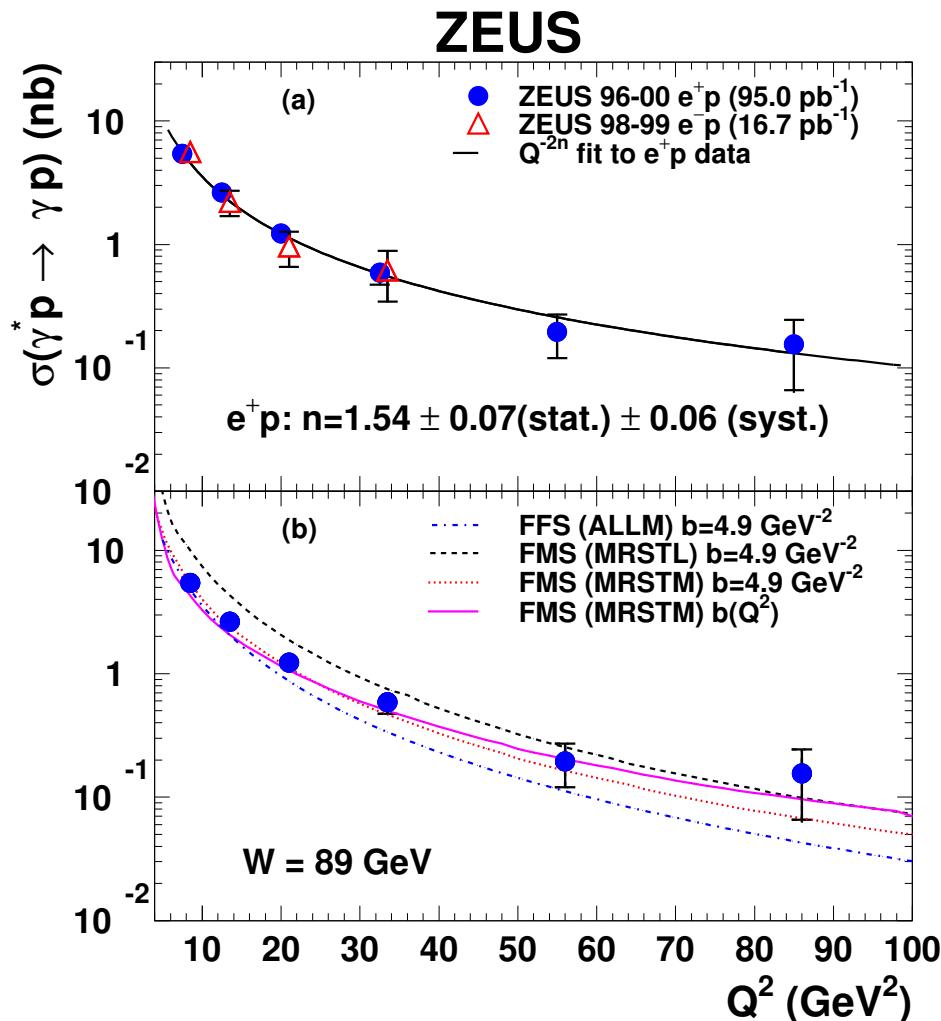


Data compatible with δ increasing with Q^2

steep rise of σ with $W \rightarrow$ DVCS is a hard process

Dependence similar as for
vector mesons

Q^2 dependence



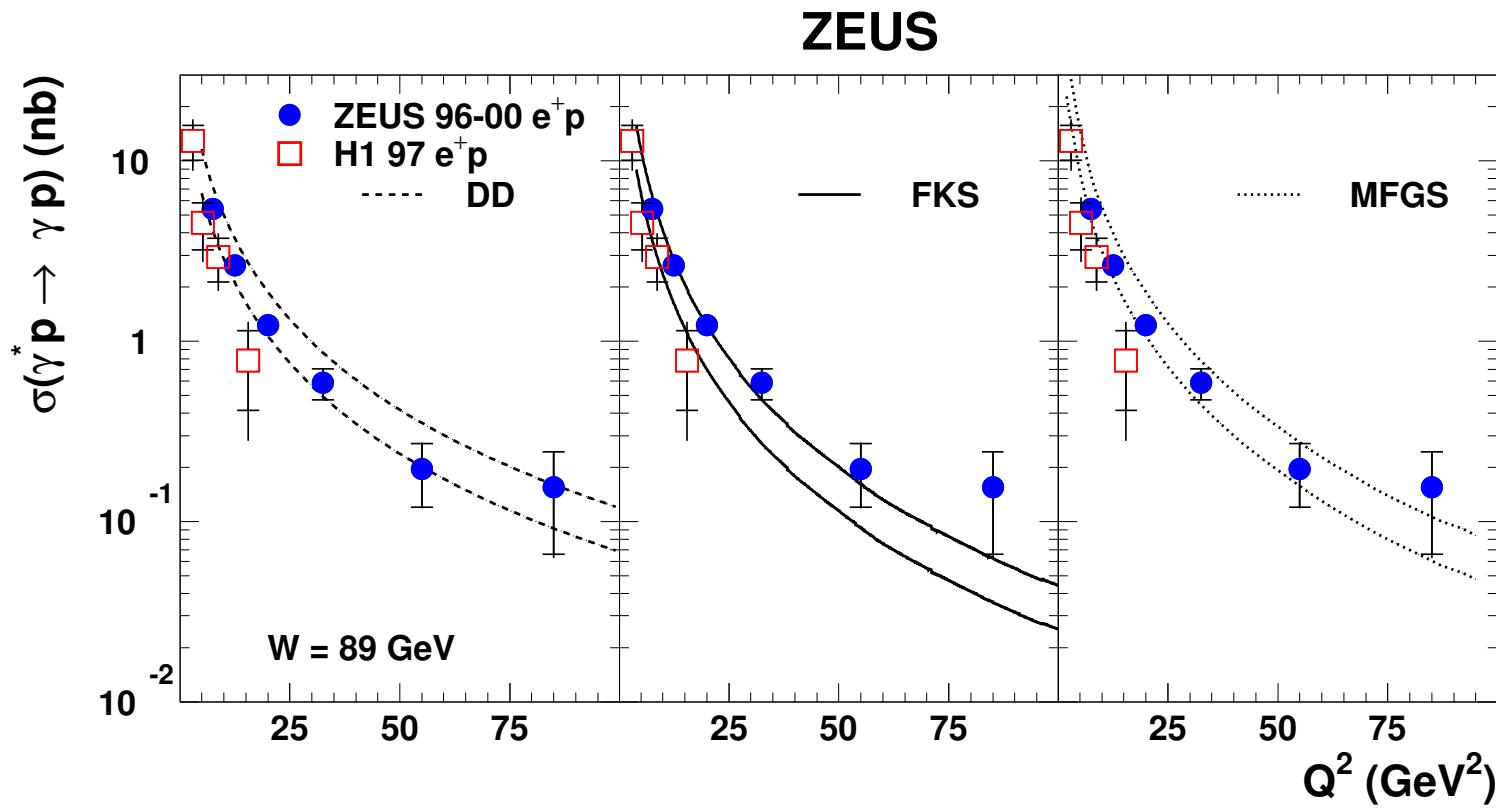
- $\sigma(\gamma^* p \rightarrow \gamma p) \sim Q^{-3}$

Freund,McDermott and Strikman
NLO calculation

$$b(Q^2) = 8[1 - 0.15\ln(Q^2/2)]\text{GeV}^{-2}$$

The best description
is given by NLO calculation

Comparison with Colour Dipole Models



upper curves $b = 4 \text{ GeV}^{-2}$

lower curves $b = 7 \text{ GeV}^{-2}$

DD (Donnachie and Dosh) - perturbative + Regge

FKS (Forshaw,Kerley and Shaw) - Regge approach

MFGS(McDermott,Frankfurt,Guzey and Strikman)- QCD colour transparency

Very different models
give similar predictions in the
measured kinematical region

Summary for DVCS

- ▶ $\sigma(\gamma^* p \rightarrow \gamma p)$ measured in the range $4 < Q^2 < 100 \text{ GeV}^2$ and $40 < W < 140 \text{ GeV}$
- ▶ No significant difference between $e^+ p$ and $e^- p$
- ▶ $\sigma(\gamma^* p \rightarrow \gamma p)$ rises steeply with $W \rightarrow$ hard process
- ▶ $\sigma(\gamma^* p \rightarrow \gamma p) \sim Q^{-3}$
- ▶ $\frac{d\sigma}{dt} = \frac{d\sigma}{dt}|_{t=0} e^{-b(Q^2)|t|}$
 Q^2 dependence of b-slope may be crucial in extracting GPDs, however
t dependence still not measured

Photon light-cone wave function - introduction

Light-Cone Wave Functions (LCWF)

$$\gamma(\vec{q}) \xrightarrow{\text{wavy line}} \begin{array}{c} l^+ \\[-1ex] \curvearrowright \\[-1ex] l^- \end{array} \rightarrow \psi_\gamma = a|\gamma_p\rangle + b|l^+l^-\rangle + (\text{other e.m.}) + c|q\bar{q}\rangle + (\text{other hadronic}) + \dots$$

$$\gamma(\vec{q}) \xrightarrow{\text{wavy line}} \begin{array}{c} q \\[-1ex] \curvearrowright \\[-1ex] \bar{q} \end{array}$$

- are the probability amplitudes to find a component with a given momentum in the momentum space
- are the solutions of LC Hamiltonian: $H_{LC}^{QCD} |\psi_\gamma\rangle = M_h^2 |\psi_\gamma\rangle$ ($H_{LC}^{QCD} = P^+P^- - P_\perp^2$)
- are usually tested through measurements of form factors
- are the best descriptions of the composite system

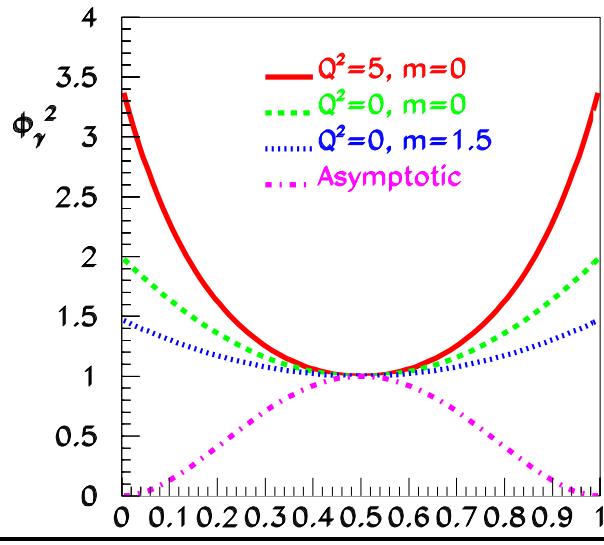
Photon light-cone wave function - QED and QCD

The LCWF for the lowest Fock states: $\psi_{\lambda_1 \lambda_2}^{\gamma} = -ee_l \frac{\bar{l}_{\lambda_1}(k) \lambda \cdot \epsilon^{\gamma} l_{\lambda_2}(q-k)}{\sqrt{u(1-u)}(Q^2 + \frac{k_{\perp}^2 + m^2}{u(1-u)})} \gamma(q)$

- longitudinal light-cone momentum fraction: $u = \frac{k_i^+}{p^+} = \frac{k_i^0 + k_i^z}{p^0 + p^z}$ $\sum_{i=1}^n u_i = 1$
- transverse momenta: $\vec{k}_{\perp i} \quad \sum_{i=1}^n \vec{k}_{\perp i} = \vec{0}_{\perp}$

The **hadronic** $|q\bar{q}\rangle$ LCWF is expected by pQCD to be the same as for **electromagnetic** $|l^+l^-\rangle$ for $k_{\perp}^2 \gg \Lambda_{QCD}^2$:

S.J.Brodsky, L.Frankfurt, J.F.Gunion, A.H.Mueller and M.Strikman, Phys.Rev.**D50**,3134(1994)



- for the transversely polarised photons:

$$\Phi_{f\bar{f}/\gamma_T^*}^2 \sim \sum_{\mu=1}^2 \frac{1}{4} Tr \psi^2 = \frac{m^2 + k_{\perp}^2 [u^2 + (1-u)^2]}{[k_{\perp}^2 + a^2]^2}$$

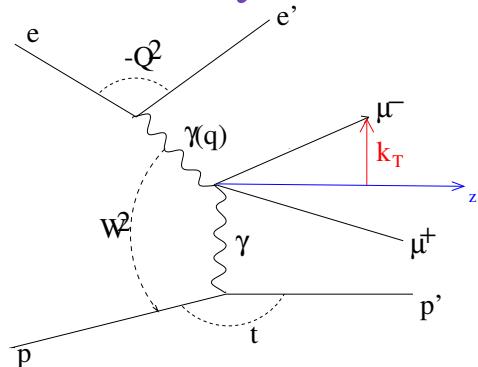
$$a^2 = m^2 + Q^2 u (1 - u)$$
- for the longitudinally polarised photons:

$$\Phi_{f\bar{f}/\gamma_L^*}^2 \sim \frac{Q^2 [u^2 (1-u)^2]}{[k_{\perp}^2 + a^2]^2}$$

LCWF - measurement of QED component

ZEUS e^+p sample 55.4 pb^{-1}

$30 < W < 170 \text{ GeV}, Q^2 \sim 0$



$$\gamma\gamma \rightarrow \mu^+\mu^-$$

- ➡ Proton and electron undetected
- ➡ Only 2 μ in detector
- ➡ Diffractive (small t)
- ➡ $4 < M_{\mu\mu} < 15 \text{ GeV}$

$$* u = \frac{E_1 + p_{z'_1}}{E_1 + E_2 + p_{z'_1} + p_{z'_2}} \quad 0 < u < 1$$

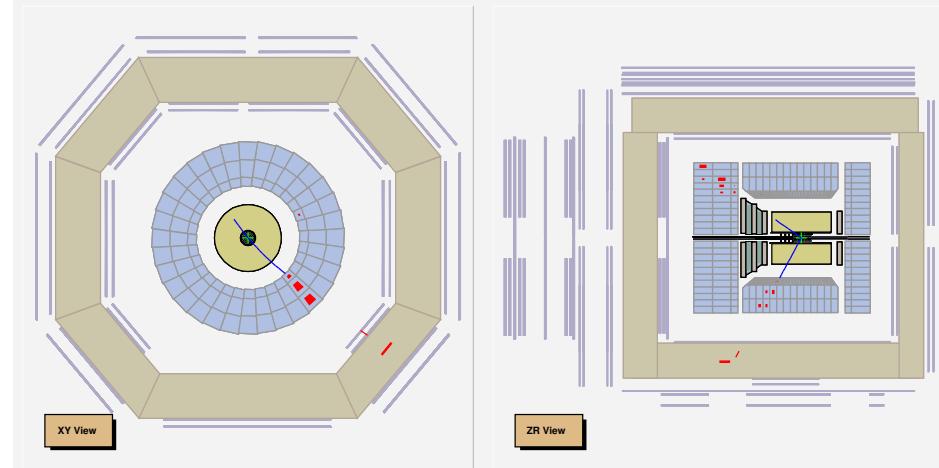
$$* \vec{k}_{T1} + \vec{k}_{T2} = \vec{0}$$

$$* W^2 = (q + P)^2 \approx 2E_p \sum_i (E_i - p_{zi})$$

$$* t = (P - P')^2 \approx -p_{T\mu\mu}^2$$

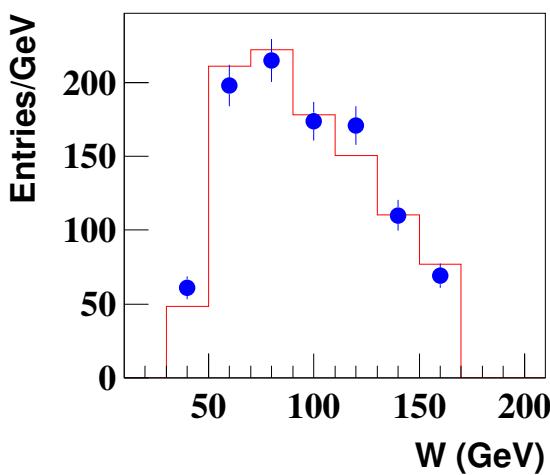
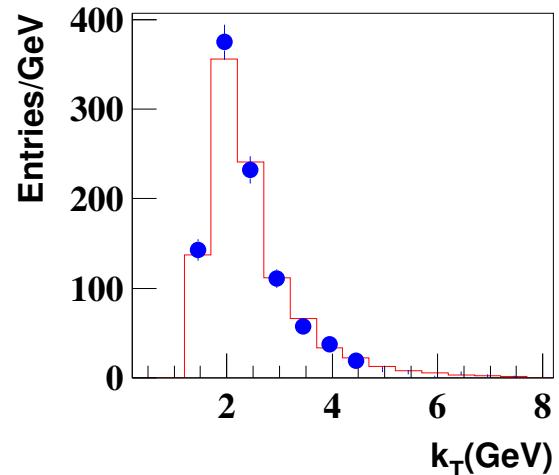
$$* M_{\mu\mu}^2 = (p_1 + p_2)^2$$

Zeus Run 31846 Event 1246					date: 19-01-1999 time: 21:36:15	
$E_e = 5.81 \text{ GeV}$	$E_\gamma = 3.81 \text{ GeV}$	$E_t = 1.61 \text{ GeV}$	$E_b = 3.51 \text{ GeV}$	$E_o = 2.29 \text{ GeV}$		
$E_\nu = 0.00 \text{ GeV}$	$p_\nu = 0.46 \text{ GeV}$	$p_t = 0.36 \text{ GeV}$	$p_b = 0.28 \text{ GeV}$	$p_o = 4.20 \text{ GeV}$		
$\phi_e = 0.66$	$t_e = -2.60 \text{ ns}$	$t_t = -1.92 \text{ ns}$	$t_b = -100.00 \text{ ns}$	$t_o = -2.40 \text{ ns}$		

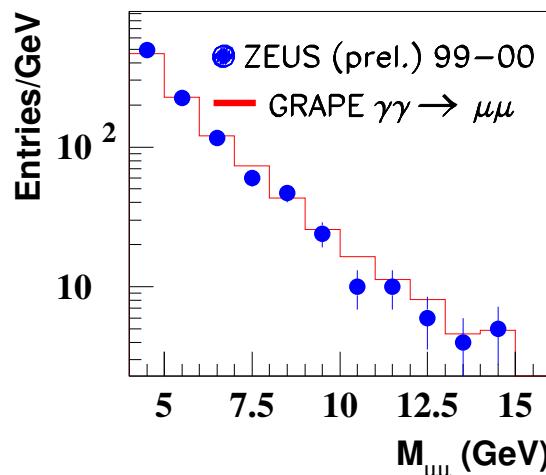
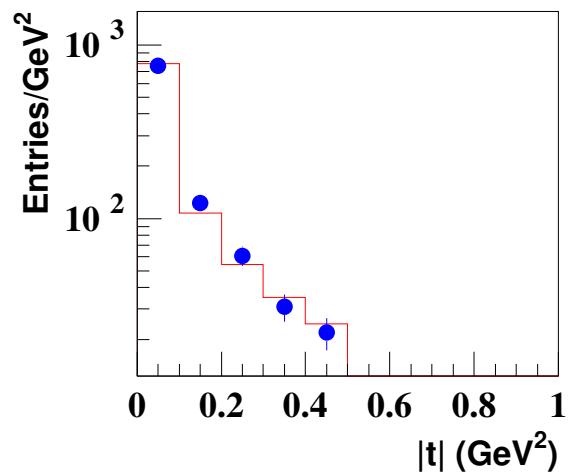


Kinematical variable distributions for the LCWF

ZEUS

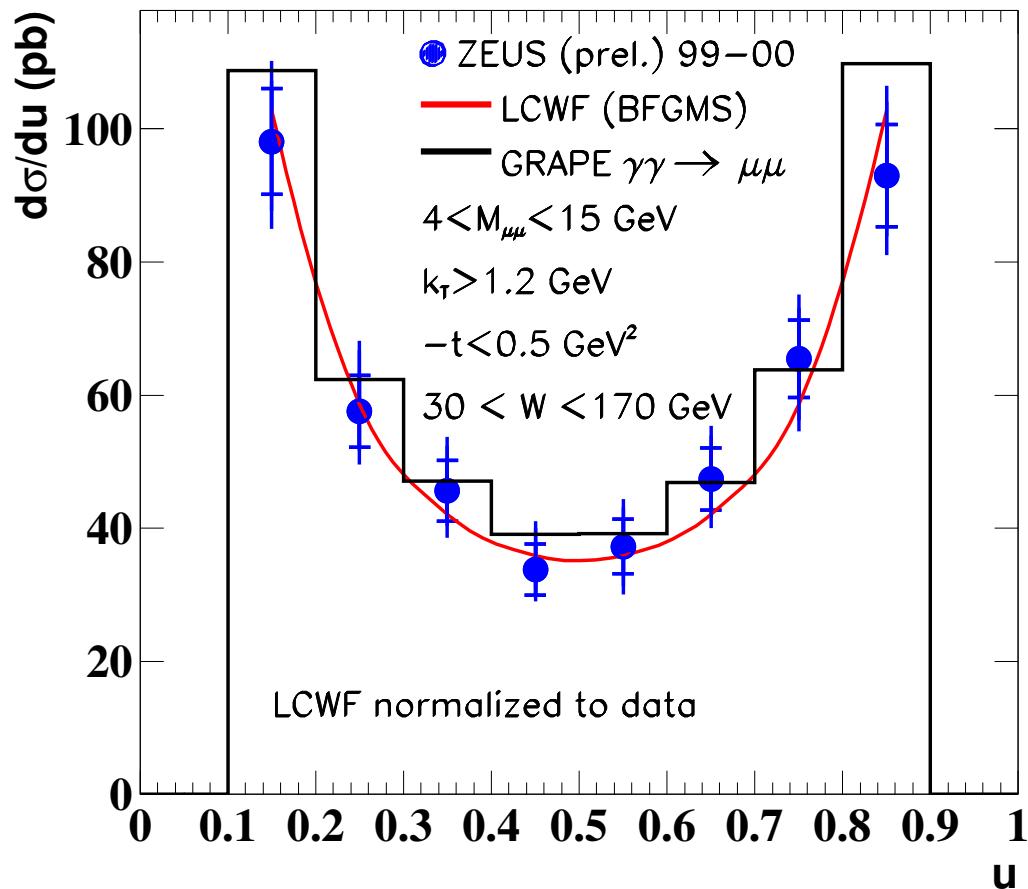


DATA and MC
in good agreement



LCWF - measurement of QED component

ZEUS



Calculations: Brodsky, Frankfurt, Gu-
nion, Müller, Strikman (BFGMS)

- Electromagnetic LCWF of the photon is measured
- BFGMS is in agreement with QED and data
- Results demonstrate the first proof that diffractive dissociation of particles can be reliably used to measure their LCWF

Summary for Light Cone Wave Function

- ▶ Photon electromagnetic LCWF is measured and is in agreement with QED
- ▶ This demonstrate the first proof that diffractive dissociation of particles can be reliably used to measure their LCWF
- ▶ Measured EM LCWF gives support for the method used in previous measurements of the pion LCWF and possible future applications