Low x Physics at HERA

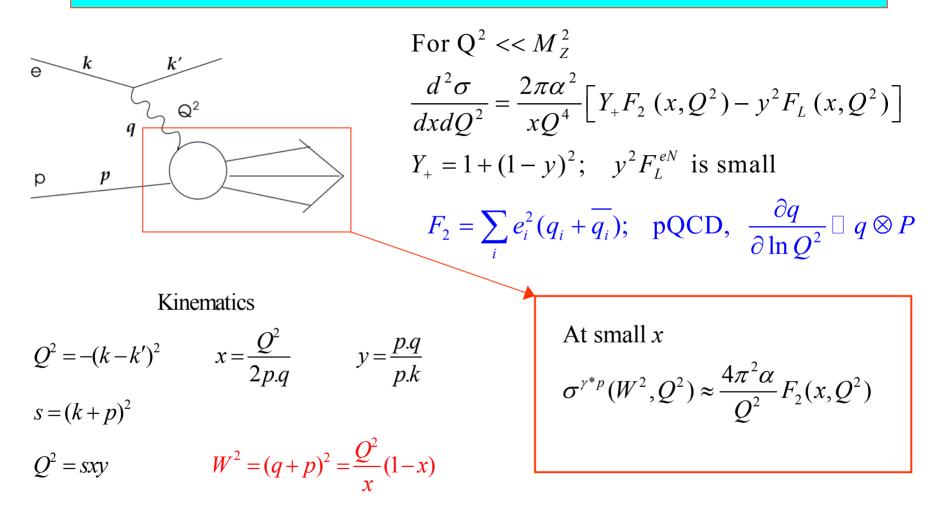
Robin Devenish (Oxford) for H1 and ZEUS

Outline

x? - is deep inelastic scattering Bjorken x

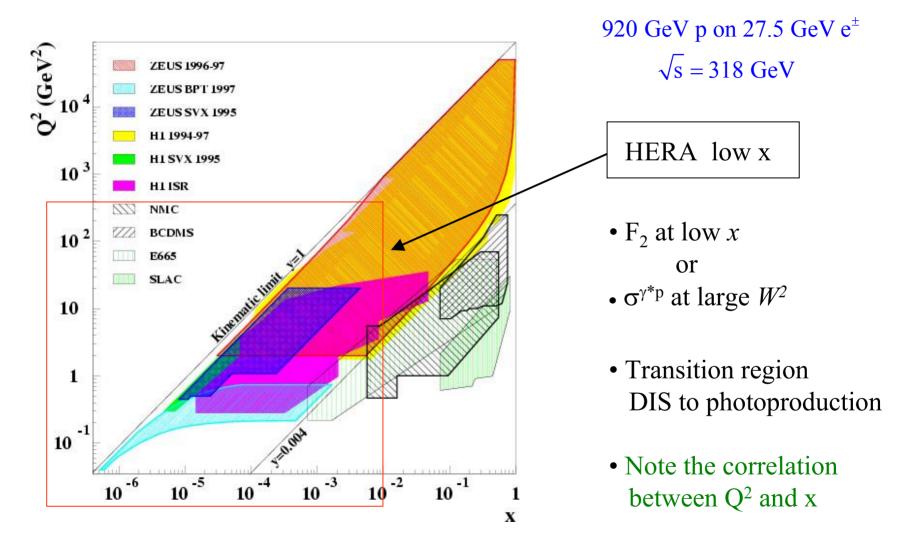
- Formalism and phase space
- F_2 at low x
- Contexts and Pictures
- More details on F_2 at low x
- F_2 at very low Q^2 & transition to photoproduction
- Universality at low *x*?
- Diffractive Processes
- Proton rest frame & dipole models
- Summary

Inelastic Scattering Formalism



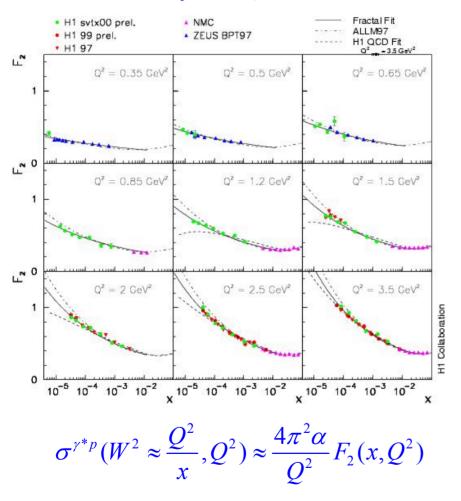
inelasticity, y = (E - E')/E in a fixed target frame

HERA & DIS kinematic regions



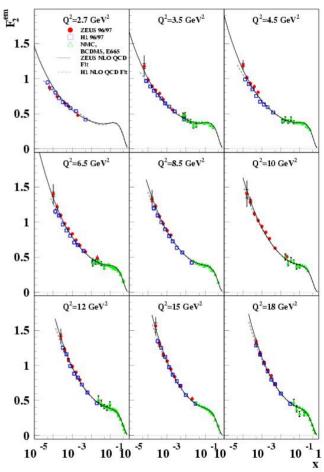
The rise of F_2 at low x

very low Q²



medium Q²

ZEUS+H1



 $\sigma(\gamma^*p)$ rises more rapidly with W² as Q² increases

Contexts

$\log x \leftrightarrow \log W^2$

low x

- pQCD
- DGLAP NLO sufficient?
- ln(1/x) summations
- BFKL
 - -perturbative Pomeron
- gluon dominance
- universality?
- vary 'size' of γ^*
- saturation

Hadron-hadron

- optical theorem
 - $\sigma_{tot} \sim Im {<} \gamma {*} p |A| \gamma {*} p {>} / w^2$
- forward elastic amplitude
- Regge theory

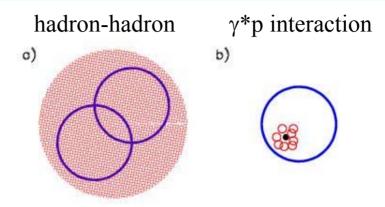
 $\sigma_{tot} \sim const.(w^2)^{\alpha(0)-1}$ Pomeron $\alpha(0)-1 = 0.08$

- quasi-elastic scattering
- diffraction
- unitarity limit

Colour Dipole models

Interaction pictures

Transverse views of particle interactions

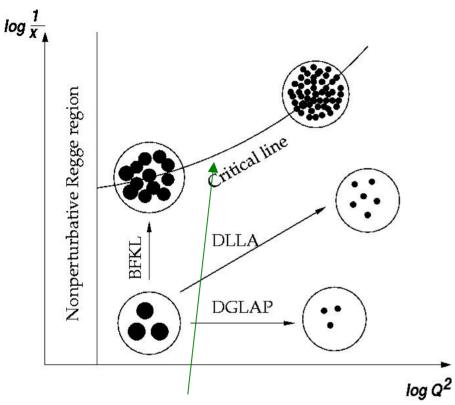


Bartels & Kowalski

- a) Diffuse gluon 'radiation cloud' drives the interaction and size of interaction region, which is larger than the hadrons, grows slowly with energy
- b) γ^* with small transverse dimensions, d, interacting with a proton also with a radiation cloud but more intense because of limited size calculable using pQCD

At HERA the size of the photon can be varied from that of hadron (photoproduction) to much smaller, since $d \sim 1/Q$

High gluon density and saturation



• Critical line indicates region above which saturation will occur – gluons overlap – nonlinear evolution Gluon dynamics dominates but how rapidly does F_2 increase?

• DGLAP dominated by gluon splitting function $P_{gg} \sim 1/x$

• DLLA

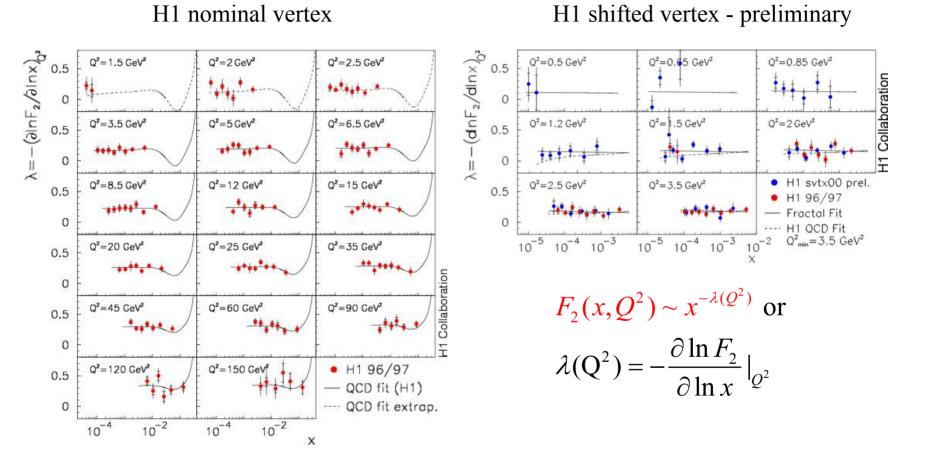
 $\exp\sqrt{\left(12\alpha_s/\pi\right)\ln 1/x\ln Q^2/Q_0^2}$

• BFKL

~ $x^{-\lambda}$ with λ as large as 0.5

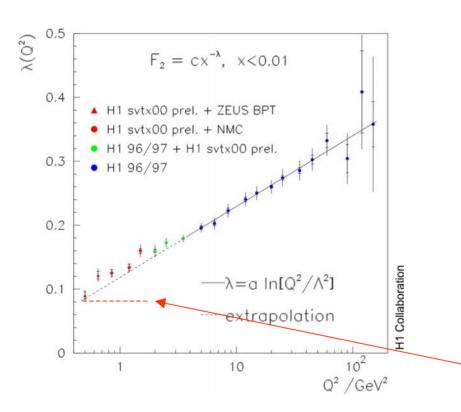
- Other summations:
 - CCFM (angular ordering)
 - those from Thorne; Altarelli et al
- on the edge or just outside the reach of HERA?

Model independent study of F₂ at low x



- characterise the rise of F_2 – taking full account of errors - for Q^2 fixed and x < 0.01, λ roughly constant

$\lambda(Q^2)$ vs Q^2

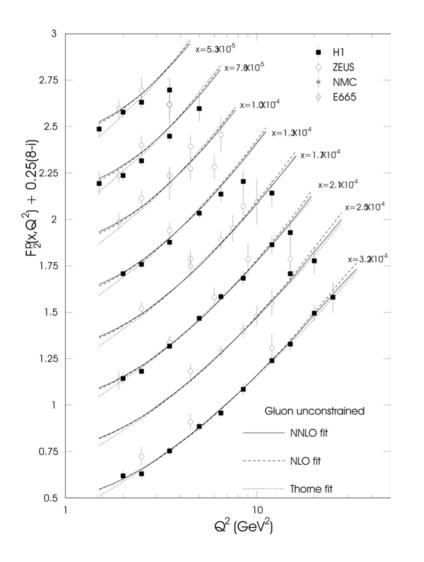


Fit $\lambda(Q^2) = a \ln(Q^2 / \Lambda^2)$ gives $\Lambda = 292 \pm 20(\text{stat}) \pm 51(\text{sys})$ $F_2(x,Q^2) \approx c(Q^2) x^{-\lambda(Q^2)}$ with $c(Q^2) \approx 0.18$ indep. of Q^2

• no sign of rise slowing at large Q^2 and small x as might be expected from saturation

• at very small Q^2 the value of λ is consistent with that expected from hadron-hadron scattering: $\lambda \sim 0.08$

Beyond standard NLO evolution?



From the MRST team

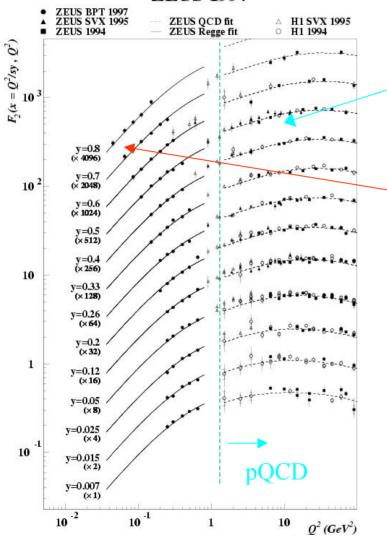
 F_2 at low *x* fit using DGLAP NLO, NNLO and some resummation of ln(1/x)terms (Thorne fit)

All give acceptable fits parton densities are different but need other observables (eg F_L) to distinguish.

Standard NLO DGLAP fits describe data at low x down to $Q^2 \sim 1.5 \text{ GeV}^2$ - more in Milstead's talk

F₂ as **Q**² tends to zero

ZEUS 1997



- NLO pQCD describes F_2 down to $Q^2 \sim 1.5$ GeV²
- At very small Q², EM current conservation requires

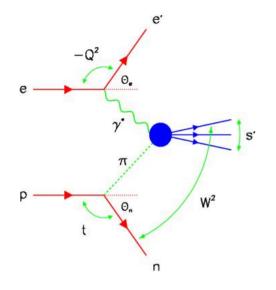
 $F_2(x,Q^2) \rightarrow Q^2 \times const.$ as $Q^2 \rightarrow 0$

- Data shows a smooth transition in Q²
- Many models describe the transition region
- Regge based approaches
- Generalised Vector Dominance
- Colour dipoles (more later)
- Self-similarity

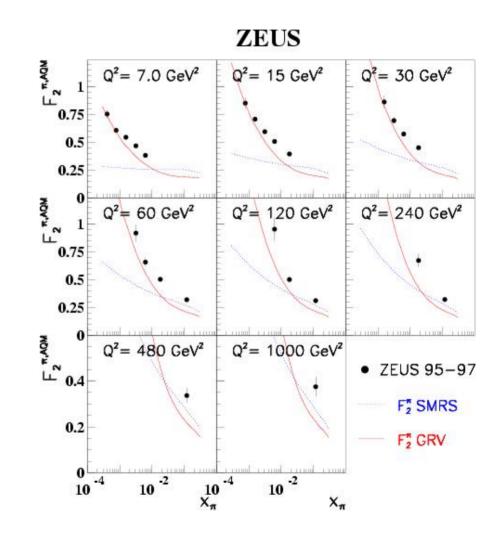
Universality at low x

- At small *x* the dynamics of Q^2 evolution is dominated by gluon splitting
- Far from the valence region in *x*, the identity of the parent particle becomes unimportant
- Recently both H1 and ZEUS have published measurements of deep inelastic scattering with an identified forward neutron in the final state
- Is there evidence of similar rapid growth at small *x* in other structure functions?
- At small p-n momentum transfers ('t'), single pion exchange dominates and the pion structure function can be isolated

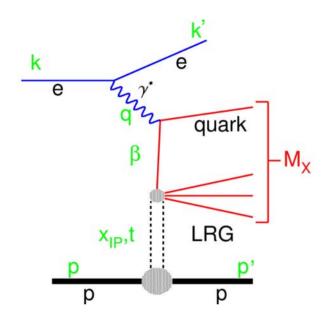
\mathbf{F}^{π} at low x



Although there is uncertainty in the normalisation, there is no doubt that F^{π} is rising steeply at low x



Diffraction at HERA



Identify diffractive events either using leading proton spectrometer or LRG in main detector

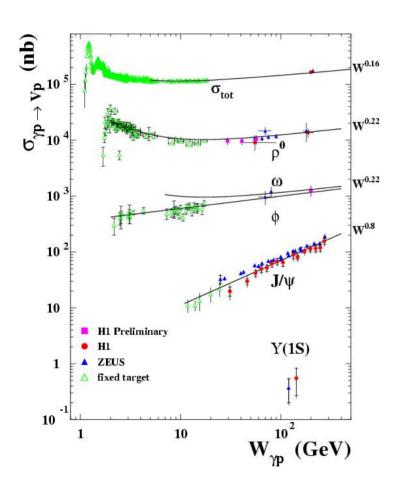
$$\gamma^* p \rightarrow X p$$

- involves vacuum q. no. exchange
- X is a vector meson or a hadronic system separated from the proton by a large rapidity gap (LRG)
- At HERA
 - vector mesons vs Q^2
 - inclusive diffraction vs $M_X \& Q^2$
 - hard diffraction and jets

 W^2 dependence of all of these

Concentrate on the W^2 dependence of diffractive cross-sections

Vector meson photoproduction



 $\sigma(\gamma p \rightarrow V p)$ vs W for real photoproduction of vector mesons compared with $\sigma_{tot}(\gamma p)$

lines indicate a power law fit $\sigma \sim W^{\delta}$

to data with W > 10 GeV

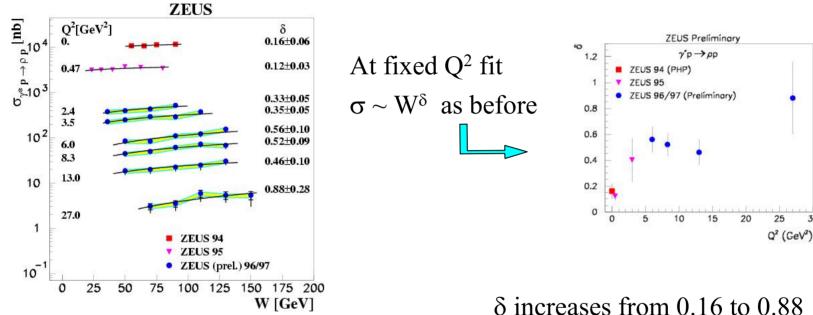
for ρ , ω , $\phi \quad \delta \sim 0.22$ comparable to $\delta \sim 0.16$ for $\sigma_{tot}(\gamma p)$

for J/ ψ $\delta \sim 0.8$

Faster rise if hard scale is set by large M_V

ρ^0 production vs Q²

 $\sigma(\gamma^* p \rightarrow \rho^0 p)$ measured as a function of W and Q^2



NB use $\lambda \sim \delta/4$ to compare rates of rise in diffraction $(|A|^2)$ and $\sigma_{tot}(\gamma * p)$ (ImA)

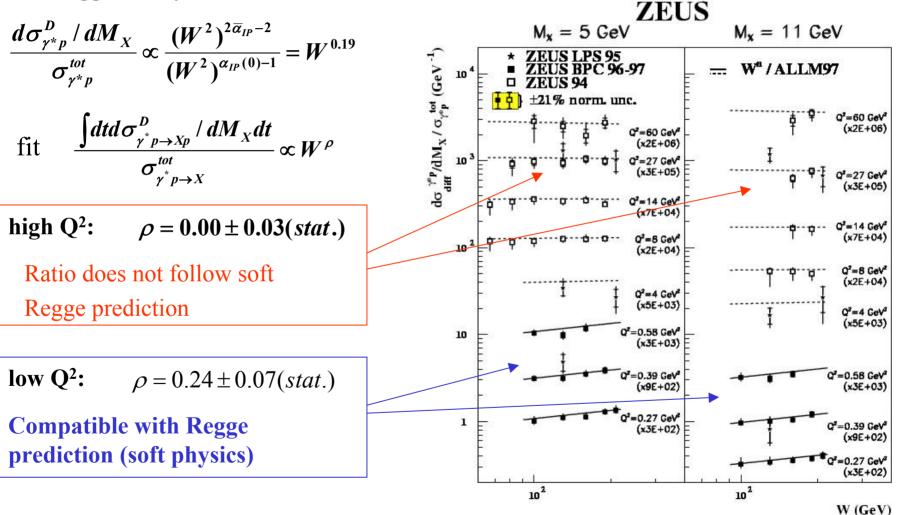
 δ increases from 0.16 to 0.88 as Q^2 increases from 0 to 27 GeV²

faster rise as γ^* provides the hard scale

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Inclusive diffraction/total vs W

Regge theory:



Putting it all together

Results so far:

• low *x* rise of $F_2 \rightarrow \sigma_{tot}(\gamma^* p)$ increases faster with W² than expected from hadron-hadron scattering

• similar features seen in other quasi-elastic processes where there is a hard scale

• dominance of gluon dynamics & some evidence for universality at low x

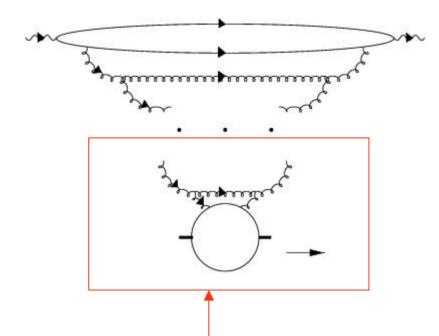
• smooth transition from DIS to photoproduction at $Q^2 = 0$ many models to describe this

• the hard scale is associated with small transverse size of the probe (γ^*) or final state particle (vector meson)

• is there a framework in which all this can be put together?

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Proton rest frame view



proton rest frame view of the $\gamma^* p$ interaction

 $\gamma^* \rightarrow q\bar{q}$ followed by multiple gluon radiation

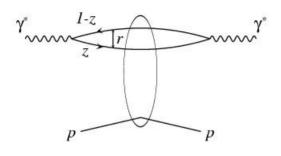
The q \bar{q} dipole forms a long time t ~ 1/(m_px) before interacting with the proton

As the cascade grows the energy and $k_{\rm T}$ of the gluons decreases – their transverse size grows – no longer calculable perturbatively

colour dipole models

Nikolaev & Zakharov; Golec-Biernat & Wuesthoff; Mueller; Gotsman, Levin & Maor; Buchmuller & Hebecker..; Forshaw, Shaw, McDermott...

Colour Dipole Models with saturation



r transverse separation, conjugate to k_T

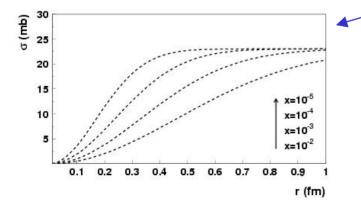
z longitudinal photon momentum fraction

known wave function

$$\sigma_{\gamma^*P} = \int d^2 r \, dz \, \Psi_{\gamma^*}^*(r, z, Q^2) \, \sigma_{qq}(x, r) \, \Psi_{\gamma^*}(r, z, Q^2)$$

model unknown dipole cross-section σ_{qq} (e.g. Golec-Biernat & Wuesthoff)

$$\sigma_{qq}(x,r) = \sigma_0 \left\{ 1 - \exp(-r^2/4R_0^2(x)) \right\}; \quad R_0(x) = (x/x_0)^{\lambda/2}$$

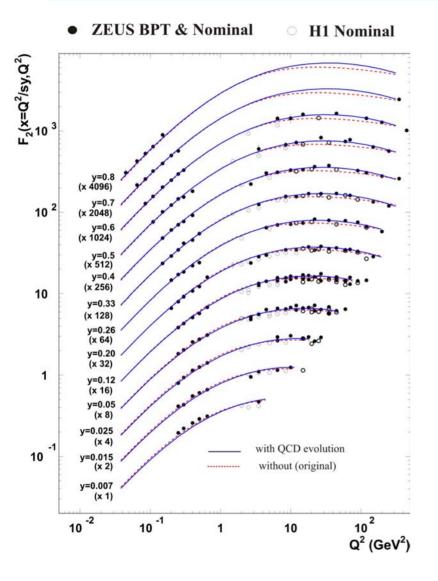


$$r \ll R_0$$
 $\sigma_{qq} \propto r^2 x^{-\lambda}$

$$r >> R_0 \qquad \sigma_{qq} \propto \sigma_0$$

unitarity bound built in & approach controlled by $R_0(x)$

Colour Dipole Model fitted to inclusive data



$$\sigma_{qq}(x,r) = \sigma_0 \left\{ 1 - \exp(-r^2 / 4R_0^2(x)) \right\}$$
$$R_0(x) = (x / x_0)^{\lambda/2}$$

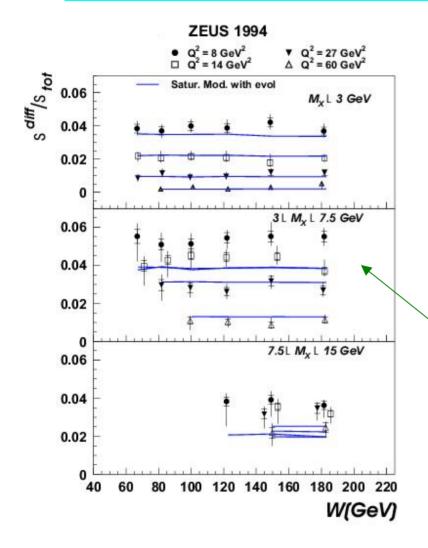
Three parameters: λ , x_0 , σ_0 to be fitted.

Extension (G-B, Bartels &Kowlaski - DIS02) include QCD evolution by requiring

 $\sigma_{qq}(r,x) = \frac{\pi^2}{3} r^2 \alpha_s xg(x,C/r^2), \text{ at small } r$ xg evolves - 5 parameter fit - improves description of high Q^2 data

• Cannot use this agreement as verifying saturation at HERA, as many other models give similar agreement, including non-saturating dipole models.

Hard diffraction in the colour dipole model



• Dipole models provide a natural framework for hard diffraction – with the same σ_{qq} and parameters as determined from inclusive data

$$\frac{d\sigma_{\gamma^{*\mathrm{P}}}^{\mathrm{Diffr}}}{dt}\Big|_{t=0} \propto \left|\Psi_{\gamma^{*}}(Q^{2},r)\right|^{2} \otimes \sigma_{\mathrm{qq}}^{2}(x,r)$$

• While not perfect, the dipole model does give a constant ratio for $\sigma^{\text{Diffr}}/\sigma^{\text{All}}$ in DIS – a key feature of the data that the soft Regge theory does not reproduce

 \bullet Also describes behaviour of δ for

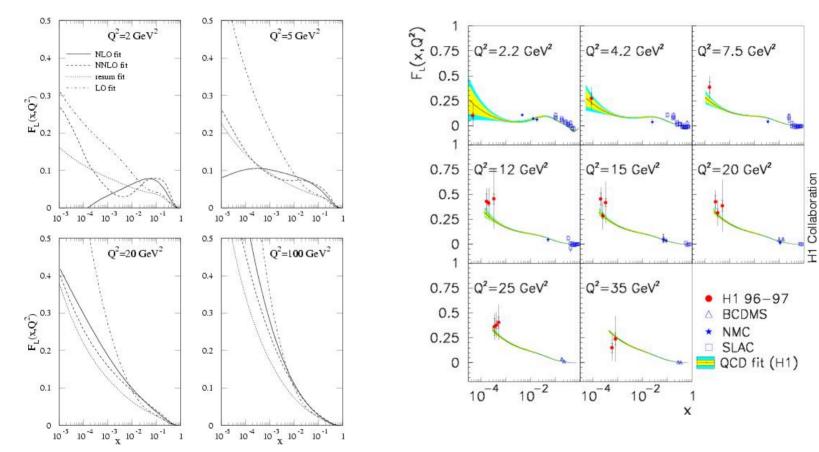
 $\sigma(\gamma * p \to \rho^0 p) = const \times W^{\delta(Q^2)}$

Summary

- HERA has provided high precision data on F_2 at low x and hard diffractive scattering
- Rise of F_2 at low x mirrored in other processes when appropriate hard scale is present
- Aspects of universality in the low *x* dynamics hinted at
- Colour dipole models are promising but saturation not proven at HERA
- HERA has opened up new avenues in strong interaction physics
 - high density perturbative gluon dynamics
 - deepened the relationship between diffractive scattering and the physics driving rising total cross-sections
- Essential input for physics at the Tevatron, RHIC & the LHC

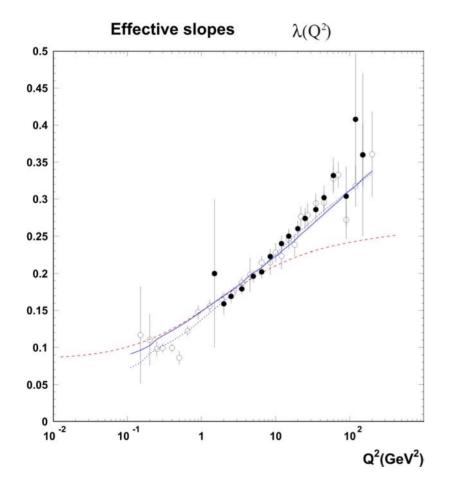
Thanks to many colleagues on H1, ZEUS and in the 'HERA low x club', for real and virtual help in preparing this talk.

F_L – models and data



F_L LO , NLO and NNLO

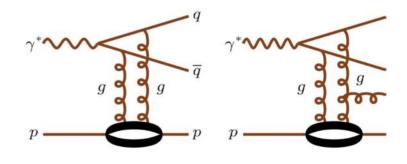
Dipole model description of $\lambda(Q^2)$



ZEUS & H1 data

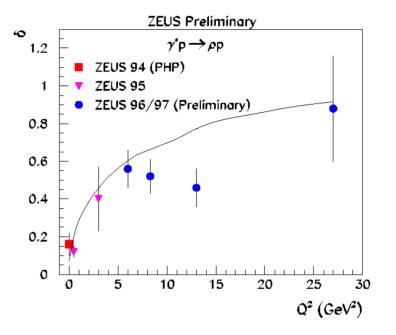
 $F_2(x,Q^2) \sim x^{-\lambda(Q^2)}$

Clearly need the inclusion of QCD evolution to get a reasonable description of λ above medium Q² values



Dipole models for Vector Mesons

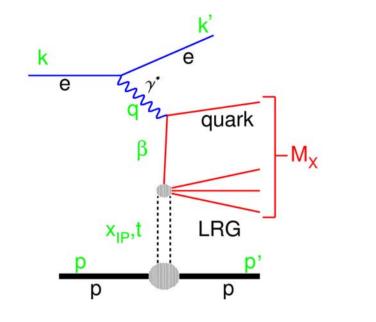
This is being studied by a number of groups – the results are encouraging. References: Munier hep-ph/0206117 Caldwell & Soares hep-ph/0101085 Forshaw, Kerley & Shaw Phys Rev D60 074012 (1999)



Plot shows the calculation from Caldwell & Soares of $\delta(Q^2)$, where

$$\sigma(\gamma * p \to \rho^0 p) = const \times W^{\delta(Q^2)}$$

Variables in diffractive DIS



Identify diffractive events either using leading proton spectrometer or LRG in main detector. If the forward proton is not detected then there will be up to 20%

background from proton dissociation.

DIS

$$\begin{vmatrix} x = Q^2/2p \cdot q \\ W^2 = (q+p)^2 \\ M_X^2 = (q+P-P')^2 \end{vmatrix}$$

 $Q^2 = -q^2$

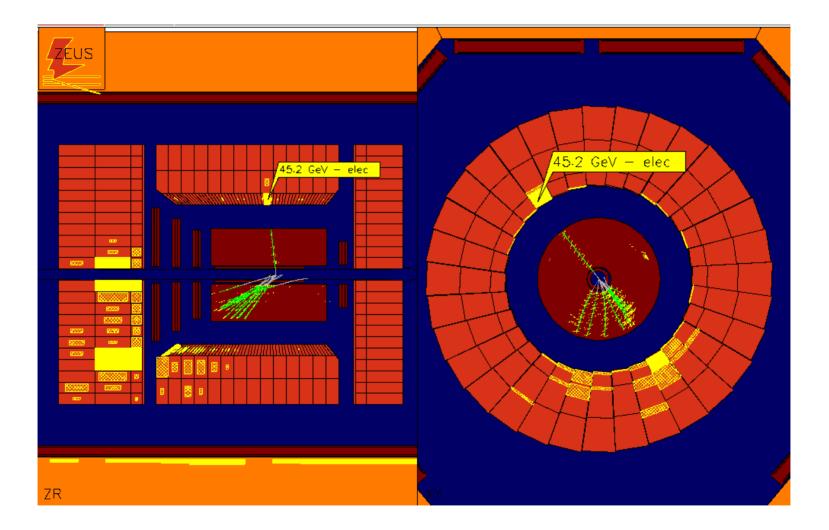
$$t = (p - p')^2 \approx -\frac{p_t^2}{x_L}$$

$$x_{L} = \frac{p'_{z}}{p_{z}} \approx 1 - x_{IP}$$

$$x_{IP} = \frac{q \cdot (p - p')}{q \cdot p} \approx \frac{Q^{2} + M_{X}^{2}}{Q^{2} + W^{2}}$$

$$\beta = \frac{Q^{2}}{2q \cdot (p - p')} \approx \frac{Q^{2}}{Q^{2} + M_{X}^{2}}$$

Standard NC event in ZEUS



Diffractive DIS event in ZEUS

